

Basics of high-power ultrafast fiber amplifiers

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Principles of fiber amplifiers

- Signal guiding, amplification
- Pump guiding and double-clad fibers
- Thermal aspects

Ultrafast fiber amplifiers

- Limiting effects in short pulse amplification
- LMA fibers and CPA amplification
- State of the art

Parabolic fiber amplifiers

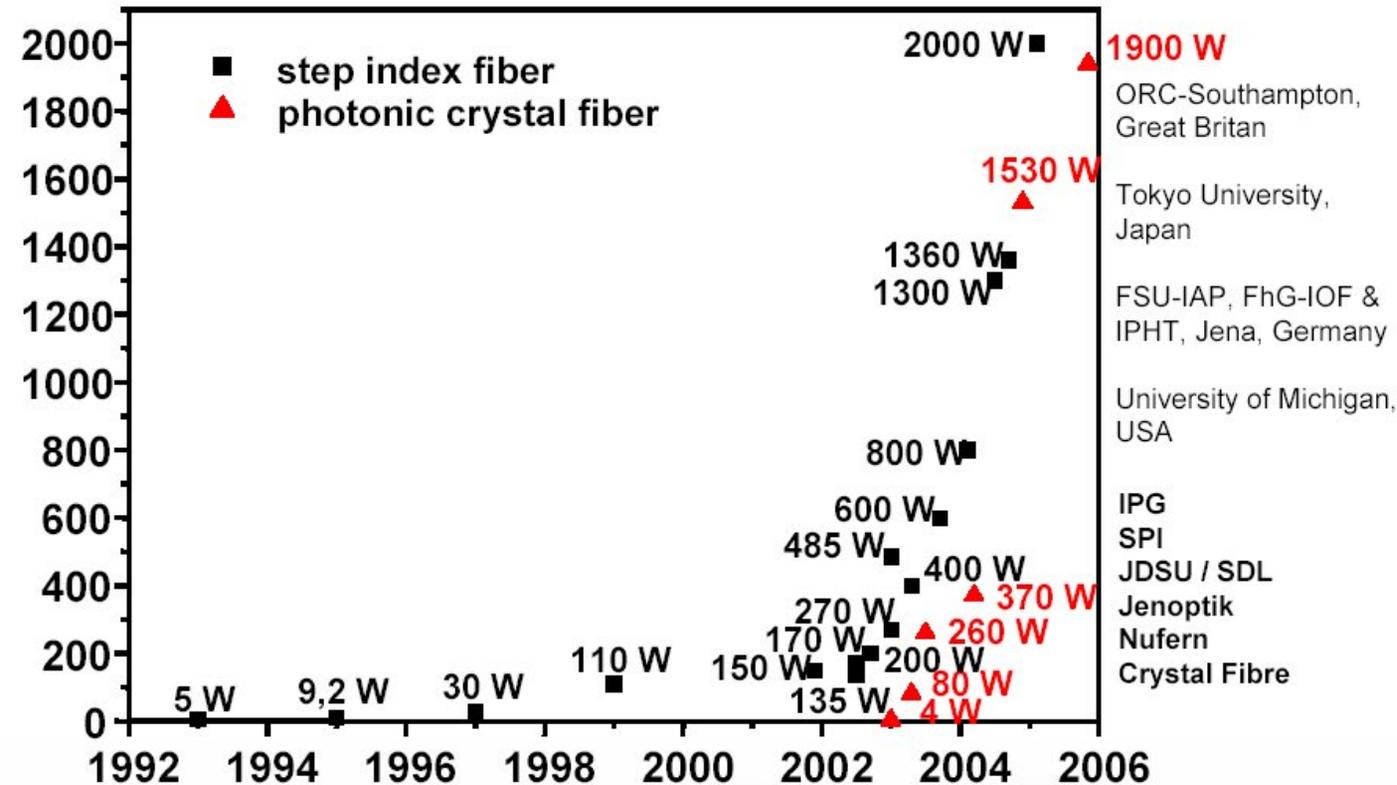
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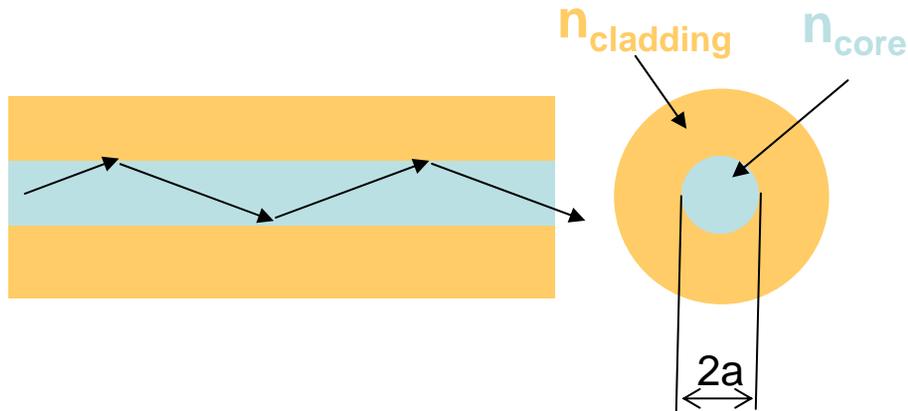
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CW fiber laser output power

