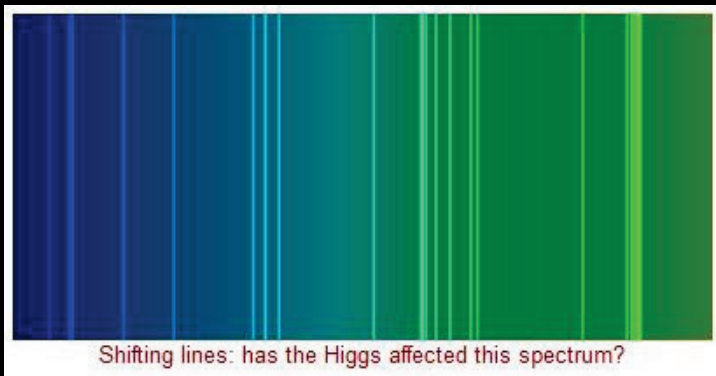


# Probing the Atomic Higgs force



CD, R. Ozeri, G. Perez, Y. Soreq  
*hep-ph: 1601.05087 + in progress*  
CD, Y. Soreq *hep-ph:1602.04838*

Cédric Delaunay  
CNRS/LAPTh  
Annecy-le-Vieux  
France



LAL Seminar  
1-04-2016 | Orsay

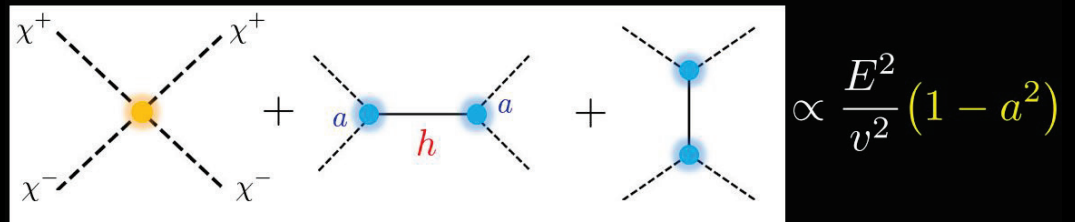
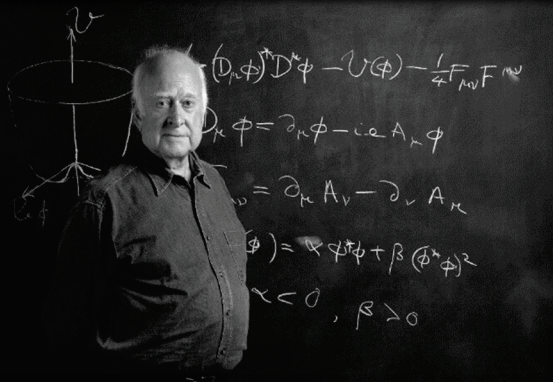
# *Outline*

1. The Higgs Mechanism and the Flavor Puzzle
2. Higgs Force in Atoms
3. Probing Higgs Couplings with Isotope Shift
4. The Weak Force and New Physics

*The Higgs Mechanism  
and the Flavor Puzzle*

# The Higgs Mechanism

- breaks EW symmetry:  $SU(2)_L \times U(1)_Y \rightarrow U(1)_{\text{QED}}$



ATLAS+CMS:  $|a - 1| \lesssim \mathcal{O}(10\%)$

ATLAS-CONF-2015-044

- provides charged fermion masses:

in the SM:  $m_f = y_f \times v$

# The flavor Puzzle

- Charged fermion masses are highly hierarchical:

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- The origin of this hierarchy is unknown, despite a host of precision flavor measurements.
- Within the SM, it is assumed to originate from hierarchical Higgs-to-fermion couplings:

$$y_f^{\text{SM}} \propto m_f$$

*How well can we test?*

# Alternative Approaches

- Froggatt-Nielsen like: lighter fermions couple to higher powers of the Higgs

Giudice-Lebedev '08  
Bauer-Carena-Gemmler '15

$$\frac{m_f}{v} \propto \frac{v^{n_f}}{\Lambda^{n_f}} \quad y_f \propto (1 + n_f) \frac{m_f}{v}$$

solving the flavor puzzle:  $n_f \sim \log(m_f/v)$

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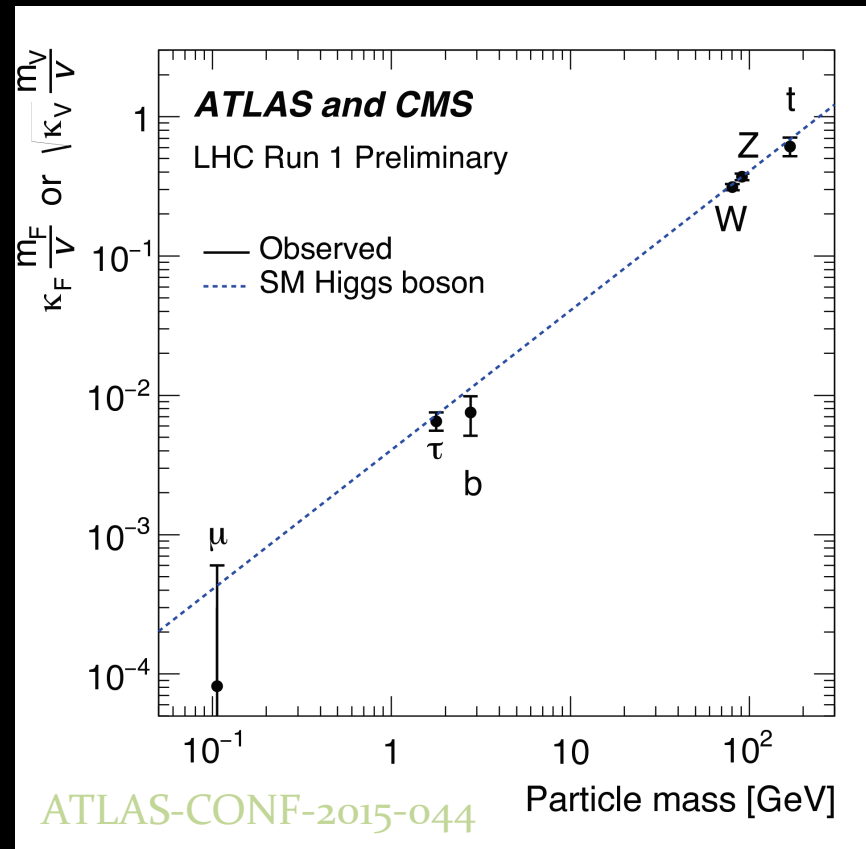
solving the flavor puzzle:  $n_f \sim \log(m_f/v)$

- First two generation masses do not come from the Higgs mechanism at all!  
(technicolor?)

