#### Higgs Hunting 2013, July 25-27, 2013, Orsay, France



Aurelio Juste ICREA/IFAE, Barcelona

For the CDF and DØ Collaborations





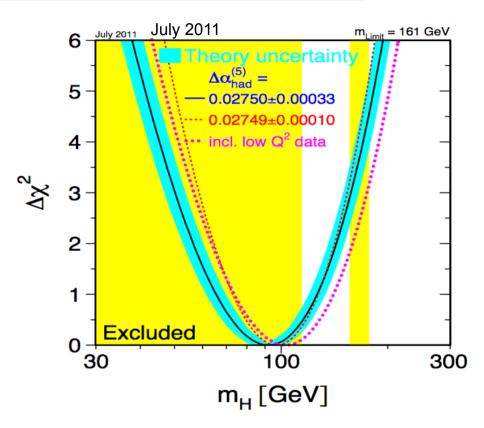




### **Historical Perspective**

- Before direct constraints from hadron colliders all we knew about the SM Higgs boson came from:
  - direct searches at LEP m<sub>H</sub>>114.4 GeV @ 95% CL
  - indirect constraints from precision electroweak observables (\*)

- (\*) using latest  $m_t$  and  $M_W$  measurements from the Tevatron
- → The Tevatron is expected to be sensitive precisely in this mass range!
- → Searches for the SM Higgs boson have been a priority at the Tevatron.



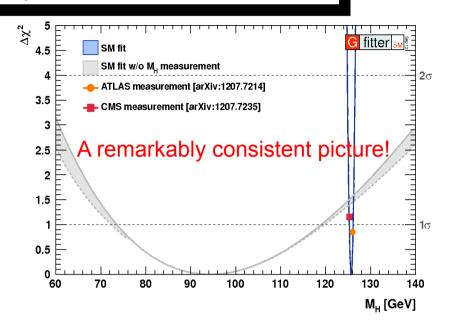
 Over the last decade the outstanding performance of the accelerator and CDF and DØ detectors, together with continuous improvements to the searches, allowed to start placing direct limits beyond LEP and eventually reach SM sensitivity in the preferred mass range.

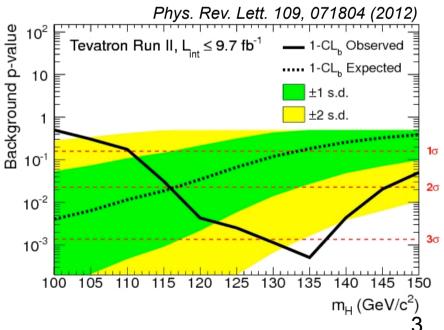
### **Historical Perspective**

 Over the last two years, the LHC has excluded most of the remaining allowed mass range and observed a new boson with mass ~125 GeV, primarily in γγ and ZZ→4I final states.

The Tevatron experiments independently reported evidence for a new particle decaying into bb with compatible mass.

- Since the observation, Higgs studies at the Tevatron have focused on the measurement of Higgs properties (rates in different decay modes, couplings, J<sup>P</sup>).
- In addition, Tevatron measurements of key electroweak parameters (m<sub>t</sub> and M<sub>W</sub>) will play an increasingly important role to help unravel the nature of the discovered boson and point to potential new physics effects.





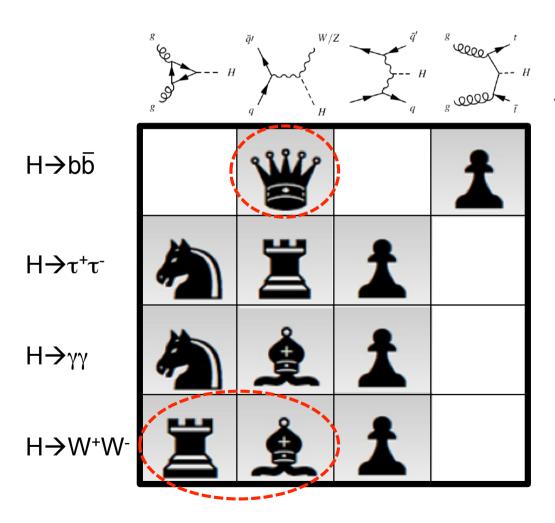
# **Today's Presentation**

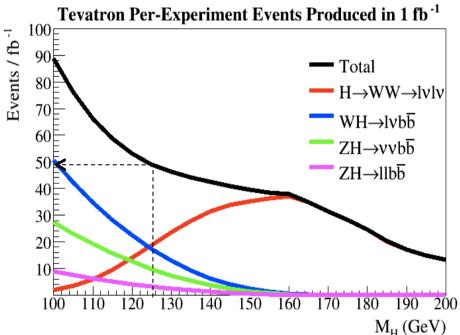
- Introduction
- Direct constraints on the Higgs boson
  - → Here will just give a quick overview of SM Higgs boson searches. For more details see talks by S. Shalhout and G. Davies.
  - → Will not discuss BSM Higgs boson searches. For an overview see talk by J. Hays.
- Indirect constraints on the Higgs boson: m<sub>t</sub> and m<sub>W</sub>
  - → Main focus of this talk
- Summary and outlook

# **Direct Constraints**

# **Search Strategies**

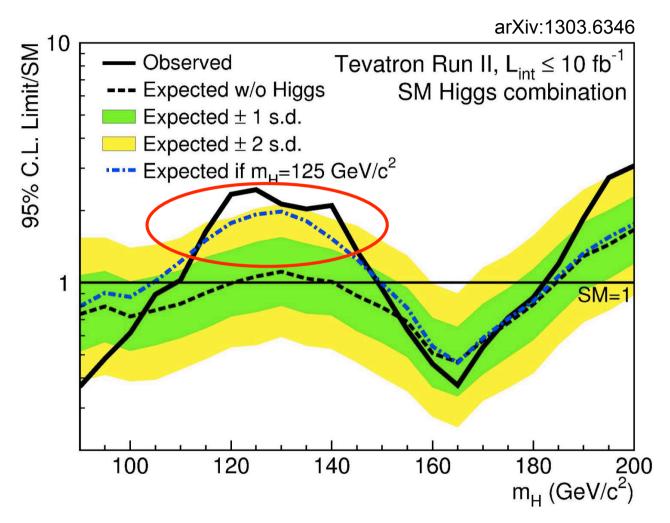
 Defined by a combination of theoretical and experimental considerations (large σxBR but experimentally feasible: trigger, backgrounds....).





For m<sub>H</sub>=125 GeV, ~1000 Higgs events produced at the Tevatron in the main search channels with 10 fb<sup>-1</sup>!

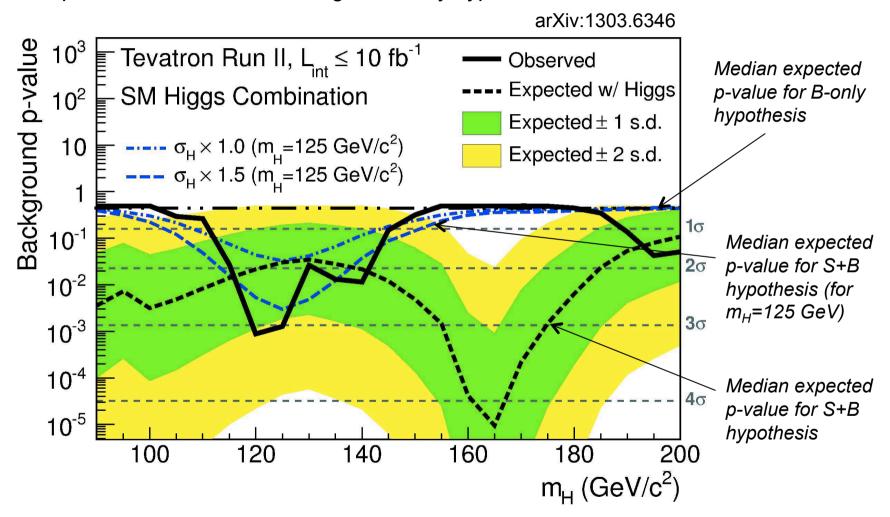
#### **Tevatron Combined Results**



- Expected exclusion:  $90 < m_H < 120$  GeV,  $140 < m_H < 184$  GeV Observed exclusion:  $90 < m_H < 109$  GeV,  $149 < m_H < 182$  GeV
- 95% CL limit at m<sub>H</sub>=125 GeV: 1.06xSM (expected), 2.44xSM (observed)

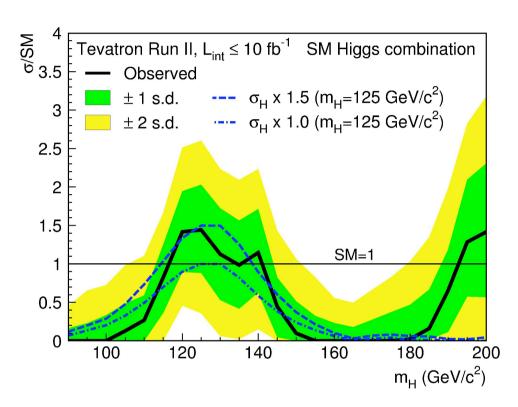
# Quantifying the Excess: p-values

• Local p-value distribution for background-only hypothesis:



• Minimum local p-value at  $m_H=120$  GeV:  $3.1\sigma$  ( $2.0\sigma$  expected) p-value at  $m_H=125$  GeV:  $3.0\sigma$  ( $1.9\sigma$  expected)

# Quantifying the Excess: Best Fit Signal Rate



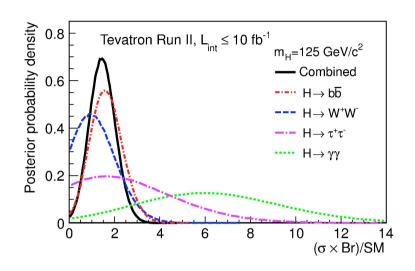
- Maximum likelihood fit to data with signal rate as free parameter.
- Best-fit signal rate at m<sub>H</sub>=125 GeV:

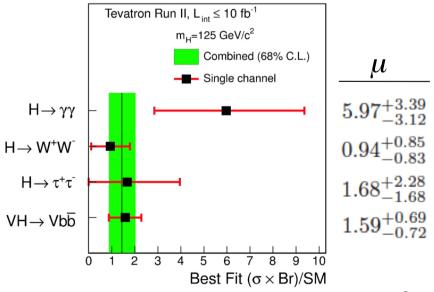
$$\mu = 1.44^{+0.59}_{-0.56}$$

Consistent with SM Higgs.

Reasonably consistent across channels.

arXiv:1303.6346





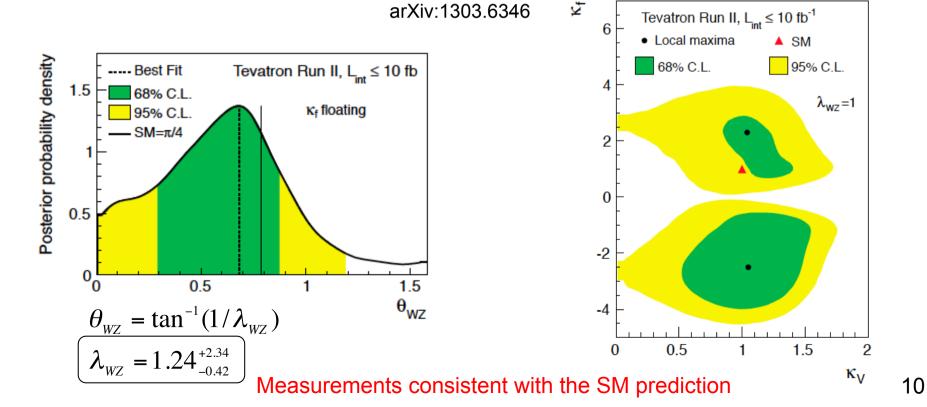
# **Probing Higgs Boson Couplings**

#### Benchmark I:

• Probe SU(2)<sub>V</sub> custodial symmetry by measuring the ratio  $\lambda_{WZ} = \kappa_W / \kappa_Z$ . Assume all fermion couplings are scaled by a common free parameter ( $\kappa_F$ ).

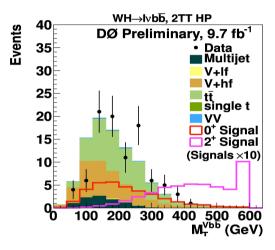
#### Benchmark II:

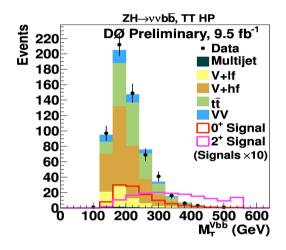
- Consider two independent multiplicative factors: common to all couplings to vector bosons ( $\kappa_V$ ) and common to all couplings to fermions ( $\kappa_f$ ). Assume  $\lambda_{WZ}$ =1.
- Measure  $\kappa_f$  and  $\kappa_V$  simultaneously.

