

# A METHOD FOR COMPUTATION OF RADIATION EMITTED BY ELECTRONS AND POSITRONS IN STRAIGHT AND BENT CRYSTALS

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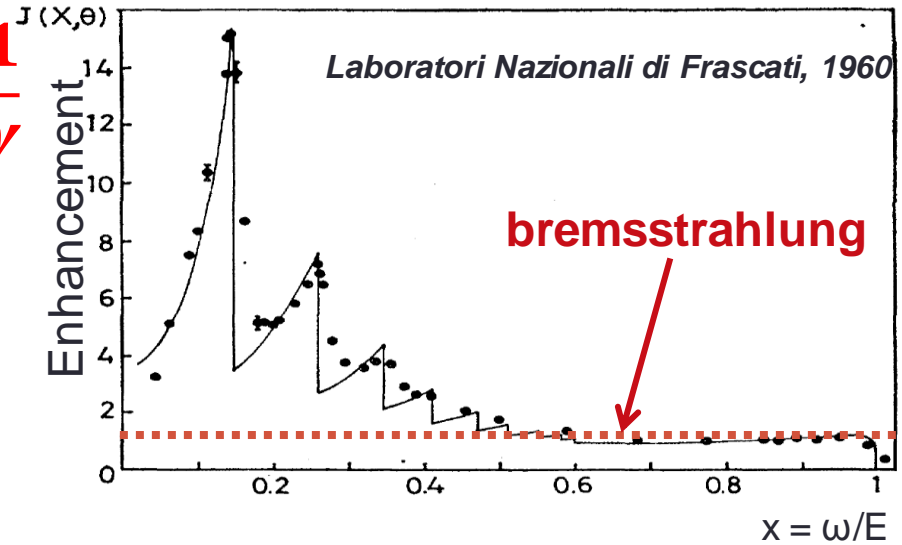
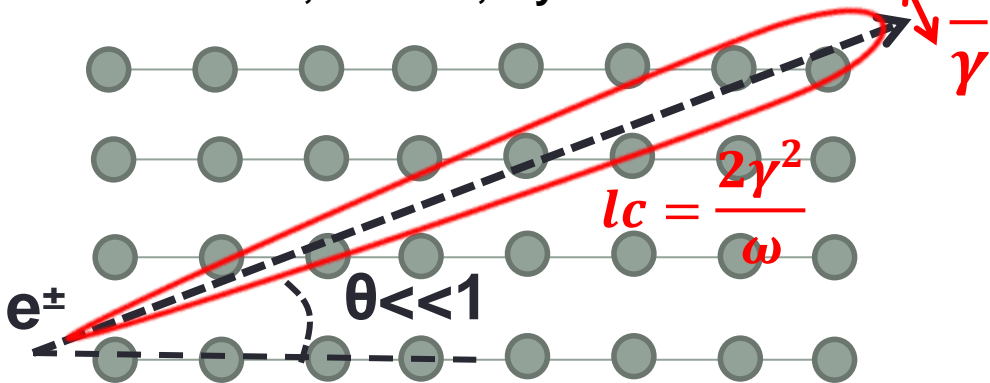


# Outlook

- Introduction of radiation processes in crystals;
- Development of an algorithm to compute the radiation generation in oriented crystals based on the Baier and Katkov method;
- Comparison with experiments at intermediate energies (1 GeV  $e^\pm$ );
- Comparison with experiments at ultra-high energies (100 GeV  $e^\pm$ );

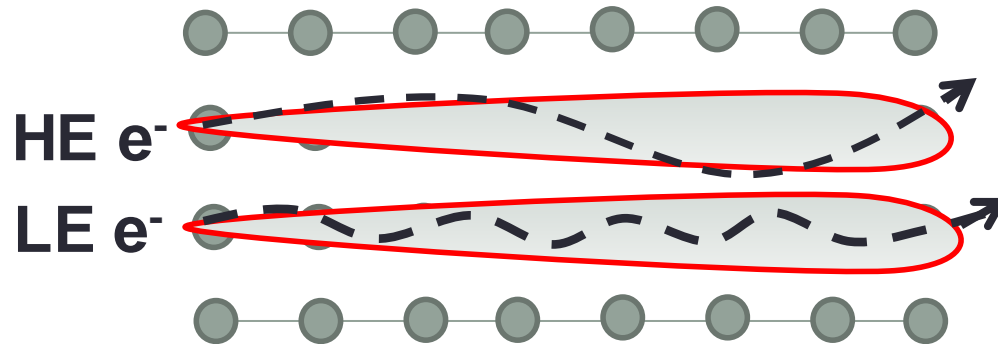
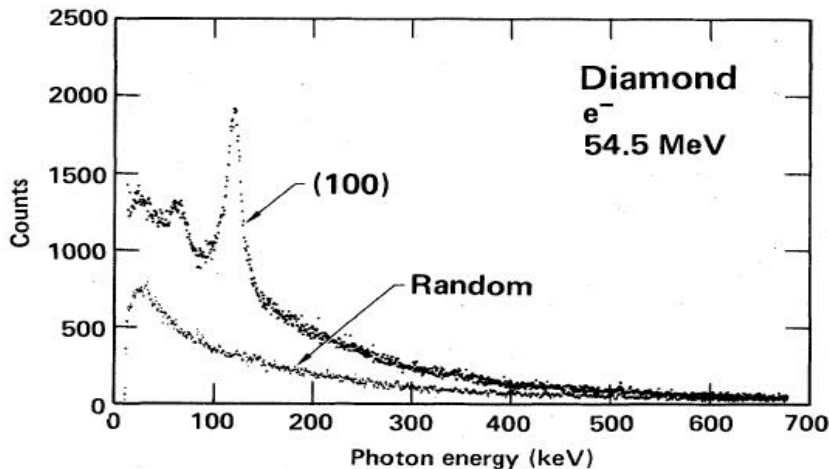
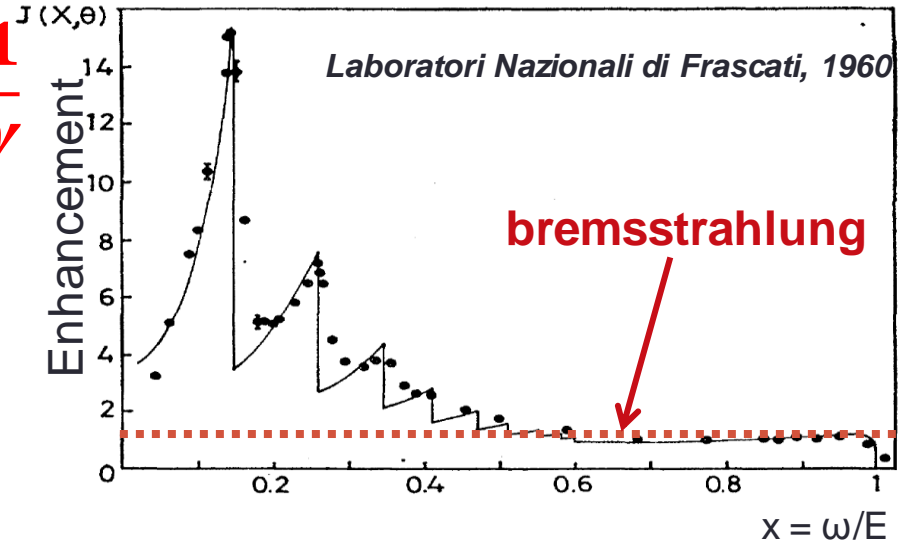
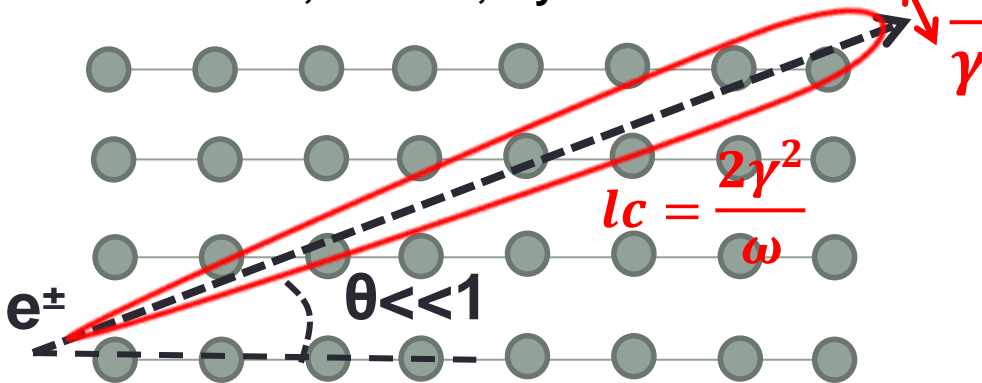
# Enhancement of bremsstrahlung radiation in aligned crystals

