



LHCb Combination of the CKM angle γ

Matthew Kenzie CERN

LAL Seminar

May 3, 2016

OP violation and the CKM matrix

2 The LHCb Experiment

3 CKM angle γ

4 LHCb Combination

5 Conclusion and Prospects



CP violation



- We live in a matter (and photon) dominated universe
- How does baryogenesis lead to a matter / antimatter asymmetry?
- CP violation is a crucial ingredient to this problem (Sacharov)
- CKM matrix is the one place in the SM with CP violation
- \blacktriangleright CPV in the SM $(\sim 10^{-20})$ does not nearly account for the observed baryon-photon ratio($\sim 10^{-10})$
- ▶ New sources of CP violation would be a clear indiction of New Physics (NP)



CKM matrix



 \blacktriangleright Quark mixing in the SM is described by the 3 \times 3 unitary CKM matrix



The matrix elements determine the transition probability



• Parameterised by three mixing angles (θ_{12} , θ_{13} , θ_{23}) and a CP violating phase (δ)



CKM matrix

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The CKM matrix exhibits a clear hierachy, sin(θ₁₃) << sin(θ₂₃) << sin(θ₁₂) << 1, so often expressed in Wolfenstein parameterisation (A, λ, ρ, η)</p>

Wolfenstein parametrisation

$$V = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} = \begin{pmatrix} 1 - \lambda^2/2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \lambda^2/2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix} + \mathcal{O}(\lambda^4)$$

- Hierachy gives very distinctive behaviour to the flavour sector of the SM which gives strong constraints on NP
- CKM matrix gives the only source of CP violation in the SM ($m_{\nu} = \theta_{QCD} = 0$)







CKM picture is now well verified

- Any discrepancies would be of great importance
- CKM angle γ is the *least well known* constraint





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The Ultimate Test



 \blacktriangleright Not just via direct / indirect disagreement but many constraints from new physics in neutral mixing require input of γ



The Ultimate Test



- \blacktriangleright LHCb expected precision in 2029 $\sim\pm1^\circ$
- \blacktriangleright Belle II expected precision in 2023 $\sim\pm2^\circ$



2. The LHCb Experiment



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LHC, CERN, Geneva





LHCb Detector



A single arm forward spectrometer



2. The LHCb Experiment

LHCb Detector

- A single arm forward spectrometer
- A factory for beauty and charm decays
- Acceptance range $2 < \eta < 5$
- ► 100K $b\overline{b}$ pairs produced per second ($10^4 \times B$ factories)
- $\sigma(b\overline{b}) = 284 \pm 54 \mu b$
- $\sigma(c\overline{c}) \approx 20 \times \sigma(b\overline{b})$

LHCb performance paper - [arXiv:1412.6352]

- IP resolution $\approx 20 \mu \mathrm{m}$
- *p* resolution $\approx 0.5\%$
- au resolution pprox 45 fs
- Calorimeter ID for γ , e, π^0
- Particle ID $\epsilon(K) \sim 95\%$ with 5% $\pi \to K$ mis-id
- Muons $\epsilon(\mu)\sim$ 97% with $(1-3)\%\pi
 ightarrow\mu$ mis-id





2. The LHCb Experiment

LHCb Trigger





- Allow detector alignment and calibration in real time!
- In turn means online and offline reconstruction are identical
- Allows performing of many analyses online
- Allows high readout rate
- High efficiency for a broad range of topics





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$$\gamma = \arg \left(- \frac{\textit{V}_{ud} \textit{V}_{ub}^*}{\textit{V}_{cd} \textit{V}_{cb}^*} \right)$$

- $\blacktriangleright \gamma$ is known very well
- Can be determined entirely from tree decays
 - Unique property among all CP violation parameters
 - Hadronic parameters can be determined from data
- Neglible theoretical uncertainty (Zupan and Brod 2013)

Theory uncertainty on γ

 $\delta\gamma/\gamma pprox \mathcal{O}(10^{-7})$ - [arXiv:1308.5663]