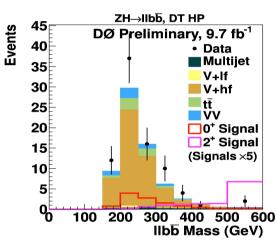


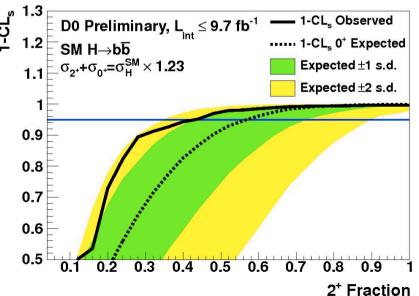
# Spin/Parity

- LHC results in bosonic decay modes consistent with JP=0+ hypothesis.
- Tevatron can check J<sup>P</sup> in VH(H→bb) production (J. Ellis et al., JHEP 1211, 134 (2012)).
   So far results from DØ on 0<sup>+</sup> vs 2<sup>+</sup> (graviton-like couplings). 0<sup>+</sup> vs 0<sup>-</sup> and Tevatron combination upcoming.









DØ Note 6387-CONF

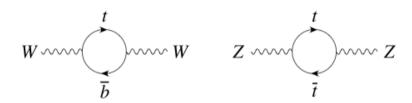
For  $\mu$ =1.23 (DØ measured rate):

- Exclude 2<sup>+</sup> at 99.9% CL (in favor of 0<sup>+</sup>).
- Assuming the excess results from a superposition of a 0<sup>+</sup> and 2<sup>+</sup> particle:
   Exclude a 2<sup>+</sup> fraction f<sub>2+</sub>>0.42 at 95% CL (in favor of pure 0<sup>+</sup>).

# **Indirect Constraints**

 Experimental and theoretical precision achieved allow to test the electroweak theory at the quantum level.

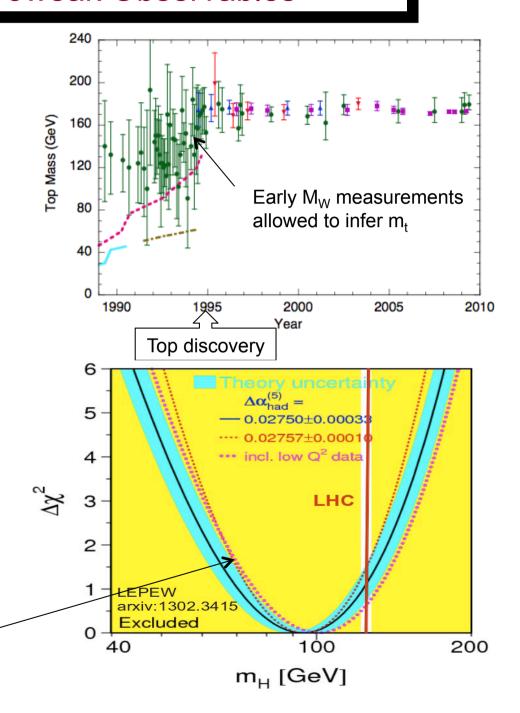
$$\label{eq:mww} {\color{blue}m_W^2} \left(1 - \frac{m_W^2}{m_Z^2}\right) = \frac{\pi \alpha}{\sqrt{2} G_F} (1 + \Delta r)$$



$$\Delta r_{\rm top} = -\frac{3\alpha}{16\pi} \frac{\cos^2 \theta_W}{\sin^4 \theta_W} \frac{m_t^2}{m_W^2}$$

$$\Delta r_{\rm Higgs} = + \frac{11\alpha}{48\pi \sin^2 \theta_W} \log \frac{m_H^2}{m_W^2}$$

Precision EW observables, along with  $m_t$ , allow to infer  $M_H$  ( $M_W$  is one of the most sensitive measurements)



 Experimental and theoretical precision achieved allow to test the electroweak theory at the quantum level.

$$m_W^2 \left(1 - \frac{m_W^2}{m_Z^2}\right) = \frac{\pi \alpha}{\sqrt{2}G_F} (1 + \Delta r)$$

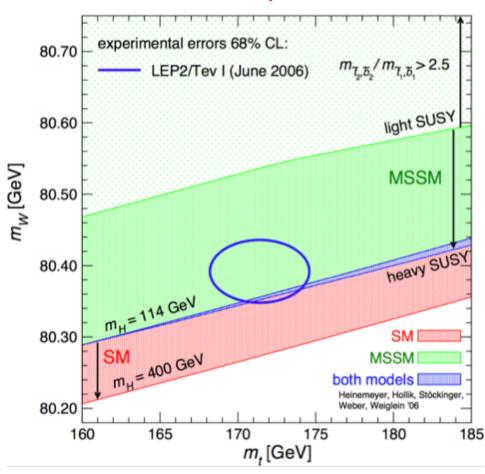
$$W \sim W \qquad Z \sim Z$$

$$\Delta r_{\text{top}} = -\frac{3\alpha}{16\pi} \frac{\cos^2 \theta_W}{\sin^4 \theta_W} \frac{m_t^2}{m_W^2}$$

$$h \qquad \qquad + \qquad \qquad M$$

$$\Delta r_{\rm Higgs} = +\frac{11\alpha}{48\pi \sin^2 \theta_W} \log \frac{m_H^2}{m_W^2}$$

#### Status early in Run II



 $\Delta m_t \sim 2$  GeV (from early Tevatron Run II meas.)  $\Delta M_W \sim 30$  MeV (dominated by LEP2)

 Experimental and theoretical precision achieved allow to test the electroweak theory at the quantum level.

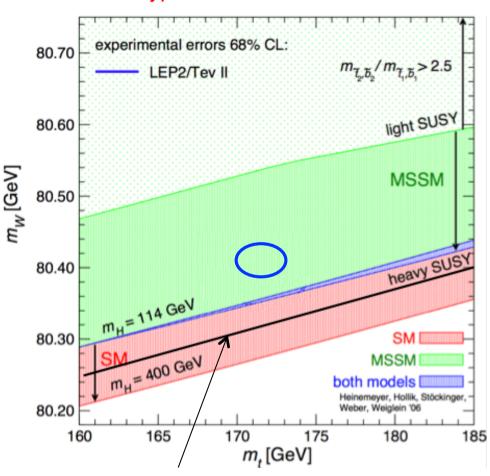
$$m_W^2 \left(1 - \frac{m_W^2}{m_Z^2}\right) = \frac{\pi \alpha}{\sqrt{2}G_F} (1 + \Delta r)$$

$$W \sim W \qquad Z \sim Z$$

$$\Delta r_{\text{top}} = -\frac{3\alpha}{16\pi} \frac{\cos^2 \theta_W}{\sin^4 \theta_W} \frac{m_t^2}{m_W^2}$$

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#### A Hypothetical Scenario in 2012



Higgs boson with M<sub>H</sub>=160 GeV discovered at Tevatron or LHC

∆m<sub>t</sub>~1 GeV (~ultimate Tevatron precision)

 $\Delta M_W \sim 15$  MeV (dominated by Tevatron) 15

 Experimental and theoretical precision achieved allow to test the electroweak theory at the quantum level.

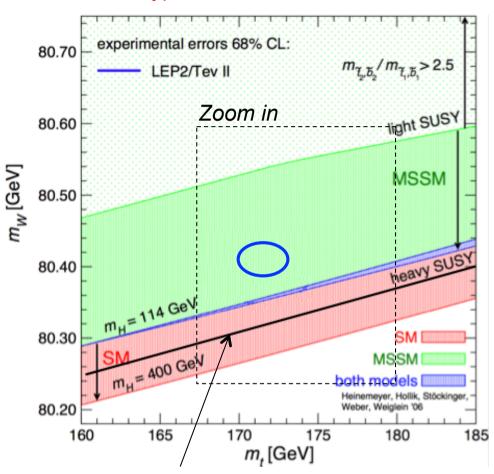
$$m_W^2 \left(1 - \frac{m_W^2}{m_Z^2}\right) = \frac{\pi \alpha}{\sqrt{2}G_F} (1 + \Delta r)$$



$$\Delta r_{\text{top}} = -\frac{3\alpha}{16\pi} \frac{\cos^2 \theta_W}{\sin^4 \theta_W} \frac{m_t^2}{m_W^2}$$

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#### A Hypothetical Scenario in 2012



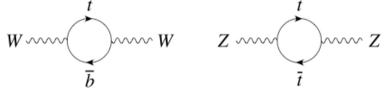
Higgs boson with  $M_H$ =160 GeV discovered at Tevatron or LHC

 $\Delta m_t \sim 1$  GeV ( $\sim$ ultimate Tevatron precision)

 $\Delta M_W \sim 15$  MeV (dominated by Tevatron) 16

 Experimental and theoretical precision achieved allow to test the electroweak theory at the quantum level.

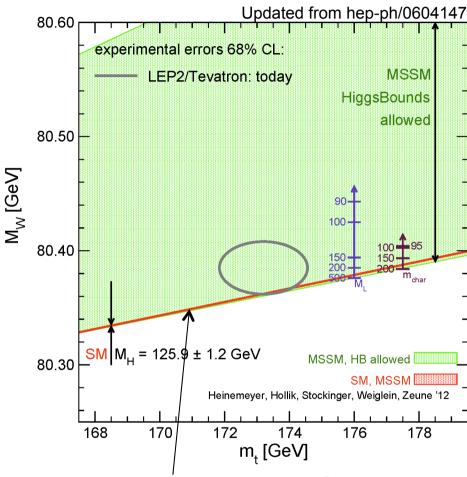
$$m_W^2 \left(1 - \frac{m_W^2}{m_Z^2}\right) = \frac{\pi \alpha}{\sqrt{2}G_F} (1 + \Delta r)$$



$$\Delta r_{\rm top} = -\frac{3\alpha}{16\pi} \frac{\cos^2 \theta_W}{\sin^4 \theta_W} \frac{m_t^2}{m_W^2}$$

$$\Delta r_{\rm Higgs} = +\frac{11\alpha}{48\pi \sin^2 \theta_W} \log \frac{m_H^2}{m_W^2}$$

#### Actual Scenario in 2012

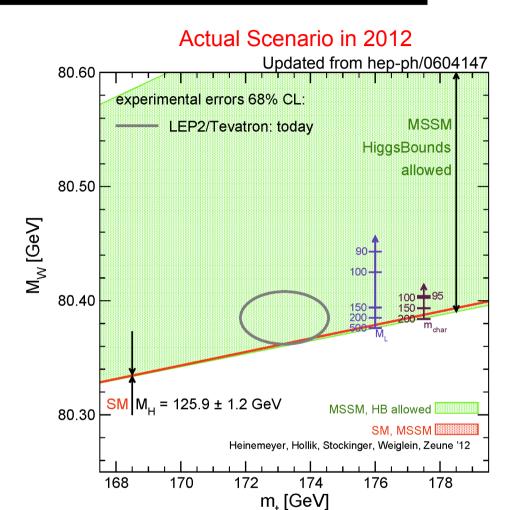


Higgs boson with  $M_H$ ~125 GeV discovered at LHC

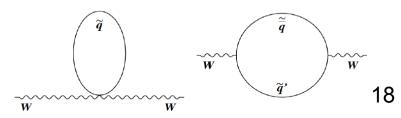
Δm<sub>t</sub>~0.94 GeV (still improving)

 $\Delta M_W \sim 15 \text{ MeV (still improving)}$ 

 Experimental and theoretical precision achieved allow to test the electroweak theory at the quantum level.

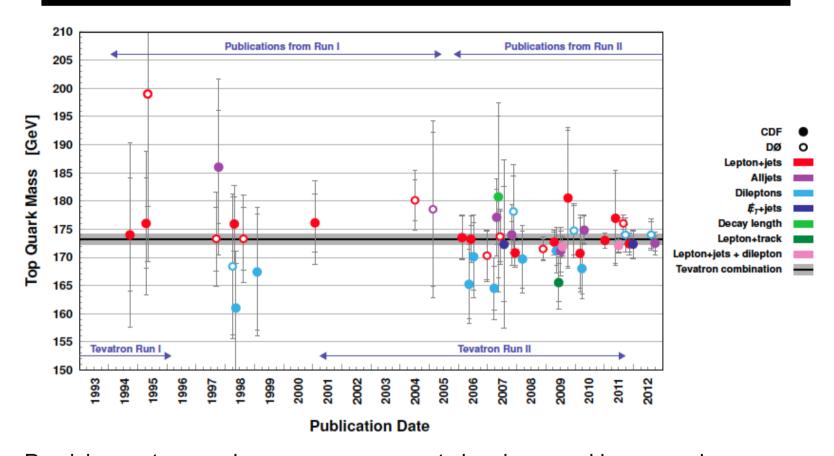


#### SUSY loops can contribute ~100 MeV to M<sub>W</sub>!



# Top Quark Mass

# **Top Quark Mass**



- Precision on top quark mass measurements has improved by one order or magnitude since Run I:
  - Best single measurement in Run I: Δm<sub>t</sub>/m<sub>t</sub>~7%
  - Best single measurement in Run II: ∆m<sub>t</sub>/m<sub>t</sub>~0.7%
- This has been possible, not only to the much large available statistics, but also to the development of novel experimental techniques, most of which have been adopted by LHC experiments.

