

# Robust Determination of the Higgs Couplings: Power to the Data

Juan González Fraile

Universitat de Barcelona

Tyler Corbett, O. J. P. Éboli, J. G-F and M. C. Gonzalez-Garcia

arXiv:1207.1344, 1211.4580, 1304.1151

<http://hep.if.usp.br/Higgs>

# Overview

Discovery of a  $\simeq 125$  GeV "Higgs-like" particle  $\rightarrow$  EWSB direct exploration:

- Spin
- Parity
- EWSB connected new states
- Couplings

Bottom-up model-independent effective Lagrangian approach:

$$\mathcal{L}_{\text{eff}} = \sum_n \frac{f_n}{\Lambda^2} \mathcal{O}_n$$

- $\mathcal{L}_{\text{eff}}$ : describe the low energy effects of new physics in the couplings of this observed new state in the coefficients of dimension-6 operators.
- Assume observed state is light electroweak doublet scalar and that  $SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$  is linearly realized in the effective theory.
- Choice of operators and basis  $\rightarrow$  **Driven by the data**
- Complementarity of experimental searches  $\rightarrow$  TGV  $\leftrightarrow$  Higgs

Determine coefficients of operators using all available data: Tevatron, LHC, TGV, EWPD.

# Overview

Discovery of a  $\simeq 125$  GeV "Higgs-like" particle  $\rightarrow$  EWSB direct exploration:

- Spin
- Parity
- EWSB connected new states
- **Couplings**

Bottom-up model-independent effective Lagrangian approach:

$$\mathcal{L}_{\text{eff}} = \sum_n \frac{f_n}{\Lambda^2} \mathcal{O}_n$$

- $\mathcal{L}_{\text{eff}}$ : describe the low energy effects of new physics in the couplings of this observed new state in the coefficients of dimension-6 operators.
- Assume observed state is light electroweak doublet scalar and that  $SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$  is linearly realized in the effective theory.
- Choice of operators and basis  $\rightarrow$  **Driven by the data**
- Complementarity of experimental searches  $\rightarrow$  TGV  $\leftrightarrow$  Higgs

Determine coefficients of operators using all available data: Tevatron, LHC, TGV, EWPD.

# Effective Lagrangian

$$\mathcal{L}_{\text{eff}} = \sum_n \frac{f_n}{\Lambda^2} \mathcal{O}_n$$

Our assumptions are:

- The observed state belongs to a  $SU(2)$  doublet.
- The state is CP-even as in SM.
- Narrow resonance and no overlapping resonances.
- $SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$  SM local symmetry, C and P even, lepton and baryon number conservation

59 dimension-6 operators are enough...<sup>1</sup>

Set reduced by considering only C and P even and EOM to eliminate/choose the basis

$$2\mathcal{O}_{\Phi,2} - 2\mathcal{O}_{\Phi,4} = \sum_{ij} \left( y_{ij}^e \mathcal{O}_{e\Phi,ij} + y_{ij}^u \mathcal{O}_{u\Phi,ij} + y_{ij}^d (\mathcal{O}_{d\Phi,ij})^\dagger + \text{h.c.} \right) ,$$

$$2\mathcal{O}_B + \mathcal{O}_{WB} + \mathcal{O}_{BB} + g'^2 \left( \mathcal{O}_{\Phi,1} - \frac{1}{2} \mathcal{O}_{\Phi,2} \right) = -\frac{g'^2}{2} \sum_i \left( -\frac{1}{2} \mathcal{O}_{\Phi L,ii}^{(1)} + \frac{1}{6} \mathcal{O}_{\Phi Q,ii}^{(1)} - \mathcal{O}_{\Phi e,ii}^{(1)} + \frac{2}{3} \mathcal{O}_{\Phi u,ii}^{(1)} - \frac{1}{3} \mathcal{O}_{\Phi d,ii}^{(1)} \right)$$

$$2\mathcal{O}_W + \mathcal{O}_{WB} + \mathcal{O}_{WW} + g^2 \left( \mathcal{O}_{\Phi,4} - \frac{1}{2} \mathcal{O}_{\Phi,2} \right) = -\frac{g^2}{4} \sum_i \left( \mathcal{O}_{\Phi L,ii}^{(3)} + \mathcal{O}_{\Phi Q,ii}^{(3)} \right) .$$

<sup>1</sup>Buchmuller & Wyler; Grzadkowski et al. arXiv: 1008.4884

# Effective Lagrangian

$$\mathcal{L}_{\text{eff}} = \sum_n \frac{f_n}{\Lambda^2} \mathcal{O}_n$$

Our assumptions are:

- The observed state belongs to a  $SU(2)$  doublet.
- The state is CP-even as in SM.
- Narrow resonance and no overlapping resonances.
- $SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$  SM local symmetry, C and P even, lepton and baryon number conservation

**59** dimension-6 operators are enough...<sup>1</sup>

Set reduced by considering only C and P even and EOM to eliminate/choose the basis

$$2\mathcal{O}_{\Phi,2} - 2\mathcal{O}_{\Phi,4} = \sum_{ij} \left( y_{ij}^e \mathcal{O}_{e\Phi,ij} + y_{ij}^u \mathcal{O}_{u\Phi,ij} + y_{ij}^d (\mathcal{O}_{d\Phi,ij})^\dagger + \text{h.c.} \right) ,$$

$$2\mathcal{O}_B + \mathcal{O}_{WB} + \mathcal{O}_{BB} + g'^2 \left( \mathcal{O}_{\Phi,1} - \frac{1}{2} \mathcal{O}_{\Phi,2} \right) = -\frac{g'^2}{2} \sum_i \left( -\frac{1}{2} \mathcal{O}_{\Phi L,ii}^{(1)} + \frac{1}{6} \mathcal{O}_{\Phi Q,ii}^{(1)} - \mathcal{O}_{\Phi e,ii}^{(1)} + \frac{2}{3} \mathcal{O}_{\Phi u,ii}^{(1)} - \frac{1}{3} \mathcal{O}_{\Phi d,ii}^{(1)} \right)$$

$$2\mathcal{O}_W + \mathcal{O}_{WB} + \mathcal{O}_{WW} + g^2 \left( \mathcal{O}_{\Phi,4} - \frac{1}{2} \mathcal{O}_{\Phi,2} \right) = -\frac{g^2}{4} \sum_i \left( \mathcal{O}_{\Phi L,ii}^{(3)} + \mathcal{O}_{\Phi Q,ii}^{(3)} \right) .$$

<sup>1</sup>Buchmuller & Wyler; Grzadkowski et al. arXiv: 1008.4884

# Effective Lagrangian

$$\mathcal{L}_{\text{eff}} = \sum_n \frac{f_n}{\Lambda^2} \mathcal{O}_n$$

Our assumptions are:

- The observed state belongs to a  $SU(2)$  doublet.
- The state is CP-even as in SM.
- Narrow resonance and no overlapping resonances.
- $SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$  SM local symmetry, C and P even, lepton and baryon number conservation

**59** dimension-6 operators are enough...<sup>1</sup>

Set reduced by considering only C and P even and EOM to eliminate/choose the basis

$$2\mathcal{O}_{\Phi,2} - 2\mathcal{O}_{\Phi,4} = \sum_{ij} \left( y_{ij}^e \mathcal{O}_{e\Phi,ij} + y_{ij}^u \mathcal{O}_{u\Phi,ij} + y_{ij}^d (\mathcal{O}_{d\Phi,ij})^\dagger + \text{h.c.} \right) ,$$

$$2\mathcal{O}_{\mathcal{B}} + \mathcal{O}_{WB} + \mathcal{O}_{BB} + g'^2 \left( \mathcal{O}_{\Phi,1} - \frac{1}{2} \mathcal{O}_{\Phi,2} \right) = -\frac{g'^2}{2} \sum_i \left( -\frac{1}{2} \mathcal{O}_{\Phi L,ii}^{(1)} + \frac{1}{6} \mathcal{O}_{\Phi Q,ii}^{(1)} - \mathcal{O}_{\Phi e,ii}^{(1)} + \frac{2}{3} \mathcal{O}_{\Phi u,ii}^{(1)} - \frac{1}{3} \mathcal{O}_{\Phi d,ii}^{(1)} \right)$$

$$2\mathcal{O}_W + \mathcal{O}_{WB} + \mathcal{O}_{WW} + g^2 \left( \mathcal{O}_{\Phi,4} - \frac{1}{2} \mathcal{O}_{\Phi,2} \right) = -\frac{g^2}{4} \sum_i \left( \mathcal{O}_{\Phi L,ii}^{(3)} + \mathcal{O}_{\Phi Q,ii}^{(3)} \right) .$$