Coherent Bremsstrahlung (1950s) Ter-Mikaelian, Ferretti, Dyson-Uberall



# Enhancement of radiation in aligned crystals

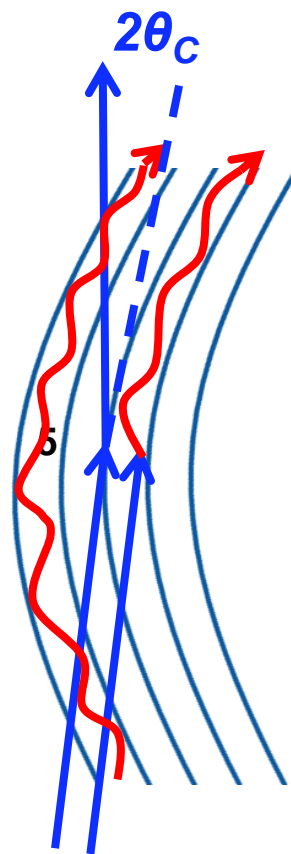
Coherent Bremsstrahlung (1950s) Ter-Mikaelian, Ferretti, Dyson-Uberall



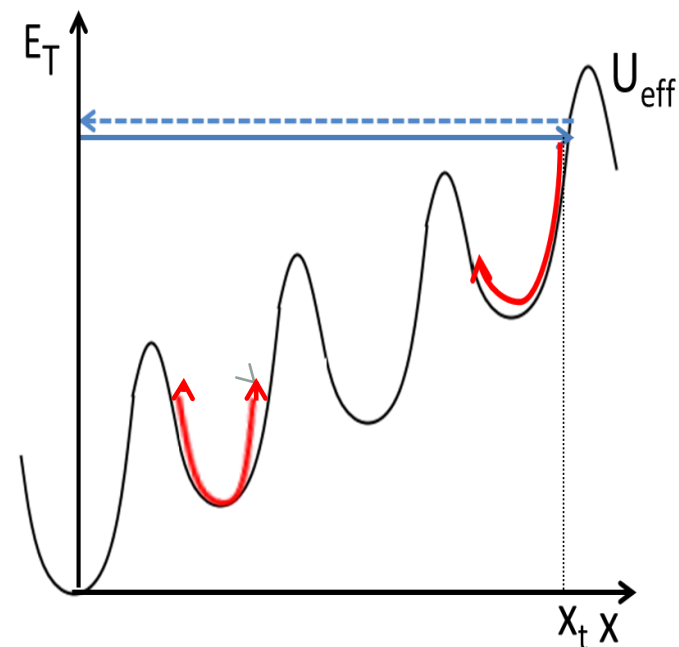
Channeling radiation ( Kumakhov)

# Channeling and related effects in a bent crystal

- ✓ **Tsyganov (1976):** channeling in bent crystals;
- ✓ **Taratin and Vorobiov (1987):** Volume Reflection of over-barrier particles.

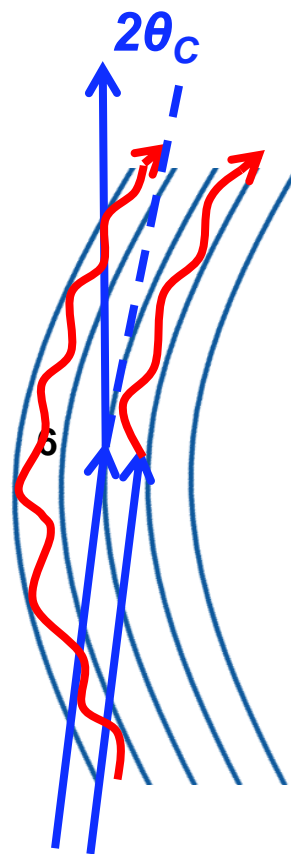


$$U_{eff}(x) = U(x) + \frac{pv}{R}x$$

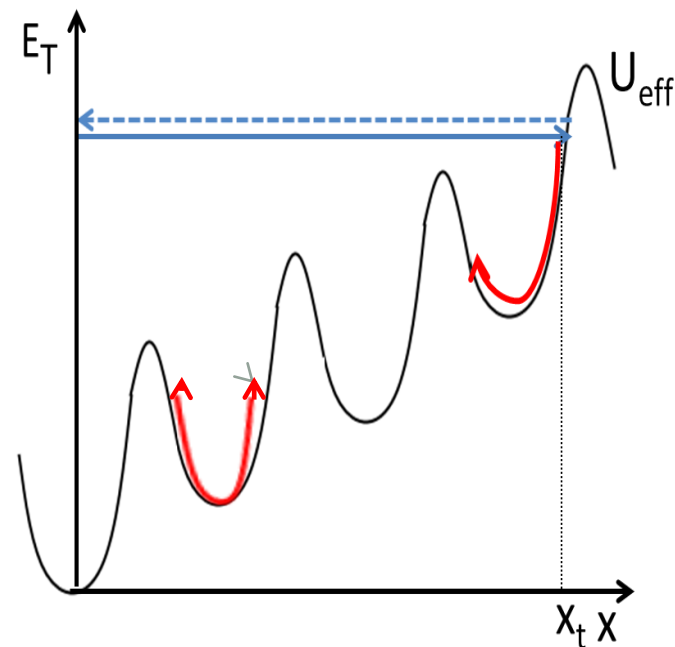


# Channeling and related effects in a bent crystal

- ✓ **Tsyganov (1976)**: channeling in bent crystals;
- ✓ **Taratin and Vorobiov (1987)**: Volume Reflection of over-barrier particles.
- Under VR the angle between the particle trajectory and crystalline planes changes during the motion
- Need of a general method for radiation computation



$$U_{eff}(x) = U(x) + \frac{pv}{R}x$$



# Baier-Katkov quasiclassical operator method (1967-1968)

General method for calculation of radiation generated by  $e^\pm$  in an external field

The electromagnetic radiated energy is evaluated with the BK formula:

$$\frac{dE}{d^3k} = \omega \frac{dN}{d^3k} \frac{\alpha}{4\pi^2} \iint dt_1 dt_2 \frac{[(E^2 + E'^2)(v_1 v_2 - 1) + \omega^2 / \gamma^2]}{2E'^2} e^{-ik'(x_1 - x_2)} \quad (1)$$

where the integration is made over the classical trajectory.

The generality of the Baier-Katkov operator method permits to simulate the electromagnetic radiation emitted by  $e^\pm$  in very different cases, e. g., straight, bent and periodically bent crystals, and for different beam energy range, from sub-GeV to TeV.

# Baier-Katkov quasiclassical operator method (1967-1968)

General method for calculation of radiation generated by  $e^\pm$  in an external field

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where the integration is made over the classical trajectory.

## Why classical trajectory?

2 types of quantum effects :

- the quantization of particle motion  $\sim \hbar\omega_0/E$   
In crystals: **negligible for electron/positron energy >10-100 MeV**
- the **quantum recoil** of the particle when it radiates a photon with energy  $\hbar\omega \sim E$   
**NOT negligible for electron/positron energy >50 GeV**



# An algorithm for radiation in crystals

## Integration of the BK formula

**SMALL ANGLE APPROXIMATION:** Since the angle between particle trajectories and crystal planes or axes is small and at ultrarelativistic energies the radiation angle  $1/\gamma$  is much smaller than unity the particle velocity  $\mathbf{v}$  and photon momentum  $\mathbf{k}$  can be represented in the form :

$$\begin{aligned}\mathbf{v}(t) &\simeq \mathbf{v}_{\perp}(t) + \mathbf{e}_z [1 - 1/2\gamma^2 - v_{\perp}^2(t)/2], \\ \mathbf{k} = \mathbf{n}\omega &\simeq \mathbf{e}_{\perp}\omega\theta + \mathbf{e}_z\omega(1 - \theta^2/2),\end{aligned}$$

where the angle  $\theta \ll 1$  represents the radiation angle. The formula (1) can be rewritten as:

$$\frac{dE}{d^3k} \sim \frac{\alpha}{8\pi^2} \frac{\varepsilon^2 + \varepsilon'^2}{\varepsilon'^2} \omega^2 C, \quad (2)$$

$$\text{where } C = |\mathbf{I}_{\perp}|^2 + \gamma^{-2} \frac{\omega^2}{\varepsilon^2 + \varepsilon'^2} |J|^2 \quad (3)$$