Transverse monomode



$$V_{SI} = \frac{2\pi}{\lambda} \cdot a \cdot \sqrt{n_{\text{core}}^2 - n_{\text{cladding}}^2}$$

single-mode condition: $V_{SI} < 2.405$

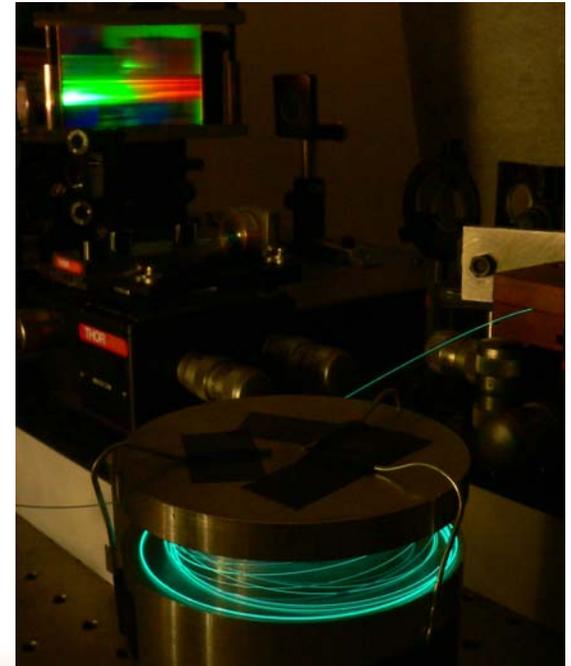


- Guiding by **Total Internal Reflection**
- Monomode operation : **index contrast, mode diameter**

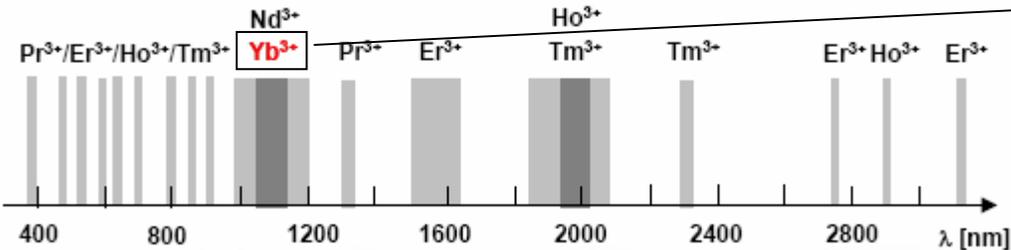


Advantages of fiber amplifiers

- High gain (interaction length and confinement)
- Immune against thermo-optical problems
- Excellent beam quality
- Integration in all-fiber systems



