### Charmonium from B decays

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### Charmonium production at hadron machines

- The first measurement of direct  $J/\psi$  and  $\psi'$  production at CDF in '97: striking discrepancy from theoretical expectation
- NRQCD: double expansion in terms of  $\alpha_s$  and v (velocity): an addition of "colour-octet" term was proposed. *Bodwin, Braaten, Lepage, PRD51 ('95)*
- Tremendous efforts have been made to obtain more precise theoretical prediction for the charmonium production at hadron machines (computation of the higher order corrections, extracting the matrix element of NRQCD, Color-singlet approach etc). 

   after many debates, still the situation is unclear!

### More investigation is needed!

### New observables to help the situation ???

So far, the study has been limited to  $J/\psi$ ,  $\psi'$ ,  $\chi_{Jc}$ 

 Prompt/secondary production of charmonium states such as η<sub>c</sub> or h<sub>c</sub> have never been done at hadron machines although they could be useful for clarify some issues (using spin symmetry)!

Charmonium from B decays to extract the matrix elements?

 Secondary charmonium production is experimentally cleaner than the prompt production. Theoretically, it is less clean (e.g. issues in the NLO estimate of the singlet contribution) but can't we still learn something?

<--- Universality of matrix elements!</pre>





$$H_{eff} = \frac{G_F}{\sqrt{2}} \sum_{q=s,d} \left\{ V_{cb}^* V_{cq} \left[ \frac{1}{3} C_{[1]}(\mu) \mathcal{O}_1(\mu) + C_{[8]}(\mu) \mathcal{O}_8(\mu) \right] - V_{tb}^* V_{tq} \sum_{i=3}^6 C_i(\mu) \mathcal{O}_i(\mu) \right\}$$
$$\mathcal{O}_1 = \left[ \bar{c} \gamma_\mu (1 - \gamma_5) c \right] \left[ \bar{b} \gamma^\mu (1 - \gamma_5) q \right]$$
$$\mathcal{O}_8 = \left[ \bar{c} T^A \gamma_\mu (1 - \gamma_5) c \right] \left[ \bar{b} T^A \gamma^\mu (1 - \gamma_5) q \right]$$

It has been pointed out by several authors that the singlet term is too small to explain the experimental data.

NLO computation

Beneke, Maltoni, Rothstein PRD59 ('99)

$$\Gamma[n] = \Gamma_0 \left[ C_{[1,8]}^2 f[n](\eta) \left(1 + \delta_P[n]\right) \right] + \frac{\alpha_s(\mu)}{4\pi} \left( C_{[1]}^2 g_1[n](\eta) + 2C_{[1]} C_{[8]} g_2[n](\eta) + C_{[8]}^2 g_3[n](\eta) \right) \right] \langle \mathcal{O}^H[n] \rangle,$$



The singlet term has a large renormalization running effect which makes it large negative (unphysical) at mb scale.

Improved NLO result

Beneke, Maltoni, Rothstein PRD59 ('99)

# $Br(B \to J/\psi X) = \underbrace{0.0754 \times 10^{-2} \langle \mathcal{O}_{1}^{\psi}(^{3}S_{1}) \rangle + 0.195 \langle \mathcal{O}_{8}^{\psi}(^{3}S_{1}) \rangle + 0.342 \underbrace{\left[ \langle \mathcal{O}_{8}^{\psi}(^{1}S_{0}) \rangle + 3.1/m_{c}^{2} \langle \mathcal{O}_{8}^{\psi}(^{3}P_{0}) \rangle \right]}_{\mathcal{M}_{1,k}^{\psi}(^{1}S_{0}^{(8)},^{3}P_{0}^{(8)})}$ $Br(B \to \eta_{c}X) = \underbrace{0.250 \times 10^{-2} \mathcal{O}_{1}^{\eta_{c}}(^{1}S_{0}) \rangle + 0.342 \langle \mathcal{O}_{8}^{\eta_{c}}(^{1}S_{0}) \rangle + 0.195 \underbrace{\left[ \langle \mathcal{O}_{8}^{\eta_{c}}(^{3}S_{1}) \rangle - 0.24/m_{c}^{2} \langle \mathcal{O}_{8}^{\eta_{c}}(^{1}P_{1}) \rangle \right]}_{\mathcal{M}_{1,0.24}^{\eta_{c}}(^{3}S_{1}^{(8)},^{1}P_{1}^{(8)})}$ Caveat: A large uncertainty (factor two?)

Beneke, Maltoni, Rothstein PRD59 ('99)