Ghosh-Gupta-Perez '15  
Altmannshofer-Gori-Kagan-Silvestrini-Zupan '15



# The Flavor Puzzle at LHC

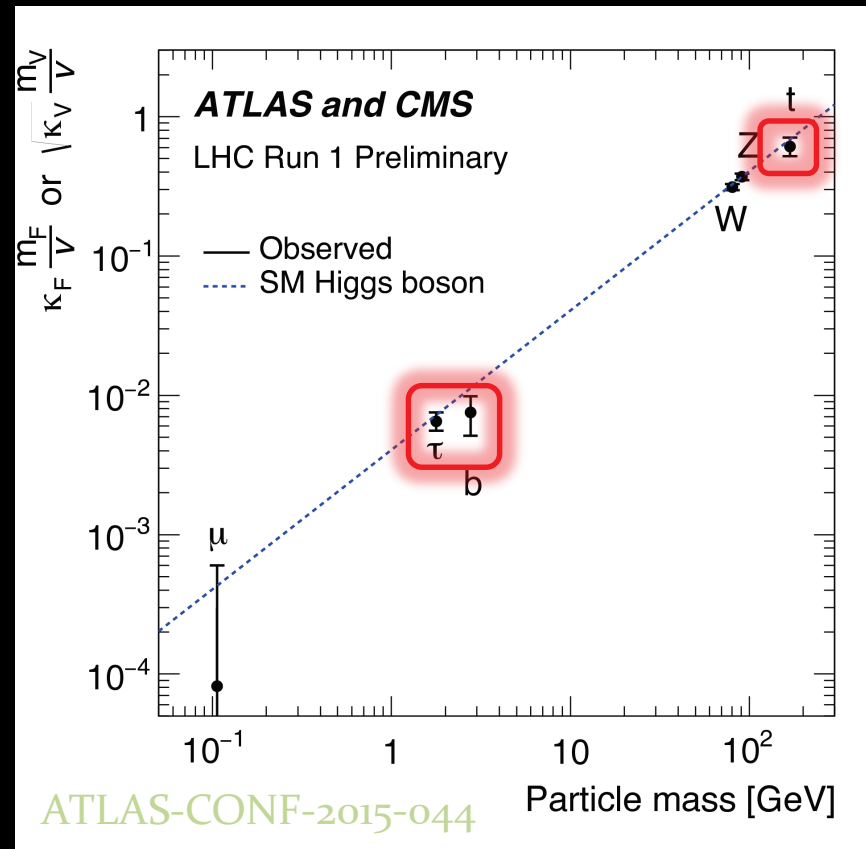


# The Flavor Puzzle at LHC

ATLAS+CMS  
Higgs-signal/SM:

$\mu^{\tau\tau}$	$1.12^{+0.25}_{-0.23}$
$\mu^{bb}$	$0.69^{+0.29}_{-0.27}$

$\mu_{ttH}$	$2.3^{+0.7}_{-0.6}$
-------------	---------------------



→ the Higgs mechanism is likely to be  
the dominant source of 3<sup>rd</sup> generation masses

# The Flavor Puzzle at LHC

There is an opportunity to probe  $c$ -coupling directly, thanks to charm-tagging:

in  $VH$  production

Perez-Soreq-Stamou-Tobioka '15

in  $Hc$  production

Isidori-Goertz '15

Other probes exist:

- $h \rightarrow J/\psi\gamma$

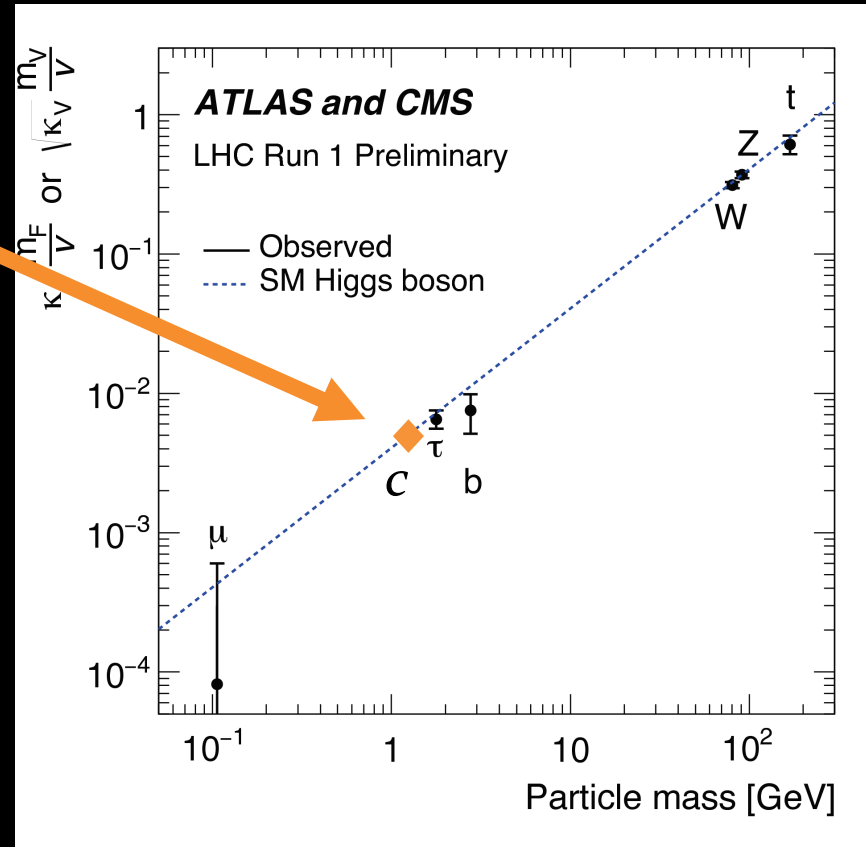
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- global fits

