After the discovery of a Higgs boson(-like) particle – a theorist's perspective

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February 21, 2014

Discoveries at the LHC

Expectations (2008):



't Hooft: "A Higgs, or more" "A super world"

Gross:



Veltman:



"The unexpected"

Discoveries at the LHC

Expectations (2008 it Hooft: "A Higgs,	3): Gross: or more" "A supe	ve vr world" "T	Itman: Itman: ne unexpected"
Nov. 2011 HCP 2011:	13 Dec. 2011	4 July 2012	8 Oct. 2013 Nobel prize:
Exclusion of a wide Higgs mass range, some theorists' thought: "complete exclusion until the end of 2011"	ATLAS & CMS report an excess of events: Too early to draw conclusions	ATLAS & CMS announce the discovery of a Higgs-like particle	for the theoret. discovery of a mechanism that contributes to our understanding of the origin of mass

How was it discovered?

Discovery channels (LHC: p p collider, p = proton):

 $pp \rightarrow H \rightarrow ZZ^* \rightarrow 4\ell$:

(here: $\ell = \mu$, μ = muon)





 $pp \rightarrow H \rightarrow \gamma \gamma$:





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How was it discovered?

Discovery channels (LHC: p p collider, p = proton):

 $pp \rightarrow H \rightarrow ZZ^* \rightarrow 4\ell$: (here: $\ell = \mu, \mu = muon$) $pp \rightarrow H \rightarrow \gamma \gamma + 2$ jets:



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• Mass: free parameter in the Standard Model

expectation from precision measurements: O(100 GeV)(e.g. mass of the *W* boson)

Moriond '13: CMS: $m_H = 125.7 \pm 0.3$ (stat) ± 0.3 (syst) GeV

ATLAS: $m_H = 125.5 \pm 0.2$ (stat) $^{+0.5}_{-0.6}$ (syst) GeV

• **Spin?** Landau-Yang theorem:

Massive spin-1 particle cannot decay into two photons:

Decay into photons observed \Rightarrow spin \neq 1

Moriond '13: spin = 2: Excluded with > 99% confidence level spin = 0: compatible model dependent

• CP? Moriond '13: CP-even: compatible spin = 0 CP-odd: Exclusion with $\gtrsim 98\%$ confidence level

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- Couplings? so far compatible with the Standard Model:
 - Measurement of further production und decay channels:

 $pp \rightarrow H \rightarrow WW$ (compatible with SM)

$$pp \rightarrow H \rightarrow \tau \tau$$
 (Evidence!)

• • •

- still relatively large errors (\sim 20 %)
- not all couplings accessible

Signal strengths:



- Mass: free parameter in the Stand Model expectation from pr models: Other mass by W boson) Moriond '13: C Higgs be given eters. ATLA other parameter 0.3 (stat) ± 0.3 (syst) GeV
- Spin? Landau-Yang theorem:

Massive spin-1 particle cannot decay into two photons:

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- still relatively large errors (\sim 20 %)
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Further Results of the LHC

Supersymmetric partner particles:

- not found yet
- strongest constraints: colour charged particles: gluinos and first generation of squarks signature: jets and missing energy
- much less constrained:
 - ⋆ top squarks
 - purely electroweak particles signature: three charged leptons

Note specific assumptions: simplified models, ...

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Further Results of the LHC

Further particles:

not found yet



*Only a selection of the available mass limits on new states or phenomena shown

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Reasons for further investigations

- Main question: What is the underlying theory?
- Reasons to search for beyond the Standard Model Physics

From experiments:

* dark matter

- * matter-antimatter asymmetry in the universe
- * neutrino oscillations

From theory:

- * grand unification
- * embedding of gravity
- * Higgs mass M_H : sensitive to physics at high energy scales Λ quantum corrections: $\delta M_H^2 \sim \Lambda^2$ (hierarchy problem)

Minimal Supersymmetric Standard Model (MSSM)

MSSM: * Extension of the Standard Model (SM)

★ Further symmetry:

Supersymmetry (SUSY):