# An algorithm for radiation in crystals

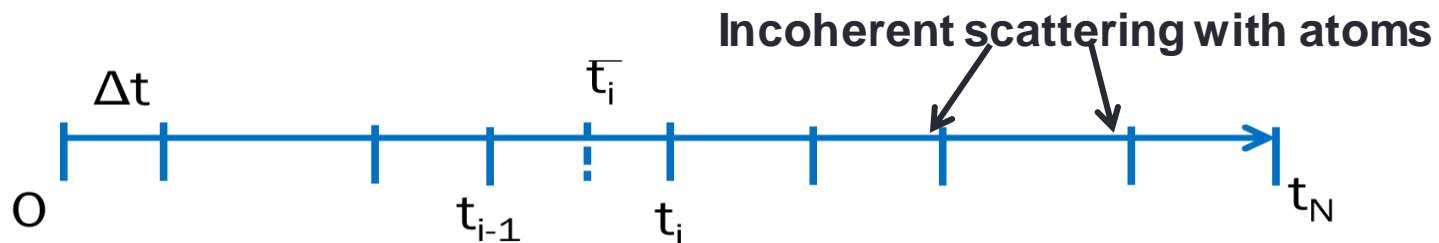
## Integration of the BK formula

**SMALL ANGLE APPROXIMATION:** the integrals of eq. (1) can be represented as follows:

$$I_{\perp} = \int_{-\infty}^{+\infty} dt (v_{\perp} - \theta) e^{-i\phi(t)} \quad (4)$$

being  $\phi(t) = \frac{\omega'}{2} \int_{-\infty}^t dt' [\gamma^{-2} + (v_{\perp}(t') - \theta)^2]$  and  $\omega' = \omega\varepsilon/\varepsilon'$ .

## ACCOUNT OF INCOHERENT SCATTERING:

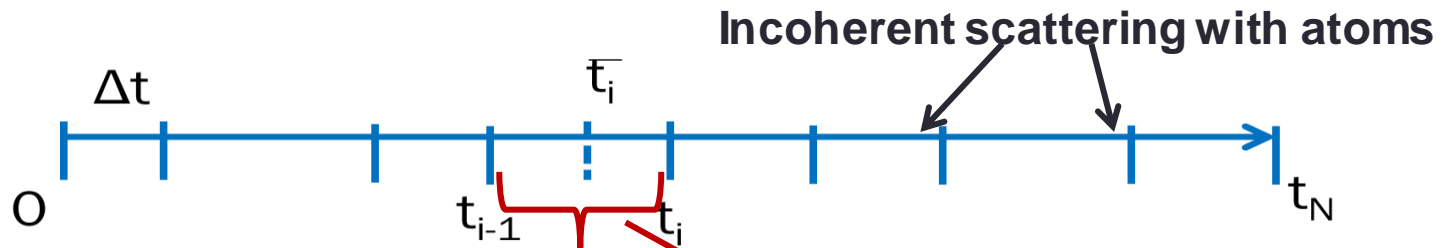


The **particle trajectory** is then divided in **N small steps**, within which the particle trajectory is calculated through the integration of equation of motion in the continuous potential. **At the end of each step the scattering by nuclei and electrons is sampled** and the transverse velocity for the *i*-step becomes

$$\mathbf{v}_{\perp,i} \rightarrow \mathbf{v}_{\perp,i} + \theta_{s,i}$$

# An algorithm for radiation in crystals

## Integration of the BK formula



In order to improve the convergence of its integration over  $t$  and  $\theta$  (photon emission angle), the integrals of eq. 4 are computed as follows after an integration by parts:

$$J \approx i \sum_{i=1}^N \left\{ \exp [i\phi(t_i)] \left[ \frac{1}{\phi_{t_i+0}} - \frac{1}{\phi_{t_i-0}} \right] - \exp [i\phi(\bar{t}_i)] \left[ \frac{2\ddot{\phi}}{\dot{\phi}^3} \right]_{t_i} \sin ([\phi(t_i - 0) - \phi(t_{i-1} + 0)]/2) \right\}$$

**If incoherent scattering is switched off, it is go to zero.**

$$\dot{\phi}(t < t_i) = \frac{\omega'}{2} [1/\gamma^2 + (\mathbf{v}_{\perp}(t) - \boldsymbol{\theta})^2],$$

$$\ddot{\phi}(t) = \omega' (\mathbf{v}_{\perp}(t) - \boldsymbol{\theta}) \dot{\mathbf{v}}_{\perp}(t),$$

$$\dot{\phi}(t_i + 0) = \frac{\omega'}{2} [1/\gamma^2 + (\mathbf{v}_{\perp}(t) + \boldsymbol{\theta}_{s,i} - \boldsymbol{\theta})^2], \quad \dot{\mathbf{v}}_{\perp} = -\frac{1}{\varepsilon} \frac{\partial U(\mathbf{r})}{\partial \mathbf{r}_{\perp}}, \quad U(\mathbf{r}) \text{ being the continuous potential.}$$

The contributions of the trajectory ends are not taken into account, thus neglecting the soft contribution of transition radiation.

**The integration over  $\theta$  leads to the radiation spectral intensity,  $\omega dN/d\omega$ .**

# COMPARISON WITH EXPERIMENTS

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GeV and sub-GeV energy range:

- No quantum correction;
- Thin crystals  $\rightarrow$  single photon emission.

# RADCHARM++

**The algorithm for direct integration of the BK formula has been included in the RADCHARM++ routine [1] , which is an expansion of the DYNECHARM++ code [2] (see E. Bagli talk)**

- The electrical characteristic of the crystal are evaluated by using the atomic form factors from x-ray diffraction data;
- Numerical integration of the classical equation of motion of particle trajectories under the continuum potential approximation;
- At the end of each step the multiple and single scattering by nuclei and electrons is sampled.

**DYNECHARM++ has already been implemented in Geant4 [3].  
The RADCHARM++ routine can also be implemented to include  
the bremsstrahlung radiation enhancement in crystals.**

[1] L. Bandiera, et al., Nucl. Instrum. Methods Phys. Res., Sect. B 355, 44 (2015).