CD-Golling-Perez-Soreq '13

- $\Gamma_h \leq 1.7 \text{ GeV}$

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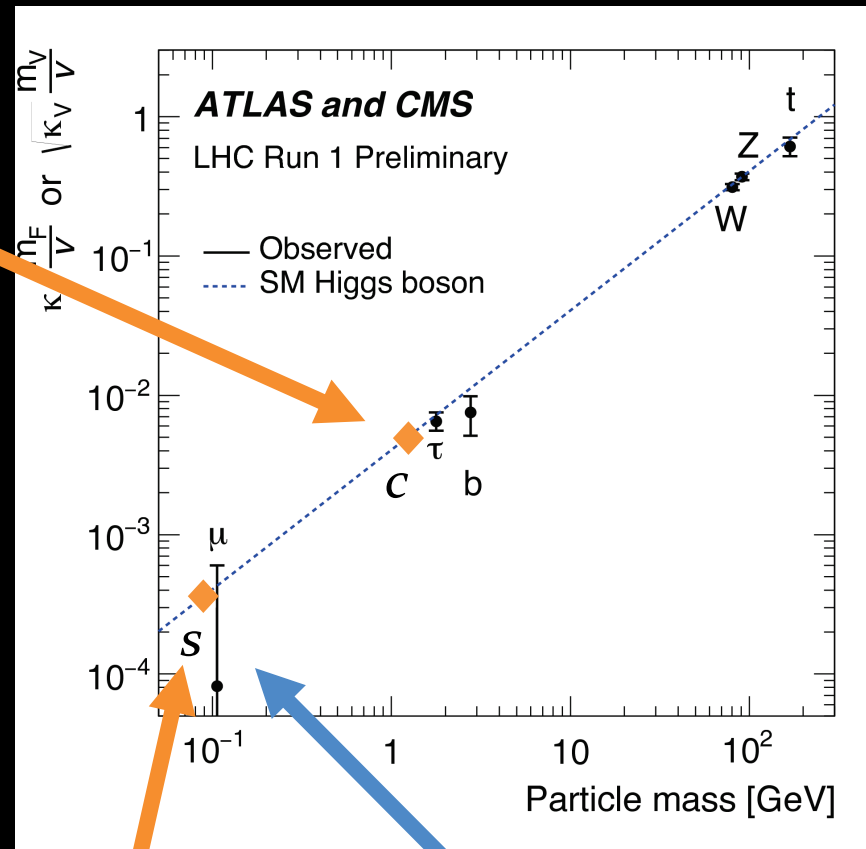
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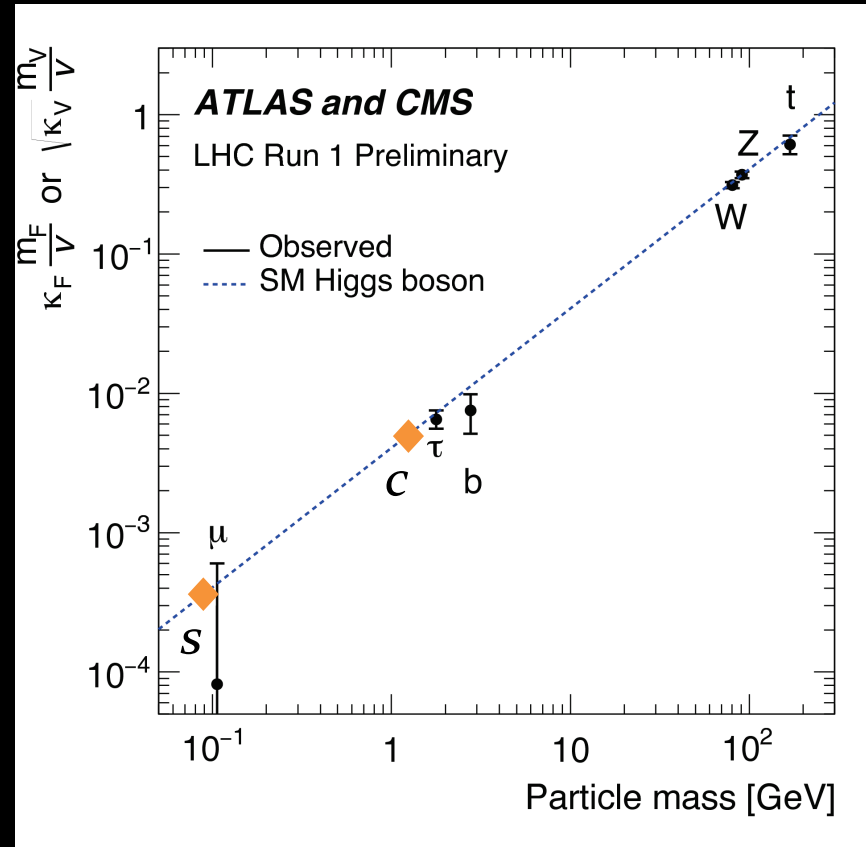
Kagan et al. '14