<sup>1</sup>Buchmuller & Wyler; Grzadkowski et al. arXiv: 1008.4884

# The right of choice

Higgs interactions with gauge bosons<sup>2</sup>:

$$\begin{aligned}
 \mathcal{O}_{GG} &= \Phi^\dagger \Phi G_{\mu\nu}^a G^{a\mu\nu} , & \mathcal{O}_{WW} &= \Phi^\dagger \hat{W}_{\mu\nu} \hat{W}^{\mu\nu} \Phi , & \mathcal{O}_{BB} &= \Phi^\dagger \hat{B}_{\mu\nu} \hat{B}^{\mu\nu} \Phi , \\
 \mathcal{O}_{BW} &= \Phi^\dagger \hat{B}_{\mu\nu} \hat{W}^{\mu\nu} \Phi , & \mathcal{O}_W &= (D_\mu \Phi)^\dagger \hat{W}^{\mu\nu} (D_\nu \Phi) , & \mathcal{O}_B &= (D_\mu \Phi)^\dagger \hat{B}^{\mu\nu} (D_\nu \Phi) , \\
 \mathcal{O}_{\Phi,1} &= (D_\mu \Phi)^\dagger \Phi \Phi^\dagger (D^\mu \Phi) , & \mathcal{O}_{\Phi,2} &= \frac{1}{2} \partial^\mu (\Phi^\dagger \Phi) \partial_\mu (\Phi^\dagger \Phi) , & \mathcal{O}_{\Phi,4} &= (D_\mu \Phi)^\dagger (D^\mu \Phi) (\Phi^\dagger \Phi) ,
 \end{aligned}$$

Higgs interactions with fermions:

$$\begin{aligned}
 \mathcal{O}_{e\Phi,ij} &= (\Phi^\dagger \Phi) (\bar{L}_i \Phi e_{Rj}) & \mathcal{O}_{\Phi L,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu \Phi) (\bar{L}_i \gamma^\mu L_j) & \mathcal{O}_{\Phi L,ij}^{(3)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu^a \Phi) (\bar{L}_i \gamma^\mu \sigma_a L_j) \\
 \mathcal{O}_{u\Phi,ij} &= (\Phi^\dagger \Phi) (\bar{Q}_i \Phi u_{Rj}) & \mathcal{O}_{\Phi Q,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu \Phi) (\bar{Q}_i \gamma^\mu Q_j) & \mathcal{O}_{\Phi Q,ij}^{(3)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu^a \Phi) (\bar{Q}_i \gamma^\mu \sigma_a Q_j) \\
 \mathcal{O}_{d\Phi,ij} &= (\Phi^\dagger \Phi) (\bar{Q}_i \Phi d_{Rj}) & \mathcal{O}_{\Phi e,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu \Phi) (\bar{e}_{Ri} \gamma^\mu e_{Rj}) \\
 & & \mathcal{O}_{\Phi u,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu \Phi) (\bar{u}_{Ri} \gamma^\mu u_{Rj}) \\
 & & \mathcal{O}_{\Phi d,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu \Phi) (\bar{d}_{Ri} \gamma^\mu d_{Rj}) \\
 & & \mathcal{O}_{\Phi ud,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu \Phi) (\bar{u}_{Ri} \gamma^\mu d_{Rj})
 \end{aligned}$$

In the absence of theoretical prejudice chose a basis where the operators are more directly related to the existing data

$Z$  properties,  $W$  decays, low energy  $\nu$  scattering, atomic  $P$ , FCNC, Moller scattering  $P$  and  $e^+e^- \rightarrow f\bar{f}$  at LEP2.

EWPD at tree level, TGV

$$^2 D_\mu \Phi = \left( \partial_\mu + i \frac{1}{2} g' B_\mu + i g \frac{\sigma_a}{2} W_\mu^a \right) \Phi, \quad \hat{B}_{\mu\nu} = i \frac{g'}{2} B_{\mu\nu}, \quad \hat{W}_{\mu\nu} = i \frac{g}{2} \sigma^a W_{\mu\nu}^a$$

# The right of choice

Higgs interactions with gauge bosons<sup>2</sup>:

$$\begin{aligned}
 \mathcal{O}_{GG} &= \Phi^\dagger \Phi G_{\mu\nu}^a G^{a\mu\nu} , & \mathcal{O}_{WW} &= \Phi^\dagger \hat{W}_{\mu\nu} \hat{W}^{\mu\nu} \Phi , & \mathcal{O}_{BB} &= \Phi^\dagger \hat{B}_{\mu\nu} \hat{B}^{\mu\nu} \Phi , \\
 \mathcal{O}_{BW} &= \Phi^\dagger \hat{B}_{\mu\nu} \hat{W}^{\mu\nu} \Phi , & \mathcal{O}_W &= (D_\mu \Phi)^\dagger \hat{W}^{\mu\nu} (D_\nu \Phi) , & \mathcal{O}_B &= (D_\mu \Phi)^\dagger \hat{B}^{\mu\nu} (D_\nu \Phi) , \\
 \mathcal{O}_{\Phi,1} &= (D_\mu \Phi)^\dagger \Phi \Phi^\dagger (D^\mu \Phi) , & \mathcal{O}_{\Phi,2} &= \frac{1}{2} \partial^\mu (\Phi^\dagger \Phi) \partial_\mu (\Phi^\dagger \Phi) , & \mathcal{O}_{\Phi,4} &= (D_\mu \Phi)^\dagger (D^\mu \Phi) (\Phi^\dagger \Phi) ,
 \end{aligned}$$

Higgs interactions with fermions:

$$\begin{aligned}
 \mathcal{O}_{e\Phi,ij} &= (\Phi^\dagger \Phi) (\bar{L}_i \Phi e_{Rj}) & \mathcal{O}_{\Phi L,ij}^{(1)} &= \Phi^\dagger \overleftrightarrow{\mathcal{D}}_\mu \Phi (\bar{L}_i \gamma^\mu L_j) & \mathcal{O}_{\Phi L,ij}^{(3)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu^a \Phi) (\bar{L}_i \gamma^\mu \sigma_a L_j) \\
 \mathcal{O}_{u\Phi,ij} &= (\Phi^\dagger \Phi) (\bar{Q}_i \Phi u_{Rj}) & \mathcal{O}_{\Phi Q,ij}^{(1)} &= \Phi^\dagger \overleftrightarrow{\mathcal{D}}_\mu \Phi (\bar{Q}_i \gamma^\mu Q_j) & \mathcal{O}_{\Phi Q,ij}^{(3)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu^a \Phi) (\bar{Q}_i \gamma^\mu \sigma_a Q_j) \\
 \mathcal{O}_{d\Phi,ij} &= (\Phi^\dagger \Phi) (\bar{Q}_i \Phi d_{Rj}) & \mathcal{O}_{\Phi e,ij}^{(1)} &= \Phi^\dagger (D_\mu \Phi) (\bar{e}_{Ri} \gamma^\mu e_{Rj}) \\
 & & \mathcal{O}_{\Phi u,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu \Phi) (\bar{u}_{Ri} \gamma^\mu u_{Rj}) \\
 & & \mathcal{O}_{\Phi d,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu \Phi) (\bar{d}_{Ri} \gamma^\mu d_{Rj}) \\
 & & \mathcal{O}_{\Phi ud,ij}^{(1)} &= \tilde{\Phi}^\dagger (\overleftrightarrow{\mathcal{D}}_\mu \Phi) (\bar{u}_{Ri} \gamma^\mu d_{Rj})
 \end{aligned}$$

In the absence of theoretical prejudice chose a basis where the operators are more directly related to the existing data

$Z$  properties,  $W$  decays, low energy  $\nu$  scattering, atomic  $P$ , FCNC, Moller scattering  $P$  and  $e^+ e^- \rightarrow f \bar{f}$  at LEP2.

EWPD at tree level, TGV

$$^2 D_\mu \Phi = \left( \partial_\mu + i \frac{1}{2} g' B_\mu + i g \frac{\sigma_a}{2} W_\mu^a \right) \Phi, \quad \hat{B}_{\mu\nu} = i \frac{g'}{2} B_{\mu\nu}, \quad \hat{W}_{\mu\nu} = i \frac{g}{2} \sigma^a W_{\mu\nu}^a$$



# The right of choice

Higgs interactions with gauge bosons<sup>2</sup>:

$$\begin{aligned}
 \mathcal{O}_{GG} &= \Phi^\dagger \Phi G_{\mu\nu}^a G^{a\mu\nu} , & \mathcal{O}_{WW} &= \Phi^\dagger \hat{W}_{\mu\nu} \hat{W}^{\mu\nu} \Phi , & \mathcal{O}_{BB} &= \Phi^\dagger \hat{B}_{\mu\nu} \hat{B}^{\mu\nu} \Phi , \\
 \mathcal{O}_{BW} &= \Phi^\dagger \hat{B}_{\mu\nu} \hat{W}^{\mu\nu} \Phi , & \mathcal{O}_W &= (D_\mu \Phi)^\dagger \hat{W}^{\mu\nu} (D_\nu \Phi) , & \mathcal{O}_B &= (D_\mu \Phi)^\dagger \hat{B}^{\mu\nu} (D_\nu \Phi) , \\
 \mathcal{O}_{\Phi,1} &= (D_\mu \Phi)^\dagger \Phi \Phi^\dagger (D^\mu \Phi) , & \mathcal{O}_{\Phi,2} &= \frac{1}{2} \partial^\mu (\Phi^\dagger \Phi) \partial_\mu (\Phi^\dagger \Phi) , & \mathcal{O}_{\Phi,4} &= (D_\mu \Phi)^\dagger (D^\mu \Phi) (\Phi^\dagger \Phi) ,
 \end{aligned}$$