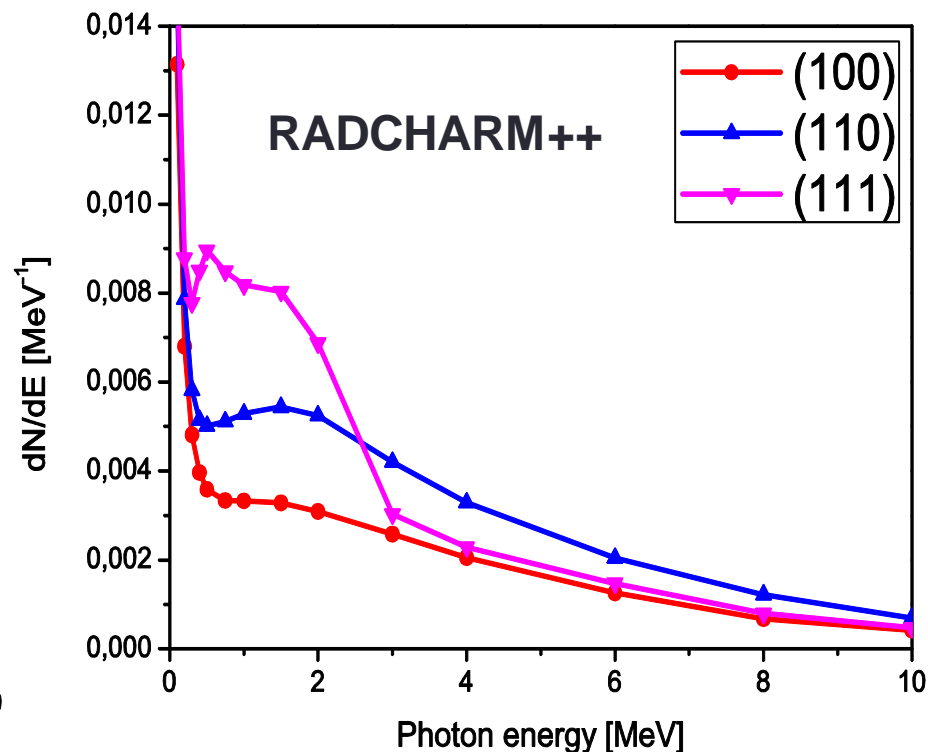
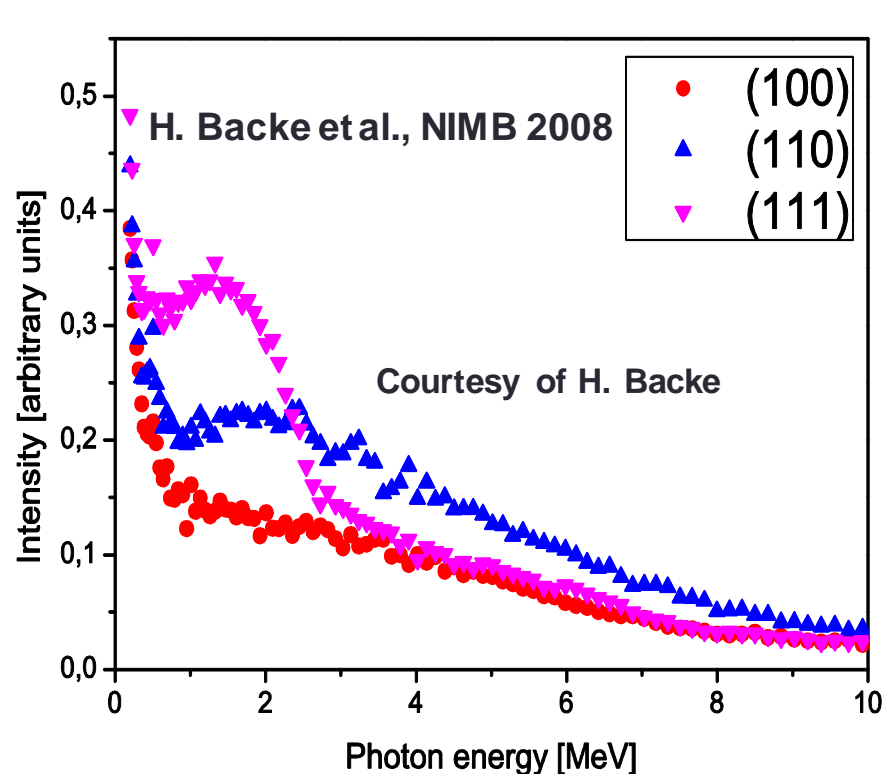
[2] E. Bagli, V. Guidi, Nucl. Instr. and Meth. in Phys. Res. Section B 309 (2013) 124

[3] E. Bagli, M. Asai, D. Brandt, et al. Eur. Phys. J. C (2014) 74: 2996.

# Comparison with previous experiments:

Simulation of e.m. radiation emitted by ultrarelativistic electrons in the field of any crystal plane

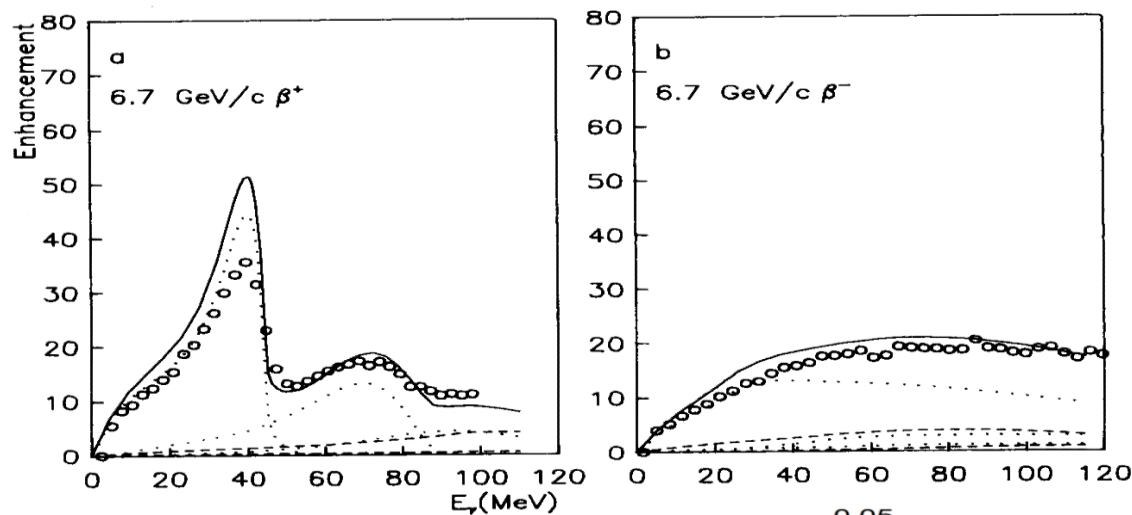
Comparison with past experiments performed at the Mainzer Mikrotron with 855 MeV electrons interacting with a 175  $\mu\text{m}$  straight Si crystal



## Comparison with previous experiments:

Simulation of e.m. radiation emitted by both positive and negative particles  
negative particles

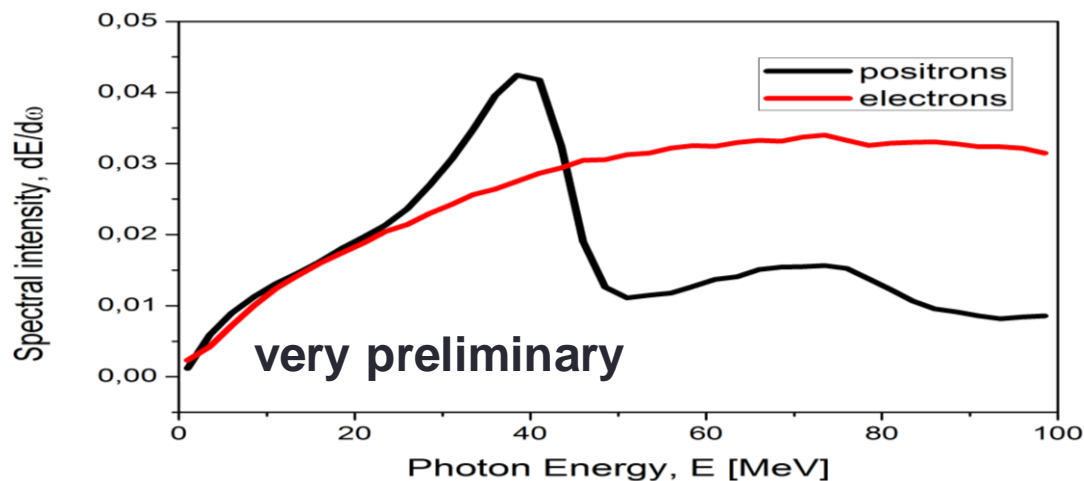
Comparison with past experiments at CERN: 6.7 GeV positrons/electrons channeling in a 0.1 mm thick Si (110)



Enhancement with respect to random

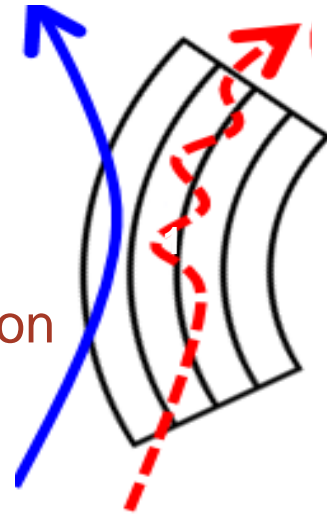
J. Bak et al.,  
Nucl.Phys. B254(1985)491.

Spectral Intensity,  
 $dE/d\omega = \omega dN/d\omega$   
(ideal case)



# Experiment with bent crystals: Motivations

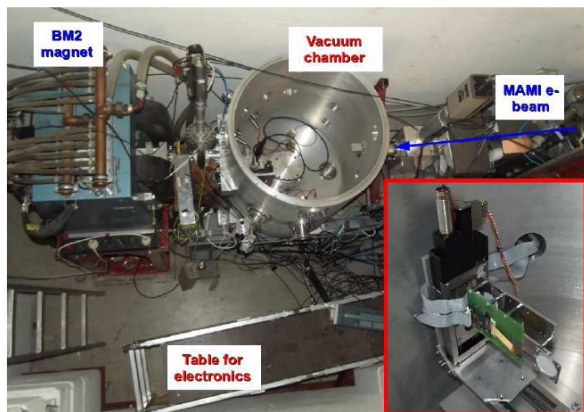
- A lot of attention is devoted to channeling effects of **electron around GeV** :
  - Interest for alternatives x-ray sources
  - Relatively large availability accelerators
- Study of the influence of the curvature on Channeling Radiation (CR) and Coherent Bremsstrahlung (CB). This experimental knowledge may be exploited to **determine with more accuracy the CR contribution in crystalline undulators**;
- Steering of sub-GeV electron trajectories through channeling in bent crystals was not possible before due to the lack of thin-enough bent crystals.