Sensitivity to muon-coupling, with high-enough luminosity

ATL-PHYS-PUB-2014-016

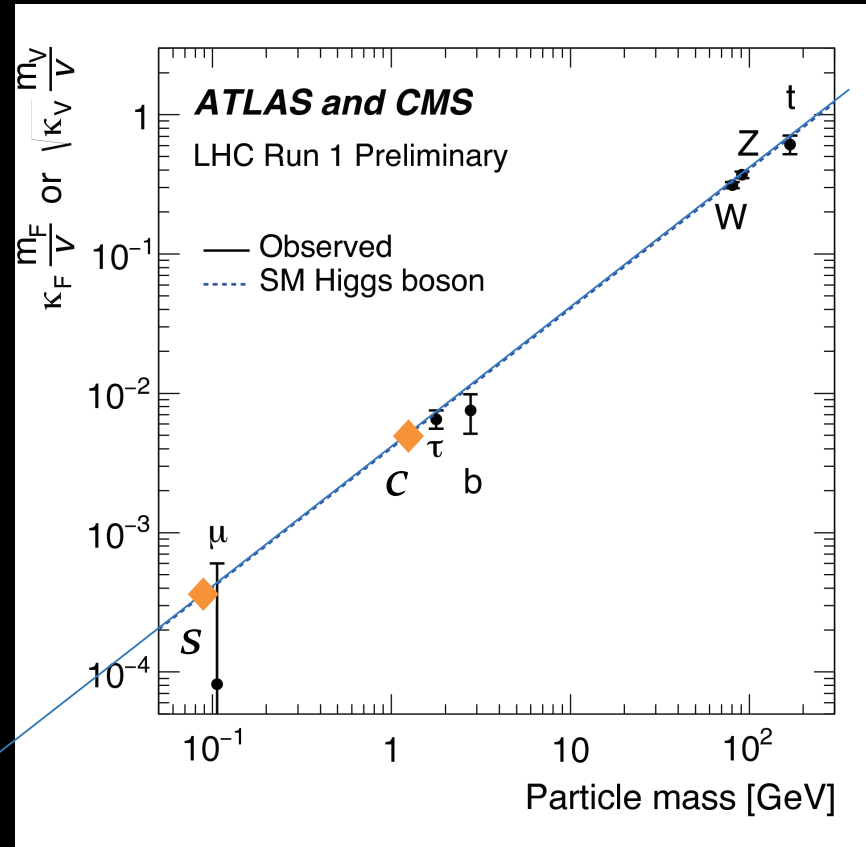
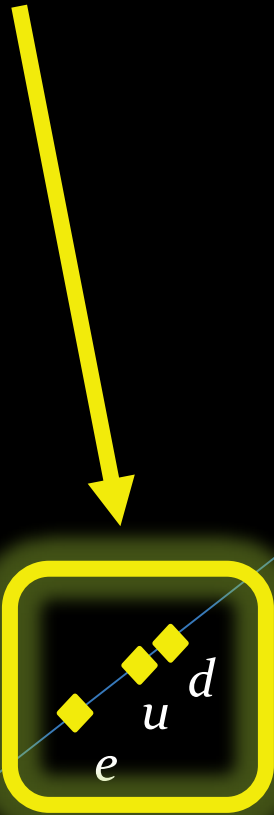
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What about  $e, u, d$ ?



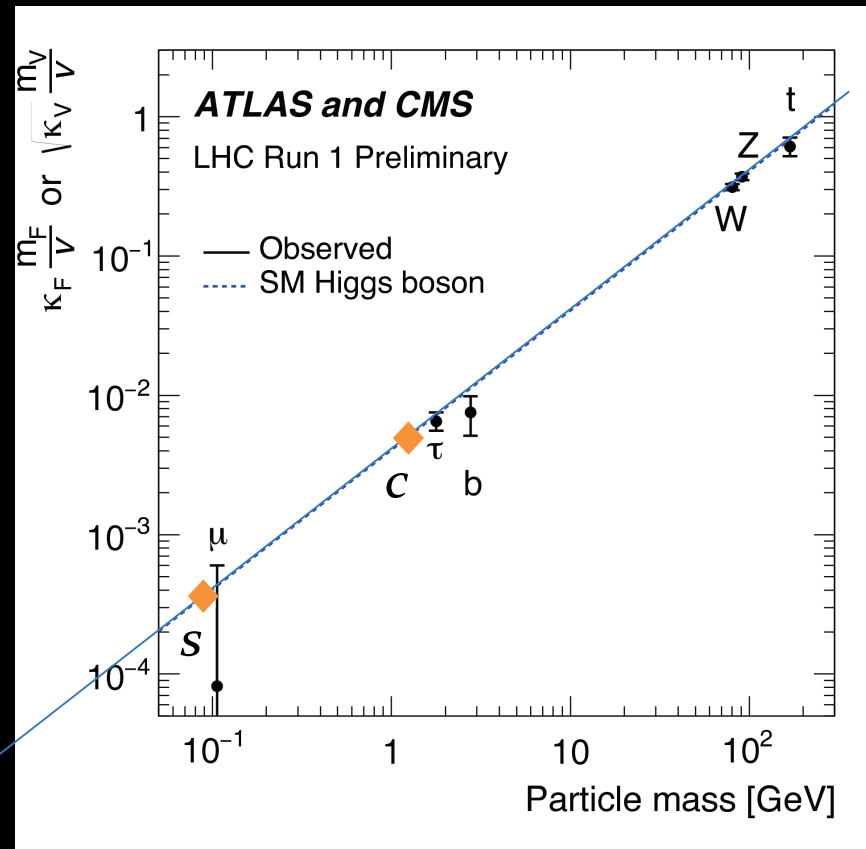
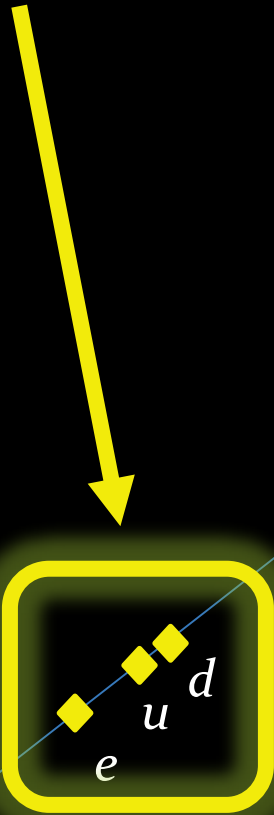
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What about  $e, u, d$ ?



[stable nuclei]  
[chemistry]

Probing the couplings to the building blocks of matter is an important test of the Higgs mechanism