Higgs interactions with fermions:

$$\begin{aligned}
 \mathcal{O}_{e\Phi,ij} &= (\Phi^\dagger \Phi) (\bar{L}_i \Phi e_{Rj}) & \mathcal{O}_{\Phi L,ij}^{(1)} &= \Phi^\dagger \overleftrightarrow{\mathcal{D}}_\mu \Phi (\bar{L}_i \gamma^\mu L_j) & \mathcal{O}_{\Phi L,ij}^{(3)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu^a \Phi) (\bar{L}_i \gamma^\mu \sigma_a L_j) \\
 \mathcal{O}_{u\Phi,ij} &= (\Phi^\dagger \Phi) (\bar{Q}_i \Phi u_{Rj}) & \mathcal{O}_{\Phi Q,ij}^{(1)} &= \Phi^\dagger \overleftrightarrow{\mathcal{D}}_\mu \Phi (\bar{Q}_i \gamma^\mu Q_j) & \mathcal{O}_{\Phi Q,ij}^{(3)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu^a \Phi) (\bar{Q}_i \gamma^\mu \sigma_a Q_j) \\
 \mathcal{O}_{d\Phi,ij} &= (\Phi^\dagger \Phi) (\bar{Q}_i \Phi d_{Rj}) & \mathcal{O}_{\Phi e,ij}^{(1)} &= \Phi^\dagger \overleftrightarrow{\mathcal{D}}_\mu \Phi (\bar{e}_{Ri} \gamma^\mu e_{Rj}) \\
 & & \mathcal{O}_{\Phi u,ij}^{(1)} &= \Phi^\dagger \overleftrightarrow{\mathcal{D}}_\mu \Phi (\bar{u}_{Ri} \gamma^\mu u_{Rj}) \\
 & & \mathcal{O}_{\Phi d,ij}^{(1)} &= \Phi^\dagger \overleftrightarrow{\mathcal{D}}_\mu \Phi (\bar{d}_{Ri} \gamma^\mu d_{Rj}) \\
 & & \mathcal{O}_{\Phi ud,ij}^{(1)} &= \Phi^\dagger \overleftrightarrow{\mathcal{D}}_\mu \Phi (\bar{u}_{Ri} \gamma^\mu d_{Rj})
 \end{aligned}$$

In the absence of theoretical prejudice chose a basis where the operators are more directly related to the existing data

$Z$  properties,  $W$  decays, low energy  $\nu$  scattering, atomic  $P$ , FCNC, Moller scattering  $P$  and  $e^+ e^- \rightarrow f \bar{f}$  at LEP2.

EWPD at tree level, [TGV](#)

$${}^2 D_\mu \Phi = \left( \partial_\mu + i \frac{1}{2} g' B_\mu + i g \frac{\sigma_a}{2} W_\mu^a \right) \Phi, \quad \hat{B}_{\mu\nu} = i \frac{g'}{2} B_{\mu\nu}, \quad \hat{W}_{\mu\nu} = i \frac{g}{2} \sigma^a W_{\mu\nu}^a$$

# The right of choice

Higgs interactions with gauge bosons<sup>2</sup>:

$$\begin{aligned}
 \mathcal{O}_{GG} &= \Phi^\dagger \Phi G_{\mu\nu}^a G^{a\mu\nu} , & \mathcal{O}_{WW} &= \Phi^\dagger \hat{W}_{\mu\nu} \hat{W}^{\mu\nu} \Phi , & \mathcal{O}_{BB} &= \Phi^\dagger \hat{B}_{\mu\nu} \hat{B}^{\mu\nu} \Phi , \\
 \mathcal{O}_{BW} &= \Phi^\dagger \hat{B}_{\mu\nu} \hat{W}^{\mu\nu} \Phi , & \mathcal{O}_W &= (D_\mu \Phi)^\dagger \hat{W}^{\mu\nu} (D_\nu \Phi) , & \mathcal{O}_B &= (D_\mu \Phi)^\dagger \hat{B}^{\mu\nu} (D_\nu \Phi) , \\
 \mathcal{O}_{\Phi,1} &= (D_\mu \Phi)^\dagger \Phi \Phi^\dagger (D^\mu \Phi) , & \mathcal{O}_{\Phi,2} &= \frac{1}{2} \partial^\mu (\Phi^\dagger \Phi) \partial_\mu (\Phi^\dagger \Phi) , & \mathcal{O}_{\Phi,4} &= (D_\mu \Phi)^\dagger (D^\mu \Phi) (\Phi^\dagger \Phi) ,
 \end{aligned}$$

Higgs interactions with fermions:

$$\begin{aligned}
 \mathcal{O}_{e\Phi,ij} &= (\Phi^\dagger \Phi) (\bar{L}_i \Phi e_{Rj}) & \mathcal{O}_{\Phi L,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu \Phi) (\bar{L}_i \gamma^\mu L_j) & \mathcal{O}_{\Phi L,ij}^{(3)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu^a \Phi) (\bar{L}_i \gamma^\mu \sigma_a L_j) \\
 \mathcal{O}_{u\Phi,ij} &= (\Phi^\dagger \Phi) (\bar{Q}_i \Phi u_{Rj}) & \mathcal{O}_{\Phi Q,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu \Phi) (\bar{Q}_i \gamma^\mu Q_j) & \mathcal{O}_{\Phi Q,ij}^{(3)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu^a \Phi) (\bar{Q}_i \gamma^\mu \sigma_a Q_j) \\
 \mathcal{O}_{d\Phi,ij} &= (\Phi^\dagger \Phi) (\bar{Q}_i \Phi d_{Rj}) & \mathcal{O}_{\Phi e,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu \Phi) (\bar{e}_{Ri} \gamma^\mu e_{Rj}) \\
 & & \mathcal{O}_{\Phi u,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu \Phi) (\bar{u}_{Ri} \gamma^\mu u_{Rj}) \\
 & & \mathcal{O}_{\Phi d,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu \Phi) (\bar{d}_{Ri} \gamma^\mu d_{Rj}) \\
 & & \mathcal{O}_{\Phi ud,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu \Phi) (\bar{u}_{Ri} \gamma^\mu d_{Rj})
 \end{aligned}$$

In the absence of theoretical prejudice chose a basis where the operators are more directly related to the existing data

$Z$  properties,  $W$  decays, low energy  $\nu$  scattering, atomic  $P$ , FCNC, Moller scattering  $P$  and  $e^+ e^- \rightarrow f \bar{f}$  at LEP2.

EWPD at tree level, TGV

$${}^2 D_\mu \Phi = \left( \partial_\mu + i \frac{1}{2} g' B_\mu + i g \frac{\sigma_a}{2} W_\mu^a \right) \Phi, \quad \hat{B}_{\mu\nu} = i \frac{g'}{2} B_{\mu\nu}, \quad \hat{W}_{\mu\nu} = i \frac{g}{2} \sigma^a W_{\mu\nu}^a$$

# The right of choice

Higgs interactions with gauge bosons<sup>2</sup>:

$$\begin{aligned}
 \mathcal{O}_{GG} &= \Phi^\dagger \Phi G_{\mu\nu}^a G^{a\mu\nu} , & \mathcal{O}_{WW} &= \Phi^\dagger \hat{W}_{\mu\nu} \hat{W}^{\mu\nu} \Phi , & \mathcal{O}_{BB} &= \Phi^\dagger \hat{B}_{\mu\nu} \hat{B}^{\mu\nu} \Phi , \\
 \mathcal{O}_{BW} &= \Phi^\dagger \hat{B}_{\mu\nu} \hat{W}^{\mu\nu} \Phi , & \mathcal{O}_W &= (D_\mu \Phi)^\dagger \hat{W}^{\mu\nu} (D_\nu \Phi) , & \mathcal{O}_B &= (D_\mu \Phi)^\dagger \hat{B}^{\mu\nu} (D_\nu \Phi) , \\
 \mathcal{O}_{\Phi,1} &= (D_\mu \Phi)^\dagger \Phi \Phi^\dagger (D^\mu \Phi) , & \mathcal{O}_{\Phi,2} &= \frac{1}{2} \partial^\mu (\Phi^\dagger \Phi) \partial_\mu (\Phi^\dagger \Phi) , & \mathcal{O}_{\Phi,4} &= (D_\mu \Phi)^\dagger (D^\mu \Phi) (\Phi^\dagger \Phi) ,
 \end{aligned}$$