#### Handles for a Precise Measurement

#### Jet Energy Scale (JES)

- Top mass measurement requires precise mapping between reconstructed jets and original partons → correct for detector, jet algorithm and physics effects.
- What's crucial is the relative energy calibration between data and MC jets:

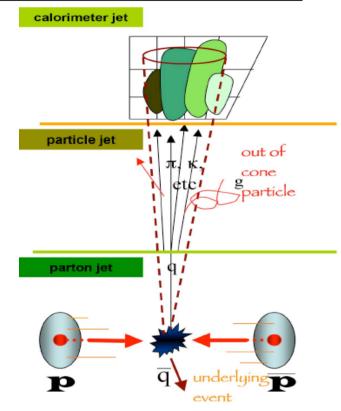
$$\Delta E_{jet}/E_{jet}\sim 1\% \Rightarrow \Delta m_t \sim 1 \text{ GeV}$$

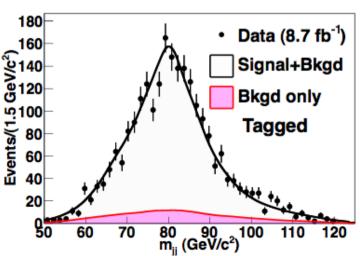
- Handles:
  - dijets, photon+jets, Z+jets
  - W mass from W→jj in top quark decays (in-situ calibration)

B-tagging: reduction of physics as well as combinatorial background

Sophisticated mass extraction techniques: maximize statistical sensitivity; minimize some systematic uncertainties (e.g. JES)

Simulation: accurate detector modeling and stateof-the-art theoretical knowledge (gluon radiation, b-fragmentation, etc) required.





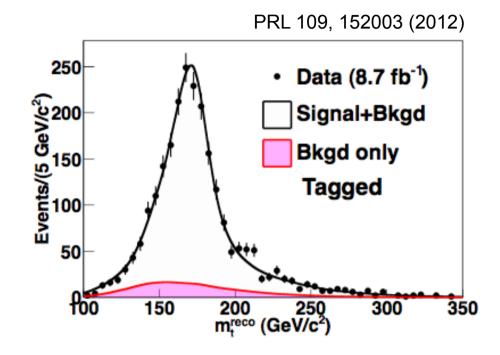
# Mass Extraction Techniques: Template Methods

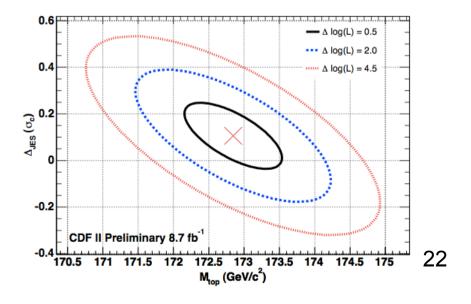
- Identify kinematical variables strongly correlated with m<sub>t</sub>. Compare data and MC with different m<sub>t</sub> hypotheses.
- Example: reconstructed m<sub>t</sub> from kinematic fit in lepton+jets channel.

$$\chi^{2} = \sum_{i=l, jets} \frac{(p_{T}^{i, fit} - p_{T}^{i, meas})^{2}}{\sigma_{i}^{2}} + \sum_{j=x, y} \frac{(UE_{j}^{fit} - UE_{j}^{meas})^{2}}{\sigma_{j}^{2}} + \frac{(M_{jj} - M_{W})^{2}}{\Gamma_{W}^{2}} + \frac{(M_{ly} - M_{W})^{2}}{\Gamma_{t}^{2}} + \frac{(M_{bly} - m_{t}^{reco})^{2}}{\Gamma_{t}^{2}} + \frac{(M_{bly} - m_{t}^{reco})^{2}}{\Gamma_{t}^{2}}$$

Usually pick solution with lowest  $\chi^2$ .

 Reduce JES systematic by using in-situ hadronic W mass in tt events: simultaneous determination of m<sub>t</sub> and JES scaling factor from reconstructed m<sub>t</sub> and M<sub>W</sub> templates.





# Mass Extraction Techniques: Dynamic Methods

- Compute event-by-event probability as a function of m<sub>t</sub> making use of all reconstructed objects in the events (integrate over unknowns). Maximize sensitivity by:
  - summing over all permutations of jets and neutrino solutions
  - allowing better measured events to contribute more.

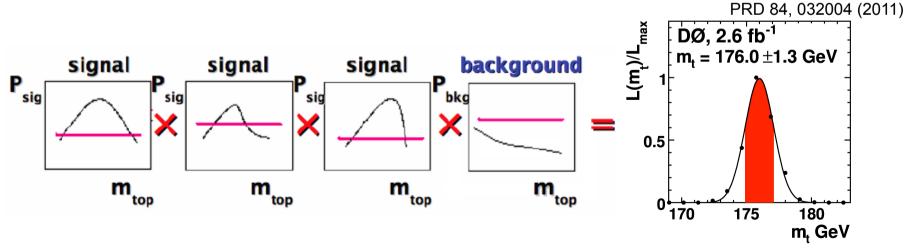
Example: matrix element method

parton distribution functions

$$P(x; m_t) = \frac{1}{\sigma} \int d^n \sigma(y; m_t) dq_1 dq_2 f(q_1) f(q_2) W(x \mid y)$$

differential cross section (LO matrix element)

transfer function: mapping from parton-level variables (y) to reconstructed-level variables (x)

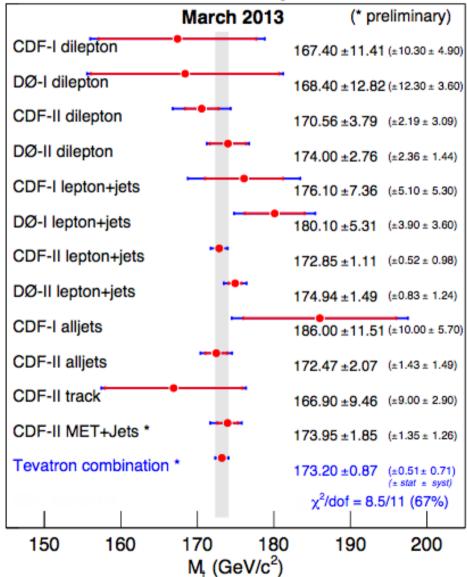


Also incorporate in-situ JES calibration.

### **Tevatron Summary**

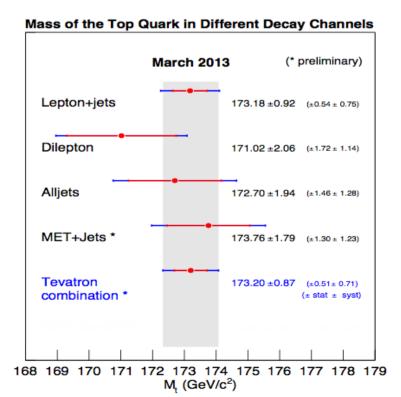
arXiv:1305.3929

#### Mass of the Top Quark



$$m_t = 173.20 \pm 0.51(stat) \pm 0.71(syst) GeV$$
  
 $X^2/NDF = 8.5/11 (prob=67\%)$ 

- 0.5% overall precision! Statistical and systematic uncertainties comparable.
- Consistent measurement across decay channels.



# **Tevatron Summary**

Uncertainty Source	∆m <sub>t</sub> (Gev)
In-situ JES	0.36
Light-jet response (1)	0.16
Light-jet response (2)	0.15
Model of b jets	0.11
Response to b/q/g jets	0.09
Out-of-cone correction	0.01
Total JES	0.44
Signal model	0.52
Jet model	0.08
Lepton model	0.05
Background data	0.13
Background MC	0.06
Multiple interactions	0.07
Method calibration	0.06
Total systematic	0.71
Statistical	0.51
Total	0.87

```
m_t = 173.20 \pm 0.51(stat) \pm 0.71(syst) GeV

X^2/NDF = 8.5/11 (prob=67\%)
```

- 0.5% overall precision! Statistical and systematic uncertainties comparable.
- Dominant systematics:
  - In-situ JES: scales like 1/√L
  - Signal modeling
- Prospects: DØ measurements being updated to the full dataset:
  - Increased statistics by x2 (dileptons),
     x3 (lepton+jets)
  - Alljets measurement for first time.
  - Reduced systematics
- Tevatron combination may reach ultimate precision ~0.7-0.8 GeV.

# **LHC Summary**

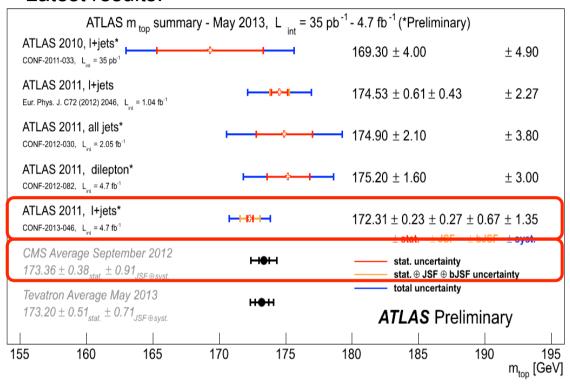
#### July 2012 LHC Combination:

 $m_t = 173.3 \pm 0.5(stat) \pm 1.3(syst) GeV$  $X^2/NDF = 2.5/6 (prob=87\%)$ 

Uncertainty source	δm <sub>top</sub> (GeV)
in-situ JES	0.38
rJES (CMS)	0.06
light-jet response (2)	0.07
model of b-jets	0.68
detector model	0.19
underlying event	0.47
radiation	0.69
color reconnection	0.55
MHI	0.25
lepton model	0.01
background data	0.16
background MC	0.01
method calibration	0.13
statistical	0.47
Total	1.40

ATLAS-CONF-2012-095 CMS PAS TOP-12-001

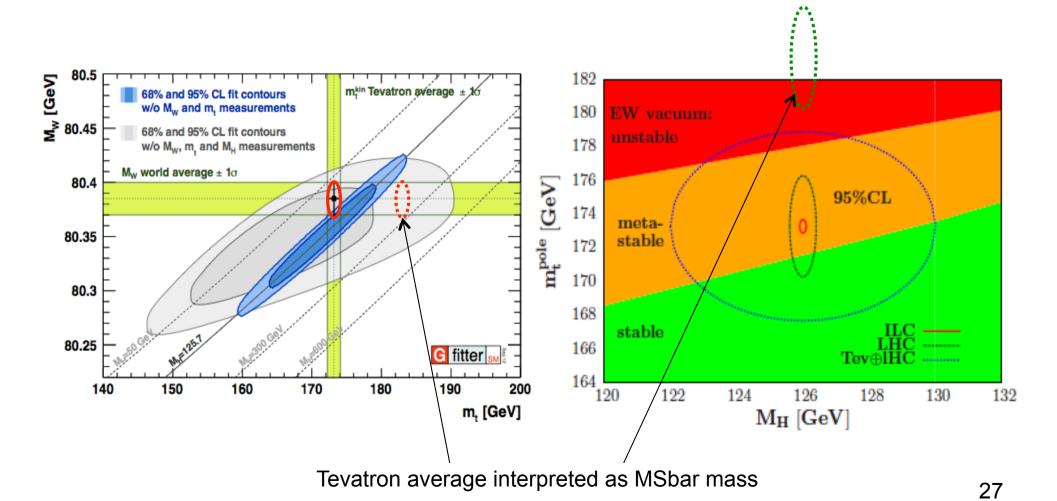
#### Latest results:



- LHC experiments are making fast progress and starting to reach precision competitive with the Tevatron.
- Promising future owing to large available samples, allowing for in-situ x-checks and reduction of systematic uncertainties (e.g. via MC tuning).
- Still, hard to imagine going below 0.5 GeV.

# What Top Quark Mass?

- What's really being measured is the top quark mass in whatever MC generator was used to extract it. But different MCs include different radiative corrections!
- Naively one expects it to be close to the pole mass.
- It's important to know since there is a ~10 GeV shift between pole and MSbar mass!



#### A Partial Check

Compare measured tt cross section

$$\sigma_{
m tar t}=7.56^{+\,0.63}_{-\,0.56}$$
 pb D0 coll. arXiv:1105.5384  $\sigma_{
m tar t}=7.50^{\,+0.48}_{\,-0.48}$  pb CDF coll. CDF-note-9913

to theoretical prediction (approx NNLO) as a function of MSbar mass.

