### Lattice result on $J/\psi \to \eta_c \gamma$ , $h_c \to \eta_c \gamma$

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In collaboration with D. Becirevic



Based on "Lattice QCD study of the radiative decays  $J/\psi \to \eta_c \gamma$  and  $h_c \to \eta_c \gamma$ " D.Becirevic and F.Sanfilippo, arXiv:1206.1445, JHEP 1301 (2013) 028

### Summary

### $J/\psi ightarrow \eta_c \gamma$ radiative decay

- $\Gamma(J/\psi \to \eta_c \gamma)$  experimental situation
- 2 Theoretical puzzle

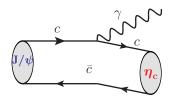
#### Two interesting tests of SM

- **1**  $\Delta_c = M_{J/\psi} M_{\eta_c}$  hyperfine splitting
- 2  $J/\psi \rightarrow e^+e^-$  decay constant

#### $h_c \to \eta_c \gamma$ radiative decay

- **1**  $\Gamma(h_c \to \eta_c \gamma)$  lattice determination
- **2** Prediction for  $\Gamma_{h_c}$

### $J/\psi \rightarrow \eta_c \gamma$ radiative decay



### Current experimental situation is unclear

 $\Gamma(J/\psi \to \eta_c \gamma)_{\rm PDG} = (1.58 \pm 0.37) \, \text{keV}$ :



- $\bullet$  KEDR (arXiv:1002.2071): (2.17  $\pm$  0.14  $\pm$  0.37) keV (preliminary) final result expected this year
- BESIII will hopefully clarify the situation

### $J/\psi \to \eta_c \gamma$ radiative decay

#### Theoretical predictions are inconclusive

- Dispersive bound from  $\Gamma(\eta_c \to 2\gamma)$ :  $\Gamma(J/\psi \to \eta_c \gamma) < 3.2 \text{ keV} [\text{M.A. Shifman, Z. Phys. C 6 ('80)}]$
- Two QCD sum rule calculations gave two different results:
  - $\sim$  (1.7  $\pm$  0.4) keV [A.Y. Khodjamirian, Sov. J. Nucl. Phys. 39 ('84)]  $\sim$  (2.6  $\pm$  0.5) keV [Beilin and Radyushkin, Nucl. Phys. B 260 ('85)]
- QCD eff. theory
  - $\bullet$  (1.5  $\pm$  1.0) keV [N.Brambilla et al, PRD73 ('06)]
  - $(2.14 \pm 0.40) \text{ keV}$  [A.Pineda, J.Segovia, arXiv:1302.3528]
- Potential Quark Models:
  - ullet ~ 3.3 keV [M.B Voloshin, Prog.Part.Nucl.Phys. 61 ('07)]
  - ullet  $\sim$  **2.85** keV [E. Eichten et al., RMP80 ('08)]

#### Lattice QCD computations

- Quenched and single lattice spacing: 2.35(10) keV[J.J Dudek et al., PRD 73 ('06)]
- Unquenched but still single lattice spacing: 2.77 (5) keV[Chen et. al, PRD 84 ('11)]

Both results obtained at large negative  $q^2$ 's, then extrapolated to  $q^2 = 0$ 

### Lattice QCD

#### Desired features

Continuum: Several lattice spacings to take continuum limit

Renormalization: Non perturbative

Momentum: Work directly at  $q^2 = 0$  to avoid the  $q^2$  extrapolation

Unquenching: Include 2 physical light, strange and charm dynamical quarks

### What we currently have...

Continuum: 4 different lattice spacings ( $a \in [0.054; 0.100] \text{ fm}$ )

Renormalization: Non perturbative (RI-MOM)

Momentum: Work at  $q^2 = 0$  using twisted boundary conditions

Unquenching: Only 2 dynamical light quarks ( $M_{\pi} \in [280; 500] \text{ MeV}$ )