Higgs interactions with fermions:

$$\begin{aligned}
 \mathcal{O}_{e\Phi,ij} &= (\Phi^\dagger \Phi) (\bar{L}_i \Phi e_{Rj}) & \mathcal{O}_{\Phi L,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu \Phi) (\bar{L}_i \gamma^\mu L_j) & \mathcal{O}_{\Phi L,ij}^{(3)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu^a \Phi) (\bar{L}_i \gamma^\mu \sigma_a L_j) \\
 \mathcal{O}_{u\Phi,ij} &= (\Phi^\dagger \Phi) (\bar{Q}_i \Phi u_{Rj}) & \mathcal{O}_{\Phi Q,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu \Phi) (\bar{Q}_i \gamma^\mu Q_j) & \mathcal{O}_{\Phi Q,ij}^{(3)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu^a \Phi) (\bar{Q}_i \gamma^\mu \sigma_a Q_j) \\
 \mathcal{O}_{d\Phi,ij} &= (\Phi^\dagger \Phi) (\bar{Q}_i \Phi d_{Rj}) & \mathcal{O}_{\Phi e,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu \Phi) (\bar{e}_{Ri} \gamma^\mu e_{Rj}) \\
 & & \mathcal{O}_{\Phi u,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu \Phi) (\bar{u}_{Ri} \gamma^\mu u_{Rj}) \\
 & & \mathcal{O}_{\Phi d,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu \Phi) (\bar{d}_{Ri} \gamma^\mu d_{Rj}) \\
 & & \mathcal{O}_{\Phi ud,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{\mathcal{D}}_\mu \Phi) (\bar{u}_{Ri} \gamma^\mu d_{Rj})
 \end{aligned}$$

In the absence of theoretical prejudice chose a basis where the operators are more directly related to the existing data

Z properties, W decays, low energy  $\nu$  scattering, atomic P, FCNC, Moller scattering P and  $e^+ e^- \rightarrow f \bar{f}$  at LEP2.

EWPD at tree level, TGV

$${}^2 D_\mu \Phi = \left( \partial_\mu + i \frac{1}{2} g' B_\mu + i g \frac{\sigma_a}{2} W_\mu^a \right) \Phi, \quad \hat{B}_{\mu\nu} = i \frac{g'}{2} B_{\mu\nu}, \quad \hat{W}_{\mu\nu} = i \frac{g}{2} \sigma^a W_{\mu\nu}^a$$

# The right of choice

Higgs interactions with gauge bosons<sup>2</sup>:

$$\begin{aligned}
 \mathcal{O}_{GG} &= \Phi^\dagger \Phi G_{\mu\nu}^a G^{a\mu\nu} , & \mathcal{O}_{WW} &= \Phi^\dagger \hat{W}_{\mu\nu} \hat{W}^{\mu\nu} \Phi , & \mathcal{O}_{BB} &= \Phi^\dagger \hat{B}_{\mu\nu} \hat{B}^{\mu\nu} \Phi , \\
 \mathcal{O}_{BW} &= \Phi^\dagger \hat{B}_{\mu\nu} \hat{W}^{\mu\nu} \Phi , & \mathcal{O}_W &= (D_\mu \Phi)^\dagger \hat{W}^{\mu\nu} (D_\nu \Phi) , & \mathcal{O}_B &= (D_\mu \Phi)^\dagger \hat{B}^{\mu\nu} (D_\nu \Phi) , \\
 \mathcal{O}_{\Phi,1} &= (D_\mu \Phi)^\dagger \Phi \Phi^\dagger (D^\mu \Phi) , & \mathcal{O}_{\Phi,2} &= \frac{1}{2} \partial^\mu (\Phi^\dagger \Phi) \partial_\mu (\Phi^\dagger \Phi) , & \mathcal{O}_{\Phi,4} &= (D_\mu \Phi)^\dagger (D^\mu \Phi) (\Phi^\dagger \Phi) ,
 \end{aligned}$$

Higgs interactions with fermions:

$$\begin{aligned}
 \mathcal{O}_{e\Phi,ij} &= (\Phi^\dagger \Phi) (\bar{L}_i \Phi e_{Rj}) & \mathcal{O}_{\Phi L,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{D}_\mu \Phi) (\bar{L}_i \gamma^\mu L_j) & \mathcal{O}_{\Phi L,ij}^{(3)} &= \Phi^\dagger (\overleftrightarrow{D}_\mu^a \Phi) (\bar{L}_i \gamma^\mu \sigma_a L_j) \\
 \mathcal{O}_{u\Phi,ij} &= (\Phi^\dagger \Phi) (\bar{Q}_i \Phi u_{Rj}) & \mathcal{O}_{\Phi Q,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{D}_\mu \Phi) (\bar{Q}_i \gamma^\mu Q_j) & \mathcal{O}_{\Phi Q,ij}^{(3)} &= \Phi^\dagger (\overleftrightarrow{D}_\mu^a \Phi) (\bar{Q}_i \gamma^\mu \sigma_a Q_j) \\
 \mathcal{O}_{d\Phi,ij} &= (\Phi^\dagger \Phi) (\bar{Q}_i \Phi d_{Rj}) & \mathcal{O}_{\Phi e,ij}^{(1)} &= \Phi^\dagger (D_\mu \Phi) (\bar{e}_{Ri} \gamma^\mu e_{Rj}) \\
 & & \mathcal{O}_{\Phi u,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{D}_\mu \Phi) (\bar{u}_{Ri} \gamma^\mu u_{Rj}) \\
 & & \mathcal{O}_{\Phi d,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{D}_\mu \Phi) (\bar{d}_{Ri} \gamma^\mu d_{Rj}) \\
 & & \mathcal{O}_{\Phi ud,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{D}_\mu \Phi) (\bar{u}_{Ri} \gamma^\mu d_{Rj})
 \end{aligned}$$

In the absence of theoretical prejudice chose a basis where the operators are more directly related to the existing data

**Z properties, W decays, low energy  $\nu$  scattering, atomic P, FCNC, Moller scattering P and  $e^+e^- \rightarrow f\bar{f}$  at LEP2.**

EWPD at tree level, TGV

$$^2 D_\mu \Phi = \left( \partial_\mu + i \frac{1}{2} g' B_\mu + i g \frac{\sigma_a}{2} W_\mu^a \right) \Phi, \quad \hat{B}_{\mu\nu} = i \frac{g'}{2} B_{\mu\nu}, \quad \hat{W}_{\mu\nu} = i \frac{g}{2} \sigma^a W_{\mu\nu}^a$$

# The right of choice

Higgs interactions with gauge bosons<sup>2</sup>:

$$\begin{aligned}
 \mathcal{O}_{GG} &= \Phi^\dagger \Phi G_{\mu\nu}^a G^{a\mu\nu} , & \mathcal{O}_{WW} &= \Phi^\dagger \hat{W}_{\mu\nu} \hat{W}^{\mu\nu} \Phi , & \mathcal{O}_{BB} &= \Phi^\dagger \hat{B}_{\mu\nu} \hat{B}^{\mu\nu} \Phi , \\
 \mathcal{O}_{BW} &= \Phi^\dagger \hat{B}_{\mu\nu} \hat{W}^{\mu\nu} \Phi , & \mathcal{O}_W &= (D_\mu \Phi)^\dagger \hat{W}^{\mu\nu} (D_\nu \Phi) , & \mathcal{O}_B &= (D_\mu \Phi)^\dagger \hat{B}^{\mu\nu} (D_\nu \Phi) , \\
 \mathcal{O}_{\Phi,1} &= (D_\mu \Phi)^\dagger \Phi \Phi^\dagger (D^\mu \Phi) , & \mathcal{O}_{\Phi,2} &= \frac{1}{2} \partial^\mu (\Phi^\dagger \Phi) \partial_\mu (\Phi^\dagger \Phi) , & \mathcal{O}_{\Phi,4} &= (D_\mu \Phi)^\dagger (D^\mu \Phi) (\Phi^\dagger \Phi) ,
 \end{aligned}$$

Higgs interactions with fermions:

$$\begin{aligned}
 \mathcal{O}_{e\Phi,33} &= (\Phi^\dagger \Phi) (\bar{L}_3 \Phi e_{R3}) & \mathcal{O}_{\Phi L,ij}^{(1)} &= \Phi^\dagger \overleftrightarrow{(D_\mu \Phi)} (\bar{L}_i \gamma^\mu L_j) & \mathcal{O}_{\Phi L,ij}^{(3)} &= \Phi^\dagger \overleftrightarrow{(D_\mu^a \Phi)} (\bar{L}_i \gamma^\mu \sigma_a L_j) \\
 \mathcal{O}_{u\Phi,ij} &= (\Phi^\dagger \Phi) (\bar{Q}_i \Phi u_{Rj}) & \mathcal{O}_{\Phi Q,ij}^{(1)} &= \Phi^\dagger \overleftrightarrow{(D_\mu \Phi)} (\bar{Q}_i \gamma^\mu Q_j) & \mathcal{O}_{\Phi Q,ij}^{(3)} &= \Phi^\dagger \overleftrightarrow{(D_\mu^a \Phi)} (\bar{Q}_i \gamma^\mu \sigma_a Q_j) \\
 \mathcal{O}_{d\Phi,33} &= (\Phi^\dagger \Phi) (\bar{Q}_3 \Phi d_{R3}) & \mathcal{O}_{\Phi e,ij}^{(1)} &= \Phi^\dagger \overleftrightarrow{(D_\mu \Phi)} (\bar{e}_{Ri} \gamma^\mu e_{Rj}) \\
 & & \mathcal{O}_{\Phi u,ij}^{(1)} &= \Phi^\dagger \overleftrightarrow{(D_\mu \Phi)} (\bar{u}_{Ri} \gamma^\mu u_{Rj}) \\
 & & \mathcal{O}_{\Phi d,ij}^{(1)} &= \Phi^\dagger \overleftrightarrow{(D_\mu \Phi)} (\bar{d}_{Ri} \gamma^\mu d_{Rj}) \\
 & & \mathcal{O}_{\Phi ud,ij}^{(1)} &= \Phi^\dagger \overleftrightarrow{(D_\mu \Phi)} (\bar{u}_{Ri} \gamma^\mu d_{Rj})
 \end{aligned}$$

In the absence of theoretical prejudice chose a basis where the operators are more directly related to the existing data

$Z$  properties,  $W$  decays, low energy  $\nu$  scattering, atomic  $P$ , **FCNC**, Moller scattering  $P$  and  $e^+e^- \rightarrow f\bar{f}$  at LEP2.