 $Q|Boson\rangle = |Fermion\rangle, \qquad Q|Fermion\rangle = |Boson\rangle$

Q = supersymmetry generator

- **Recipe:** Standard Model particles + 2nd Higgs doublet (2HDM) (Generation of fermion masses, anomaly cancelations)
 - Superpartners
 - Explicit soft SUSY-breaking ⇒ many new (complex) parameters (Else: mass_{superpartner} = mass_{2HDM-particle} ← exp. excluded)
 - R-Parity: discrete symmetry

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Minimal Supersymmetric Standard Model (MSSM)



Particles with same quantum numbers can mix:

charged Higgsinos and Gauginos \rightarrow Charginos neutral Higgsinos and Gauginos \rightarrow Neutralinos

Higgs Sector at Born Level

Higgs potential: gauge couplings H_d , H_u : Higgs doublets $V_{\text{Higgs}} = \frac{g^2 + {g'}^2}{8} (H_d^+ H_d - H_u^+ H_u)^2 + \frac{g^2}{2} |H_d^+ H_u|^2$ $+ |\mu|^2 (H_d^+ H_d + H_u^+ H_u)$ μ : coupl. betw. Higgs superfields $+(m_1^2H_d^+H_d+m_2^2H_u^+H_u)$ soft breaking terms $+ (\epsilon_{ii} | m_{12}^2 | e^{i \varphi_{m_{12}^2}} H_d^i H_{ii}^j + h.c.)$ non-vanishing phases: • one phase in the Higgs potential: $\varphi_{m_{e}^2}$ \Rightarrow maybe CP- or T-violation? (Time reversal-operator is • phase difference of Higgs doublets ξ antiunitary \Rightarrow complex conjugation

Higgs Sector at Born Level

Higgs potential: gauge couplings H_d , H_u : Higgs doublets $V_{\text{Higgs}} = \frac{g^2 + {g'}^2}{8} (H_d^+ H_d - H_u^+ H_u)^2 + \frac{g^2}{2} |H_d^+ H_u|^2$ $+ |\mu|^2 (H_A^+ H_d + H_\mu^+ H_\mu)$ μ : coupl. betw. Higgs superfields $+(m_1^2H_d^+H_d+m_2^2H_u^+H_u)$ soft breaking terms $+ (\epsilon_{ii} | m_{12}^2 | e^{i\varphi_{m_{12}^2}} H_d^i H_{ii}^j + h.c.)$

- one phase in the Higgs potential: $\varphi_{m_{12}^2}$ can be rotated away
- phase difference of Higgs doublets ξ: vanishes because of minimum condition

ho CP violation
at Born level
in the Higgs sector

Higgs Sector at Born Level

Physical mass eigenstates (at Born level):

• 5 Higgs bosons: 3 neutral H, h, A; 2 charged H^{\pm}

Masses of the Higgs bosons:

• not all independent:

often: Mass M_A or $M_{H^{\pm}}$ (and tan β) as free parameter tan $\beta = \frac{v_2}{v_1}$: ratio of the Higgs vac. expect. values

• lightest Higgs boson: h

Upper theoretical Born mass bound:

 $M_h \leq M_Z = 91 \text{ GeV}$

with quantum corrections of higher orders: $M_h \lesssim 140 \text{ GeV}$

____ dependent on the MSSM parameters

Why a precise Higgs mass prediction?

 Needed as consistent input for the calculation of cross sections and decay widths in the MSSM



A precise theoretical prediction is needed to fully exploit this constraint:

 $\Delta M_H^{ ext{exp.}} < 1 ext{ GeV}$ vs $\Delta M_H^{ ext{theory}} pprox 3 ext{ GeV}$

 In the discussion of the amount of fine-tuning of the MSSM the precise theoretical prediction of the Higgs boson mass enters.

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Contributions to the Higgs boson masses

Quantum corrections:





One-loop level $\mathcal{O}(\alpha_t)$: $\alpha_t \sim (\text{top Yukawa coupl.})^2$

Two-loop level $\mathcal{O}(\alpha_t \alpha_s)$:

Real parameters:



no mixing between CP-even and CP-odd states

 \Rightarrow Lightest Higgs boson is CP-even.

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Implications of a 125.5 GeV Higgs boson (MSSM)



- A 125.5 \pm 3 GeV mass constrains the parameter space but does not exclude the MSSM. (theory uncertainty \approx 3 GeV)
- here: no known 3-loop contributions included
 [Martin; Harlander, Kant, Mihaila, Steinhauser]

Implications of a 125.5 GeV Higgs boson (MSSM)



For parameter scans, see e.g.