- Wilson regularization of QCD with twisted mass term (tmQCD)
- QCD gauge field configurations produced by ETM collaboration



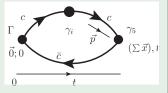
### Form factor computation

$$\Gamma(J/\psi \to \eta_c \gamma) \propto |\langle \eta_c | J^{em} | J/\psi \rangle|_{q^2=0}^2$$

#### Three points functions

$$C_{ij}^{(3)}(t) = \langle \text{Tr} \left[ S_c(y;0) \gamma_i S_c(0,x) \gamma_j S_c^{\vec{p}}(x,y) \gamma_5 \right] \rangle$$
at intermediate times:

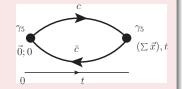
at intermediate times:
$$C_{ij}^{(3)}(t) \underset{0 \ll t \ll T}{\simeq} \frac{Z_{J\psi} Z_{\eta c} \exp \left[ \left( \frac{E_{\eta c} - M_{J/\psi}}{t} \right) t \right] \left\langle \eta_c \middle| J_j^{em} \middle| (J/\psi)_i \right\rangle$$



#### Two points functions

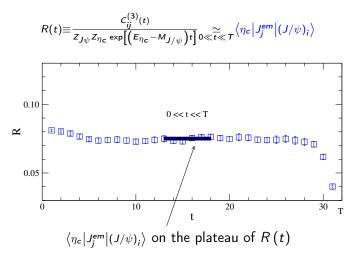
$$C^{(2)}(t) = \langle \text{Tr} \left[ S_c(0,0;\vec{x},t) \gamma_5 S_c(\vec{x},t;\vec{0},0) \gamma_5 \right] \rangle =$$
at large times:
$$\underset{t \to \infty}{\sim} Z_{nc}^2 \exp(-M_{nc} t)$$

and similarly for 
$$Z_{J/\psi}$$
,  $M_{J/\psi}$ .

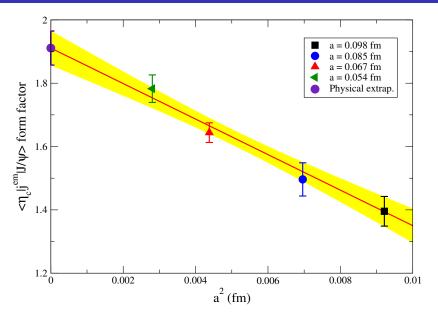


### Matrix element $\langle \eta_c | J_i^{em} | (J/\psi)_i \rangle$ (example)

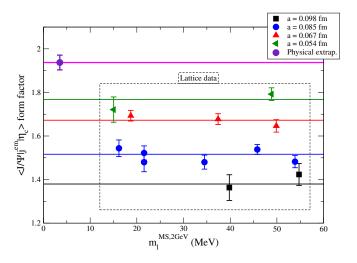
Determine Z and M from 2 points functions and combine with  $C^{(3)}$ :



## Continuum extrapolation of $\langle \eta_c | J_j^{em} | (J/\psi)_i \rangle$



### Dependence on light quark mass



Insensitive to variation of the light sea quark mass  $m_\ell^{\rm sea}$  (expected because  $m_c^{\rm val}\gg m_\ell^{\rm sea}$ )

Expected insensitivity to the dynamical strange quark  $(m_c\gg m_s\gg m_\ell^{sea})$ 

#### Numerical result

#### Our final result

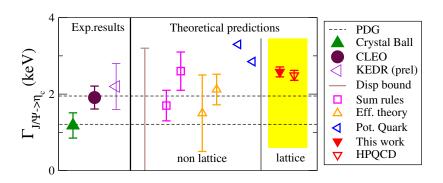
- Putting everything together we get:  $\Gamma(J/\psi \to \eta_c \gamma) = 2.58 \, (13) \, \mathrm{keV}$
- Our value is clearly:
  - larger than Crystal Ball('86) 1.18(33) keV
  - compatible with CLEO('09) 1.91(30) keV, and KEDR('10) 2.17(40) keV