EWPD at tree level, TGV

$$^2 D_\mu \Phi = \left( \partial_\mu + i \frac{1}{2} g' B_\mu + i g \frac{\sigma_a}{2} W_\mu^a \right) \Phi, \quad \hat{B}_{\mu\nu} = i \frac{g'}{2} B_{\mu\nu}, \quad \hat{W}_{\mu\nu} = i \frac{g}{2} \sigma^a W_{\mu\nu}^a$$

# The right of choice

Higgs interactions with gauge bosons<sup>2</sup>:

$$\begin{aligned}
 \mathcal{O}_{GG} &= \Phi^\dagger \Phi G_{\mu\nu}^a G^{a\mu\nu} , & \mathcal{O}_{WW} &= \Phi^\dagger \hat{W}_{\mu\nu} \hat{W}^{\mu\nu} \Phi , & \mathcal{O}_{BB} &= \Phi^\dagger \hat{B}_{\mu\nu} \hat{B}^{\mu\nu} \Phi , \\
 \mathcal{O}_{BW} &= \Phi^\dagger \hat{B}_{\mu\nu} \hat{W}^{\mu\nu} \Phi , & \mathcal{O}_W &= (D_\mu \Phi)^\dagger \hat{W}^{\mu\nu} (D_\nu \Phi) , & \mathcal{O}_B &= (D_\mu \Phi)^\dagger \hat{B}^{\mu\nu} (D_\nu \Phi) , \\
 \mathcal{O}_{\Phi,1} &= (D_\mu \Phi)^\dagger \Phi \Phi^\dagger (D^\mu \Phi) , & \mathcal{O}_{\Phi,2} &= \frac{1}{2} \partial^\mu (\Phi^\dagger \Phi) \partial_\mu (\Phi^\dagger \Phi) , & \mathcal{O}_{\Phi,4} &= (D_\mu \Phi)^\dagger (D^\mu \Phi) (\Phi^\dagger \Phi) ,
 \end{aligned}$$

Higgs interactions with fermions:

$$\begin{aligned}
 \mathcal{O}_{e\Phi,33} &= (\Phi^\dagger \Phi) (\bar{L}_3 \Phi e_{R3}) & \mathcal{O}_{\Phi L,ij}^{(1)} &= \Phi^\dagger \overleftrightarrow{(D_\mu \Phi)} (\bar{L}_i \gamma^\mu L_j) & \mathcal{O}_{\Phi L,ij}^{(3)} &= \Phi^\dagger \overleftrightarrow{(D_\mu^a \Phi)} (\bar{L}_i \gamma^\mu \sigma_a L_j) \\
 \mathcal{O}_{u\Phi,ij} &= (\Phi^\dagger \Phi) (\bar{Q}_i \Phi u_{Rj}) & \mathcal{O}_{\Phi Q,ij}^{(1)} &= \Phi^\dagger \overleftrightarrow{(D_\mu \Phi)} (\bar{Q}_i \gamma^\mu Q_j) & \mathcal{O}_{\Phi Q,ij}^{(3)} &= \Phi^\dagger \overleftrightarrow{(D_\mu^a \Phi)} (\bar{Q}_i \gamma^\mu \sigma_a Q_j) \\
 \mathcal{O}_{d\Phi,33} &= (\Phi^\dagger \Phi) (\bar{Q}_3 \Phi d_{R3}) & \mathcal{O}_{\Phi e,ij}^{(1)} &= \Phi^\dagger \overleftrightarrow{(D_\mu \Phi)} (\bar{e}_{Ri} \gamma^\mu e_{Rj}) \\
 & & \mathcal{O}_{\Phi u,ij}^{(1)} &= \Phi^\dagger \overleftrightarrow{(D_\mu \Phi)} (\bar{u}_{Ri} \gamma^\mu u_{Rj}) \\
 & & \mathcal{O}_{\Phi d,ij}^{(1)} &= \Phi^\dagger \overleftrightarrow{(D_\mu \Phi)} (\bar{d}_{Ri} \gamma^\mu d_{Rj}) \\
 & & \mathcal{O}_{\Phi ud,ij}^{(1)} &= \Phi^\dagger \overleftrightarrow{(D_\mu \Phi)} (\bar{u}_{Ri} \gamma^\mu d_{Rj})
 \end{aligned}$$

In the absence of theoretical prejudice chose a basis where the operators are more directly related to the existing data

$Z$  properties,  $W$  decays, low energy  $\nu$  scattering, atomic  $P$ , FCNC, Moller scattering  $P$  and  $e^+e^- \rightarrow f\bar{f}$  at LEP2.

**EWPD at tree level, TGV**

$$^2 D_\mu \Phi = \left( \partial_\mu + i \frac{1}{2} g' B_\mu + i g \frac{\sigma_a}{2} W_\mu^a \right) \Phi, \quad \hat{B}_{\mu\nu} = i \frac{g'}{2} B_{\mu\nu}, \quad \hat{W}_{\mu\nu} = i \frac{g}{2} \sigma^a W_{\mu\nu}^a$$

# The right of choice

Higgs interactions with gauge bosons<sup>2</sup>:

$$\begin{aligned}
 \mathcal{O}_{GG} &= \Phi^\dagger \Phi G_{\mu\nu}^a G^{a\mu\nu} , & \mathcal{O}_{WW} &= \Phi^\dagger \hat{W}_{\mu\nu} \hat{W}^{\mu\nu} \Phi , & \mathcal{O}_{BB} &= \Phi^\dagger \hat{B}_{\mu\nu} \hat{B}^{\mu\nu} \Phi , \\
 \mathcal{O}_{BW} &= \Phi^\dagger \hat{B}_{\mu\nu} \hat{W}^{\mu\nu} \Phi , & \mathcal{O}_W &= (D_\mu \Phi)^\dagger \hat{W}^{\mu\nu} (D_\nu \Phi) , & \mathcal{O}_B &= (D_\mu \Phi)^\dagger \hat{B}^{\mu\nu} (D_\nu \Phi) , \\
 \mathcal{O}_{\Phi,1} &= (D_\mu \Phi)^\dagger \Phi \Phi^\dagger (D^\mu \Phi) , & \mathcal{O}_{\Phi,2} &= \frac{1}{2} \partial^\mu (\Phi^\dagger \Phi) \partial_\mu (\Phi^\dagger \Phi) , & \mathcal{O}_{\Phi,4} &= (D_\mu \Phi)^\dagger (D^\mu \Phi) (\Phi^\dagger \Phi) ,
 \end{aligned}$$

Higgs interactions with fermions:

$$\begin{aligned}
 \mathcal{O}_{e\Phi,33} &= (\Phi^\dagger \Phi) (\bar{L}_3 \Phi e_{R3}) & \mathcal{O}_{\Phi L,ij}^{(1)} &= \Phi^\dagger \overleftrightarrow{(D_\mu \Phi)} (\bar{L}_i \gamma^\mu L_j) & \mathcal{O}_{\Phi L,ij}^{(3)} &= \Phi^\dagger \overleftrightarrow{(D_\mu^a \Phi)} (\bar{L}_i \gamma^\mu \sigma_a L_j) \\
 \mathcal{O}_{u\Phi,ij} &= (\Phi^\dagger \Phi) (\bar{Q}_i \Phi u_{Rj}) & \mathcal{O}_{\Phi Q,ij}^{(1)} &= \Phi^\dagger \overleftrightarrow{(D_\mu \Phi)} (\bar{Q}_i \gamma^\mu Q_j) & \mathcal{O}_{\Phi Q,ij}^{(3)} &= \Phi^\dagger \overleftrightarrow{(D_\mu^a \Phi)} (\bar{Q}_i \gamma^\mu \sigma_a Q_j) \\
 \mathcal{O}_{d\Phi,33} &= (\Phi^\dagger \Phi) (\bar{Q}_3 \Phi d_{R3}) & \mathcal{O}_{\Phi e,ij}^{(1)} &= \Phi^\dagger \overleftrightarrow{(D_\mu \Phi)} (\bar{e}_{Ri} \gamma^\mu e_{Rj}) \\
 & & \mathcal{O}_{\Phi u,ij}^{(1)} &= \Phi^\dagger \overleftrightarrow{(D_\mu \Phi)} (\bar{u}_{Ri} \gamma^\mu u_{Rj}) \\
 & & \mathcal{O}_{\Phi d,ij}^{(1)} &= \Phi^\dagger \overleftrightarrow{(D_\mu \Phi)} (\bar{d}_{Ri} \gamma^\mu d_{Rj}) \\
 & & \mathcal{O}_{\Phi ud,ij}^{(1)} &= \Phi^\dagger \overleftrightarrow{(D_\mu \Phi)} (\bar{u}_{Ri} \gamma^\mu d_{Rj})
 \end{aligned}$$

In the absence of theoretical prejudice chose a basis where the operators are more directly related to the existing data

$Z$  properties,  $W$  decays, low energy  $\nu$  scattering, atomic  $P$ , FCNC, Moller scattering  $P$  and  $e^+e^- \rightarrow f\bar{f}$  at LEP2.