Estimated MSbar mass:

$$m_t^{\overline{\rm MS}}(m_t) = 163.3 \pm 2.7 \,{\rm GeV}$$

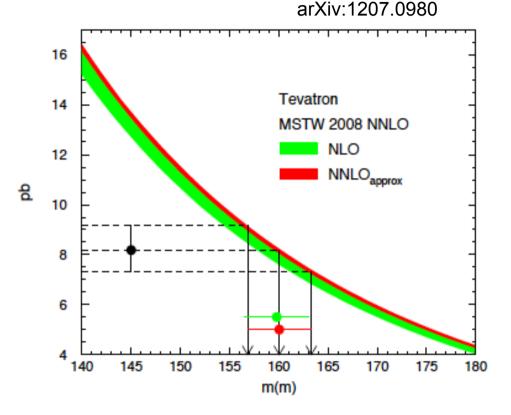
translated to pole mass (conversion known to 3-loops in QCD):

$$m_t^{\rm pole} = 173.3 \pm 2.8 \,{\rm GeV}$$

→ In good agreement with Tevatron average for m<sub>t</sub>!

$$\underline{m}_{\underline{t}}$$
 = 173.20 ± 0.51(stat) ± 0.71(syst) GeV  
X<sup>2</sup>/NDF = 8.5/11 (prob=67%)

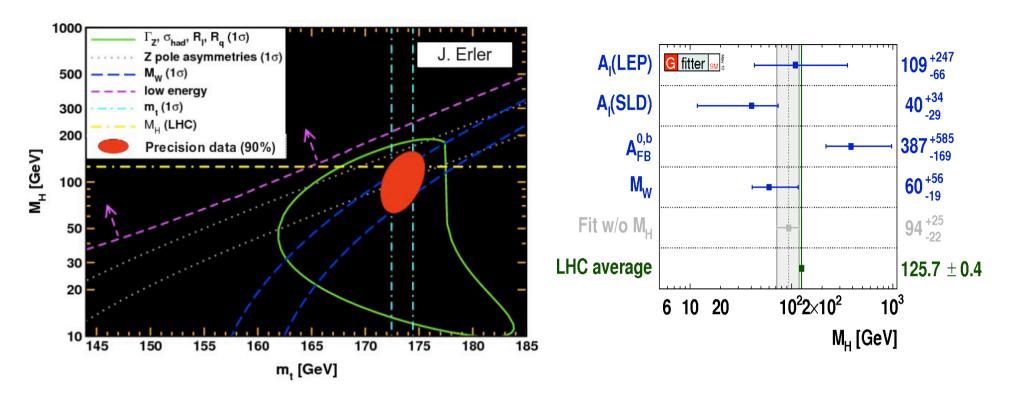
→ supports direct measurements close to pole mass definition but comparison only good to O(3 GeV).



	ABM11	JR09	MSTW08	NN21
$m_t^{\overline{ ext{MS}}}(m_t)$	$162.0^{+2.3}_{-2.3}{}^{+0.7}_{-0.6}$	$163.5^{+2.2}_{-2.2}{}^{+0.6}_{-0.2}$	$163.2^{+2.2+0.7}_{-2.2-0.8}$	$164.4^{+2.2}_{-2.2}{}^{+0.8}_{-0.2}$
$m_t^{ m pole}$	$171.7^{+2.4}_{-2.4}{}^{+0.7}_{-0.6}$	$173.3^{+2.3}_{-2.3}{}^{+0.7}_{-0.2}$	$173.4^{+2.3}_{-2.3}^{+0.8}_{-0.8}$	$174.9^{+2.3}_{-2.3}^{+0.8}_{-0.3}$

# W Boson Mass

#### **Motivation**

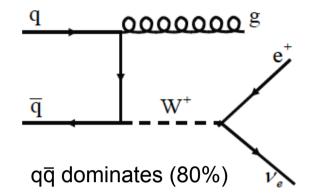


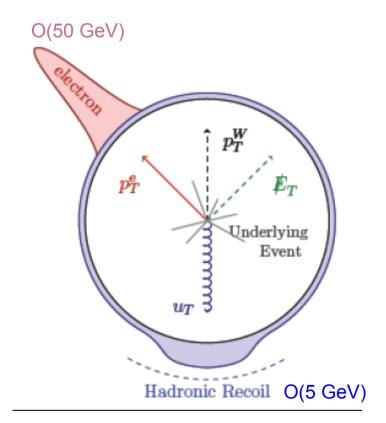
- Most sensitivity to  $M_H$  comes from  $M_W$  and  $\sin^2\theta_{eff}$  (via Z-pole asymmetries).
- For equal contribution to the Higgs mass uncertainty need: ΔM<sub>W</sub>=0.006Δm<sub>t</sub>
  - Latest Tevatron combination: Δm<sub>t</sub>~0.9 GeV → would need ΔM<sub>W</sub>~5 MeV!
  - Before Run II had:  $\Delta M_W \sim 30$  MeV (driven by LEP2 measurements).
- $\rightarrow$  Given current precision on m<sub>t</sub>, progress on  $\Delta M_W$  would have the biggest impact on precision EW fit.

# Measurement: Basic Strategy

#### In principle just need to:

- measure:
  - charged lepton
  - hadronic recoil
- build from MC templates of sensitive kinematic variables:
  - W transverse mass
  - lepton p<sub>T</sub>
  - missing p<sub>T</sub>
- perform a binned likelihood fit to data as a function of M<sub>W</sub>





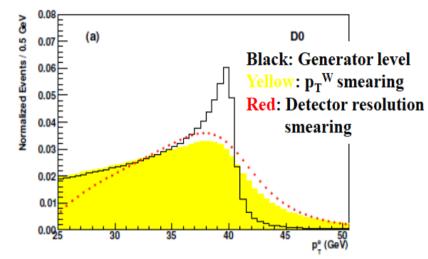
# Measurement: Basic Strategy

O(50 GeV)

Electron shower

#### In practice need to:

- measure:
  - charged lepton to 0.01% precision
  - hadronic recoil to <1% precision
- understand in great detail energy flow in event, driven by multiple physics and instrumental effects that affect each other
- employ state-of-art theoretical predictions for W boson production and decay, including radiative effects
- develop a parameterized fast MC implementing all relevant effects
- build from MC templates of sensitive kinematic variables:
  - W transverse mass
  - lepton p<sub>⊤</sub>
  - missing p<sub>T</sub>
- perform a binned likelihood fit to data as a function of M<sub>w</sub>



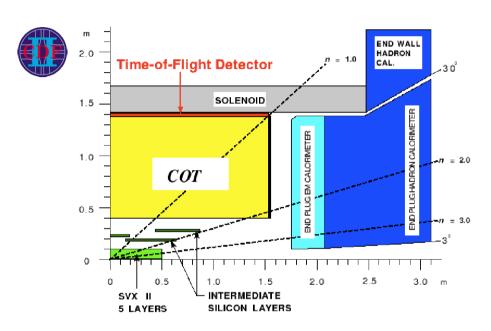
Energy under the electron cone O(100 MeV) In-cone FSR  $p_T^W$ Underlying event O(200 MeV) leakage O(50 MeV) OoC FSR Soft Recoil Min Bias Zero Bias

Hard Recoil O(5 GeV)

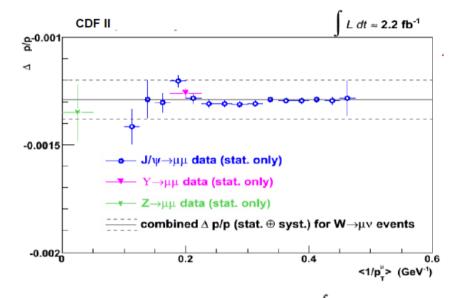
32

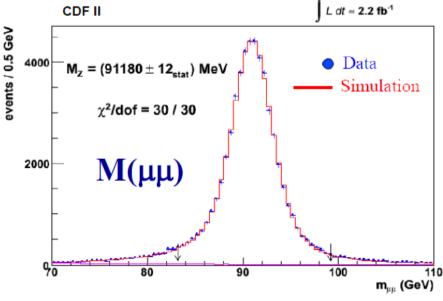
### **Measurement: Different Tactics**

 CDF and DØ employ different tactics, especially in terms of fast MC simulation and calibration, capitalizing on the strengths of their detectors.



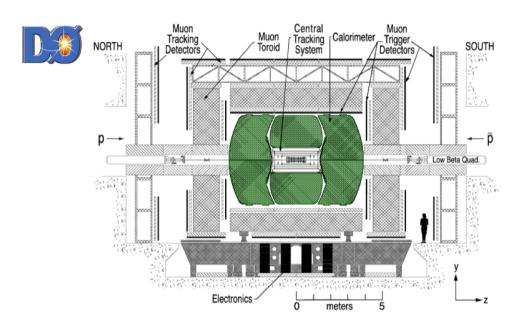
- Central tracking provides very good lepton momentum measurement
  - Muon  $\Delta p_T/p_T \sim 3.2\%$  at  $p_T = 45$  GeV
- Detailed tracker model
- Focus on momentum scale calibration
  - Using  $J/\Psi \rightarrow \mu\mu$ ,  $\Upsilon \rightarrow \mu\mu$ ,  $Z \rightarrow \mu\mu$
- Use E/p peak from W→ev to calibrate EM calorimeter
- → Measurement on W→ev and W→µv events



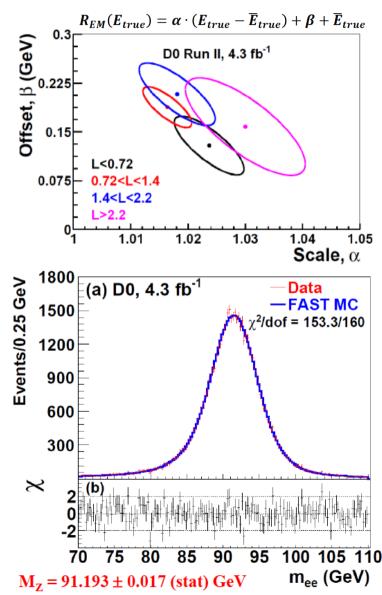


### **Measurement: Different Tactics**

 CDF and DØ employ different tactics, especially in terms of fast MC simulation and calibration, capitalizing on the strengths of their detectors.



- EM calorimeter provides very good electron energy measurement
  - Muon  $\Delta E/E \sim 3.3\%$  at  $E_T = 45$  GeV
- Detailed calorimeter, E-flow model
- Focus on calorimeter energy scale calibration
  - Using Z→ee
- → Measurement on W→ev events



# **Systematic Uncertainties**

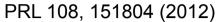
• Breakdown of systematic uncertainties for m<sub>T</sub> measurement:

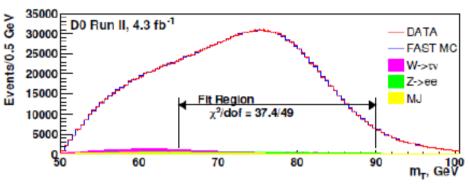
Source	CDF $m_T(\mu, \nu)$	$CDF\ m_T(e, \nu)$	$D Ø \ m_T(e, \nu)$	
Experimental – Statistical power of the calibration sample.				
Lepton Energy Scale	7	10	16	
Lepton Energy Resolution	1	4	2	
Lepton Energy Non-Linearity			4	
Lepton Energy Loss			4	
Recoil Energy Scale	5	5		
Recoil Energy Resolution	7	7		
Lepton Removal	2	3		
Recoil Model			5	
Efficiency Model			1	
Background	3	4	2	
W production and decay model - Not statistically driven.				
PDF	10	10	11	
QED	4	4	7	
Boson $p_T$	3	3	2	
W statistics	12		13	

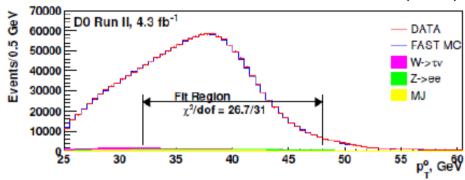
2.2 fb<sup>-1</sup>

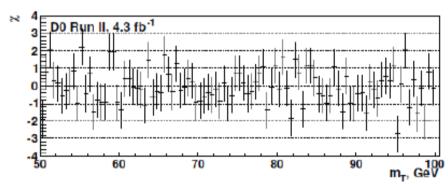
5.3 fb<sup>-1</sup>

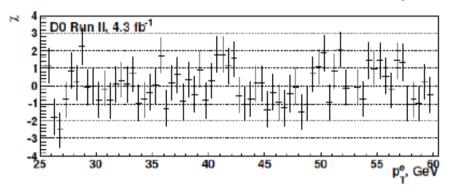
# DØ Results







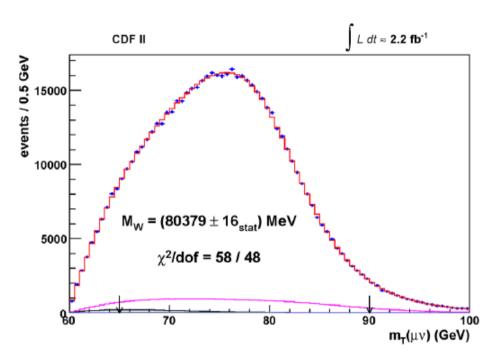


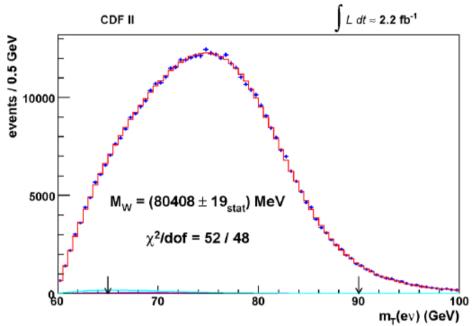


Method $(4.3 fb^{-1})$	$M_{W}$ (MeV)
$m_T(e,  u)$	$80371 \pm 13 ( extsf{stat})$
$p_T(e)$	$80343 \pm 14 (\mathrm{stat})$
$ \not\!\!E_T(e, \nu) $	$80355 \pm 15 ( ext{stat})$
Combination $m_T \oplus p_T \ (4.3  fb^{-1})$	$80367 \pm 26 (\mathrm{syst} + \mathrm{stat})$
Combination $(5.3 fb^{-1})$	$80375 \pm 23 (syst + stat)$