#### Recent development

- Recently HPQCD collab. (PRD86 (2012) 094501) reported result:
  - using Staggered quarks (HISQ regularization)
  - including also dynamical strange quark
- They show that  $\langle \eta_c | J_i^{em} | (J/\psi)_i \rangle$  does not depend on  $m_s^{sea}$
- Excellent agreement with our result in the continuum limit:

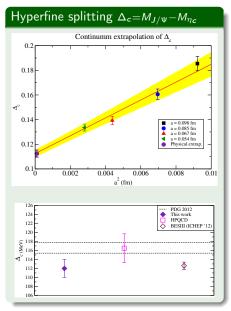
$$\Gamma_{J/\psi \to \eta_c \gamma} = 2.49 \, (19) \, \text{keV}$$

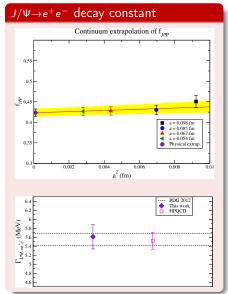
### Is the $J/\psi \to \eta_c \gamma$ puzzle solved?



- Two different lattice approaches give the same result in the continuum
- On the theory side the problem is (almost) fully solved
- This becomes a precision test of QCD
- The experimental situation needs to be clarified

### Other two precise tests of SM

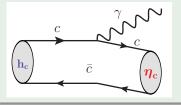




### $h_c o \eta_c \gamma$ radiative decay

### $h_c$ : charmonium $J^{PC} = 1^{+-}$ state

- $h_c$  only recently observed at CLEO (2005)
- Br  $(h_c \to \eta_c \gamma) = 53(7) \%$  at BESIII 2010

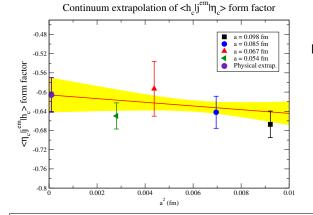


#### Open questions on the initial state

- Lifetime of  $h_c$  not measured yet
- No QCD based estimate: no QCDSR, no effective theory

As before, we can provide the first unquenched LQCD result in the continuum:  $\Gamma(h_c \to \eta_c \gamma)$  is a prediction!

### $h_c \to \eta_c \gamma$ radiative decay



We get:

$$\Gamma(h_c \to \eta_c \gamma) = 0.72(5) \text{ MeV}.$$

BESIII measured:

$$Br(h_c \to \eta_c \gamma) = (53 \pm 7) \%$$

We can predict  $h_c$  lifetime:

$$\Gamma_{h_e} = \frac{\Gamma(h_e \to \eta_e \gamma)}{\text{Br}(h_e \to \eta_e \gamma)}$$

$$1.37(23) \text{ MeV}.$$

Very recently (Confinement '12) BESIII presented preliminary results for  $\Gamma_{h_c}$ :

$$\Gamma^{incl}_{h_c} = 0.73(45)(28)\,\mathrm{MeV}, \qquad \Gamma^{excl}_{h_c} = 0.70(28)(22)\,\mathrm{MeV}$$

### Conclusions

### Main message from Lattice QCD side

First full determination of  $J/\psi \to \eta_c \gamma$  and  $h_c \to \eta_c \gamma$  form factors:

- high statistics unquenched simulations
- continuum extrapolation under control
- non-perturbatively renormalized

#### Results

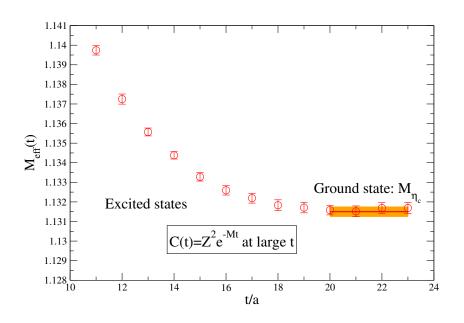
- Assessed theoretical estimate of  $\Gamma(J/\psi \to \eta_c \gamma)$
- Provided a prediction for  $h_c$  lifetime