EWPD at tree level, TGV

$$^2 D_\mu \Phi = \left( \partial_\mu + i \frac{1}{2} g' B_\mu + i g \frac{\sigma_a}{2} W_\mu^a \right) \Phi, \quad \hat{B}_{\mu\nu} = i \frac{g'}{2} B_{\mu\nu}, \quad \hat{W}_{\mu\nu} = i \frac{g}{2} \sigma^a W_{\mu\nu}^a$$

# The right of choice

Higgs interactions with gauge bosons<sup>2</sup>:

$$\begin{aligned}
 \mathcal{O}_{GG} &= \Phi^\dagger \Phi G_{\mu\nu}^a G^{a\mu\nu} , & \mathcal{O}_{WW} &= \Phi^\dagger \hat{W}_{\mu\nu} \hat{W}^{\mu\nu} \Phi , & \mathcal{O}_{BB} &= \Phi^\dagger \hat{B}_{\mu\nu} \hat{B}^{\mu\nu} \Phi , \\
 \mathcal{O}_{BW} &= \Phi^\dagger \hat{B}_{\mu\nu} \hat{W}^{\mu\nu} \Phi , & \mathcal{O}_W &= (D_\mu \Phi)^\dagger \hat{W}^{\mu\nu} (D_\nu \Phi) , & \mathcal{O}_B &= (D_\mu \Phi)^\dagger \hat{B}^{\mu\nu} (D_\nu \Phi) , \\
 \mathcal{O}_{\Phi,1} &= (D_\mu \Phi)^\dagger \Phi \Phi^\dagger (D^\mu \Phi) , & \mathcal{O}_{\Phi,2} &= \frac{1}{2} \partial^\mu (\Phi^\dagger \Phi) \partial_\mu (\Phi^\dagger \Phi) , & \mathcal{O}_{\Phi,4} &= (D_\mu \Phi)^\dagger (D^\mu \Phi) (\Phi^\dagger \Phi) ,
 \end{aligned}$$

Higgs interactions with fermions:

$$\begin{aligned}
 \mathcal{O}_{e\Phi,33} &= (\Phi^\dagger \Phi) (\bar{L}_3 \Phi e_{R3}) & \mathcal{O}_{\Phi L,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{D}_\mu \Phi) (\bar{L}_i \gamma^\mu L_j) & \mathcal{O}_{\Phi L,ij}^{(3)} &= \Phi^\dagger (\overleftrightarrow{D}_\mu^a \Phi) (\bar{L}_i \gamma^\mu \sigma_a L_j) \\
 \mathcal{O}_{u\Phi,ij} &= (\Phi^\dagger \Phi) (\bar{Q}_i \Phi u_{Rj}) & \mathcal{O}_{\Phi Q,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{D}_\mu \Phi) (\bar{Q}_i \gamma^\mu Q_j) & \mathcal{O}_{\Phi Q,ij}^{(3)} &= \Phi^\dagger (\overleftrightarrow{D}_\mu^a \Phi) (\bar{Q}_i \gamma^\mu \sigma_a Q_j) \\
 \mathcal{O}_{d\Phi,33} &= (\Phi^\dagger \Phi) (\bar{Q}_3 \Phi d_{R3}) & \mathcal{O}_{\Phi e,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{D}_\mu \Phi) (\bar{e}_{Ri} \gamma^\mu e_{Rj}) \\
 & & \mathcal{O}_{\Phi u,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{D}_\mu \Phi) (\bar{u}_{Ri} \gamma^\mu u_{Rj}) \\
 & & \mathcal{O}_{\Phi d,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{D}_\mu \Phi) (\bar{d}_{Ri} \gamma^\mu d_{Rj}) \\
 & & \mathcal{O}_{\Phi ud,ij}^{(1)} &= \Phi^\dagger (\overleftrightarrow{D}_\mu \Phi) (\bar{u}_{Ri} \gamma^\mu d_{Rj})
 \end{aligned}$$

In the absence of theoretical prejudice chose a basis where the operators are more directly related to the existing data

Z properties, W decays, low energy  $\nu$  scattering, atomic P, FCNC, Moller scattering P and  $e^+e^- \rightarrow f\bar{f}$  at LEP2.

EWPD at tree level, TGV

$$^2 D_\mu \Phi = \left( \partial_\mu + i \frac{1}{2} g' B_\mu + i g \frac{\sigma_a}{2} W_\mu^a \right) \Phi, \quad \hat{B}_{\mu\nu} = i \frac{g'}{2} B_{\mu\nu}, \quad \hat{W}_{\mu\nu} = i \frac{g}{2} \sigma^a W_{\mu\nu}^a$$



# Effective Lagrangian for Higgs Interactions

$$\mathcal{L}_{eff} = -\frac{\alpha_s v}{8\pi} \frac{f_g}{\Lambda^2} \mathcal{O}_{GG} + \frac{f_{WW}}{\Lambda^2} \mathcal{O}_{WW} + \frac{f_W}{\Lambda^2} \mathcal{O}_W + \frac{f_B}{\Lambda^2} \mathcal{O}_B + \frac{f_{\text{bot}}}{\Lambda^2} \mathcal{O}_{d\Phi,33} + \frac{f_\tau}{\Lambda^2} \mathcal{O}_{e\Phi,33} \\ \text{Unitary gauge:} \quad \left( + \frac{f_{BB}}{\Lambda^2} \mathcal{O}_{BB} \right)$$

$$\begin{aligned} \mathcal{L}_{\text{eff}}^{\text{HVV}} &= g_{Hgg} H G_{\mu\nu}^a G^{a\mu\nu} + g_{H\gamma\gamma} H A_{\mu\nu} A^{\mu\nu} + g_{HZ\gamma}^{(1)} A_{\mu\nu} Z^\mu \partial^\nu H + g_{HZ\gamma}^{(2)} H A_{\mu\nu} Z^{\mu\nu} \\ &+ g_{HZZ}^{(1)} Z_{\mu\nu} Z^\mu \partial^\nu H + g_{HZZ}^{(2)} H Z_{\mu\nu} Z^{\mu\nu} + g_{HWW}^{(1)} \left( W_{\mu\nu}^+ W^{-\mu} \partial^\nu H + \text{h.c.} \right) \\ &+ g_{HWW}^{(2)} H W_{\mu\nu}^+ W^{-\mu\nu} \\ \mathcal{L}_{\text{eff}}^{Hff} &= g_{Hij}^f \bar{f}_L^i f_R^j H + \text{h.c.} \end{aligned}$$

$$\begin{aligned} g_{Hgg} &= -\frac{\alpha_s}{8\pi} \frac{f_g v}{\Lambda^2} & , g_{H\gamma\gamma} &= -\left( \frac{g^2 v s^2}{2\Lambda^2} \right) \frac{f_{WW} + f_{BB}}{2} , \\ g_{HZ\gamma}^{(1)} &= \left( \frac{g^2 v}{2\Lambda^2} \right) \frac{s(f_W - f_B)}{2c} & , g_{HZ\gamma}^{(2)} &= \left( \frac{g^2 v}{2\Lambda^2} \right) \frac{s[2s^2 f_{BB} - 2c^2 f_{WW}]}{2c} , \\ g_{HZZ}^{(1)} &= \left( \frac{g^2 v}{2\Lambda^2} \right) \frac{c^2 f_W + s^2 f_B}{2c^2} & , g_{HZZ}^{(2)} &= -\left( \frac{g^2 v}{2\Lambda^2} \right) \frac{s^4 f_{BB} + c^4 f_{WW}}{2c^2} , \\ g_{HWW}^{(1)} &= \left( \frac{g^2 v}{2\Lambda^2} \right) \frac{f_W}{2} & , g_{HWW}^{(2)} &= -\left( \frac{g^2 v}{2\Lambda^2} \right) f_{WW} \end{aligned}$$

$$g_{Hij}^f = -\frac{m_i^f}{v} \delta_{ij} + \frac{v^2}{\sqrt{2}\Lambda^2} f_{f\Phi,ij}^f$$