## **CDF** Results

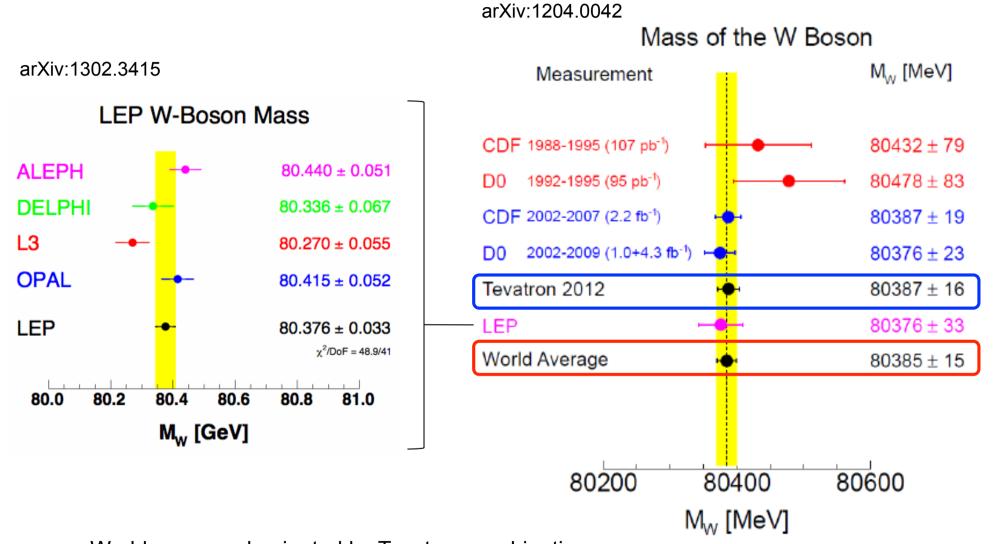
PRL 108, 151803 (2012)





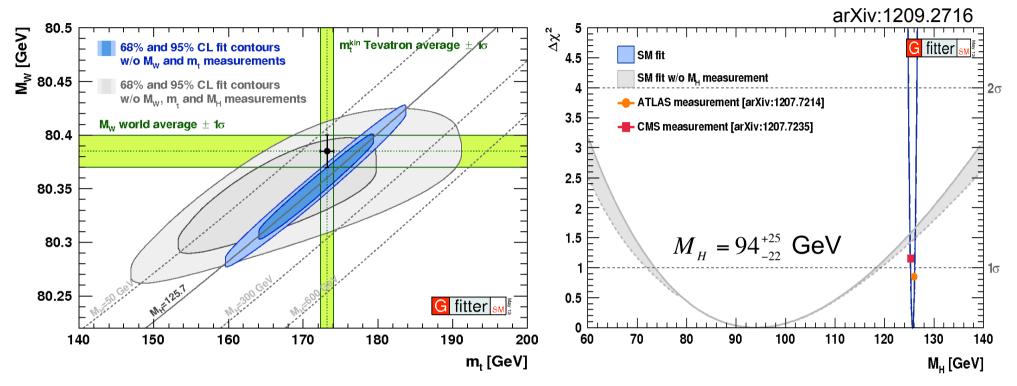
Method $(2.2 fb^{-1})$	$M_W$ (MeV)	Method $(2.2 fb^{-1})$	$M_W$ (MeV)	
$m_T(\mu, \nu)$	$80379 \pm 16 ({\rm stat})$	$m_T(e, \nu)$	$80408 \pm 19 (stat)$	
$p_T(\mu)$	$80348 \pm 18 ({\rm stat})$	$p_T(e)$	$80393 \pm 21 (stat)$	
$E_T(\mu, \nu)$	$80406 \pm 22 ({\sf stat})$	$E_T(e, \nu)$	$80431 \pm 25 (\text{stat})$	
Combination $(2.2 fb^{-1})$		$80387 \pm 19MeV  ext{(syst} +  ext{stat)}$		

#### **New World Average**



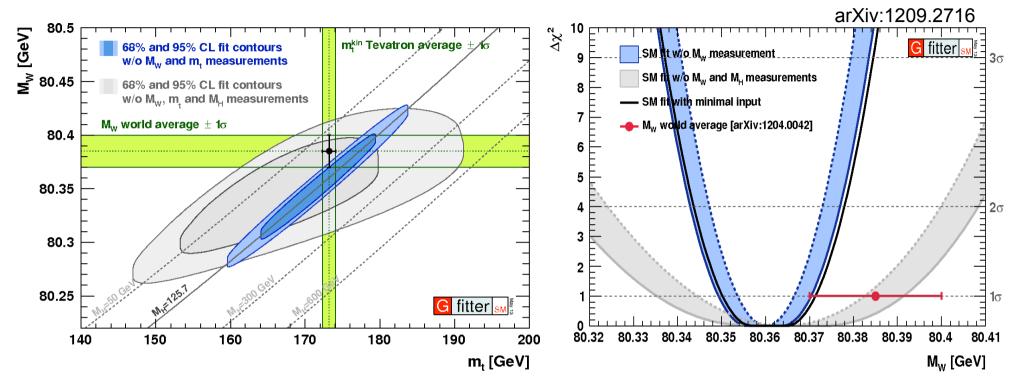
- World average dominated by Tevatron combination.
- PDF uncertainties are now the main limitation.

## Constraints on the Higgs Boson Mass



- Good consistency between direct and indirect measurements of M<sub>W</sub> and m<sub>t</sub> (w/ and w/o direct M<sub>H</sub> measurement).
- Indirect  $M_H$  determination consistent (1.3 $\sigma$ ) with direct measurement at LHC (taken to be  $M_H$ =125.7±0.4 GeV).

#### Constraints on the Higgs Boson Mass



- Good consistency between direct and indirect measurements of M<sub>W</sub> and m<sub>t</sub> (w/ and w/o direct M<sub>H</sub> measurement).
- Conversely, the test can be "turned around" and use electroweak fit, together with direct M<sub>H</sub> measurement, to determine M<sub>W</sub>:

#### **Indirect measurement:**

$$\begin{split} M_{\scriptscriptstyle W} = 80.3593 \pm (0.0056(m_{\scriptscriptstyle t}) \pm 0.0026(M_{\scriptscriptstyle Z}) \pm 0.0018(\Delta\alpha_{\scriptscriptstyle had}) \\ &\pm 0.0017(\alpha_{\scriptscriptstyle S}) \pm 0.0002(M_{\scriptscriptstyle H}) \pm (0.0040(theo)) \text{GeV} \\ = 80.359 \pm (0.011 \text{ GeV}) \end{split}$$

#### **Direct measurement:**

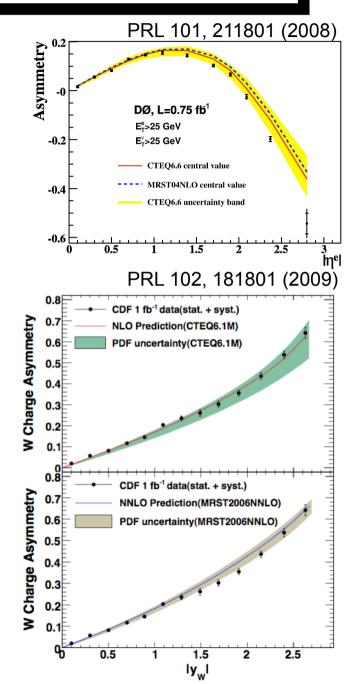
$$M_W = 80.385 \pm 0.015 \text{ GeV}$$

Desirable to further improve it!

#### **Future Prospects**

#### Main improvements expected:

- Use of full 10 fb<sup>-1</sup> dataset (so far 2.2 fb<sup>-1</sup> used by CDF and 5.3 fb<sup>-1</sup> used by DØ).
- Reduction of PDF uncertainties:
  - include in PDF fits measurements at Tevatron and LHC that can constrain PDF uncertainties:
     W charge asymmetry, Z rapidity,...
  - extend η coverage to end-cap calorimeters (DØ): sensitivity to PDF uncertainties introduced through η acceptance cuts.
    - Will also reduce correlation on PDF uncertainty between CDF and DØ.



#### **Future Prospects**

#### Main improvements expected:

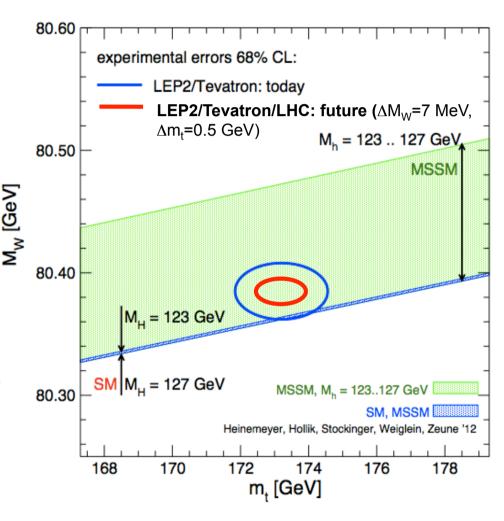
- Use of full 10 fb<sup>-1</sup> dataset (so far 2.2 fb<sup>-1</sup> used by CDF and 5.3 fb<sup>-1</sup> used by DØ).
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     W charge asymmetry, Z rapidity,...
  - extend η coverage to end-cap calorimeters (DØ): sensitivity to PDF uncertainties introduced through η acceptance cuts.
    - Will also reduce correlation on PDF uncertainty between CDF and DØ.

#### **Projected M<sub>W</sub> uncertainty:**

CDF: 10 MeV

DØ: 15 MeV

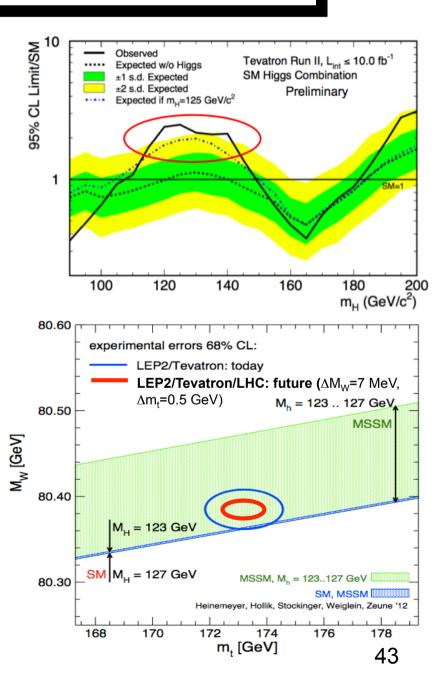
Combination: <10 MeV</li>



## **Summary and Outlook**

- Latest Tevatron results based on full Run II dataset in most search channels.
- Tevatron has achieved SM sensitivity over most of the accessible mass range.
- Excess in 115<m<sub>H</sub><140 GeV region with local p-value at m<sub>H</sub>=125 GeV corresponding to 3.0σ significance. So far emerging picture consistent with discovered boson at the LHC.
- Early results on Higgs couplings and spin/parity are consistent with a SM Higgs boson. Tevatron combination on spin/parity tests upcoming.
- Exciting prospects for precise measurements of m<sub>t</sub> and M<sub>W</sub> at the Tevatron, which will be key to capitalizing on precision measurement in the Higgs sector.

# Stay tuned!

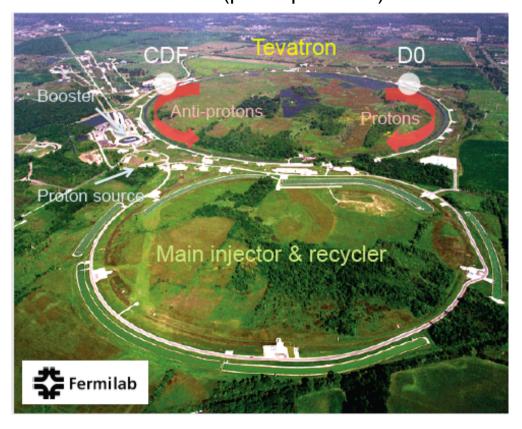


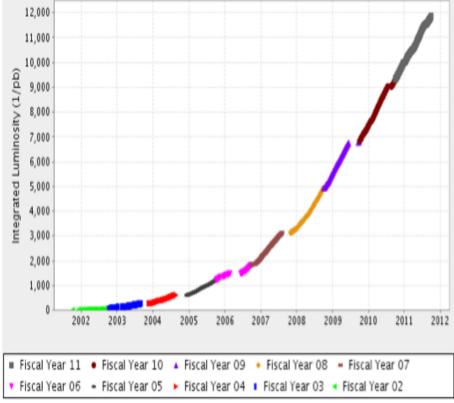
# Backup

#### The Tevatron

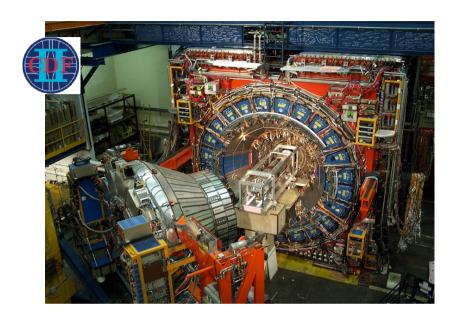
- Proton-antiproton collider at √s=1.96 TeV
  - Tevatron accelerator: 6.5 km circumference
  - Two general-purpose experiments: CDF and DØ
  - 10-year long Run II ended Sept. 30<sup>th</sup>, 2011
  - Total integrated luminosity delivered in Run II:
     ~12 fb<sup>-1</sup> (per experiment)