[Heinemeyer, Stål, Weiglein, arXiv:1112.3026;

Arbey, Battaglia, Djouadi, Mahmoudi, Quevillon, arXiv:1112.3028]

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Higgs boson mass and CP-violating phases



- The Higgs mass does depend on the squark mixing phase φ_{X_t} .
- For $\varphi_{X_t} \neq n\pi, n \in \mathbb{N}_0$, h_1 is not a CP-eigenstate.

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Zwirner;

(no resummation)

Dedes, Degrassi, Slavich]

Higgs boson mass and CP-violating phases



- The Higgs mass does depend of the squark mixing phase φ_{X_t} .
- For $\varphi_{X_t} \neq n\pi, n \in \mathbb{N}_0$, h_1 is not a CP-eigenstate. To do: Implementation of $\mathcal{O}(\alpha_t^2)$ contr. [Hollik, Passehr, arXiv:1401.8275]

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Zwirner;

Prediction obtained via Feynman diagrammatic approach:

- + all log and non-log terms are taken into account at a certain order of perturbation theory
- possible appearance of large logs:

$$\Delta M_H \sim \log rac{M_S}{m_t}$$

M_S: SUSY particle mass scale *m_t*: top quark mass

- $\Rightarrow \bullet$ good prediction for lower SUSY mass scales
 - no reliable prediction for large SUSY mass scales

Other approach: Renormalization Group Equation (RGE) approach:

* assume: all SUSY particles are heavy of order $\sim M_S$:



 \star evolve λ to lower scale using Standard Model running (RGE)

 \star the Higgs mass² is then $M_h^2(m_t) = 2\lambda(m_t)v^2$ $v \approx 174 \text{ GeV}$

 \Rightarrow logs resummed to all orders: good prediction for large SUSY masses

 \rightarrow Combine both approaches (now:only for real parameters)

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Similar findings for large SUSY masses (pure RGE approach):

[Draper, Lee, Wagner, arXiv:1312.5743]



2-loop NNLL result 3-loop NNLL result 4-loop NNLL result resummed result

Differences:

- 3-loop running for λ
- 2-loop matching
- not yet implemented into a computer code

Deviations from the Standard Model?



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Needed:

- precise predictions of the signal cross sections in the SM Calculation of higher-order corrections necessary
- precise determination of the background
- good error estimate

(e.g. including possible errors originating from a contamination of control regions by events of unknown particles

[Feigl, H.R., Zeppenfeld, arXiv:1205.3468])

Deviations from the Standard Model (SM)?

Where can they originate from?

· changes in loop-induced couplings due to unknown particles

- changes of couplings due to mixing effects in an enlarged Higgs sector
- · changes due to additional decay into invisible particles
- changes due to two degenerate Higgs bosons

Next-tominimal Supersymmetric Standard Model (NMSSM)

In the MSSM: Parameter μ in the superpotential W, $W = \mu H_1 H_2 + ...$ A priori: arbitrary value But: Order of the electroweak/SUSY-breaking scale

In the NMSSM: Additional Higgs superfield singlet S:

 μ is generated via a vacuum expectation value of the scalar Higgs singlet, $W = \lambda SH_1H_2 + ..., \lambda =$ new coupling

- 7 Higgs bosons: *h*₁, *h*₂, *h*₃, *A*₁, *A*₂, *H*[±]
- CP-violation in the Higgs sector possible already at Born level (complex parameters)
- Mass of the light CP-even MSSM-like Higgs boson can be larger than in the MSSM
- 5 neutralinos (4 neutralinos in the MSSM)

In the NMSSM:

2 Higgs bosons could be nearly degenerate with masses of \sim 125 GeV

> [Gunion, Jiang, Kraml, arXiv:1207.1545]

For illustration (one possibility):



Illustration: Cross-over region of H_1 and H_2 at $A_{\kappa} \approx -210$ GeV:

- Masses $M_{H_i} \sim 125 \text{ GeV}$
- trilinear, SUSY-breaking singlet coupling • H_1 and H_2 interchange their role