- γ can probe for new physics at extrememly high energy scales (Zupan)
 - (N)MFV new physics scenarios: $\sim \mathcal{O}(10^2) \,\, {
 m TeV}$
 - gen. FV new physics scenarios: $\sim O(10^3)$ TeV

γ from experiment

- γ is NOT known very well
- It is quite challenging to measure
- The decay rates are small

Branching ratio for suppressed γ mode BR(B⁻ \rightarrow DK⁻, D \rightarrow π K) \approx 2 \times 10⁻⁷

- \blacktriangleright Small interference effect typically $\sim 10\%$
- Fully hadronic decays hard to trigger on
- Many channels have a $K_{\rm S}^0$ in the final state low efficiency
- Many channels have a π^0 in the final state very hard at LHCb
- Many different decay channels, many observables and many hadronic unknowns make it statistically challenging



Methods to measure γ

Reconstruct the D^0/\overline{D}^0 in a final state accesible to both to acheieve interference



GLW method

- CP eigenstates e.g. $D \rightarrow KK$
- Gronau, London, Wyler (1991)
- ADS method
 - CF or DCS decays e.g. $D \rightarrow K\pi$
 - Atwood, Dunietz, Soni (1997,2001)
- GGSZ method
 - 3-body final states e.g. $D \to K^0_{
 m S} \pi \pi$
 - Giri, Grossman, Soffer, Zupan (2003)

- [Phys. Lett. B253 (1991) 483]
- [Phys. Lett. B265 (1991) 172]
- [Phys. Rev. D63 (2001) 036005]
- [Phys. Rev. Lett. 78 (1997) 3257]
 - [Phys. Rev. D68 (2003) 054018]

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The cartesian coordinates



An example GLW analysis - $B^{\pm} \rightarrow D^0 \ K^{\pm}$, $D^0 \rightarrow K^+ \ K^-$



$$A_{h}^{f} = \frac{\Gamma(B^{-} \to [f]_{D}h^{-}) - \Gamma(B^{+} \to [f]_{D}h^{+})}{\Gamma(B^{-} \to [f]_{D}h^{-}) + \Gamma(B^{+} \to [f]_{D}h^{+})}$$

$$R^f_{K/\pi} = rac{\Gamma(B^\pm o [f]_D K^\pm)}{\Gamma(B^\pm o [f]_D \pi^\pm)}$$

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An example ADS analysis - $B^{\pm} \rightarrow D^0 \ K^{\pm}$, $D^0 \rightarrow K^{\pm} \ \pi^{\pm}$

Favoured mode



[arXiv:1603.08993]

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An example ADS analysis - $B^{\pm} \rightarrow D^0 \ K^{\pm}$, $D^0 \rightarrow K^{\pm} \ \pi^{\pm}$

Suppressed mode



LHCb CER

HC-D



An example ADS analysis - $B^\pm o D^0$ K^\pm , $D^0 o K^\pm$ π^\pm

- Define observables as yield ratios (many systematics cancel)
- Along with the GLW observables build a system of equations to overconstrain the parameters

ADS ratios of favoured to suppressed

$$R_{ADS}^{\bar{f}} = \frac{\Gamma(B^- \to [\bar{f}]_D h^-) + \Gamma(B^+ \to [f]_D h^+)}{\Gamma(B^- \to [f]_D h^-) + \Gamma(B^+ \to [\bar{f}]_D h^+)}$$

Corresponding charge asymmetries

$$A_{\rm ADS}^{\bar{f}} = \frac{\Gamma(B^- \to [\bar{f}]_D h^-) - \Gamma(B^+ \to [f]_D h^+)}{\Gamma(B^- \to [\bar{f}]_D h^-) + \Gamma(B^+ \to [f]_D h^+)}$$