#### CDF and DØ Detectors





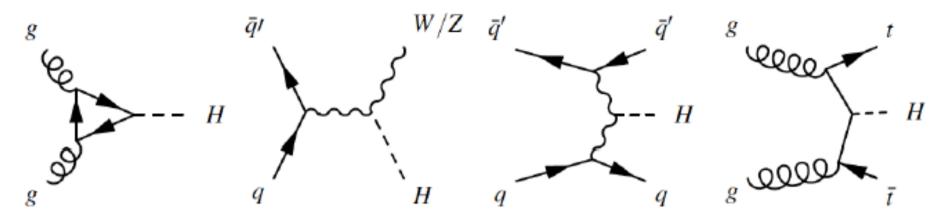
- Central tracking system embedded in a solenoidal magnetic field:
  - Silicon vertex detector
  - Tracking chamber (CDF)
     Fiber tracker (DØ)
- Preshowers
- Electromagnetic and hadronic calorimeters
- Muon system

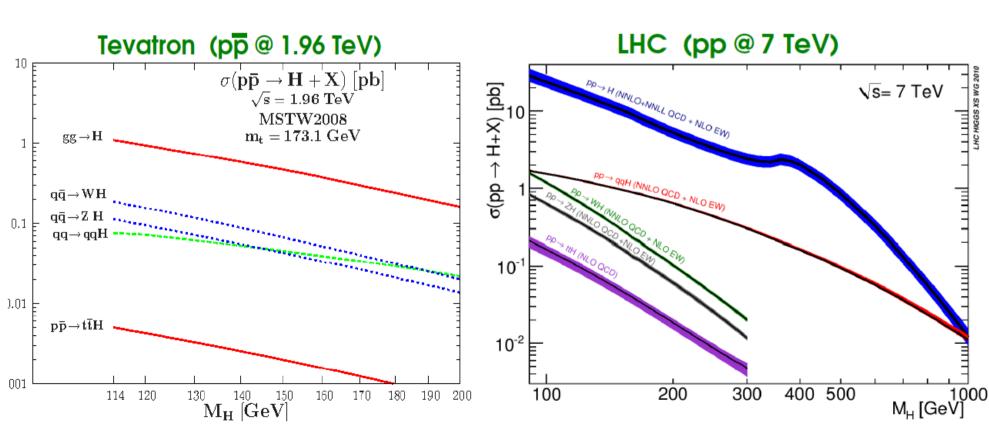


- Include multi-level trigger systems to efficiently select events with topologies of interest.
- Data taking efficiency: ~90%.
- Recorded up to 10 fb<sup>-1</sup> in Run II.

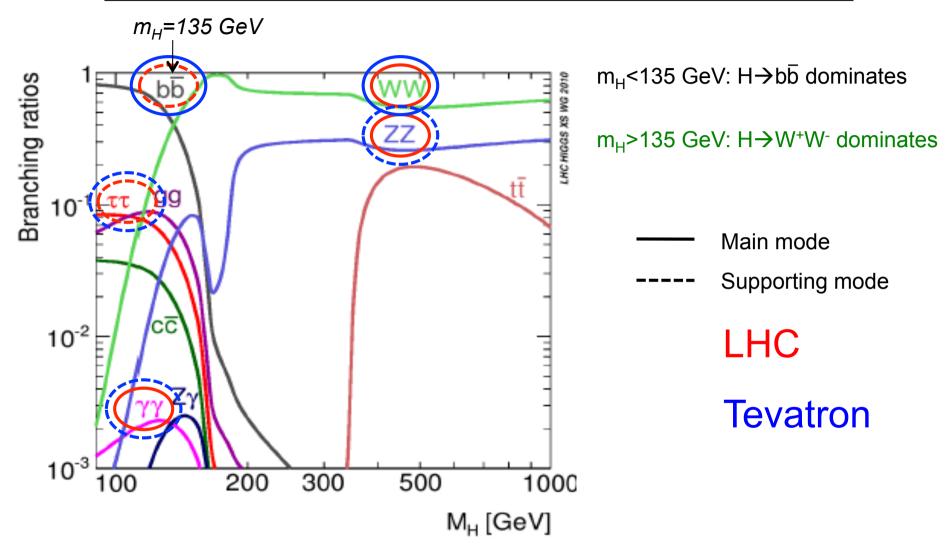
All main search channels used full dataset

## SM Higgs Production at Hadron Colliders





## **SM Higgs Decay Modes**

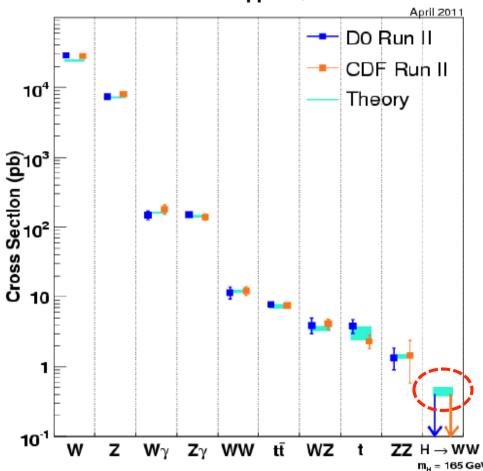


→ Many decay modes being explored to increase the sensitivity of the search to the SM Higgs boson, but also to a non-SM one!

## The Stairway to the Higgs

Higgs boson searches at the Tevatron are background-dominated.

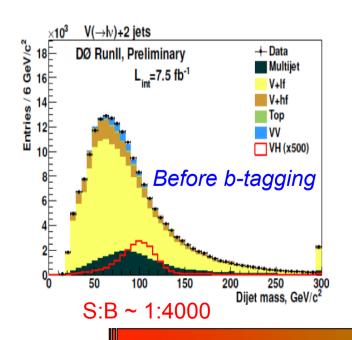


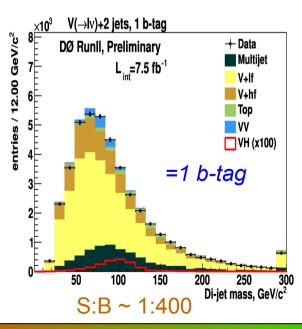


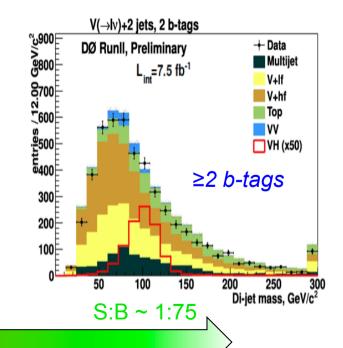
- Instrumental backgrounds: measured directly from data
  - QCD multijet production with mismeasured jets leadings to missing transverse energy or jets misidentified as leptons.
- Physics backgrounds: estimated using simulation and state-of-art theoretical predictions, and further calibrated to data whenever possible
  - W/Z+jets production (w/ real or misidentified heavy flavor jets)
  - Diboson production
  - Double and single top quark production

# Searching for H→bb̄

- Highest sensitivity channel at the Tevatron for m<sub>H</sub><130 GeV.</li>
- Identify events consistent with leptonic W/Z decays in association with jets. Require b-tagging to significantly improve the S/B ratio.
- Main backgrounds: W+jets, tt̄
- Most important discriminant variable: dijet mass.

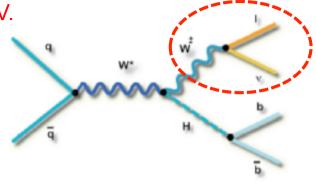




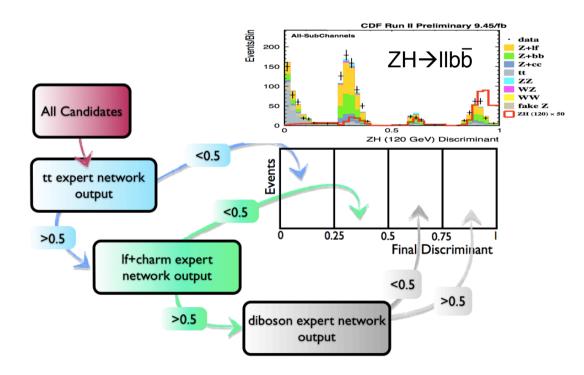


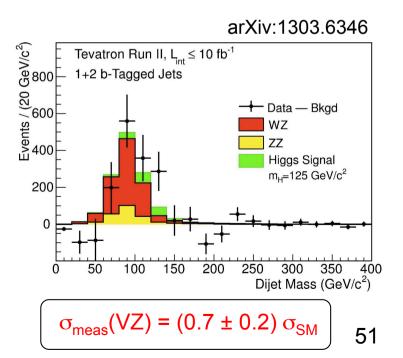
# Searching for H→bb̄

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- Identify events consistent with leptonic W/Z decays in association with jets. Require b-tagging to significantly improve the S/B ratio.
- Main backgrounds: W+jets, tt̄
- Most important discriminant variable: dijet mass.
- Improve sensitivity via sophisticated multivariate techniques.
- Validate search strategy by measuring WZ/ZZ (Z→bb̄).



See talk by XXX

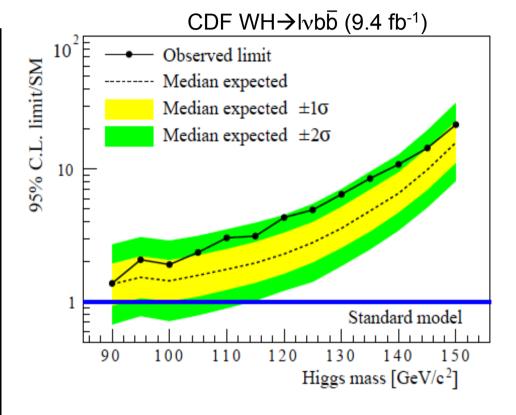




# Summary of H→bb̄ Results

95% CL Limits at  $m_H = 125 \text{ GeV}$ 

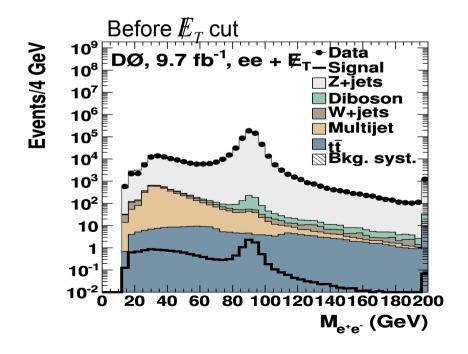
	Channel	Exp/obs Limit (σ/SM)
	WH→I√bb̄ (9.4 fb-1)	2.8/4.9
-	ZH→vvb̄b (9.4 fb⁻¹)	3.3/3.1
	ZH→I <sup>+</sup> I <sup>-</sup> b̄b̄ (9.4 fb <sup>-1</sup> )	3.6/7.2
	WH→l√bb̄ (9.7 fb <sup>-1</sup> )	4.7/5.2
<b>)</b>	ZH→vvbb̄ (9.5 fb-1)	3.9/4.3
	ZH→I <sup>+</sup> I <sup>-</sup> b̄b̄ (9.7 fb <sup>-1</sup> )	5.1/7.1
	VH/VBF→jjbb̄ (9.4 fb <sup>-1</sup> )	11.0/9.0
	ttH→I+jets (9.4 fb <sup>-1</sup> )	12.4/17.6
	ttH→jets (5.7 fb <sup>-1</sup> )	26.2/36.2

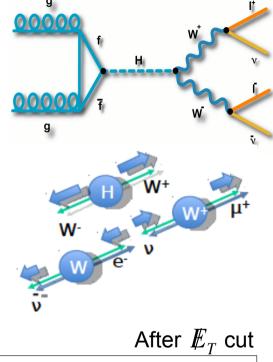


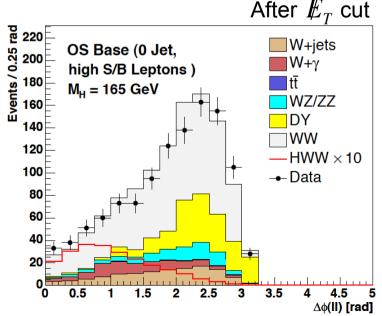
- Limits from individual VH, H→bb̄ channels at ~2-3xSM at m<sub>H</sub>=125 GeV and quickly degrading towards high mass.
- Important to consider additional channels with different mass dependence.