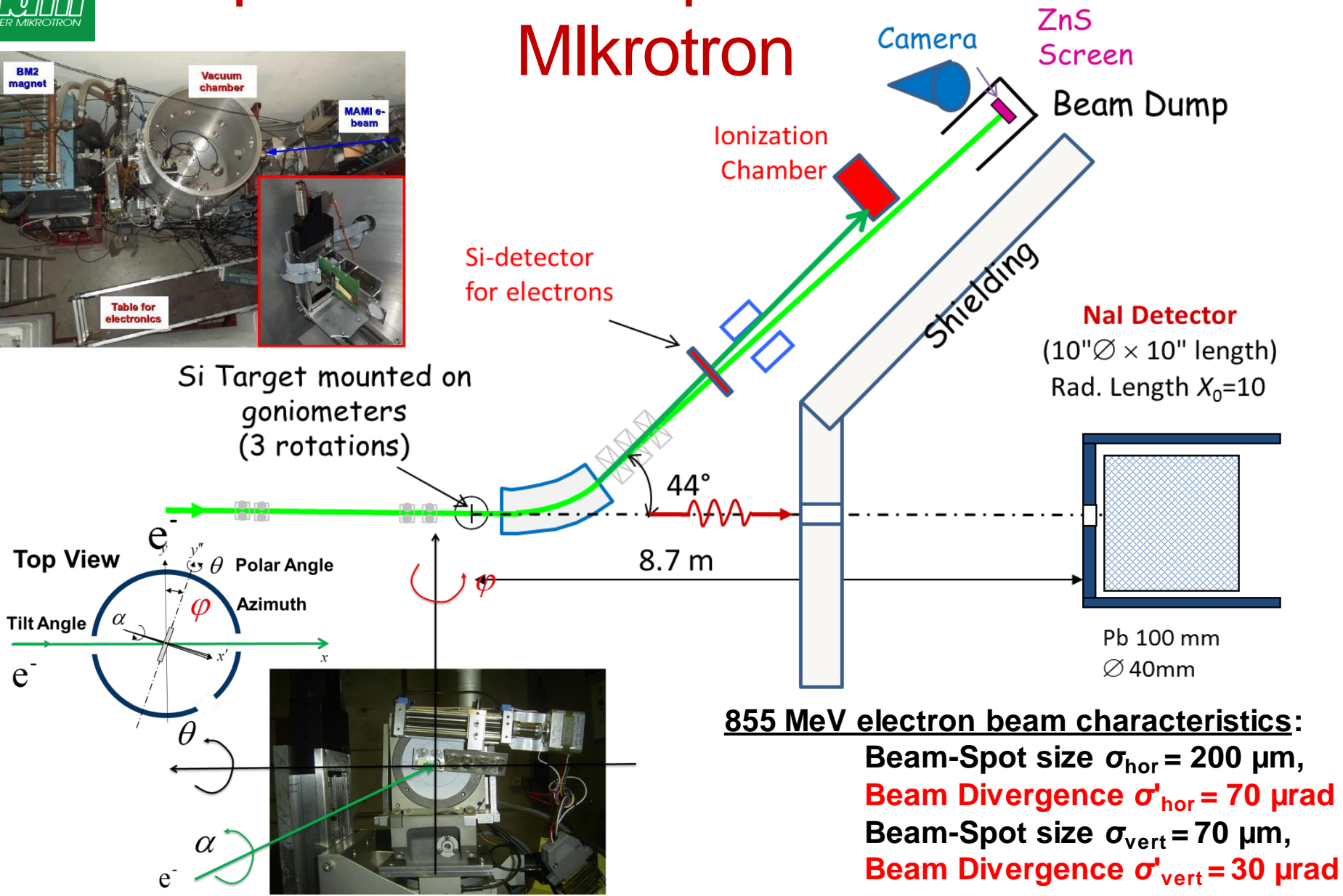




# Experimental setup at the MAInzer Mikrotron



## Mikrotron



### 855 MeV electron beam characteristics:

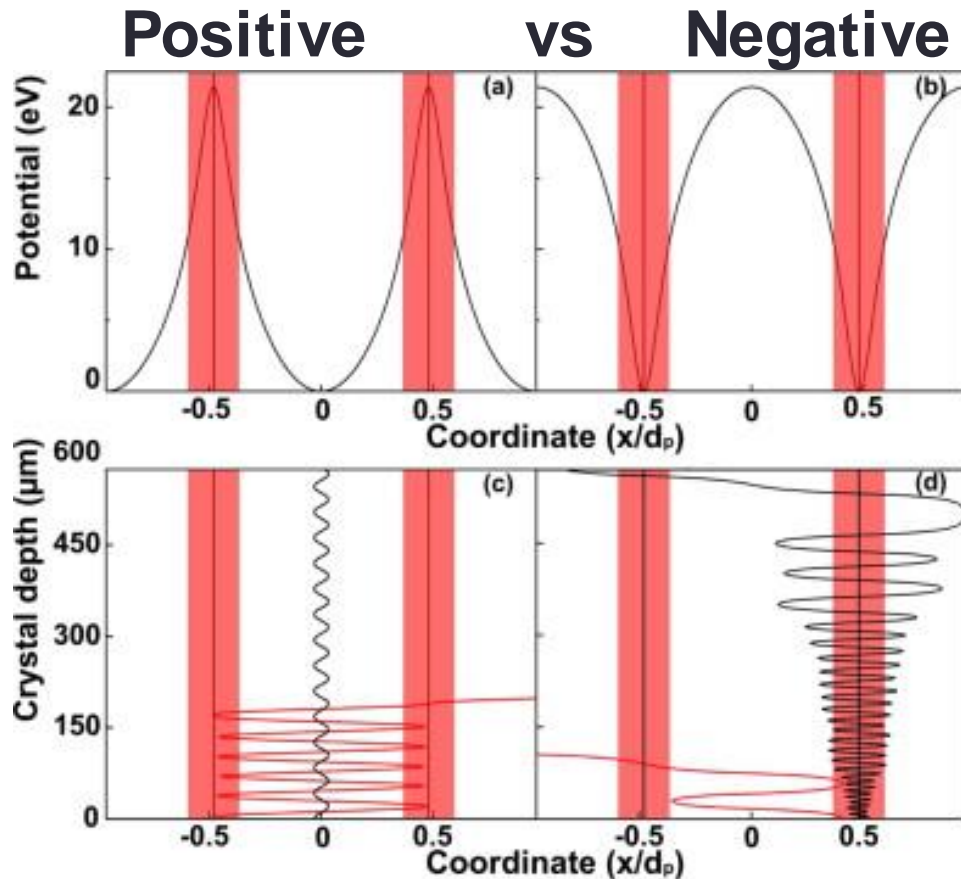
Beam-Spot size  $\sigma_{\text{hor}} = 200 \mu\text{m}$ ,

Beam Divergence  $\sigma'_{\text{hor}} = 70 \mu\text{rad}$

Beam-Spot size  $\sigma_{\text{vert}} = 70 \mu\text{m}$ ,

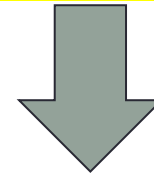
Beam Divergence  $\sigma'_{\text{vert}} = 30 \mu\text{rad}$

# Dechanneling of positive and negative particles



Channeled negative particles are dechanneled faster than positive ones due to higher probability to suffer nuclear incoherent scattering;

Ultra thin bent crystals are required for efficient deflection of negative particles