▶ Relatively trivial extension to multibody *D* decays $(D \rightarrow 4\pi, D \rightarrow K3\pi, D \rightarrow KK\pi^0, D \rightarrow \pi\pi\pi^0, D \rightarrow K\pi\pi^0)$, multibody *B* decays $(B^{\pm} \rightarrow DK^{\pm}\pi^{+}\pi^{-})$ and other initial *B* states $(B^0 \rightarrow DK^{*0})$

An example GGSZ analysis

- ▶ Requires a self-conjugate 3-body final state $(D^0 \rightarrow K^0_S \pi^- \pi^+, D^0 \rightarrow K^0_S K^- K^+)$
- The basic idea is to perform a GLW/ADS type analysis in each bin of the D decay phase space
- Compare Dalitz distribution for B⁺ and B⁻
 - \blacktriangleright Model dependent: use a Dalitz model describing all the intermediate resonances and fit for $x_{\pm}, \, y_{\pm}$
 - Model independent: define bins which maximise sensitivity to x_{\pm} , y_{\pm}



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An example GGSZ analysis



An example GGSZ analysis

- GGSZ analyses have excellent standalone sensitivity with a single solution
- \blacktriangleright Can trivially extend the methodology for neutral $B^0
 ightarrow D^0 K^{*0}$ decays

There are also other methods

- \blacktriangleright Time-dependent method using $B^0_s \to D^-_s K^+$
 - Large interference occurs via B_s^0 mixing (requires knowledge of $2\beta_s$)
 - Time dependent, flavour tagged analysis unique to LHCb [arXiv:1407.6127]

- GLS method
 - Grossman, Ligeti, Soffer (2003) [Phys. Rev. D67 (2003) 071301]
 - Uses ADS-like method with singly Cabibbo suppressed D decays (e.g. $D^0 o K^0_S K \pi$)
 - Poor sensitivity with current statistics

4. LHCb Combination

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LHCb γ combination inputs

B decay D d		D decay	Туре	$\int \mathcal{L}$	Ref.
LHCb Inputs	$B^+ \rightarrow DK^+$	$D \rightarrow hh$	GLW/ADS	$3{\rm fb}^{-1}$	[arXiv:1603.08993]
	$B^+ ightarrow DK^+$	$D ightarrow h\pi\pi\pi$	GLW/ADS	$3{\rm fb}^{-1}$	[arXiv:1603.08993]
	$B^+ \rightarrow DK^+$	$D ightarrow hh\pi^0$	GLW/ADS	$3 \mathrm{fb}^{-1}$	[arXiv:1504.05442]
	$B^+ \rightarrow DK^+$	$D ightarrow K^0_{ m S} hh$	GGSZ	$3 {\rm fb}^{-1}$	[arXiv:1405.2797]
	$B^+ \rightarrow DK^+$	$D ightarrow K_{ m S}^0 K \pi$	GLS	$3 {\rm fb}^{-1}$	[arXiv:1402.2982]
	$B^0 \rightarrow D^0 K^{*0}$	$D \to K \pi$	ADS	$3 {\rm fb}^{-1}$	[arXiv:1407.3186]
	$B^+ \rightarrow DK^+\pi\pi$	$D \rightarrow hh$	GLW/ADS	$3 {\rm fb}^{-1}$	[arXiv:1505.07044]
	$B_s^0 \rightarrow D_s^{\mp} K^{\pm}$	$D_s^+ ightarrow hhh$	TD	$1{\rm fb}^{-1}$	[arXiv:1407.6127]
	$B^0 ightarrow D^0 K^+ \pi^-$	D ightarrow hh	GLW-Dalitz	$3 \mathrm{fb}^{-1}$	[arXiv:1602.03455]
	$B^0 ightarrow D^0 K^{*0}$	$D ightarrow K^0_{ m S} \pi \pi$	GGSZ	$3 \mathrm{fb}^{-1}$	[arXiv:1604.01525]
Decay Parameters		Source		Ref.	
Auxilliary Inputs	$D^0 - \overline{D}^0$ mixing		HFAG	-	[arXiv:1412.7515]
	$D \to K \pi \pi \pi$	$(\delta_D, \kappa_D, r_D)$	CLEO+LHCb	-	[arXiv:1602.07430]
	$D \rightarrow \pi \pi \pi \pi$	(F^+)	CLEO	-	[arXiv:1504.05878]
	$D \rightarrow K \pi \pi^0$	$(\delta_D, \kappa_D, r_D)$	CLEO+LHCb	-	[arXiv:1602.07430]
	$D \rightarrow hh\pi^0$	(F^+)	CLEO	-	[arXiv:1504.05878]
	$D \rightarrow K^0_S K \pi$	(δ_D, κ_D)	CLEO	-	[arXiv:1203.3804]
	$D \rightarrow K_{S}^{0}K\pi$	(<i>r</i> _D)	CLEO	-	[arXiv:1203.3804]
	$D \rightarrow K^0_S K \pi$	(<i>r</i> _D)	LHCb	-	[arXiv:1509.06628]
	$B^0_{\circ} \rightarrow D^0 K^{*0}$	$(\kappa_B, \bar{R}_B, \bar{\Delta}_B)$	LHCb	-	[arXiv:1602.03455]
	$ B_s^0 \to D_s^+ K^-$	(ϕ_s)	LHCb	-	[arXiv:1411.3104]
Com	bination:			[LHCb-CONF-2016-001]	