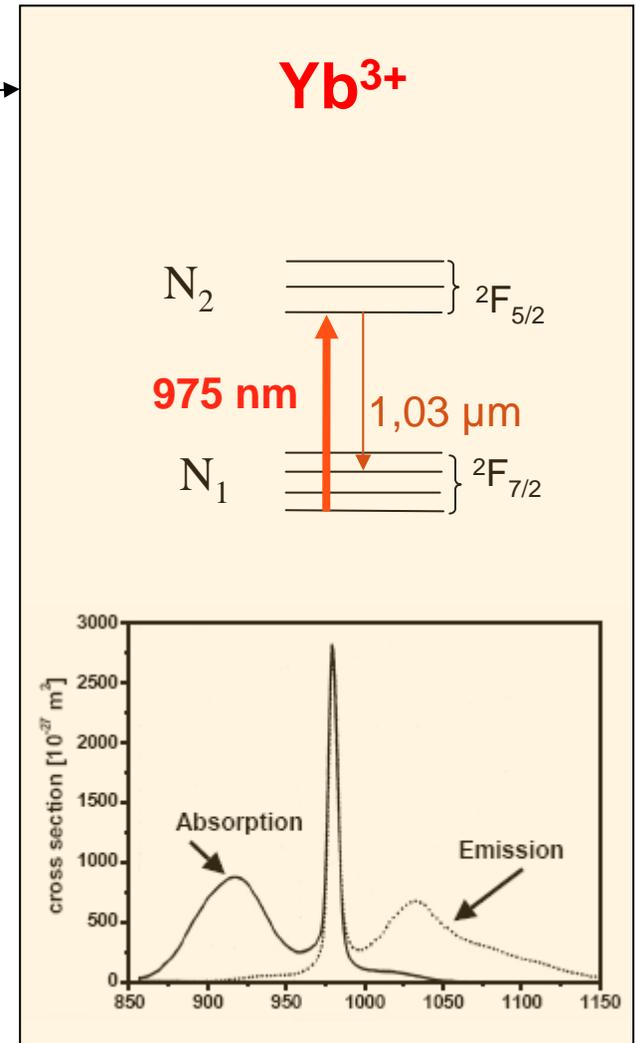
Emission bands of rare-earth dopants



Advantages of the Yb dopant

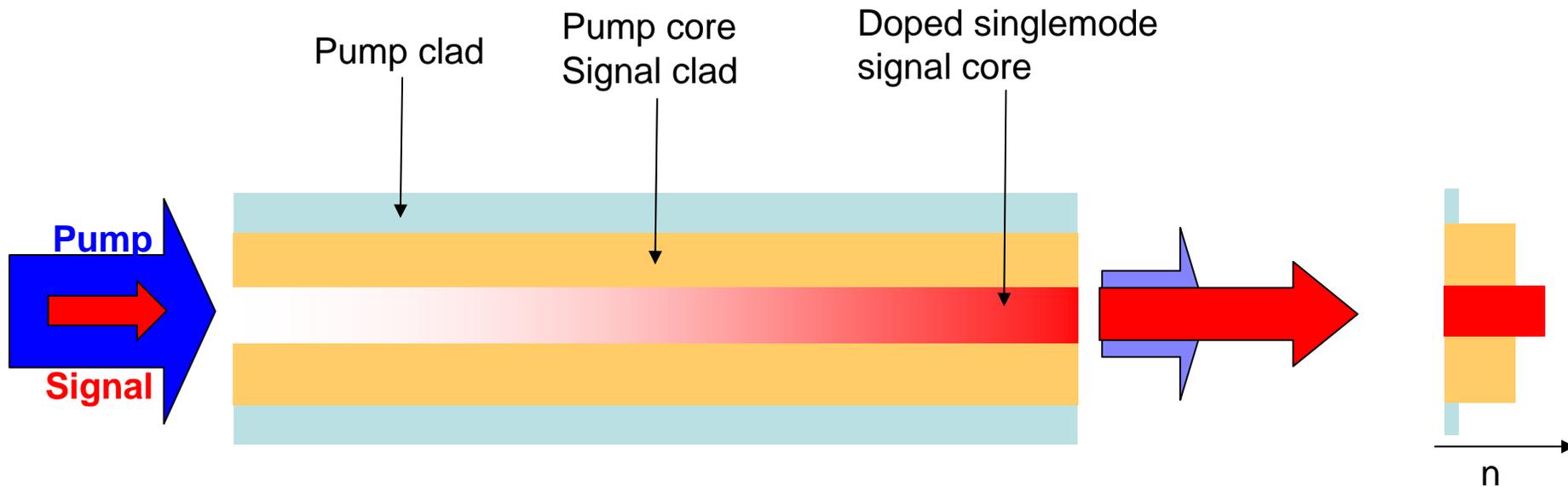
- No excited-state absorption
- Pumped by high-power diodes at 980 nm
- Broad emission bandwidth \rightarrow short pulses
- Low quantum defect \rightarrow high efficiency

