# Charmonium from $B$ decays 

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Workshop: New results on Charmonium production and decays
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## Charmonium production at hadron machines

- The first measurement of direct J/ $\Psi$ and $\Psi^{\prime}$ production at CDF in '97: striking discrepancy from theoretical expectation

CDF, PRL79 ('97)

- NRQCD: double expansion in terms of $\alpha_{s}$ and $v$ (velocity): an addition of "colour-octet" term was proposed.
- Tremendous efforts have been made to obtain more precise theoretical prediction for the charmonium production at hadron machines (computation of the higher order corrections, extracting the matrix element of NRQCD, Color-singlet approach etc). $\Rightarrow$ after many debates, still the situation is unclear!

More investigation is needed!

## New observables to help the situation ???

So far, the study has been limited to $J / \Psi, \psi^{\prime}, X_{J c}$

- Prompt/secondary production of charmonium states such as $\eta_{c}$ or $h_{c}$ have never been done at hadron machines although they could be useful for clarify some issues (using spin symmetry)!

Charmonium from $B$ decays to extract the matrix elements?

- Secondary charmonium production is experimentally cleaner than the prompt production. Theoretically, it is less clean (e.g. issues in the NLO estimate of the singlet contribution) but can't we still learn something?
<--- Universality of matrix elements!


# Revisiting the NRQCD computation of inclusive $B$ decaying into charmonium 



Beneke, Maltoni,
Rothotein
PRD59 ('99)

$$
\begin{gathered}
H_{e f f}=\frac{G_{F}}{\sqrt{2}} \sum_{q=s, d}\left\{V_{c b}^{*} V_{c q}\left[\frac{1}{3} C_{[1]}(\mu) \mathcal{O}_{1}(\mu)+C_{[8]}(\mu) \mathcal{O}_{8}(\mu)\right]-V_{t b}^{*} V_{t q} \sum_{i=3}^{6} C_{i}(\mu) \mathcal{O}_{i}(\mu)\right\} \\
\mathcal{O}_{1}=\left[\bar{c} \gamma_{\mu}\left(1-\gamma_{5}\right) c\right]\left[\bar{b} \gamma^{\mu}\left(1-\gamma_{5}\right) q\right] \\
\mathcal{O}_{8}=\left[\bar{c} T^{A} \gamma_{\mu}\left(1-\gamma_{5}\right) c\right]\left[\bar{b} T^{A} \gamma^{\mu}\left(1-\gamma_{5}\right) q\right]
\end{gathered}
$$

It has been pointed out by several authors that the singlet term is too small to explain the experimental data.

## Revisiting the NRQCD computation of inclusive $B$ decaying into charmonium

NLO computation

$$
\Gamma[n]=\Gamma_{0}\left[C_{[1,8]}^{2} f[n](\eta)\left(1+\delta_{P}[n]\right)\right.
$$

$$
\left.+\frac{\alpha_{s}(\mu)}{4 \pi}\left(C_{[1]}^{2} g_{1}[n](\eta)+2 C_{[1]} C_{[8]} g_{2}[n](\eta)+C_{[8]}^{2} g_{3}[n](\eta)\right)\right]\left\langle\mathcal{O}^{H}[n]\right\rangle
$$



The singlet term has a large renormalization running effect which makes it large negative (unphysical) at mb scale.

# Revisiting the NRQCD computation of inclusive $B$ decaying into charmonium 

## Improved NLO result



## Revisiting the NRQCD computation of inclusive $B$ decaying into charmonium

$$
\begin{aligned}
& \operatorname{Br}(B \rightarrow J / \psi X)=\quad \text { exp. }(1.094 \pm 0.032) \times 10^{-2} \\
& \quad 0.0754 \times 10^{-2}\left\langle\mathcal{O}_{1}^{\psi}\left({ }^{3} S_{1}\right)\right\rangle+0.195\left\langle\mathcal{O}_{8}^{\psi}\left({ }^{3} S_{1}\right)\right\rangle+0.342 \underbrace{\left[\left\langle\mathcal{O}_{8}^{\psi}\left({ }^{1} S_{0}\right)\right\rangle+3.1 / m_{c}^{2}\left\langle\mathcal{O}_{8}^{\psi}\left({ }^{3} P_{0}\right)\right\rangle\right]} \\
& \operatorname{Br}\left(B \rightarrow \eta_{c} X\right)=\text { exp. }\left\langle 0.9 \times 10^{-2}<- \text { update from LHCb? } \mathcal{M}_{1, k}^{\psi}{ }^{1} S_{0}^{(8)},{ }^{3} P_{0}^{(8)}\right)
\end{aligned} \quad \underbrace{0.250 \times 10^{-2}\left\langle\mathcal{O}_{1}^{\eta_{c}}\left({ }^{1} S_{0}\right)\right\rangle+0.342\left\langle\mathcal{O}_{8}^{\eta_{c}}\left({ }^{1} S_{0}\right)\right\rangle+0.195 \underbrace{\left[\left\langle\mathcal{O}_{8}^{\eta_{c}}\left({ }^{3} S_{1}\right)\right\rangle-0.24 / m_{c}^{2}\left\langle\mathcal{O}_{8}^{\eta_{c}}\left({ }^{1} P_{1}\right)\right\rangle\right]}_{8}}_{\left.\mathcal{M}_{1,0.24}^{\eta_{c}}{ }^{3} S_{1}^{(8)},{ }^{1} P_{1}^{(8)}\right)} \begin{aligned}
& \text { Spin symmetry }
\end{aligned}
$$

Can we use this result to extract some information on the octet matrix elements?