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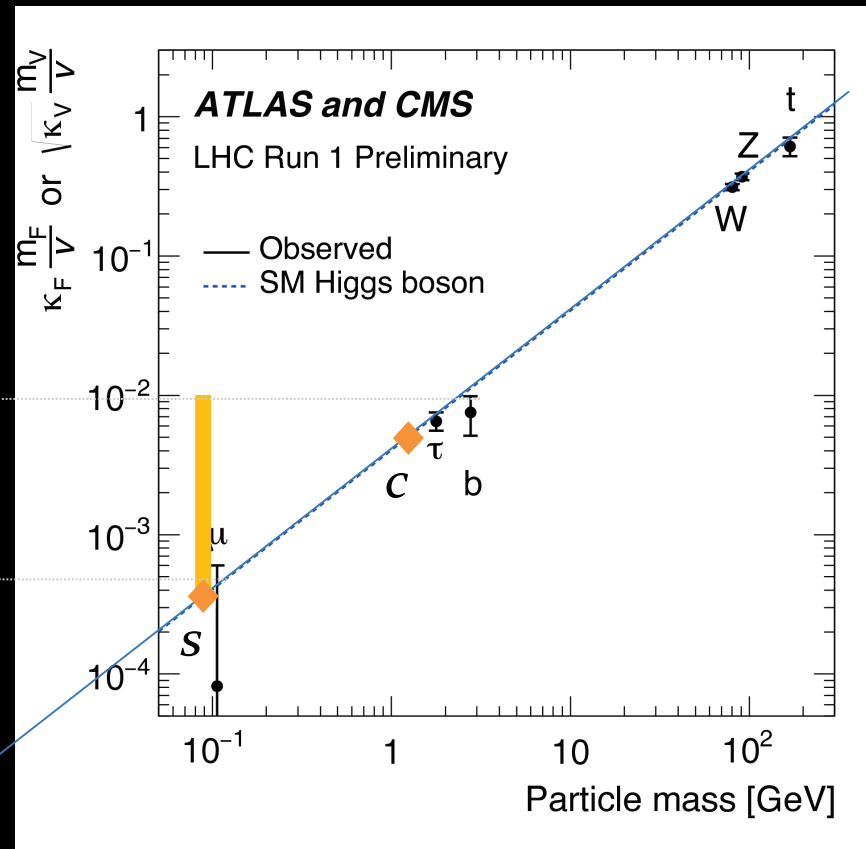
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Higgs-to-light-fermion couplings could be much larger than the SM prediction. LHC is and will remain weak in bounding them.



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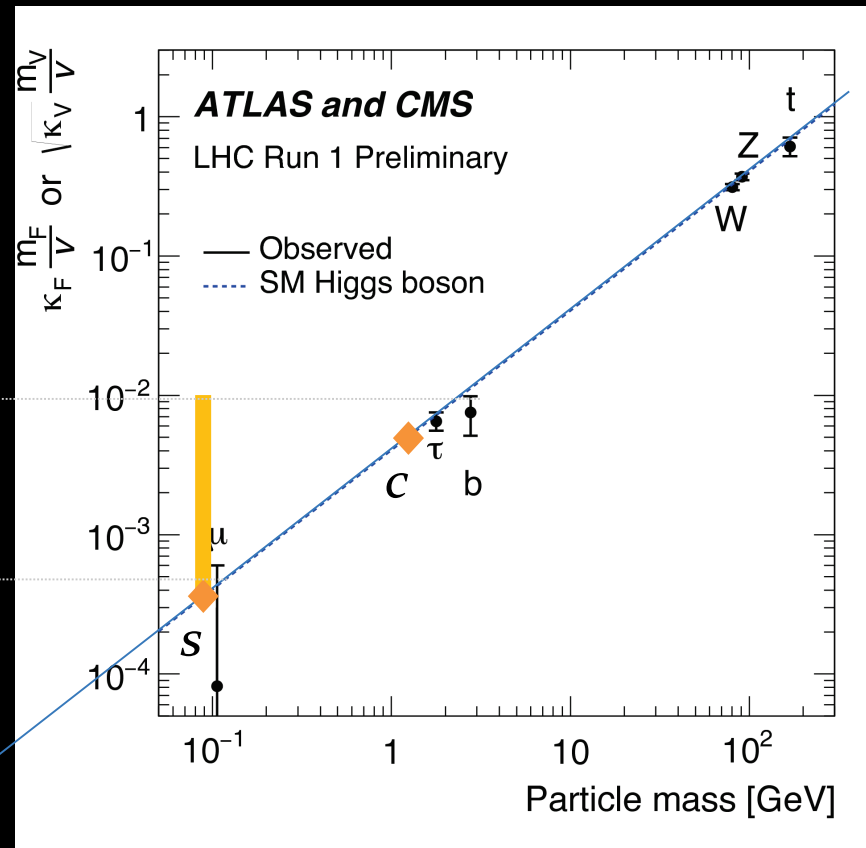
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*The Higgs force  
in Optical Clock Transitions*

# The Atomic Higgs Force

- The Higgs results in an attractive force between nuclei and their bound electrons (à la Yukawa):

$$V_{\text{Higgs}}(r) = -\frac{y_e y_A}{4\pi} \frac{e^{-m_h r}}{r} \approx -\frac{y_e y_A}{4\pi m_h^2} \frac{\delta(r)}{r^2}$$

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- $y_A = Z y_p + (A - Z) y_n$  with: Shifman-Vainshtein-Zakharov '78  
+ nuclear data, see e.g. micrOmegas

$$y_n \approx 7.7 y_u + 9.4 y_d + 0.75 y_s + 2.6 \times 10^{-4} c_g$$

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- $c_g$  constrained by LHC, weaker sensitivity to s-coupling

# Higgs Force Strength

- Under current LHC constraints:

$$y_{n,p} \lesssim 3 \text{ (0.2)} \quad \text{and} \quad y_e \lesssim 1.3 \times 10^{-3}$$

Higgs width (direct) →

← global fit (indirect)

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- This shifts transition frequencies by:

$$\Delta\nu_{nS \rightarrow n'D,F}^{\text{Higgs}} \approx 1 \text{ Hz} \times A \frac{y_e y_{n,p}}{0.004} \frac{|\psi(0)|^2}{4n^3 a_0^{-3}}$$

↙ electron-density at the nucleus  
 ↘ Bohr radius  $(\alpha m_r)^{-1}$

# Electron Density in Nuclei

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ion charge 


See e.g. Budker-Kimball-DeMille: Atomic Physics

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- Using non-relativistic hydrogen-like wavefunction:

$$|\psi(0)|^2 \simeq \frac{4.2Z}{a_0^3} (1+n_e)^2$$

# Optical Atomic Clocks

- State-of-the-art accuracy at the  $10^{-18}$  level

Bloom et al., Nature 506, 71-76 (2014)