low thermal load

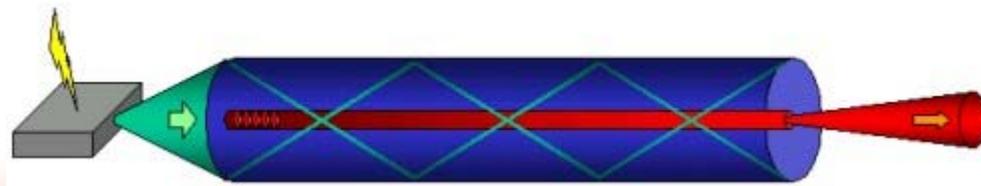


Problem : Pump must be diffraction-limited to be injected in singlemode core

Solution : Double clad fiber structure

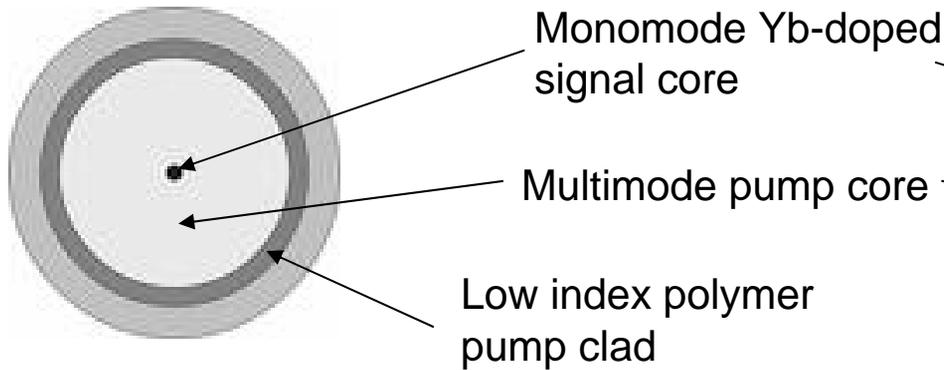


Highly efficient brightness conversion through stimulated emission

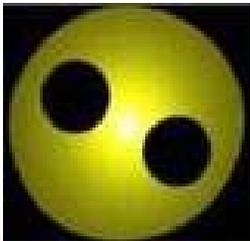


Geometry can be implemented through **doping** or **microstructuring**

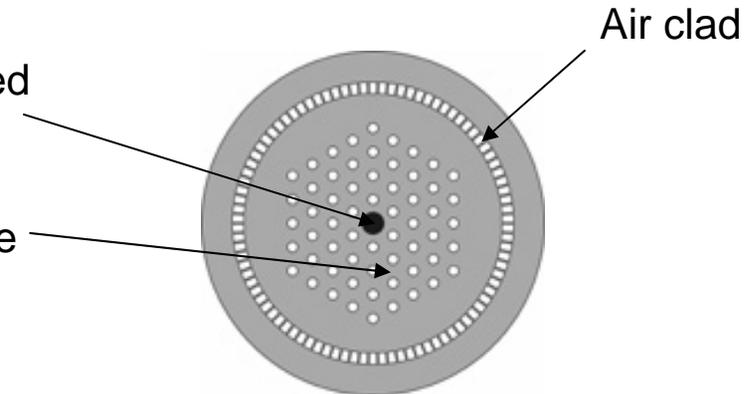
Double-clad fiber



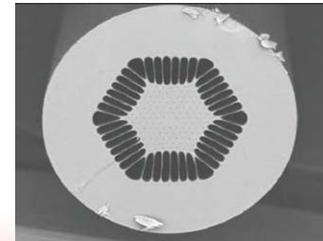
- Polymer pump clad
- Polarization-maintaining

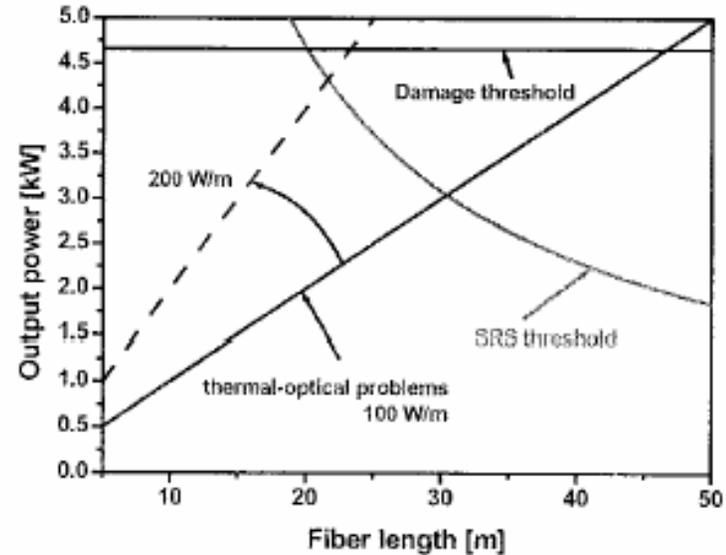
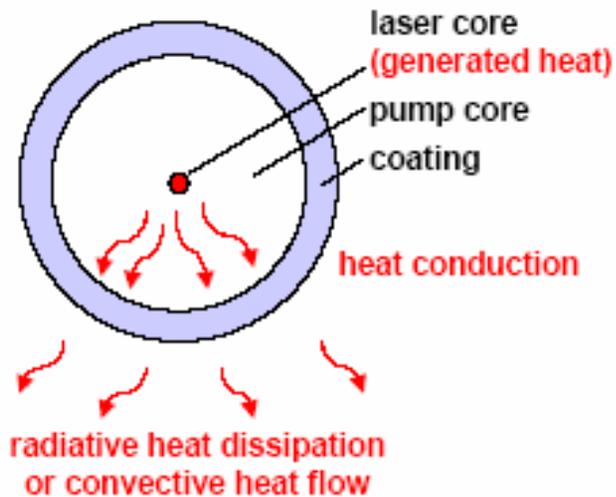