# Effective Lagrangian for Higgs Interactions

$$\mathcal{L}_{eff} = -\frac{\alpha_s v}{8\pi} \frac{f_g}{\Lambda^2} \mathcal{O}_{GG} + \frac{f_{WW}}{\Lambda^2} \mathcal{O}_{WW} + \frac{f_W}{\Lambda^2} \mathcal{O}_W + \frac{f_B}{\Lambda^2} \mathcal{O}_B + \frac{f_{\text{bot}}}{\Lambda^2} \mathcal{O}_{d\Phi,33} + \frac{f_\tau}{\Lambda^2} \mathcal{O}_{e\Phi,33} \\ \text{Unitary gauge:} \quad \left( + \frac{f_{BB}}{\Lambda^2} \mathcal{O}_{BB} \right)$$

$$\mathcal{L}_{\text{eff}}^{\text{HVV}} = g_{Hgg} H G_{\mu\nu}^a G^{a\mu\nu} + g_{H\gamma\gamma} H A_{\mu\nu} A^{\mu\nu} + g_{HZ\gamma}^{(1)} A_{\mu\nu} Z^\mu \partial^\nu H + g_{HZ\gamma}^{(2)} H A_{\mu\nu} Z^{\mu\nu} \\ + g_{HZZ}^{(1)} Z_{\mu\nu} Z^\mu \partial^\nu H + g_{HZZ}^{(2)} H Z_{\mu\nu} Z^{\mu\nu} + g_{HWW}^{(1)} \left( W_{\mu\nu}^+ W^{-\mu} \partial^\nu H + \text{h.c.} \right) \\ + g_{HWW}^{(2)} H W_{\mu\nu}^+ W^{-\mu\nu}$$

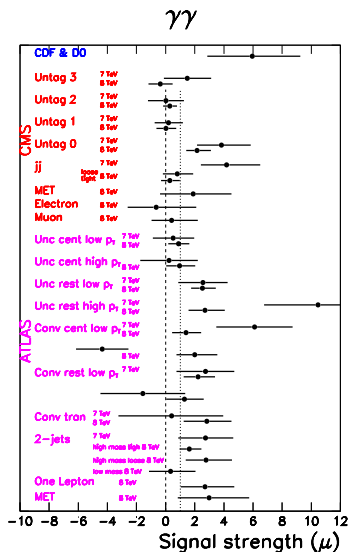
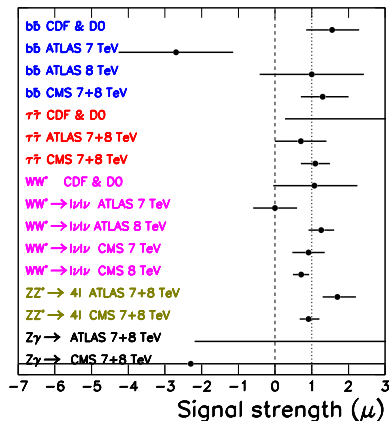
$$\mathcal{L}_{\text{eff}}^{Hff} = g_{Hij}^f \bar{f}_L f_R' H + \text{h.c.}$$

$$g_{Hgg} = -\frac{\alpha_s}{8\pi} \frac{f_g v}{\Lambda^2}, \quad g_{H\gamma\gamma} = -\left( \frac{g^2 v s^2}{2\Lambda^2} \right) \frac{f_{WW} + f_{BB}}{2}, \\ g_{HZ\gamma}^{(1)} = \left( \frac{g^2 v}{2\Lambda^2} \right) \frac{s(f_W - f_B)}{2c}, \quad g_{HZ\gamma}^{(2)} = \left( \frac{g^2 v}{2\Lambda^2} \right) \frac{s[2s^2 f_{BB} - 2c^2 f_{WW}]}{2c}, \\ g_{HZZ}^{(1)} = \left( \frac{g^2 v}{2\Lambda^2} \right) \frac{c^2 f_W + s^2 f_B}{2c^2}, \quad g_{HZZ}^{(2)} = -\left( \frac{g^2 v}{2\Lambda^2} \right) \frac{s^4 f_{BB} + c^4 f_{WW}}{2c^2}, \\ g_{HWW}^{(1)} = \left( \frac{g^2 v}{2\Lambda^2} \right) \frac{f_W}{2}, \quad g_{HWW}^{(2)} = -\left( \frac{g^2 v}{2\Lambda^2} \right) f_{WW}$$

$$g_{Hij}^f = -\frac{m_i^f}{v} \delta_{ij} + \frac{v^2}{\sqrt{2}\Lambda^2} f_{f\Phi,ij}'$$

# The statistical analysis

$$\chi^2 = \min_{\xi_{pull}} \sum_j \frac{(\mu_j - \mu_j^{\text{exp}})^2}{\sigma_j^2} + \sum_{pull} \left( \frac{\xi_{pull}}{\sigma_{pull}} \right)^2$$



# Adding TGV and EWPD

Data on triple electroweak gauge boson vertices:

$$\mathcal{L}_{WWV} = -ig_{WWV} \left\{ g_1^V \left( W_{\mu\nu}^+ W^{-\mu} V^\nu - W_\mu^+ V_\nu W^{-\mu\nu} \right) + \kappa_V W_\mu^+ W_\nu^- V^{\mu\nu} + \frac{\lambda_V}{m_W^2} W_{\mu\nu}^+ W^{-\nu\rho} V_\rho{}^\mu \right\}$$

with

$$\begin{aligned} \Delta g_1^Z &= g_1^Z - 1 = \frac{g^2 v^2}{8c^2 \Lambda^2} f_W, \\ \Delta \kappa_\gamma &= \kappa_\gamma - 1 = \frac{g^2 v^2}{8\Lambda^2} (f_W + f_B), \\ \Delta \kappa_Z &= \kappa_Z - 1 = \frac{g^2 v^2}{8c^2 \Lambda^2} (c^2 f_W - s^2 f_B). \end{aligned}$$

LEP data:

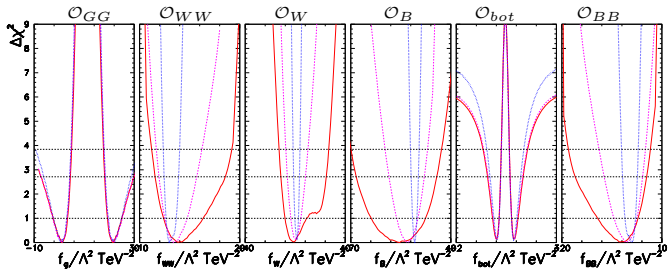
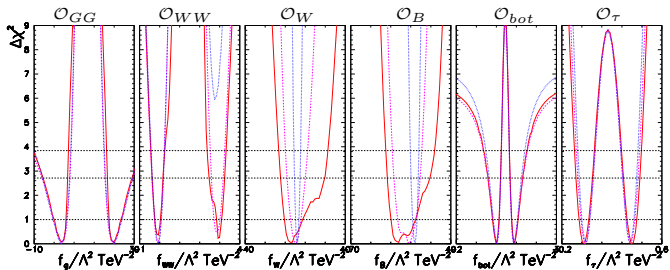
$$\begin{aligned} g_1^Z &= 0.984_{-0.049}^{+0.049} \\ \kappa_\gamma &= 1.004_{-0.025}^{+0.024} \end{aligned}$$

with a correlation factor  $\rho = 0.11$ .

Data on EWPD in terms of the S,T,U parameters:

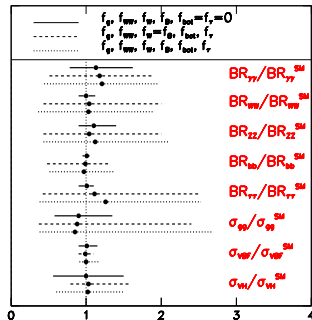
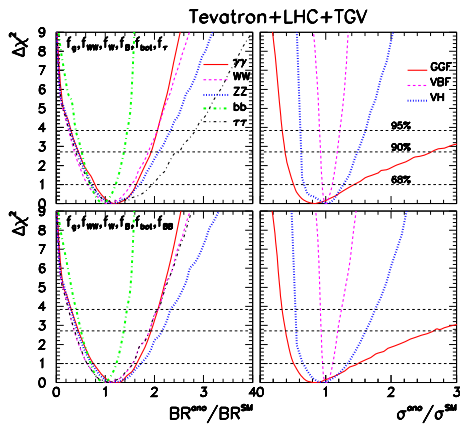
$$\begin{aligned} \Delta S &= 0.00 \pm 0.10 & \Delta T &= 0.02 \pm 0.11 & \Delta U &= 0.03 \pm 0.09 \\ \rho &= \begin{pmatrix} 1 & 0.89 & -0.55 \\ 0.89 & 1 & -0.8 \\ -0.55 & -0.8 & 1 \end{pmatrix} \end{aligned}$$