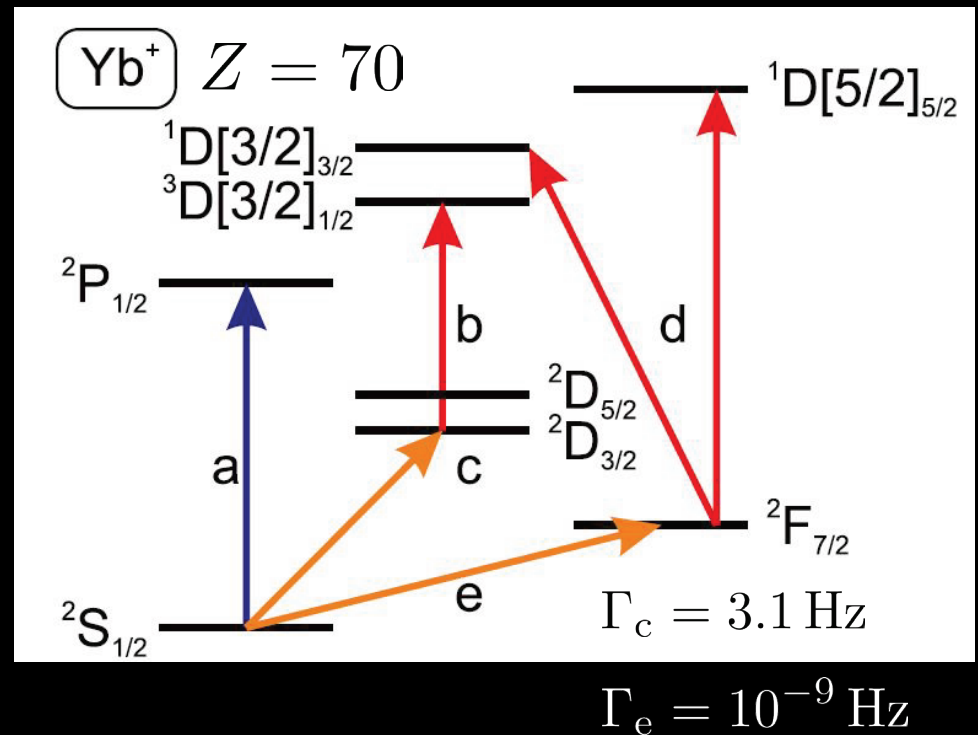
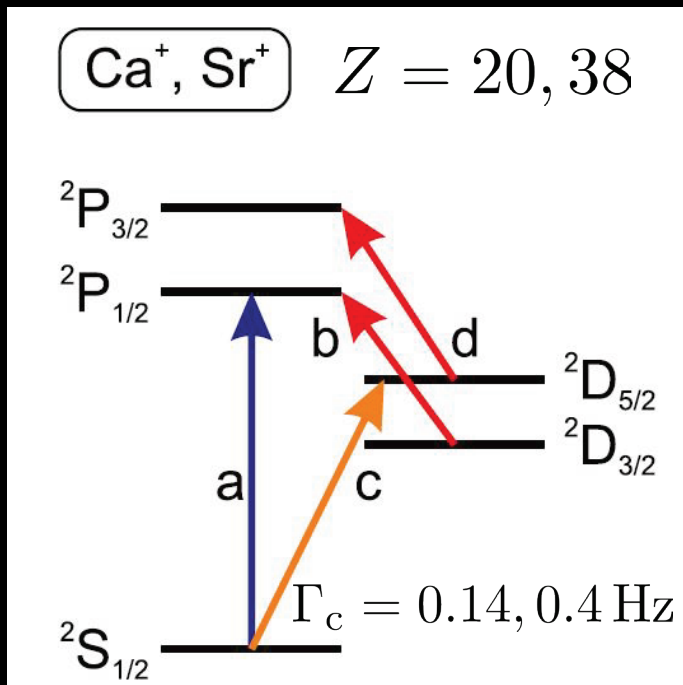
# Optical Atomic Clocks

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Bloom et al., Nature 506, 71-76 (2014)

- Narrow transitions with S-wave are needed:

Ludlow-Boyd-Ye, Rev. Mod. Phys. 87 (2015)



# Frequency Comparisons

- Experimental accuracy in  $^{40}\text{Ca}^+$ ,  $^{88}\text{Sr}^+$  is  $\sim \text{Hz}$

Dube et al., Phys. Rev. A87 (2013)  
Chwalla et al., PRL 102 (2009)

$$\nu_{E2}^{\text{Ca}^+} = 411\,042\,129\,776\,393.2(1.0)\text{Hz} \quad \sim 10^{15}\text{Hz}$$

$$\nu_{E2}^{\text{Sr}^+} = 444\,779\,044\,095\,485.5(9)\text{Hz}$$

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- Theory side is however much less promising:  
electron-electron correlations, nuclear finite-size,  
relativistic corrections, QED...  
are not accounted for at the  $10^{-15}$  level...

*Isotope Shifts  
and King plots*

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*(consider  $A' - A = 2, 4, \dots$  to avoid influence of nuclear spin)*
- The Higgs force however scales like the nuclear mass  $A$ , so there is still a net shift between isotopes!

# Isotope Shift Sources

- There are yet non-trivial IS from changes in:
  - the reduced mass:  $m_r = \frac{m_e m_A}{m_e + m_A} \simeq m_e (1 - m_e/m_A)$
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- MS/FS effects are typically in the GHz range  $\gg$  HS



# The King Plot

W. H. King,  
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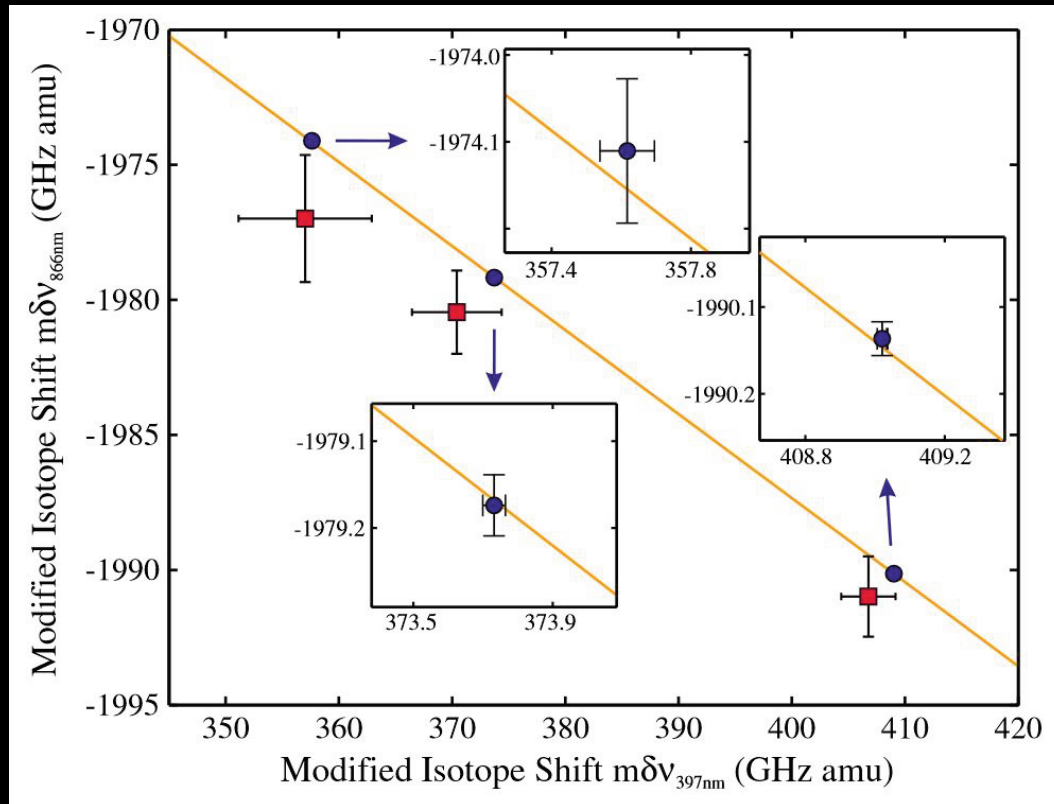
- Plot  $m\delta\nu_{AA'}^1$  vs.  $m\delta\nu_{AA'}^2$  along the isotopic chain and as long as linearity is observed,  $H_{21}$  can be bounded (unless accidentally  $m\delta\nu \propto A'$  )