Air-clad fiber



- Higher NA for pump core
- No pump radiation in the polymer clad
- Larger MFD while singlemode





- Long interaction length: Pump absorption distributed
High surface to volume ratio

→ Multi kW monomode operation is possible

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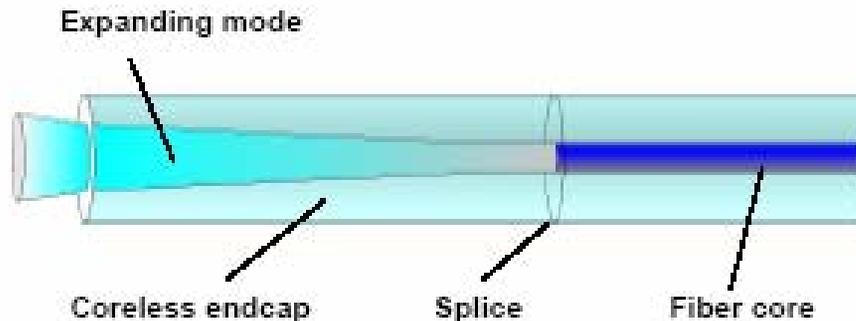
Fused silica

- Bulk damage (1 ns pulse) : 220 J/cm², **2.2 mJ** for 35 μm MFD

→ comparable to extractable energy

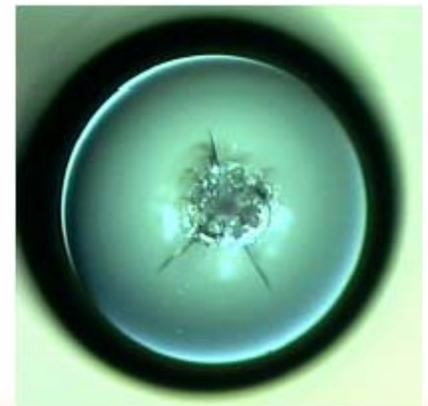
Surface damage (1 ns pulse) : 22 J/cm², **220 μJ** for 35 μm MFD

→ coreless endcaps (beam expanders)



$$\varepsilon_{sd} \approx 22 \times (\Delta\tau)^{0.4} \text{ J/cm}^2$$

(ns)
↓



- Interaction length, confinement, short pulses → **nonlinear effects**
- Spatial aspects: Self focusing (does not scale with MFD)

$$P_{crit} = \frac{\pi(0.61)^2 \lambda_0^2}{8n_0 n_2} \simeq 4\text{MW for silica @ } 1050 \text{ nm}$$

- Temporal aspects: Nonlinear Schrödinger equation with gain

$$i \frac{\partial u}{\partial z} = \frac{\beta_2}{2} \frac{\partial^2 u}{\partial t^2} - \gamma \left(\underbrace{|u|^2 u}_{\text{Self-phase modulation}} - T_R \frac{\partial |u|^2}{\partial t} \right) + i \frac{g}{2} u$$

Dispersion
Self-phase modulation
Raman effect
Gain

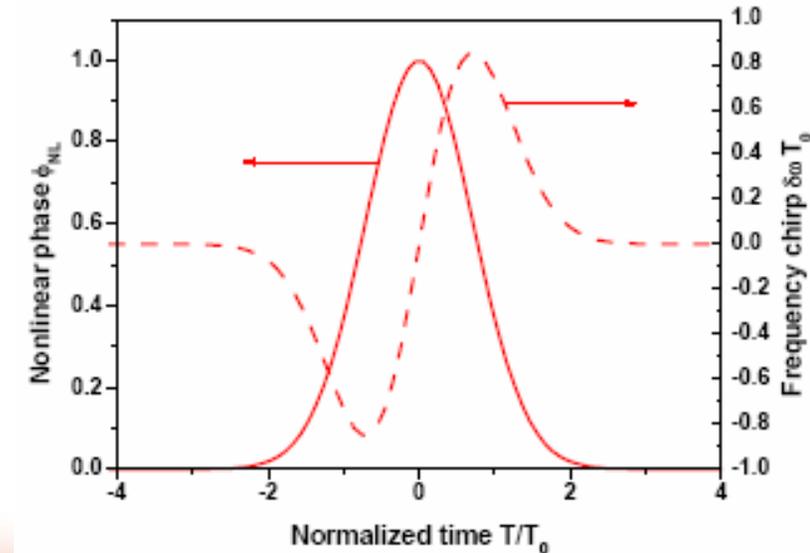
- Intensity-dependent refractive index
- Nonlinear phase-shift
- Frequency shift