## Searching for H→WW→IvIv

- Highest sensitivity channel in m<sub>H</sub>~130-200 GeV range.
- Clean dilepton+ $I\!\!E_T$ signature.
- Main backgrounds after E<sub>T</sub> cut: WW, W/Z+jets, Wγ
- After final selection expect (m<sub>H</sub>=165 GeV):
  - ~7 signal events/fb<sup>-1</sup>/experiment with S:B~1:50-1:100
- Exploit spin correlation between dibosons.
  - → Small angular separation between leptons



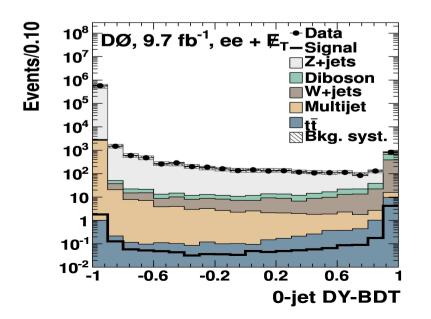


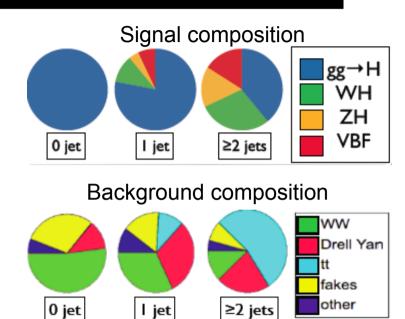


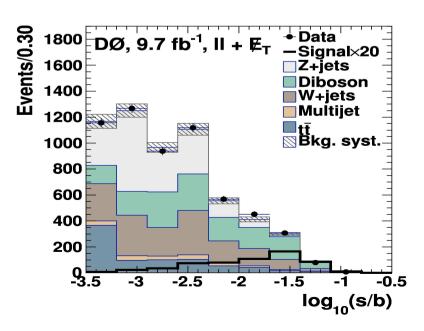
#### Searching for H→WW→IvIv

To increase the sensitivity:

- Consider all signal production modes.
- Categorize events in channels with different S:B (by jet multiplicity, lepton flavor/quality, etc) and combine at the end.
- Build multivariate discriminants against:
  - Z/γ\*→I<sup>+</sup>I<sup>-</sup> ((DØ): increase signal acceptance in event selection.
  - · Remaining background: final discriminant.







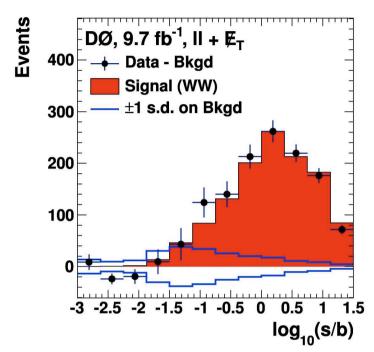
54

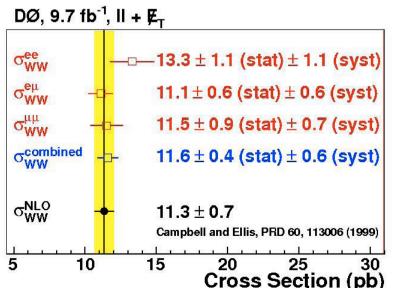
## Searching for H→WW→IvIv

To increase the sensitivity:

- Consider all signal production modes.
- Categorize events in channels with different S:B (by jet multiplicity, lepton flavor/quality, etc) and combine at the end.
- Build multivariate discriminants against:
  - Z/γ\*→I<sup>+</sup>I<sup>-</sup> ((DØ): increase signal acceptance in event selection.
  - Remaining background: final discriminant.
- Include searches for same-sign dilepton and trilepton final states.

Validate search strategy by measuring WW→lvlv.

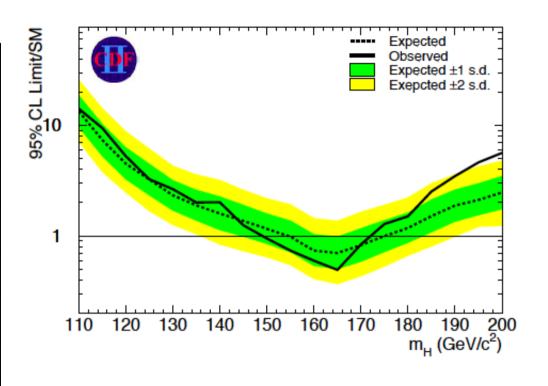




## Low Mass Results from H→WW,ττ,γγ

95% CL Limits at  $m_H = 125 \text{ GeV}$ 

Channel	Exp/obs Limit (σ/SM)
H→WW→IvIv (9.7 fb <sup>-1</sup> )	3.2/3.3
H→WW→IvIv (9.7 fb <sup>-1</sup> )	3.4/4.1
H+X→ττ+jets (8.3 fb <sup>-1</sup> )	14.8/11.7
H+X→ττjj (9.7 fb <sup>-1</sup> )	9.0/11.3
VH→ττl(I) (6.2 fb <sup>-1</sup> )	23.3/26.5
VH→ττμ (8.6 fb <sup>-1</sup> )	13.0/19.4
H→γγ (10.0 fb <sup>-1</sup> )	11.7/20.5
H→γγ (9.7 fb <sup>-1</sup> )	8.5/12.7

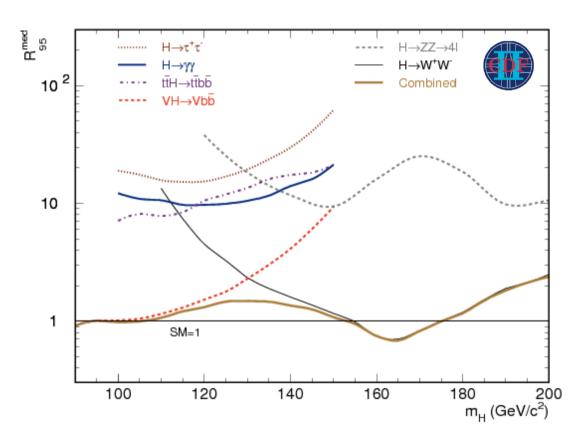


Additional channels contribute useful sensitivity at low/intermediate m<sub>H</sub>:

- $H \rightarrow WW \rightarrow l_V l_V$ : improving towards high  $m_H$ .
- H+X→ττjj, H→γγ: ~flat vs m<sub>H</sub>.

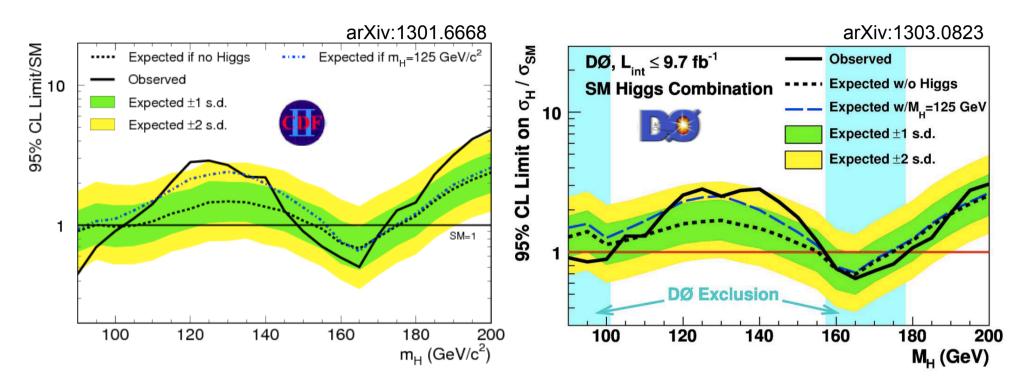
## Combined Limits on SM Higgs Production

Combination of multiple channels (and experiments!) yields the greatest sensitivity.



- Assumes SM prediction for ratio of production cross sections and branching ratios.
- More than 200 different sources of systematic uncertainties are considered (including correlations among channels and experiments), and constrained in sidebands.
- Use different techniques to cross check calculations (Bayesian, modified frequentist)
  - → results agree within ≤5%.

#### CDF and DØ Individual Results



Observed 95% CL exclusion:

 $90 < m_H < 102 \text{ GeV}, 152 < m_H < 172 \text{ GeV}$ 

At 
$$m_H = 125 \text{ GeV}$$
:

Exp. limit: 1.5 x SM Obs. limit: 2.9 x SM

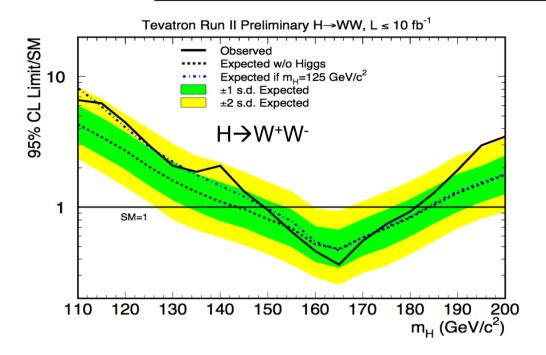
Observed 95% CL exclusion:

$$90 < m_H < 101 \text{ GeV}, 157 < m_H < 178 \text{ GeV}$$

At  $m_H = 125 \text{ GeV}$ :

Exp. limit: 1.7 x SM Obs. limit: 2.9 x SM

## Results by Channel



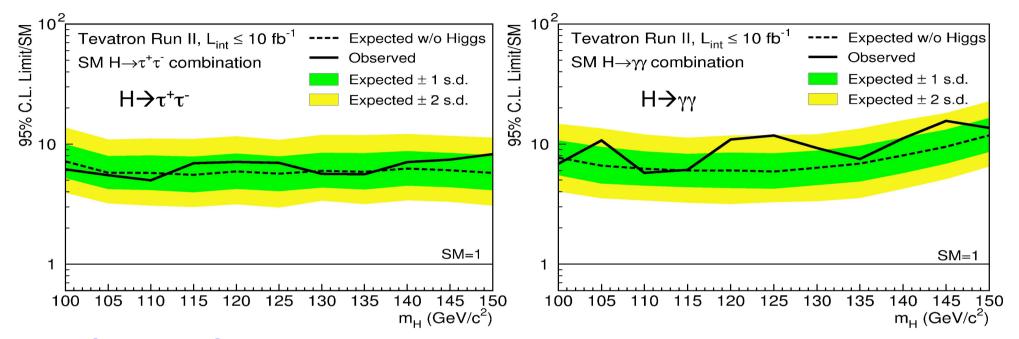
#### $H \rightarrow \tau^+\tau^-$ and $H \rightarrow \gamma\gamma$ :

- Expected sensitivity at m<sub>H</sub>~125 GeV of ~6-7xSM.
- No significant excess.

#### H→WW:

- Expected sensitivity at m<sub>H</sub>~125 GeV of ~2.1xSM.
- Very broad excess, consistent with both B-only and S+B hypotheses (m<sub>H</sub>=125 GeV).

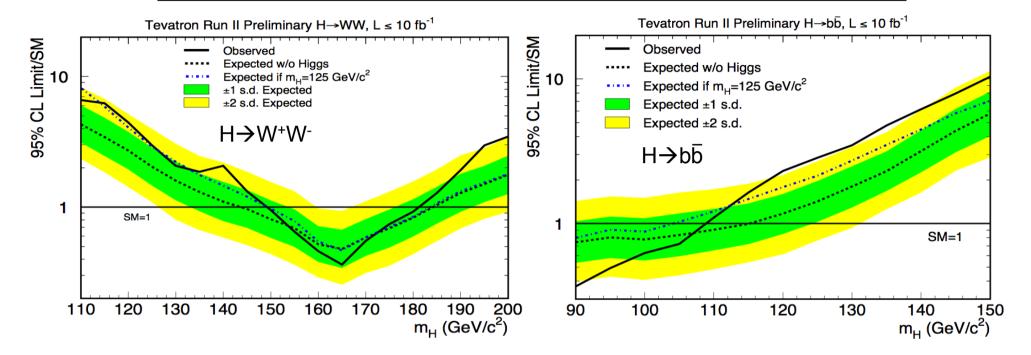
## Results by Channel



#### $H \rightarrow \tau^+\tau^-$ and $H \rightarrow \gamma\gamma$ :

- Expected sensitivity at m<sub>H</sub>~125 GeV of ~6-7xSM.
- No significant excess.

#### Results by Channel



#### $H \rightarrow \tau^+ \tau^-$ and $H \rightarrow \gamma \gamma$ :

- Expected sensitivity at m<sub>H</sub>~125 GeV of ~6-7xSM.
- No significant excess.

#### H→WW:

- Expected sensitivity at m<sub>H</sub>~125 GeV of ~2.1xSM.
- Very broad excess, consistent with both B-only and S+B hypotheses (m<sub>H</sub>=125 GeV).

#### H→bō:

- Expected sensitivity at m<sub>H</sub>~125 GeV of ~1.4xSM.
- Broad excess, somewhat above S+B hypothesis (m<sub>H</sub>=125 GeV).

## **Probing Higgs Boson Couplings**

- Several production and decay mechanisms contribute to signal rates per channel
  - → interpretation is difficult
- A better option: measure deviations of couplings from the SM prediction (*arXiv:1209.0040*). Basic assumptions:
  - there is only one underlying state at m<sub>H</sub>~125 GeV,
  - it has negligible width,
  - it is a CP-even scalar (only allow for modification of coupling strengths, leaving the Lorentz structure of the interaction untouched).