$$\begin{split} Br(B \to J/\psi X) &= \exp\left(1.094 \pm 0.032\right) \times 10^{-2} \\ & 0.0754 \times 10^{-2} \langle \mathcal{O}_{1}^{\psi}(^{3}S_{1}) \rangle + 0.195 \langle \mathcal{O}_{8}^{\psi}(^{3}S_{1}) \rangle + 0.342 \left[ \langle \mathcal{O}_{8}^{\psi}(^{1}S_{0}) \rangle + 3.1/m_{c}^{2} \langle \mathcal{O}_{8}^{\psi}(^{3}P_{0}) \rangle \right] \\ Br(B \to \eta_{c}X) &= \exp\left( \langle 0.9 \times 10^{-2} \langle - \text{ update from LHCb}\right) \right] \\ & \mathcal{M}_{1,k}^{\psi}(^{1}S_{0}^{(8)}, ^{3}P_{0}^{(8)}) \\ & 0.250 \times 10^{-2} \langle \mathcal{O}_{1}^{\eta_{c}}(^{1}S_{0}) \rangle + 0.342 \langle \mathcal{O}_{8}^{\eta_{c}}(^{1}S_{0}) \rangle + 0.195 \left[ \langle \mathcal{O}_{8}^{\eta_{c}}(^{3}S_{1}) \rangle - 0.24/m_{c}^{2} \langle \mathcal{O}_{8}^{\eta_{c}}(^{1}P_{1}) \rangle \right] \\ & \mathcal{M}_{1,0,24}^{\eta_{c}}(^{3}S_{1}^{(8)}, ^{1}P_{1}^{(8)}) \end{split}$$

#### Spin symmetry

Improved NLO result

 $\langle \mathcal{O}_8^{\eta_c}({}^{1}S_0) \rangle = \frac{1}{3} \langle \mathcal{O}_8^{J/\psi}({}^{3}S_1) \rangle,$   $\langle \mathcal{O}_8^{\eta_c}({}^{3}S_1) \rangle = \langle \mathcal{O}_8^{J/\psi}({}^{1}S_0) \rangle,$  $\langle \mathcal{O}_8^{\eta_c}({}^{1}P_1) \rangle = 3 \langle \mathcal{O}_8^{J/\psi}({}^{3}P_0) \rangle.$ 

Can we use this result to extract some information on the octet matrix elements?

### Momentum dependence of B-> J/psi X



octet matrix elements?

FIG. 10:  $p^*$  of  $J/\psi$  mesons produced directly in B decays (points). The histogram is the sum of the color-octet component from a recent NRQCD calculation [20] (dashed line) and the color-singlet  $J/\psi K^{(*)}$  component from simulation (dotted line). Beneke, Scgykerm Wolf, PRD62 ('00)

## Prediction of Br(B-> J/psi X) with the fitted matrix elements

		×10 <sup>-2</sup>	×10 <sup>-2</sup>	x10 <sup>-2</sup>	
	$\langle \mathcal{O}_1^\psi({}^3S_1)  angle$	$\langle \mathcal{O}_8^{\psi}(^3S_1) \rangle$	$\langle \mathcal{O}_8^\psi(^1S_0)  angle$	$\langle \mathcal{O}_8^\psi({}^3P_0)  angle$	$\mathcal{M}_{0,k}^\psi$
Beneke et al	1.16	1.06			2.0-2.7x10 <sup>-2</sup>
Ma et al	1.16	small?	~M <sub>Ok</sub> psi	small?	7.4 x10 <sup>-2</sup>
Chao et al	1.16	0.30±0.12	8.9±0.98	0.56±0.21	9.8x10 <sup>-2</sup>
Kniehl et al	1.32	0.17±0.05	3.04±0.35	-0.91±0.16	1.6x10 <sup>-2</sup>



## Prediction of Br(B-> etac X) with the fitted matrix elements

Assuming Spin symmetry			×10 <sup>-2</sup>	×10-2	<sup>2</sup> x10 <sup>-2</sup>					
		$\langle \mathcal{O}_1^\psi(^3S_1) \rangle$	$\langle \mathcal{O}_8^\psi(^3S_1)  angle$	$\langle \mathcal{O}_8^\psi(^1S_0)  angle$	$\langle \mathcal{O}_8^\psi({}^3P_0) \rangle$	$\mathcal{M}_{1,k}^\psi$				
	Beneke et al	0.39	0.35							
	Ma et al	0.39	small?	~M <sub>Ok</sub> <sup>psi</sup> /3	small???	2.5 x10 <sup>-2</sup>				
	Chao et al	0.39	0.10±0.04	3.0±0.32	1.68±0.63	2.8x10 <sup>-2</sup>				
	Kniehl et al	0.44	0.06±0.02	1.01±0.12	-2.73±0.48	1.3x10 <sup>-2</sup>				
Assuming, O8(1S0) for psi and O8(3S1) for etac is small 0.9×10 <sup>-2</sup> 0.8×10 <sup>-2</sup> 0.7×10 <sup>-2</sup> 0.6×10 <sup>-2</sup>					C	urrent upper limit	D.			
Spin symmetry $(\mathcal{O}^{\eta_c(1_{S_c})}) = \frac{1}{2} / \mathcal{O}^{J/\psi}({}^{3_{S_c}}))$ $(0.5 \times 10^{-2} - 0.4 \times 10^{-2} - 0.3 \times 10^{-2} $							8(150) 8(3ST) (3ST)			
\C \C \C	${}^{\eta_c}_8({}^{3}S_1)\rangle = \langle \mathcal{O}_8^{J/\psi}({}^{1}S_0)\rangle,$ ${}^{\eta_c}_8({}^{3}P_1)\rangle = \langle \mathcal{O}_8^{J/\psi}({}^{1}S_0)\rangle,$	<pre>// 0.2×10<sup>-2</sup> 0.1×10<sup>-2</sup> /</pre>			Caveat: Pr uncertain singl	roportion of th et terms are la	e arge			
	Deneke et al Nhieni et al Chao et al Experiment									