# Proof of Concept in $\text{Ca}^+$

Gebert et al. PRL 115 (2015)

$$A = 40, A' = 42, 44, 48$$

$$4S \rightarrow 3D_{5/2}$$



IS  $\sim$  1 GHz  
error  $\sim$  100 kHz

$$y_e y_n \lesssim 40$$

$$4S \rightarrow 4P_{1/2} \text{ (not-clock)}$$

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- This is ~10 times better than (comparable to) LHC8 direct (indirect) bounds, with good/better prospect for improvements!



# Higher-Order Corrections

- Need to control King's linearity at least down to:

$$\begin{array}{l} \text{Higgs force} \longrightarrow \\ \text{total IS} \longrightarrow \end{array} \frac{\text{Hz}}{\text{GHz}} \sim 10^{-9}$$

- Higher-order corrections are not trivial to compute, many-body, relativistic simulations are needed [in progress]
- Yet, IS are controlled by two small parameters:

$$\begin{aligned} \varepsilon_\mu &= m_e \mu_{AA'} \sim (A - A') 10^{-8} \\ \varepsilon_r &= \delta \langle r^2 \rangle_{AA'} / a_0^2 \sim (A - A') 10^{-11} \end{aligned}$$

- So, we can entertain NDA...

# Field Shift

- Perturbation theory: Seltzer '69  
Blundell et al. '87

$$\delta\nu_{AA'}^{\text{FS}} = -e \int d^3r_e |\psi(r_e)|^2 \delta V(r_e), \quad \delta V(r_e) = \frac{Ze}{4\pi} \int d^3r_N \frac{\delta\rho(r_N)}{|\vec{r}_e - \vec{r}_N|}$$

↑
↑

electron density
nuclear potential

nuclear charge distribution  
 ↓

- LO:  $\propto |\psi(0)|^2 \delta\langle r^2 \rangle_{AA'} \sim \mathcal{O}(\varepsilon_r)$
- NLO/LO:  $\sim \mathcal{O}(\varepsilon_\mu^2, \varepsilon_r^2, \varepsilon_\mu \varepsilon_r) / \varepsilon_r \sim 10^{-7}$
- NLO is linear up to overlap with the nucleus  $\sim \mathcal{O}(\varepsilon_r)$
- Hence, non-linearities are only of  $\mathcal{O}(\varepsilon_\mu^2) \sim 10^{-14}$

# Specific Mass Shift

- MS arises from:
  - « rescaling » Rydberg constant (normal MS)
  - electron-electron correlation, relativistic... (specific MS)
- at LO, both scale like  $m_e \mu_{AA'} \sim \mathcal{O}(\varepsilon_\mu)$
- NLO correction is parametrically: Palmer '87
$$\sim \alpha^2 m_e^2 (m_A^{-2} - m_{A'}^{-2})$$
- Hence, NLO/LO  $\sim \mathcal{O}(\alpha^2 \varepsilon_r) \sim 10^{-10}$

*Probing EW  
and BSM Physics*

# The Weak Force

$$V_{\text{weak}}(r) = -\frac{8G_{\text{F}}m_{\text{Z}}^2}{\sqrt{2}} \frac{g_{\text{e}}g_{\text{A}}}{4\pi} \frac{e^{-m_{\text{Z}}r}}{r}$$

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LEPEWWG

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Efrati-Falkowski-Soreq '14

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 [\delta g_L^{Zu}]_{ii} &= \begin{pmatrix} -0.8 \pm 3.1 \\ -0.16 \pm 0.36 \\ -0.28 \pm 3.8 \end{pmatrix} \times 10^{-2}, & [\delta g_R^{Zu}]_{ii} &= \begin{pmatrix} 1.3 \pm 5.1 \\ -0.38 \pm 0.51 \\ \times \end{pmatrix} \times 10^{-2}, \\
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- IS improved bounds:  $|\delta g_V^{Zu} + 2\delta g_V^{Zd}| \lesssim 0.018$
- still weaker than APV in Cs:  $|\delta g_V^{Zu} + 2\delta g_V^{Zd}| \lesssim 10^{-3}$



# Effective Field Theory

- Relevant operators:

$$\mathcal{O}_{eq}^V = (\bar{e}\gamma^\mu e)(\bar{q}\gamma_\mu q) \quad q = u, d$$

$$\mathcal{O}_{eq}^S = (\bar{e}e)(\bar{q}q) \quad q = u, d, s, c, b, t$$

$$\mathcal{O}_{eg} = \alpha_s(\bar{e}e)G_{\mu\nu}^2$$

operator $\mathcal{O}_i$	Upper bound on $ c_i $ ( $\Lambda = 1 \text{ TeV}$ )	Lower bound on $\Lambda_i$ [TeV] ( $c = 1$ )
$\mathcal{O}_{eu}^V$	$2.3 \times 10^{-2}$	6.6
$\mathcal{O}_{ed}^V$	$1.1 \times 10^{-2}$	9.3
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$\mathcal{O}_{ed}^S$	$2.1 \times 10^{-3}$	22
$\mathcal{O}_{es}^S$	$2.7 \times 10^{-2}$	6.1
$\mathcal{O}_{ec}^S$	0.20	2.3
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LEP<sub>2</sub>

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sensitive to scalar operators up to 20TeV!