Illustration: Cross-over region of H_1 and H_2 : Masses $\sim 125 \text{ GeV}$

 $A_{\kappa} < -210 \text{ GeV}$: H_1 singlet-like, H_2 non-singlet like $A_{\kappa} > -210 \text{ GeV}$: H_1 non-singlet like, H_2 singlet-like



Illustration: • Higher-order corrections are necessary

• Implemented in NMSSMCalc: program for evaluation of mass spectra and decay widths (allows also for complex parameters) [Baglio, Gröber, Mühlleitner, Nhung, H.R., Spira, Streicher, Walz]

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Legend: $m_{h_2} - m_{h_1} \le 1$ GeV, 1 GeV $< m_{h_2} - m_{h_1} \le 2$ GeV, 2 GeV $< m_{h_2} - m_{h_1} \le 3$ GeV

Signal ratio: $R_{gg}^{h}(XX) = \sum_{i=1}^{2} \frac{\sigma(gg \to h_i) BR(h_i \to XX)}{\sigma^{SM}(gg \to h) BR^{SM}(h \to XX)}$

with an effective SM Higgs mass

(Parameter scans with NMSSMTools [Ellwanger, Gunion, Hugonie])

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- Enhancement of the $h \rightarrow \gamma \gamma$ channel by a factor of 1.5 possible without enhancing $h \rightarrow VV$
- Enhancement of the $h \rightarrow \gamma \gamma$ channel is generally achieved by reducing the average $b\bar{b}$ coupling strength C_{bb}^{h}

[Gupta, H.R., Wells, arXiv:1206.3560, arXiv:1305.6397]

If no beyond Standard Model (SM) physics is seen at the LHC: (related to electroweak symmetry breaking)

How large can deviations from the SM Higgs couplings be?

How precisely do we need to measure the Higgs couplings at least?

Which future collider and detectors?

[Gupta, H.R., Wells, arXiv:1206.3560, arXiv:1305.6397] No completely model independent answer possible:

- MSSM
- Mixed-in singlet model = Standard Model + exotic Higgs boson singlet S: Scalar fields mix via $|H_{SM}|^2 |S|^2$ [Schabinger, Wells, hep-ph/0509209; Bowen, Cui, Wells, hep-ph/0701035]
 - ⇒ 2 CP-even mass eigenstates: SM-like *h*, heavier *H* $c_h = \cos \theta_h$ with couplings² $g_h^2 = c_h^2 g_{SM}^2$ and $g_H^2 = s_h^2 g_{SM}^2$ $\theta_h = \text{mix. angle}$
- Composite Higgs models: SM-like Higgs boson = pseudo-Goldstone: SM vector bosons and fermions + strong sector with Higgs multiplet in terms of an effective field theory for a strong interacting light Higgs (SILH) boson

Take into account:

[Gupta, H.R., Wells, arXiv:1206.3560, arXiv:1305.6397]

- discovery potential of the LHC for further particles related to electroweak symmetry breaking
- * constraints from electroweak precision tests



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[Gupta, H.R., Wells, arXiv:1206.3560, arXiv:1305.6397]

Maximal deviations from the SM Higgs couplings? (no new physics at the LHC)

	∆ <i>hVV</i> c. to gauge bosons <i>V</i>	∆ <i>hītt</i> c. to top quarks <i>t</i>	∆ <i>h̄bb</i> c. to bottom quarks <i>b</i>	∆ <i>hhh</i> triple Higgs coupl.		
MSSM	< 1%	3%	10%, 100%	2%, 15%		
Mixed-in Singlet	6%	6%	6%	18%		
Composite Higgs	8%	tens of %	tens of %	tens of %		
			$\tan \beta > 20$ all other no superpartners cases			
\rightarrow a aballance for the future						

\Rightarrow a challenge for the future!

After the discovery of a Higgs boson(-like) particle – a theorist's perspective

Summary

- Exciting times: Discovery of a Higgs boson
- Still exciting: What is its true nature?
- At the moment: compatible with the Standard Model
- Constraints on the parameter space of extension of the Standard Model (e.g. coming from the Higgs mass)
- Precise measurement of its couplings can help with the complete identification of the particle
- Program development:
 - * FeynHiggs: Higgs mass prediction for large stop masses improved

* NMSSMCALC: New program for the NMSSM Higgs masses/decay widths