# $\Delta\chi^2$ vrs $f_X$



## BRs and production CS

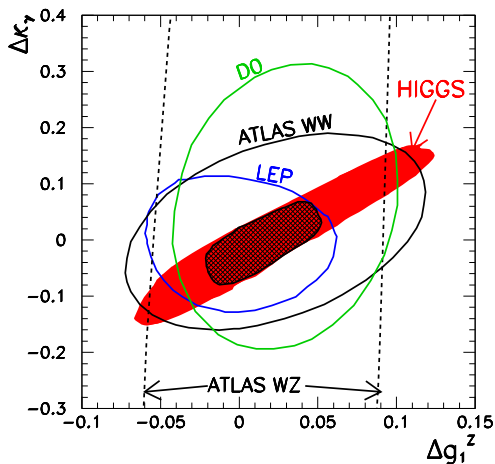
arXiv:1207.1344, 1211.4580

<http://hep.if.usp.br/Higgs>

# Determining TGV from Higgs data

arxiv:1304.1151

- Gauge Invariance  $\rightarrow$  TGV and Higgs couplings related:  $\mathcal{O}_W$  and  $\mathcal{O}_B$
- **Complementarity in experimental searches:** Higgs data bounds on  $f_W \otimes f_B \equiv \Delta\kappa_\gamma \otimes \Delta g_1^Z$



$$\Delta g_1^Z = g_1^Z - 1 = \frac{g^2 v^2}{8c^2 \Lambda^2} f_W \quad ,$$

$$\Delta \kappa_\gamma = \kappa_\gamma - 1 = \frac{g^2 v^2}{8\Lambda^2} (f_W + f_B) \quad ,$$

$$\Delta \kappa_Z = \kappa_Z - 1 = \frac{g^2 v^2}{8c^2 \Lambda^2} (c^2 f_W - s^2 f_B) \quad .$$

# Discussion and Conclusions

THANK YOU!

- **Model independent** analysis where the effects of new physics in the Higgs couplings are parametrized in  $\mathcal{L}_{eff}$ .  $SU(2)_L$  doublet  $\rightarrow SU(2)_L \times U(1)_Y$  gauge symmetry linearly realized:

$$\mathcal{L}_{eff} = \sum_n \frac{f_n}{\Lambda^2} \mathcal{O}_n \ ,$$

So far observations consistent with Higgs boson.

- Choice of basis:  
**Power to the data**  $\rightarrow$  operators whose coefficients are more easily related to existing data.
- Exploit interesting **complementarity between experimental searches**: TGV and Higgs data

arXiv:1207.1344, 1211.4580, 1304.1151  
<http://hep.if.usp.br/Higgs>



# Discussion and Conclusions

THANK YOU!

- **Model independent** analysis where the effects of new physics in the Higgs couplings are parametrized in  $\mathcal{L}_{eff}$ .  $SU(2)_L$  doublet  $\rightarrow SU(2)_L \times U(1)_Y$  gauge symmetry linearly realized:

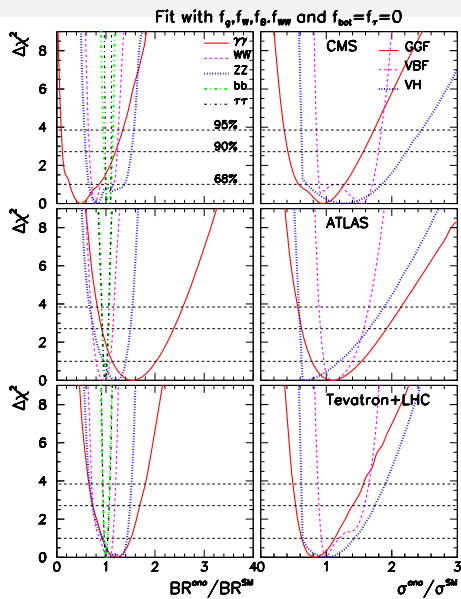
$$\mathcal{L}_{eff} = \sum_n \frac{f_n}{\Lambda^2} \mathcal{O}_n \ ,$$

So far observations consistent with Higgs boson.

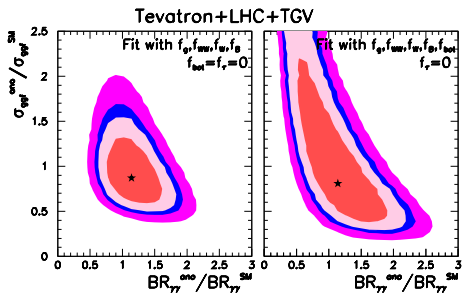
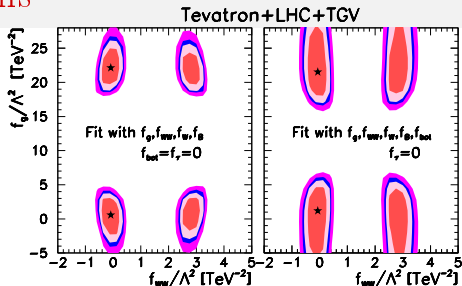
- Choice of basis:  
**Power to the data**  $\rightarrow$  operators whose coefficients are more easily related to existing data.
- Exploit interesting **complementarity between experimental searches**: TGV and Higgs data

arXiv:1207.1344, 1211.4580, 1304.1151  
<http://hep.if.usp.br/Higgs>

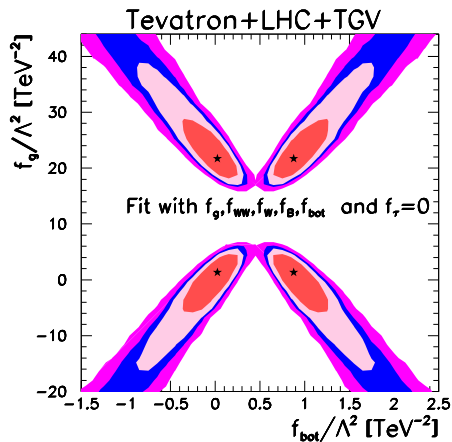
## CMS vrs ATLAS



## 2d correlations



# 2d correlations



# Best fit and ranges

|                                                  | Fit with $f_{bot} = f_\tau = 0$ |                                  | Fit with $f_{bot}$ and $f_\tau$ |                                   |
|--------------------------------------------------|---------------------------------|----------------------------------|---------------------------------|-----------------------------------|
|                                                  | Best fit                        | 90% CL allowed range             | Best fit                        | 90% CL allowed range              |
| $f_g/\Lambda^2$ (TeV <sup>-2</sup> )             | 0.64, 22.1                      | $[-1.8, 2.7] \cup [20, 25]$      | 0.71, 22.0                      | $[-6.2, 4.4] \cup [18, 29]$       |
| $f_{WW}/\Lambda^2$ (TeV <sup>-2</sup> )          | -0.083                          | $[-0.35, 0.15] \cup [2.6, 3.05]$ | -0.095                          | $[-0.39, 0.19]$                   |
| $f_W/\Lambda^2$ (TeV <sup>-2</sup> )             | 0.35                            | $[-6.2, 8.4]$                    | -0.46                           | $[-7.1, 6.5]$                     |
| $f_B/\Lambda^2$ (TeV <sup>-2</sup> )             | -5.9                            | $[-22, 6.7]$                     | -0.46                           | $[-7.1, 6.5]$                     |
| $f_{bot}/\Lambda^2$ (TeV <sup>-2</sup> )         | —                               | —                                | 0.01, 0.89                      | $[-0.34, 0.23] \cup [0.67, 1.2]$  |
| $f_\tau/\Lambda^2$ (TeV <sup>-2</sup> )          | —                               | —                                | -0.01, 0.34                     | $[-0.07, 0.05] \cup [0.28, 0.40]$ |
| $BR_{\gamma\gamma}^{ano}/BR_{\gamma\gamma}^{SM}$ | 1.13                            | $[0.78, 1.62]$                   | 1.18                            | $[0.51, 1.9]$                     |
| $BR_{WW}^{ano}/BR_{WW}^{SM}$                     | 1.00                            | $[0.9, 1.12]$                    | 1.04                            | $[0.43, 2.0]$                     |
| $BR_{ZZ}^{ano}/BR_{ZZ}^{SM}$                     | 1.10                            | $[0.9, 1.4]$                     | 1.04                            | $[0.43, 2.0]$                     |
| $BR_{bb}^{ano}/BR_{bb}^{SM}$                     | 1.01                            | $[0.95, 1.05]$                   | 0.99                            | $[0.48, 1.3]$                     |
| $BR_{\tau\tau}^{ano}/BR_{\tau\tau}^{SM}$         | 1.01                            | $[0.9, 1.1]$                     | 1.11                            | $[0.42, 2.6]$                     |
| $\sigma_{gg}^{ano}/\sigma_{gg}^{SM}$             | 0.90                            | $[0.58, 1.35]$                   | 0.88                            | $[0.37, 2.4]$                     |
| $\sigma_{VBF}^{ano}/\sigma_{VBF}^{SM}$           | 1.01                            | $[0.9, 1.15]$                    | 0.99                            | $[0.9, 1.1]$                      |
| $\sigma_{VH}^{ano}/\sigma_{VH}^{SM}$             | 1.0                             | $[0.56, 1.5]$                    | 1.03                            | $[0.79, 1.6]$                     |

Best fit values and 90% CL allowed ranges for the combination of all available Tevatron and LHC Higgs data as well as TGV.