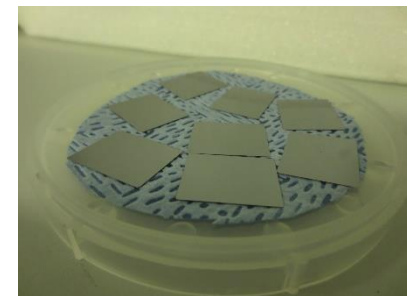
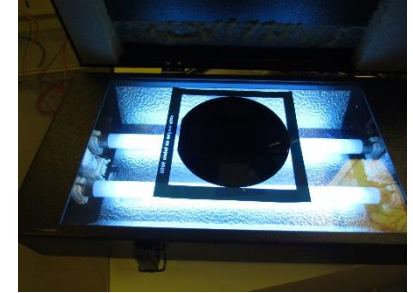
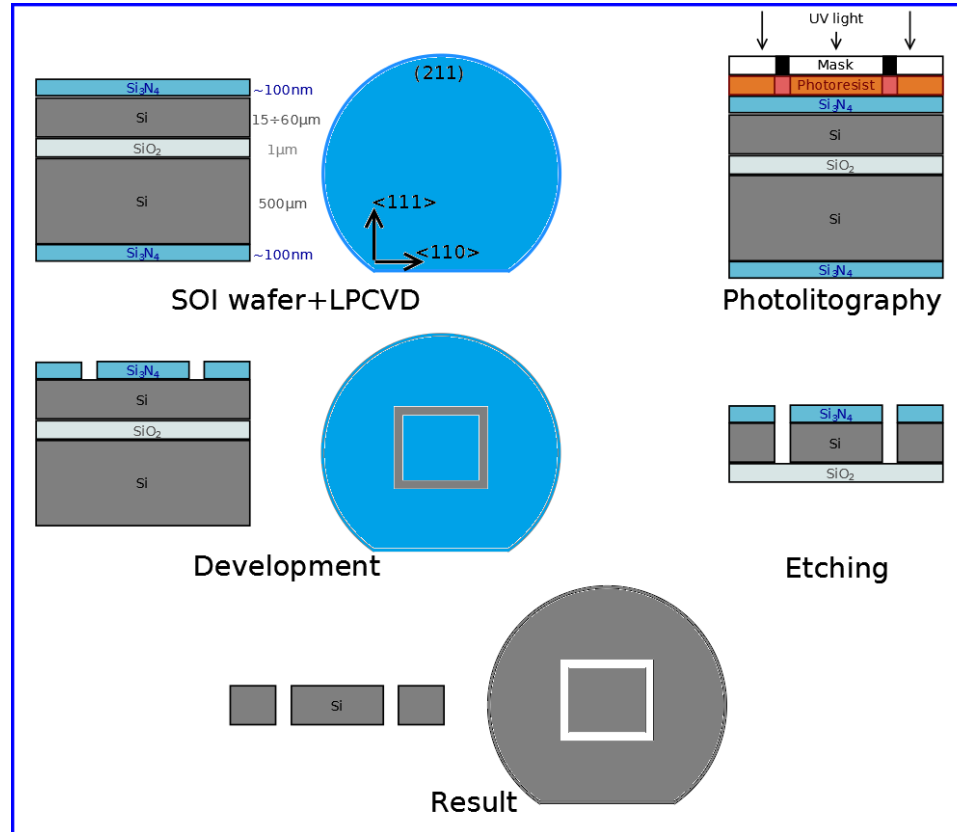
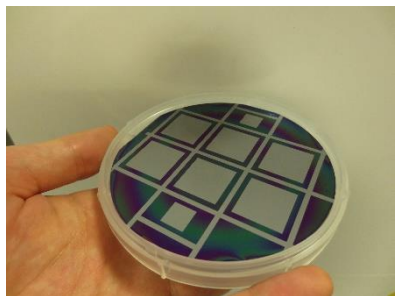
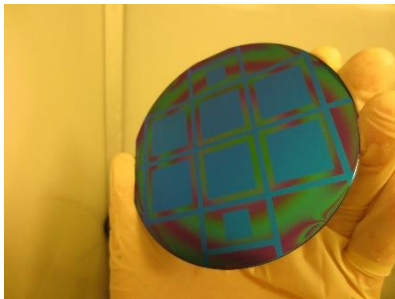


To determine how thin, one has to know the dechanneling length for negative particles!

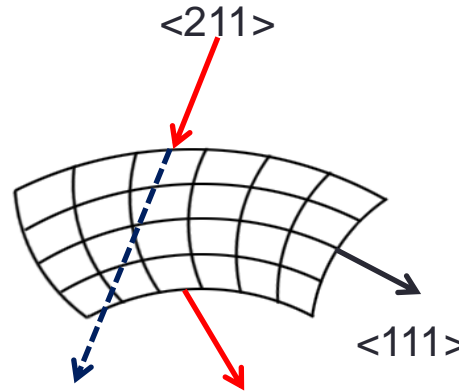
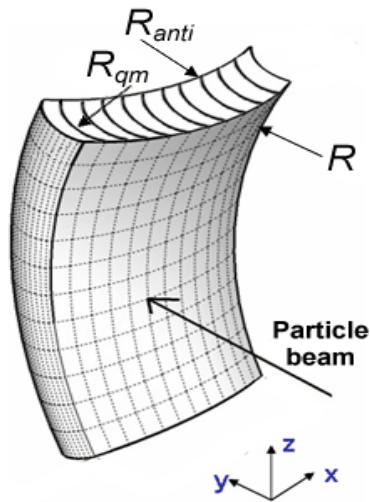
N.B.  $L_d$  decrease with energy, being some tens of microns for 1 GeV electrons in Si [1].



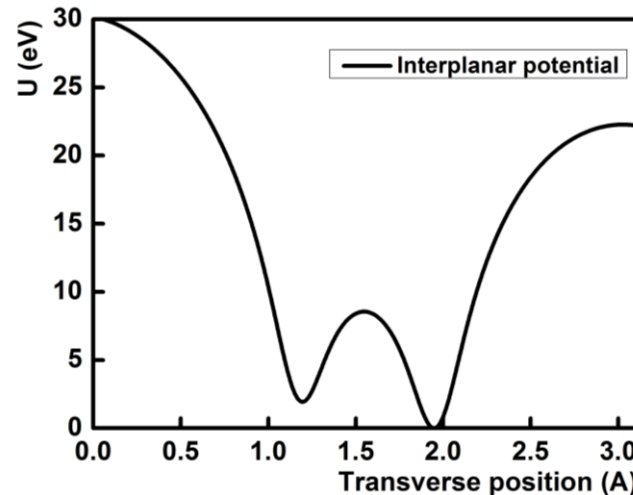
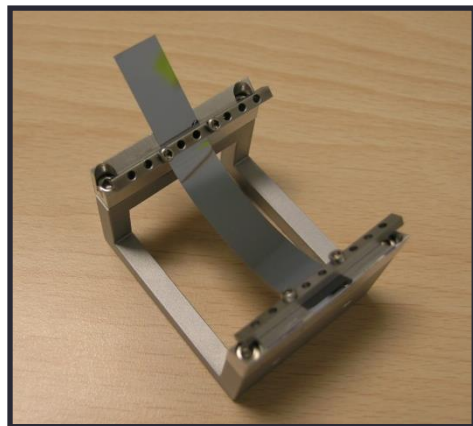
# Thin crystals manufacturing



# Crystal bending



Sketch of the **quasi-mosaic [1] crystal** highlighting the crystal crystallographic orientations. Anticlastic deformation absent due to strong bending along primary direction.