$$n(t) = n_L + n_2 I(t)$$

$$B = \int_0^L \gamma P_{peak}(z) dz$$

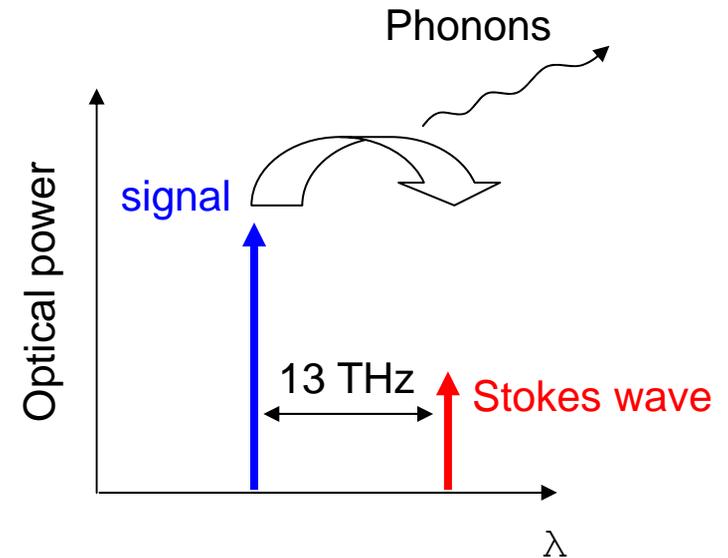
$$\delta\omega(T) = -\frac{\partial\phi_{NL}(z, T)}{\partial T} \propto \frac{L}{A_{eff}}$$

- Spectral broadening
- Nonlinear chirp, difficult to recompress



- Inelastic scattering of a signal photon to a stokes photon
- Energy transfer to acoustic modes
- Power threshold given by

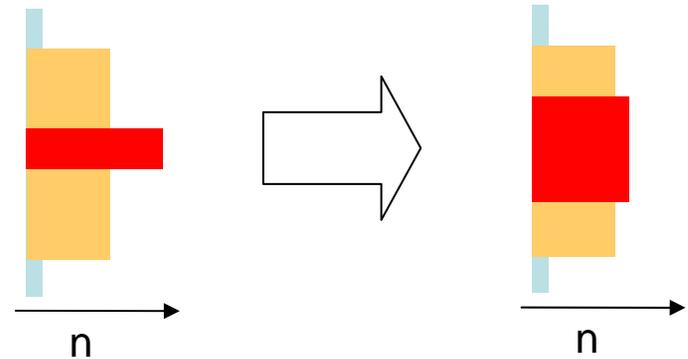
$$P_{cr} \approx \frac{16A_{eff}}{g_R L_{eff}}$$



- Temporal nonlinearity
- Optical damage

} ↓ when MFD ↑

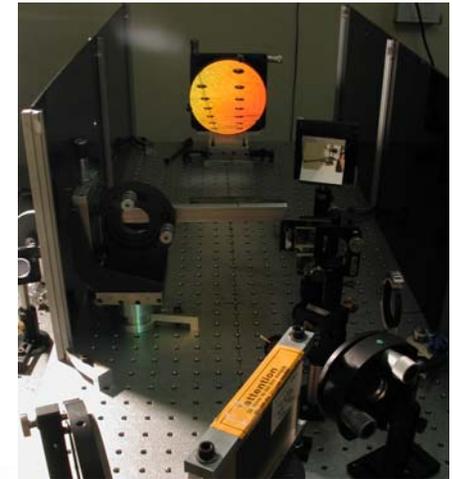
→ spread the intensity in space



- Index contrast ↓ to remain singlemode, Numerical aperture ↓
 - LMA often multimode, fundamental mode selection through coiling
 - Limited by bending loss
- MFD ~ 25 μm conventional, 35 μm for microstructured fibers

- Temporal nonlinearity
 - Optical damage
- } ↓ when peak power ↓

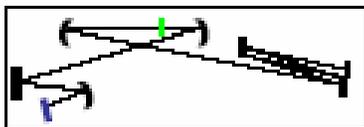
→ spread the intensity in time



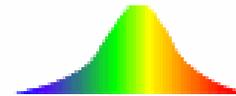
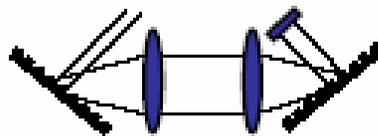
D. Strickland and G. Mourou, "Compression of amplified optical pulses,"
Opt. Comm. 56, 3, 219 (1985).