### Future perspective

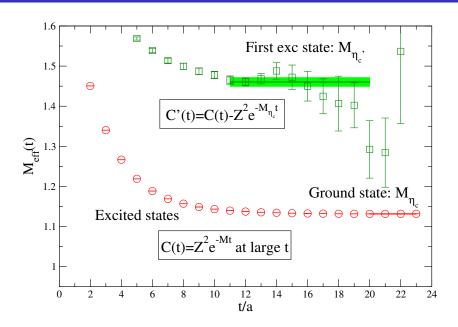
- Check insensitivity to the dynamical charm quark with lattice tmQCD
- Currently working on  $\psi' \to \eta_c \gamma$  and  $\eta'_c \to J/\psi \gamma$
- New measurement of  $\Gamma(J/\psi \to \eta_c \gamma)$  indispensable
- Improve the measurement of  $\Gamma_{h_c}$

# Backup slides

### C(t) effective mass



### C'(t) effective mass



### Can we do charm physics on current lattices?

### Some back of the envelop calculation

- Lattice spacings:  $a \sim 0.050 \div 0.100$  fm,  $1/a \sim 2 \div 4$  GeV
- ullet Charmed meson mass:  $M_{D^\pm}=1.87$  GeV,  $M_{J/\Psi}=3.1$  GeV

To study charm physics on such lattices seem questionable but...

#### Some deeper calculation

- In the free theory the cut off is given by  $p_{max} = \pi/a \sim 6 \div 12$  GeV!
- Seems to be almost good also to study b quark...

Cutoff of interacting theory is unknown: only actual computations can teach us

#### How to keep the situation under control?

Having 4 different lattice spacing, and  $\mathcal{O}(a)$  improved theory allows:

- ullet to drop coarsest lattice spacing and check for stability of a o 0 limit
- ullet to assess the convergence  $\propto a^2$  to the continuum limit:  $\Phi^{latt} = \Phi^{cont.} + \Phi' a^2$

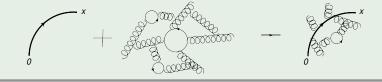
### Determination of the charm quark mass

#### Wick contraction

$$C(\tau) = \sum_{\vec{x}} \left\langle O^{\dagger}(\vec{x}, \tau) O(\vec{0}, 0) \right\rangle \underset{Wick}{=} \operatorname{Tr} \left[ \Gamma S_{I}(\vec{x}, \tau; \vec{0}, 0) \Gamma S_{c}(\vec{0}, 0; \vec{x}, \tau) \right]$$

#### Quark propagator calculation

Solving Dirac equation on gauge background provides full quark propagator  $D_q(y,x) \cdot S_q(x,0) = \delta_{y,0}$   $D_q = (\frac{1}{2\kappa} + K[U]) \mathbf{1} + i m_q \gamma_5 \tau_3$ 



### In practice

- ullet D operator and the propagator S are large matrices ( $\mathcal{O}\left(10^9\right)$  lines)
- Solving Dirac equation requires large amount of CPU resources

### Two points correlation function computation

### Combine 2 propagators with suitable Dirac structures



#### Effective mass

Where ground state dominate, correlation decays exponentially:

$$C(\tau) \xrightarrow[\tau \to \infty]{} Z_D^2 e^{-M_D \tau}$$

The derivative of its logarithm (effective mass) is a constant:

$$M_{eff}(\tau) \equiv -\partial_{\tau} \log C(\tau) \xrightarrow[\tau \to \infty]{} M_{D}$$

- Deviation from constant behavior of  $M_D\left(\tau\right)$  reveals excited states
- Derivative to be computed numerically:  $\partial_{\tau} = f(\tau + 1) f(\tau)$