## Crystal parameters

$$T = 30.5 \mu\text{m}$$

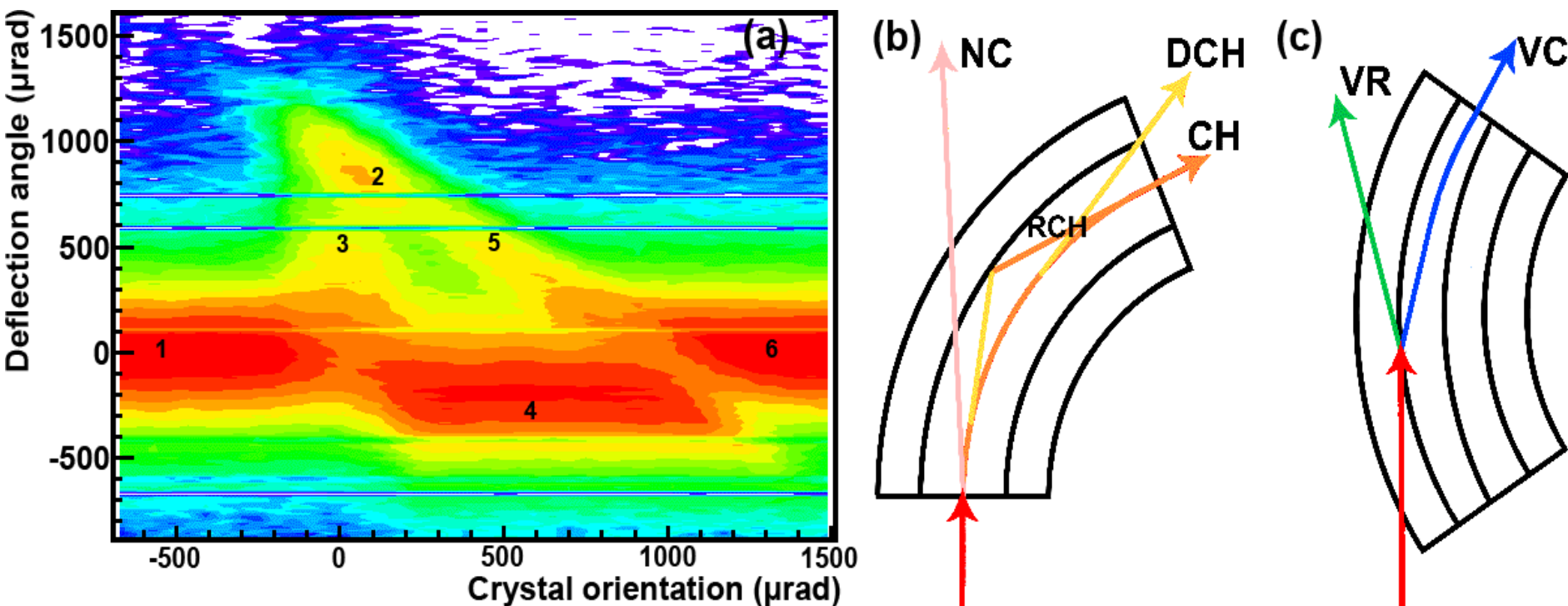
$$R_{qm} = 33.5 \text{ mm}$$

(111) Bent planes

$$U_0 = 23 \text{ eV}$$

$$\theta_c = 220 \mu\text{rad}$$

# Experimental results on beam steering

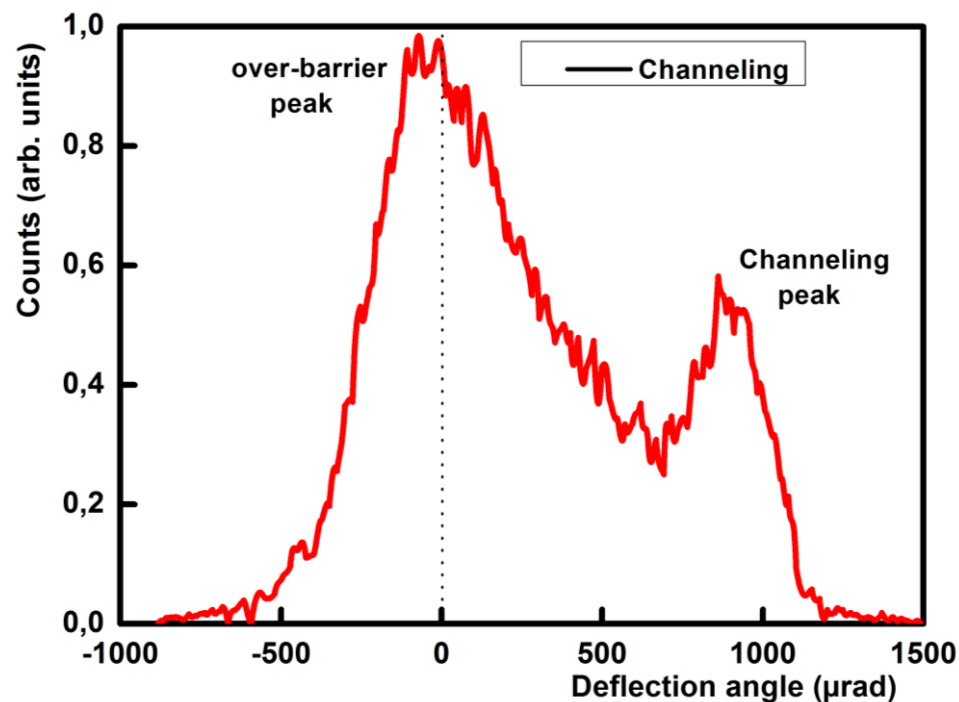


Angular scan for deflected beam distribution: (1) and (6) nonchanneling regime; (2) channeling; (3) dechanneling; (4) volume reflection; and (5) volume capture.



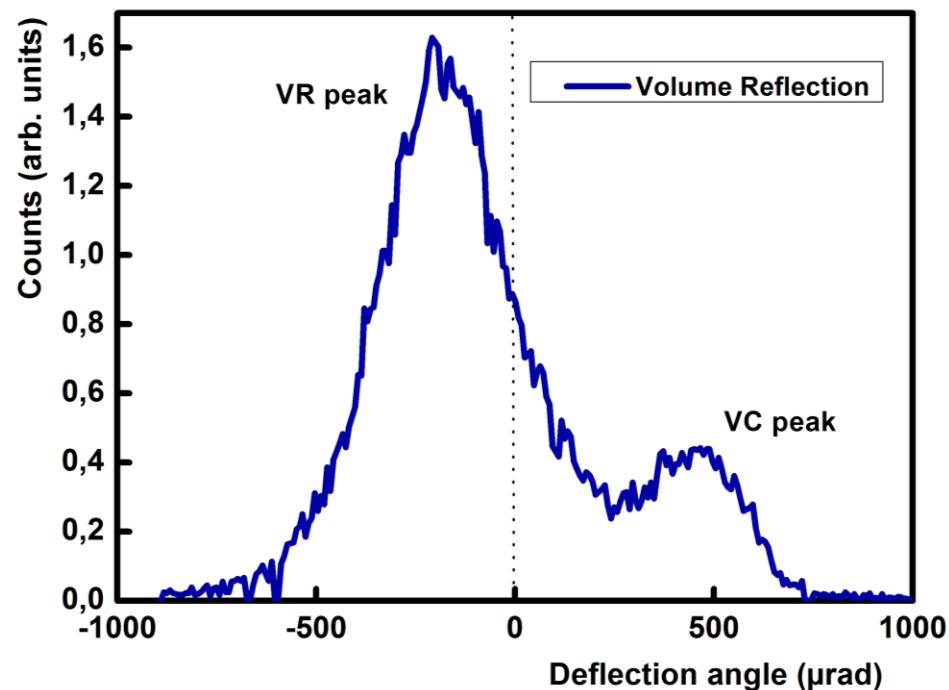
# Experimental results on beam steering

## CHANNELING



$\Theta_{\text{inc}} \sim 0 \mu\text{rad}$  Deflection =  $910 \pm 5 \mu\text{rad}$   
 Fraction of channeled is  $20.1 \pm 1.2\%$