New or updated since last combination

4. LHCb Combination

LHCb γ Combination

- Combination of all $B \rightarrow DK$ -like modes
 - [LHCb-CONF-2016-001]
- Paper to follow soon with information on $B o D\pi$ modes also
- Nominal results with a frequentist Feldman-Cousins "plugin" procedure
- 71 observables and 32 free parameters
 - $p(\chi^2, N_{\rm dof}) = 87.6\%$
 - $p(toys) = (87.0 \pm 0.2)\%$

LHCb γ Combination

- Nominal result: $\gamma = (70.9^{+7.1}_{-8.5})^{\circ}$
- Uncertainty < 10° is better than combined *B* factories
- ► The most precise single experiment measurement of *γ*
- LHCb combination paper expected later this year

LHCb γ Combination

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<u>Naive statistical treatement</u> (profile likelihood method) - plots for demonstrative purposes only

4. LHCb Combination

LHCb γ Combination

<u>Naive statistical treatement</u> (profile likelihood method) - plots for demonstrative purposes only

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Prospects

- ▶ With Run II of the LHC underway and Belle II starting soon the prospects look good
- We can reasonably expect to half the experimental uncertainty on \(\gamma\) in the next 3 years
- \blacktriangleright We can reasonably expect to have $\sim 1^\circ$ precision in the next 5-7 years
- Current systematic effects are relatively small
 - GLW/ADS
 - instrumental charge asymmetries
 - PID calibration
 - GGSZ
 - efficiency correction over the Dalitz plane
 - Time-dependent
 - Decay time resolution
 - Decay time acceptance
 - Knowledge of Δm_s , $\Delta \Gamma_s$, Γ_s
- Tree measurements of γ will not be systematically limited for a long time (not at 100 times the current dataset)

This does not include smart new ideas which people often have

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Prospects

- ▶ We are approaching the first tree-level precision measurement of the CKM triangle
- Direct measurements of $|V_{ub}|$ play a crucial role in this as well

[arXiv:1309.2293]

5. Conclusion and Prospects

Conclusions

- CKM matrix is incredibly successful description of the quark sector in the SM
- Measurements of CKM elements are becoming increasingly precise
- ▶ Finding new sources of CP violation can lead us to New Physics
- \blacktriangleright CKM angle γ is one of the only CP measurements accesible with tree-level decays
 - Theoretically very clean
 - Experimentally challenging
- LHCb has the worlds most precise single experiment measurement and dominates the world average
 - $\gamma = (70.9^{+7.1}_{-8.5})^{\circ}$
- ▶ The future looks incredibly bright with the prospect of reducing the direct measurement uncertainty by a factor of 10
 - This will compete with the indirect precision (which assumes the SM)