\* to be confirmed

# 750GeV Resonance\*

- LHC established its coupling to hadrons
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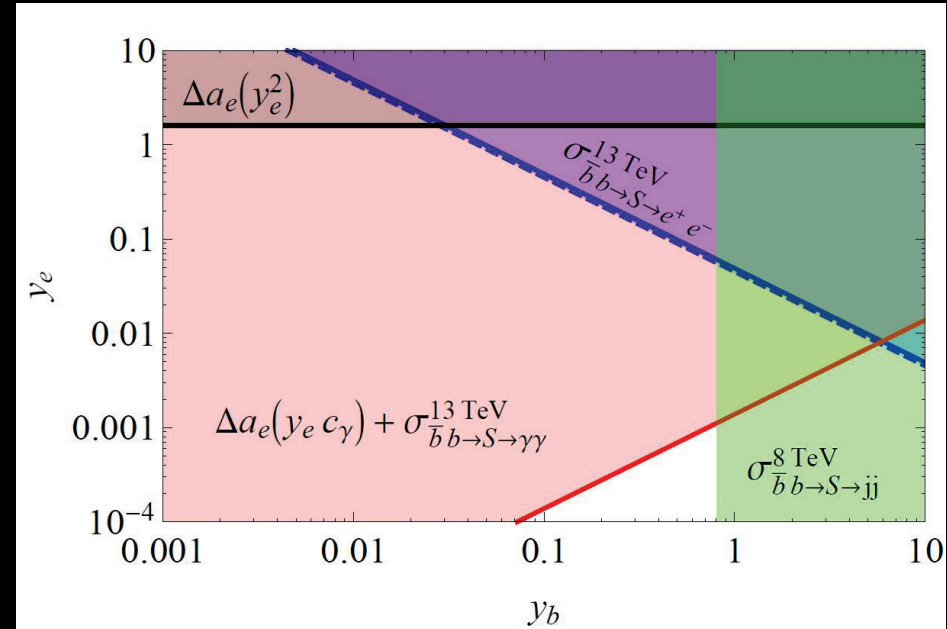
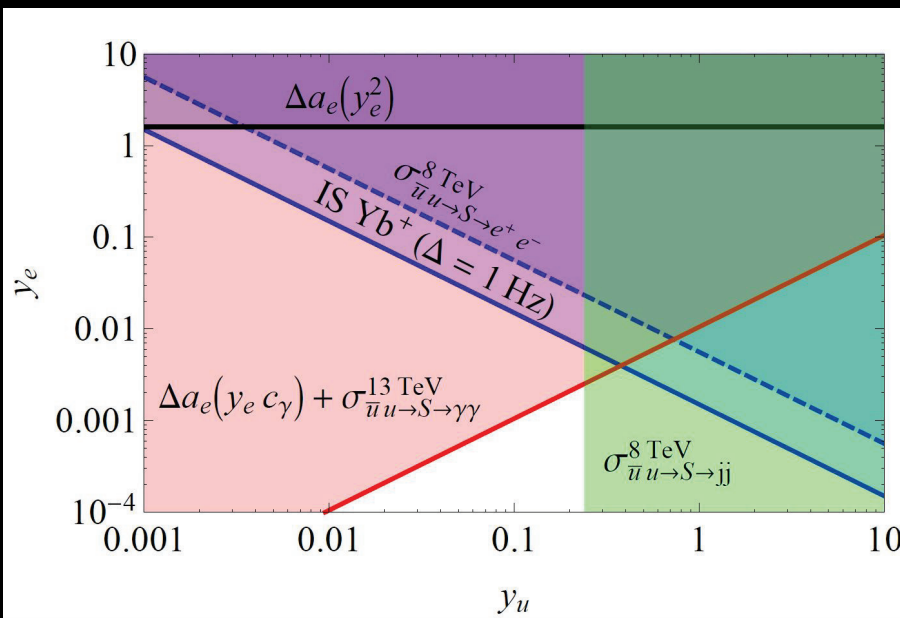
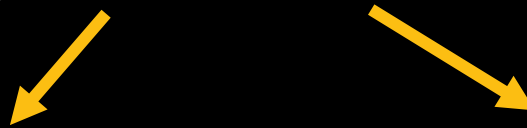
- LHC established its coupling to hadrons
- What if it further couples to electrons?
- Assuming a scalar resonance  $S$ :

Unless it's produced through gluon or heavy quark fusion, IS has more sensitivity to  $S$  couplings than LHC searches in  $e^+e^-$

$S$ couplings ( $\mu = 750$ GeV)	LHC (8, 13) bound [33, 37] ( $\Gamma_S = 45$ GeV)	IS projection ( $\Delta = 1$ Hz)
$ y_e y_u $	$(5.6, 6.0) \times 10^{-3}$	$1.5 \times 10^{-3}$
$ y_e y_d $	$(7.3, 7.8) \times 10^{-3}$	$1.2 \times 10^{-3}$
$ y_e y_s $	$(2.9, 2.5) \times 10^{-2}$	$1.5 \times 10^{-2}$
$ y_e y_c $	$(3.6, 3.0) \times 10^{-2}$	$9.6 \times 10^{-2}$
$ y_e y_b $	$(5.6, 4.5) \times 10^{-2}$	0.49
$ y_e y_t $	(0.19, 0.16)	32
$ y_e c_g $	(0.72, 0.60)	150

# 750GeV Resonance\*

- $\gamma\gamma$  signal +  $g_e - 2$  bounds the  $S\bar{e}e$  coupling
- Assuming e.g.  $u\bar{u}$  or  $b\bar{b}$  production:



*Conclusions*

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- Our method is not only sensitive to the Higgs force, but also to:
  - the weak force
  - BSM forces, as long as not coupled like charge
- Measurements in  $\text{Yb}^+$  are already underway!
- Other possibilities envisaged:  $\text{Ca}/\text{Ca}^+$ ,  $\text{Sr}/\text{Sr}^+$ ,  $\text{Dy}$