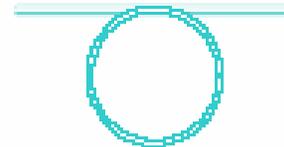
mode-locked laser



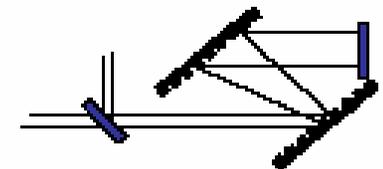
grating stretcher



amplifier

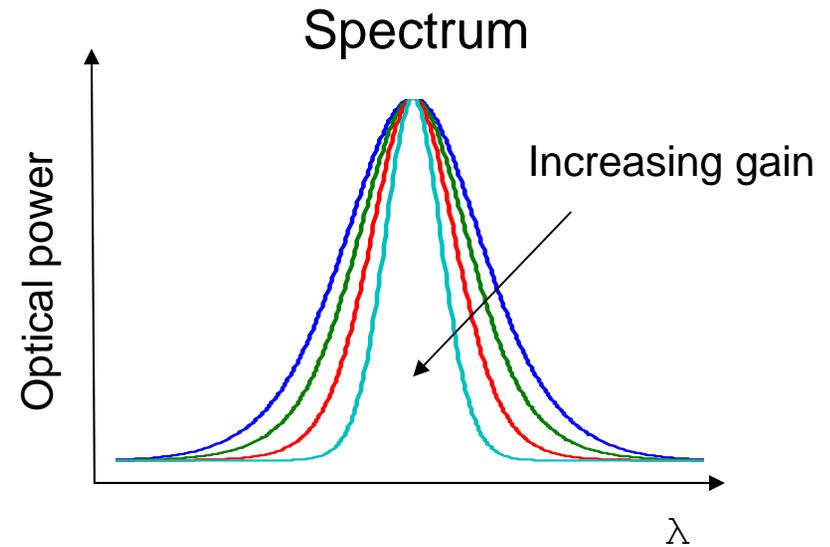


grating compressor



- Finite gain bandwidth of the medium causes gain narrowing

→ limits the pulsewidth (~100 fs for typical Yb fiber amplifier systems)



- Extractable energy $E = \varepsilon_{sat} A_{eff} \ln(G_0)$

~ mJ range for monomode LMA fibers

- **High average power** femtosecond fiber amplifier

Röser et al., Opt. Lett., vol. 30, no. 20 (2005) [Jena group \(J. Limpert\)](#)

131 W 220 fs 73 MHz

- **High energy** femtosecond fiber amplifier

Liao et al. CLEO 2006 postdeadline CPDB4 [Michigan group \(A. Galvanauskas\)](#)

500 μ J 520 fs 5 kHz

- **Review paper**

Tunnermann et al., Topics in applied physics vol. 96, pp.35-53 (2004)

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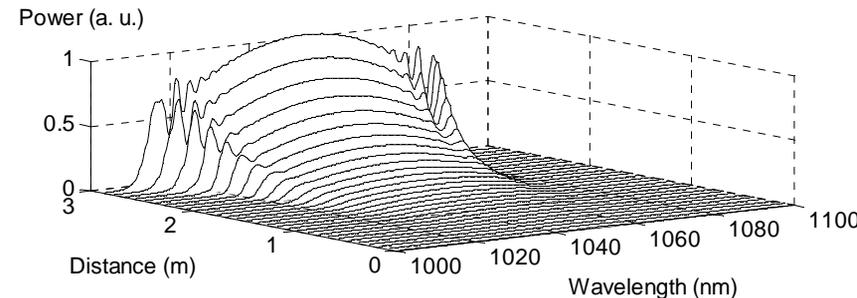
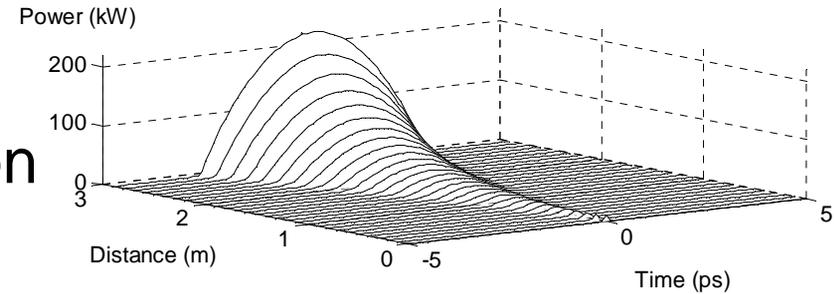
- **Nonlinear** pulse propagation with gain

$$i \frac{\partial u}{\partial z} = \frac{\beta_2}{2} \frac{\partial^2 u}{\partial t^2} - \gamma |u|^2 u + i \frac{g}{2} u$$