## Revisiting the NRQCD computation of inclusive $B$ decaying into charmonium

Momentum dependence of $B->J / p s i X$


FIG. 10: $p^{*}$ of $J / \psi$ mesons produced directly in $B$ decays (points). The histogram is the sum of the color-octet component from a recent NRQCD calculation [20] (dashed line) and the color-singlet $J / \psi K^{(*)}$ component from simulation (dotted line).

Beneke, Scgykerm Wolf, PRD62 ('00)

# Prediction of $\mathrm{Br}(\mathrm{B}->\mathrm{J} / \mathrm{psi} X)$ with the fitted matrix elements 

|  | $\left\langle\mathcal{O}_{1}^{\psi}\left({ }^{3} S_{1}\right)\right\rangle$ | $\left\langle\mathcal{O}_{8}^{\psi}\left({ }^{3} S_{1}\right)\right\rangle$ | $\left\langle\mathcal{O}_{8}^{\psi}\left({ }^{1} S_{0}\right)\right\rangle$ | $\left\langle\mathcal{O}_{8}^{\psi}\left({ }^{3} P_{0}\right)\right\rangle$ | $\mathcal{M}_{0, k}^{\psi}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Beneke et al | 1.16 | 1.06 |  |  | $2.0-2.7 \times 10^{-2}$ |
| Ma et al | 1.16 | small? | $\sim^{-2} \mathrm{Mok}^{\mathrm{psi}}$ | small? | $7.4 \times 10^{-2}$ |
| Chao et al | 1.16 | $0.30 \pm 0.12$ | $8.9 \pm 0.98$ | $0.56 \pm 0.21$ | $9.8 \times 10^{-2}$ |
| Kniehl et al | 1.32 | $0.17 \pm 0.05$ | $3.04 \pm 0.35$ | $-0.91 \pm 0.16$ | $1.6 \times 10^{-2}$ |



# Prediction of $\mathrm{Br}(\mathrm{B}->$ etac $X)$ with the fitted matrix elements 

| Assuming Spin symm | etryo.. $\times 10^{-2}$ |  | $\times 10^{-2} \times 10^{-}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\left\langle\mathcal{O}_{1}^{\psi}\left({ }^{3} S_{1}\right)\right\rangle$ | $\left\langle\mathcal{O}_{8}^{\psi}\left({ }^{3} S_{1}\right)\right\rangle$ | $\left\langle\mathcal{O}_{8}^{\psi}\left({ }^{1} S_{0}\right)\right\rangle$ | $\left\langle\mathcal{O}_{8}^{\psi}\left({ }^{3} P_{0}\right)\right\rangle$ | $\mathcal{M}_{1, k}^{\psi}$ |
| Beneke et al | 0.39 | 0.35 |  |  |  |
| Ma et al | 0.39 | small? | $\sim M_{0 k}{ }^{\text {psi }} / 3$ | small??? | $2.5 \times 10^{-2}$ |
| Chao et al | 0.39 | $0.10 \pm 0.04$ | $3.0 \pm 0.32$ | $1.68 \pm 0.63$ | $2.8 \times 10^{-2}$ |
| Kniehl et al | 0.44 | $0.06 \pm 0.02$ | $1.01 \pm 0.12$ | $-2.73 \pm 0.48$ | (1.3×10-2 |
| Assuming, $08(150)=\begin{aligned} & 0.9 \times 10^{-2} \\ & 0.8 \times 10^{-2}\end{aligned}$ <br> Current upper limit |  |  |  |  |  |
| for psi and 08(3S1) for etac is small... | $\begin{array}{r} 0.7 \times 10^{-2} \\ 0.6 \times 10^{-2} \\ 0.5 \times 10^{-2} \end{array}$ |  |  |  | Exp. |
| Spin symmetry | $\begin{aligned} & 0.4 \times 10^{-2} \\ & 0.3 \times 10^{-2} \\ & 0.2 \times 10^{-2} \end{aligned}$ |  |  |  | O8(IS0) O8(3SI) |
| $\left.\left\langle\mathcal{O}_{8}^{\eta_{c}}{ }^{1} S_{0}\right)\right\rangle=\frac{1}{3}\left\langle\mathcal{O}_{8}^{J / \psi}\left({ }^{3} S_{1}\right)\right\rangle$, |  |  |  |  | 1 |

Caveat: Proportion of the uncertain singlet terms are large...