Additional assumption made in this study:

- no additional invisible or undetected Higgs decay modes.
- Under these assumptions all production cross sections and branching ratios can be expressed in terms of a few common multiplicative factors to the SM Higgs couplings. Examples:

$$\sigma(gg \to H)BR(H \to WW) = \sigma_{SM}(gg \to H)BR_{SM}(H \to WW) \frac{\kappa_g^2 \kappa_W^2}{\kappa_H^2}$$

$$\sigma(WH)BR(H \to bb) = \sigma_{SM}(WH)BR_{SM}(H \to bb) \frac{\kappa_W^2 \kappa_b^2}{\kappa_H^2}$$

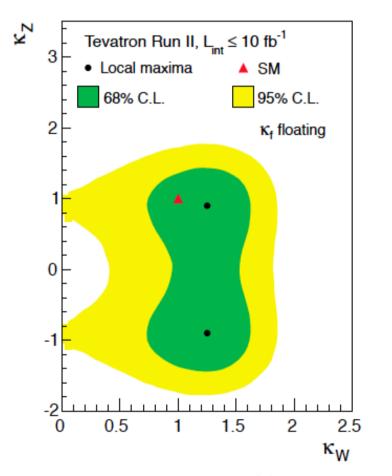
$$\kappa_g = f(\kappa_t, \kappa_b, M_H)$$

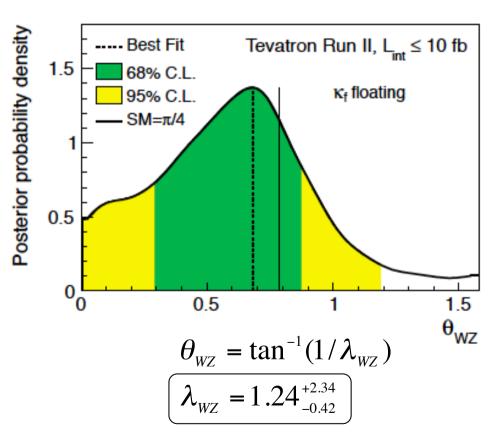
$$\kappa_H = f'(\kappa_t, \kappa_b, \kappa_\tau, \kappa_W, \kappa_Z, M_H)$$

## **Probing Higgs Boson Couplings**

#### • Benchmark I:

- Assume all fermion couplings are scaled by a common free parameter (κ<sub>F</sub>).
- Measure the parameters  $\kappa_W$  and  $\kappa_Z$  independently.
- Probe SU(2)<sub>V</sub> custodial symmetry by measuring the ratio  $\lambda_{WZ} = \kappa_W / \kappa_Z$ .

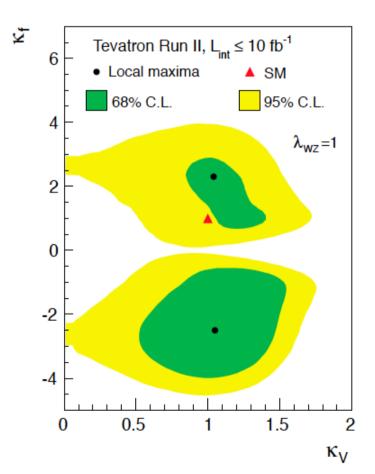


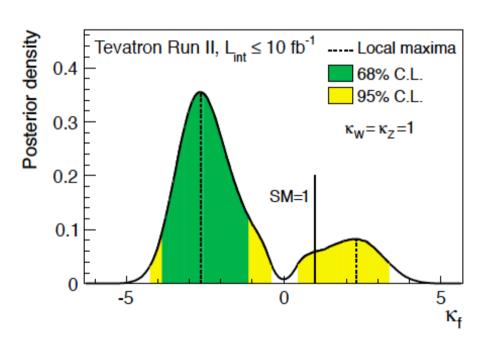


## **Probing Higgs Boson Couplings**

#### Benchmark II:

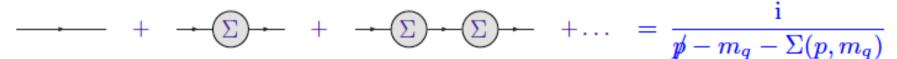
- Consider two independent multiplicative factors: common to all couplings to vector bosons ( $\kappa_V$ ) and common to all couplings to fermions ( $\kappa_f$ ). Assume  $\lambda_{WZ}$ =1.
- Measure  $\kappa_f$  and  $\kappa_V$  simultaneously.





Parameters of Lagrangian have no unique physical interpretation.
 Radiative corrections require definition of renormalization scheme.

Ex. Heavy-quark self-energy:



- Pole Mass:
  - Based on (unphysical) concept of top quark being a free parton.
  - Coincides with pole of propagator at each order:

$$\not p - m_q - \Sigma(p, m_q) \bigg|_{\not p = m_q} \rightarrow \not p - m_q^{
m pole}$$

 Definition of pole mass ambiguous up to corrections O(Λ<sub>QCD</sub>).
 Bound from lattice QCD:
 Δm<sub>t</sub>≥200 MeV

- MSbar Mass:
  - Based on 1-loop minimal subtraction

$$\delta m_q^{(1)} = m_q \frac{\alpha_s}{4\pi} \, 3C_F \, \left( \frac{1}{\epsilon} - \gamma_E + \ln 4\pi \right)$$

 MSbar scheme induces scale dependence: m(μ)

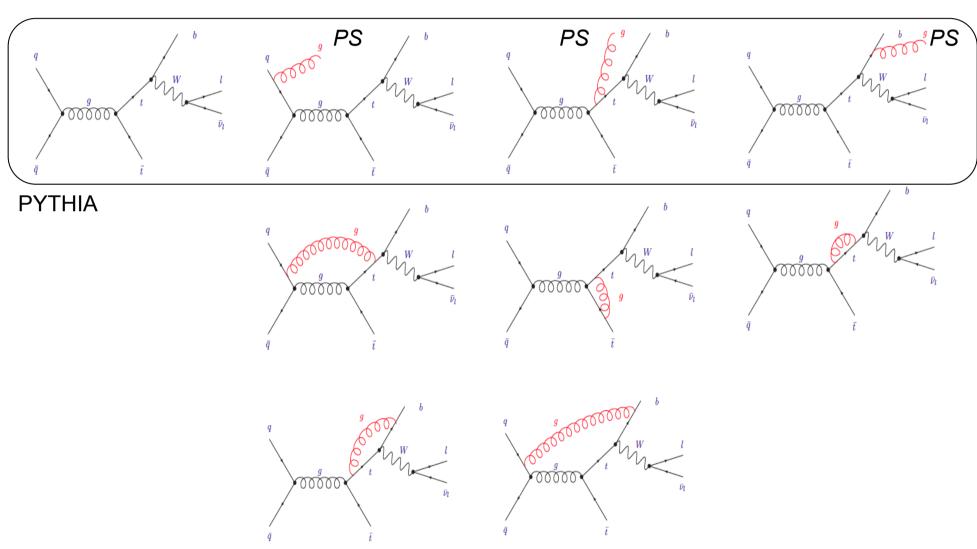
$$\left(\mu^2 \frac{\partial}{\partial \mu^2} + \beta(\alpha_s) \frac{\partial}{\partial \alpha_s}\right) m(\mu) = \gamma(\alpha_s) m(\mu)$$

Conversion between pole and MSbar schemes known to 3-loops in QCD.

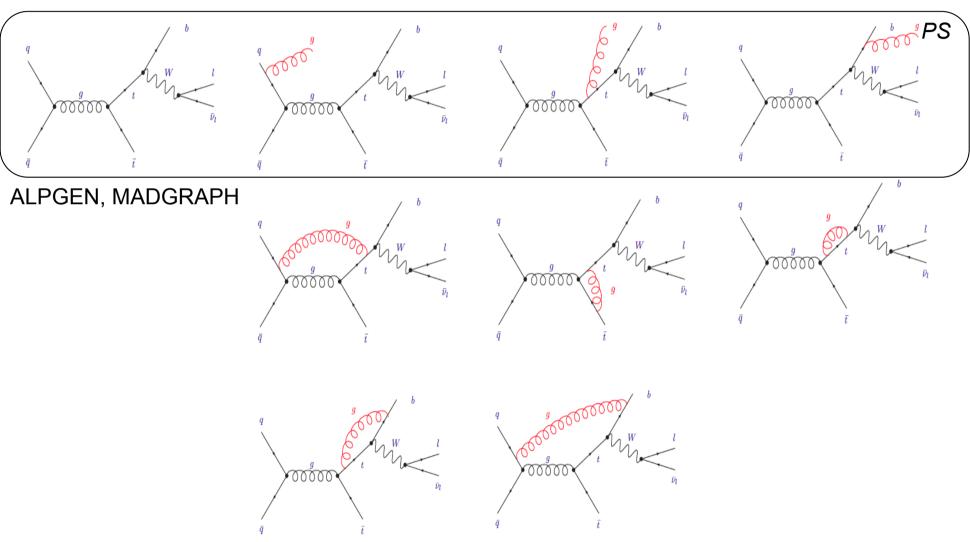
Example: 1-loop QCD

$$m^{\text{pole}} = m(\mu) \left\{ 1 + \frac{\alpha_s(\mu)}{4\pi} \left( \frac{4}{3} + \ln\left(\frac{\mu^2}{m(\mu)^2}\right) \right) + \dots \right\}$$
 65

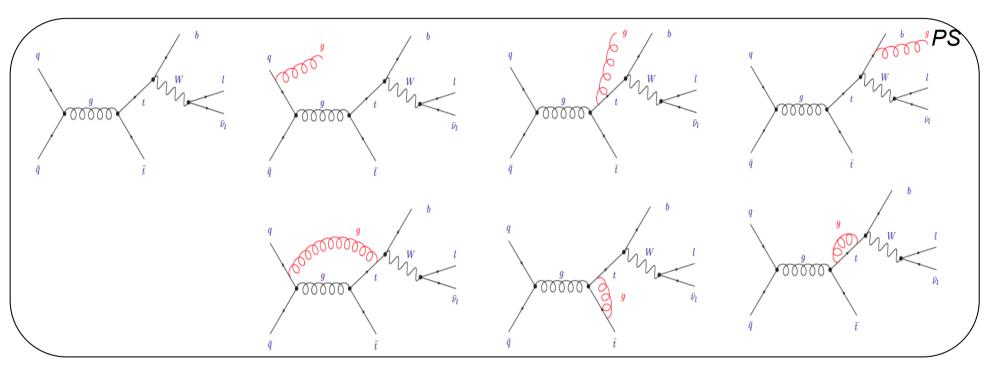
• What's really being measured is the top quark mass in whatever MC generator was used to extract it. But different MCs include different radiative corrections!



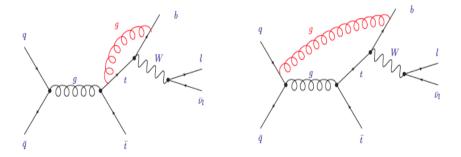
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MC@NLO, POWHEG



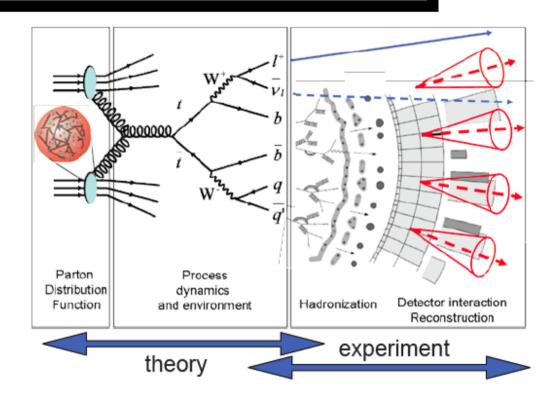
## **Systematic Uncertainties**

#### Theoretical:

- Signal simulation (PDFs, MC generator, hadronization model)
- Event modeling and environment (underlying event, color reconnection, QCD radiation, pileup)

#### Experimental:

- Physics objects and detector modeling (jet reconstruction/ resolution/scale, b-tagging, E<sub>T</sub><sup>miss</sup>,...)
- Background contamination
- Can exploit data to verify proper modeling and/or further constrain size of systematic uncertainties.



# Projected DØ Uncertainty

Source	Public. 2009 (1.0 fb <sup>-1</sup> )	Public. 2012 $(4.3 \text{ fb}^{-1})$	Proj. 10 fb <sup>-1</sup>	Proj. $10 \text{ fb}^{-1} \text{ improv.}$	Proj. 10 fb <sup>-1</sup> improv. $+$ EC
Statistical	23	13	9	9	8
Experimental syst.					
Electron energy scale	34	16	11	11	10
Electron energy resolution	2	<b>2</b>	2	2	2
EM shower model	4	4	4	2	2
Electron energy loss	4	4	4	2	2
Hadronic recoil	6	5	3	3	2
Electron ID efficiency	5	1	1	1	1
Backgrounds	2	2	2	2	2
Subtotal experimental syst.	35	18	13	12	11
W production					
and decay model					
PDF	9	11	11	11	5
$\mathbf{Q}$ ED	7	7	7	3	3
boson $p_T$	2	2	2	2	2
Subtotal $W$ model	12	13	13	12	6
Total systematic uncert.	37	22	19	17	13
Total	44	26	21	19	15

combination: 23

# Projected CDF Uncertainty

Source	0.2/fb (M eV)	2.2/fb (MeV)	10/fb (MeV)
Lepton energy scale	23	7	3
Lepton energy resolution	4	2	1
Recoil energy scale	8	4	2
Lepton removal	6	2	1
Backgrounds	6	3	2
pT(W) model	4	5	2
PDFs	11	10	5
QED radiation	10	4	4
Total systematics	34	15	8
W statistics	34	12	6
Total	48	19	10