- $\beta_2 > 0$, $g > 0$ → **Parabolic asymptotic solution**
- Also known as « **similariton** » because of its self-similar propagation properties
- Limited by higher-order effects : Raman, Finite gain bandwidth, Third-order dispersion

Advantages of parabolic amplifiers

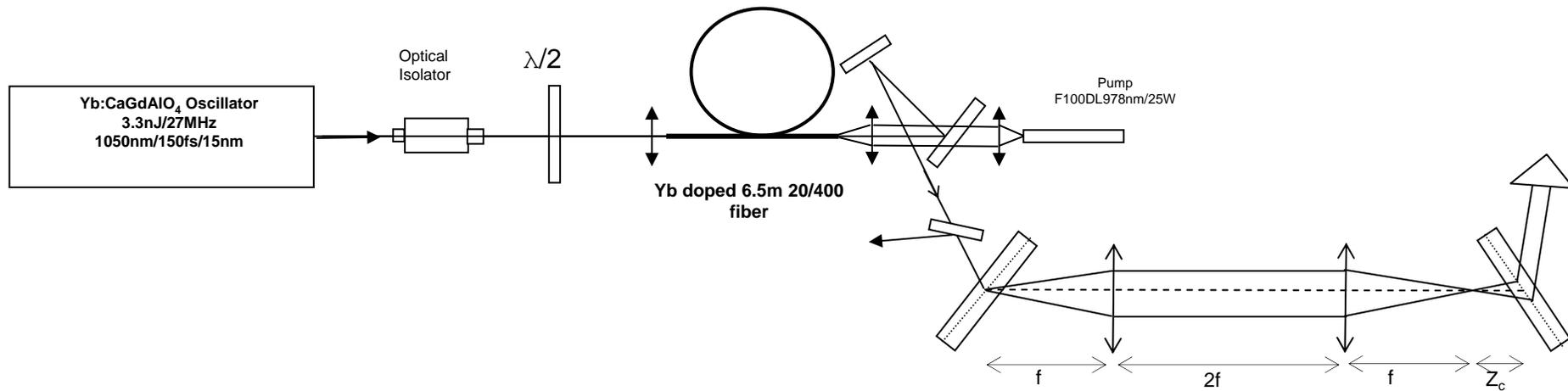
- Stable **nonlinear** pulse propagation
- Asymptotic regime
- Spectrum is generated **and**
- Linear (recompressible) chirp
- No stretcher



Drawbacks

- Limited to $\sim 1 \mu\text{J}$ with current fibers

Input : gaussian 100 fs 1 nJ
 Output : parabolic 5 ps 1 μJ
 (ideally recompressible down to 50 fs)



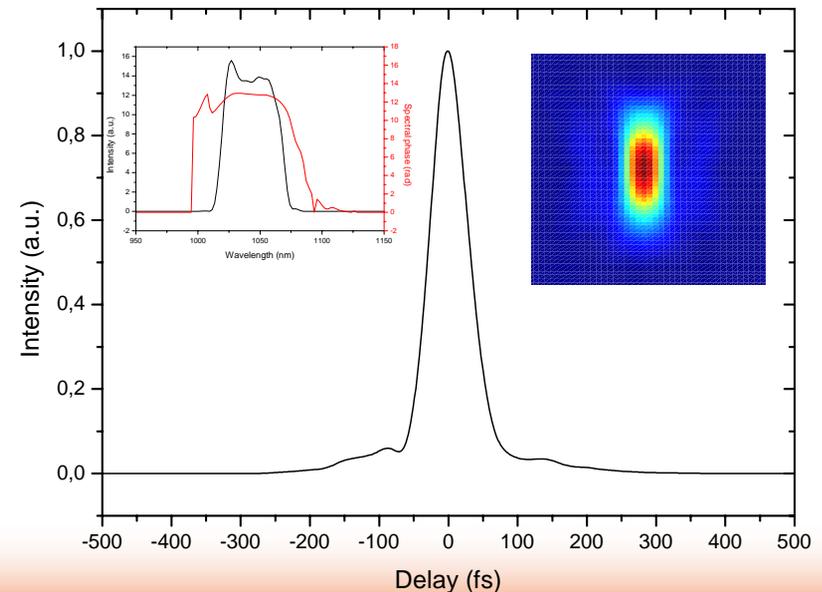
FWHM pulsewidth : **63 fs**

Pulse energy : 0.28 μ J

Peak power : **4 MW**

Average power : **7.5 W**

Rep rate : 27 MHz



- **High-harmonic generation** with a high repetition rate femtosecond **fiber-based** source