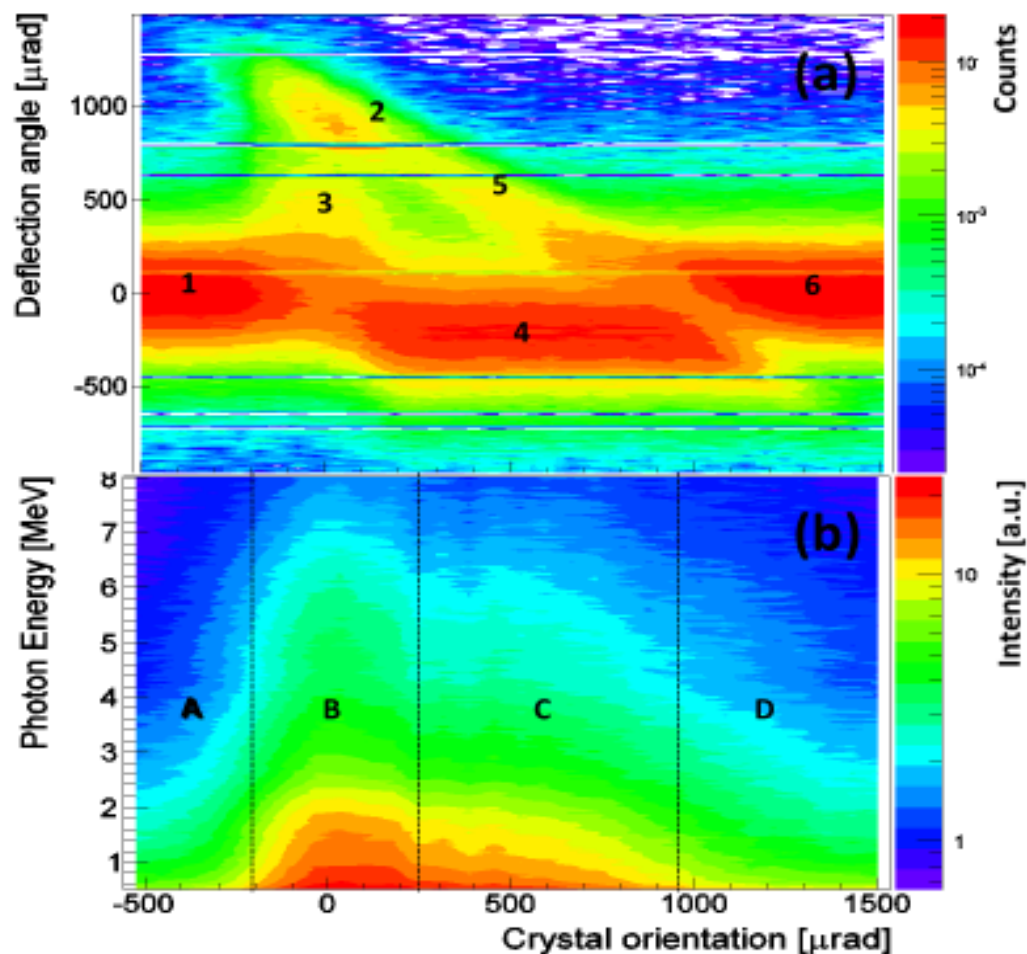
## VOLUME REFLECTION



$\Theta_{\text{inc}} \sim 450 \mu\text{rad}$  Deflection =  $191 \pm 10 \mu\text{rad}$   
 Fraction of reflected is  $76.7 \pm 1.1\%$

**First experimental observation of channeling and VR of negative particles in the sub-GeV energy range**

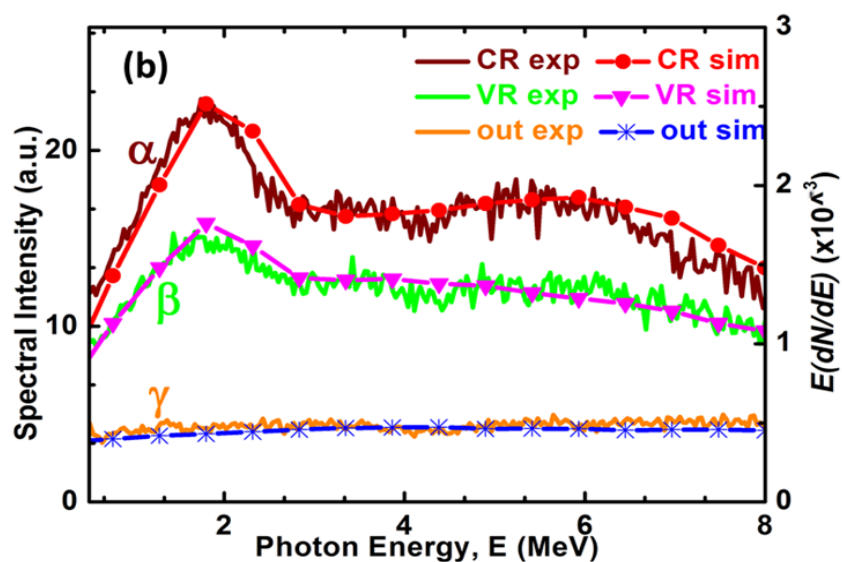
# Experimental results on radiation emission



**Observation: in VR orientation, emitted radiation seems to remain soft and intense as for channeling**

# Simulation of e.m. radiation emitted by ultrarelativistic electrons in a bent crystal

Comparison with experiment performed at the Mainzer Mikrotron with 855 MeV electrons interacting with a 30.5  $\mu\text{m}$  bent Si crystal along the (111) planes

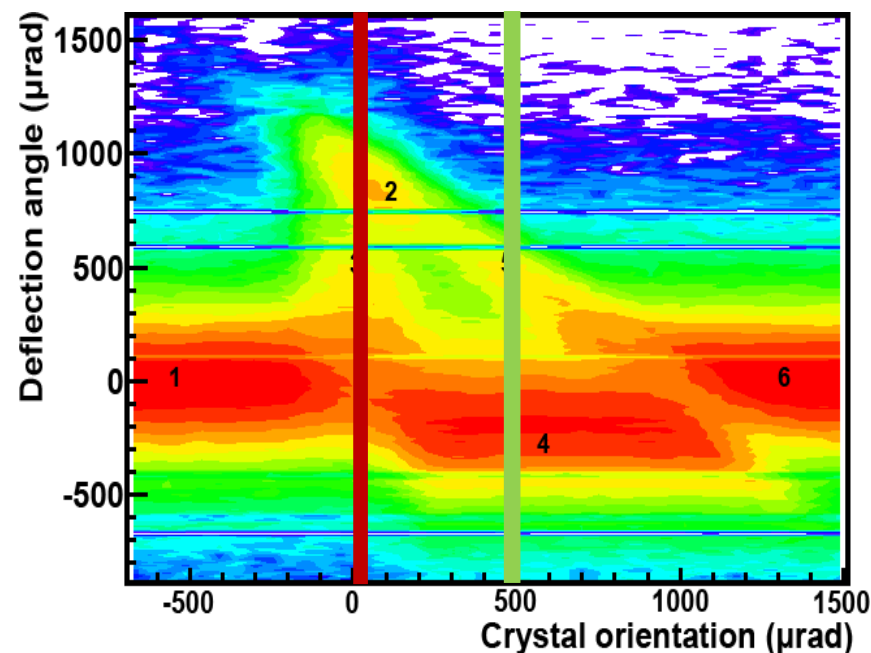


Channeling at <zero> position

Volume reflection at +490  $\mu\text{rad}$

OUT: not aligned at 8 mrad

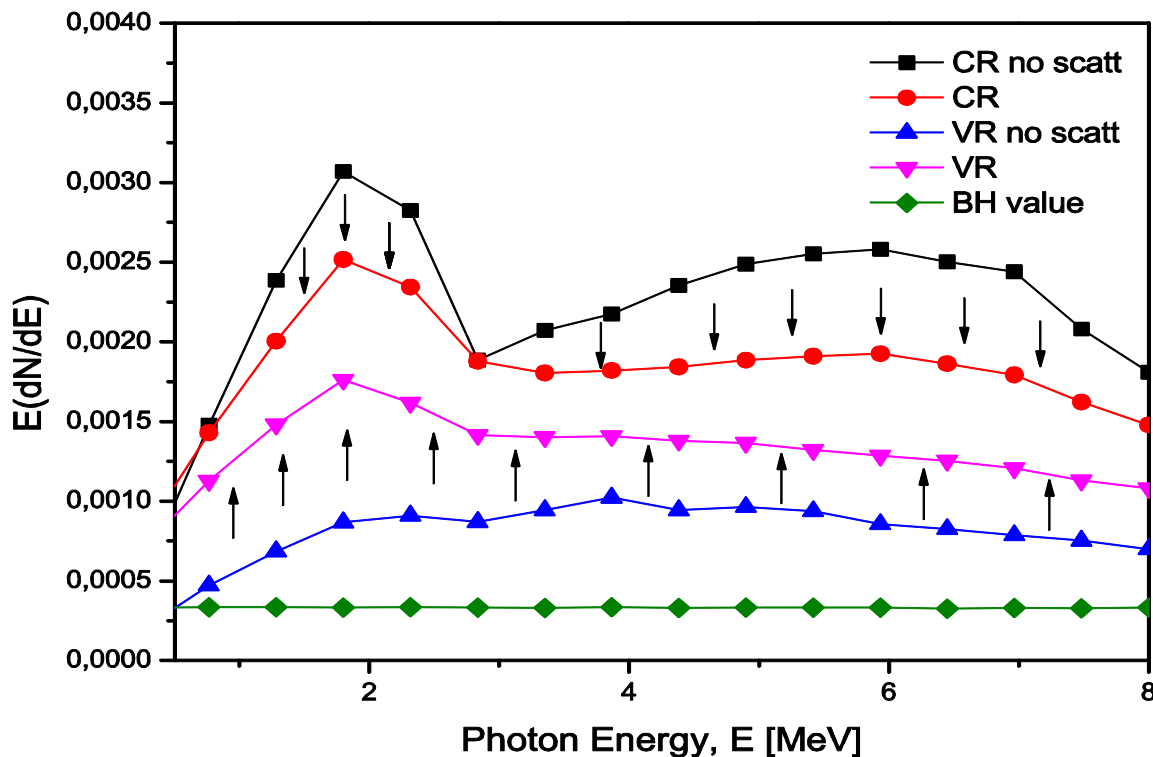
RADCHARM++



Channeling peak at  $E_\gamma \sim 1.8$  MeV  
for 855 MeV electrons  
in (111) Si bent planes

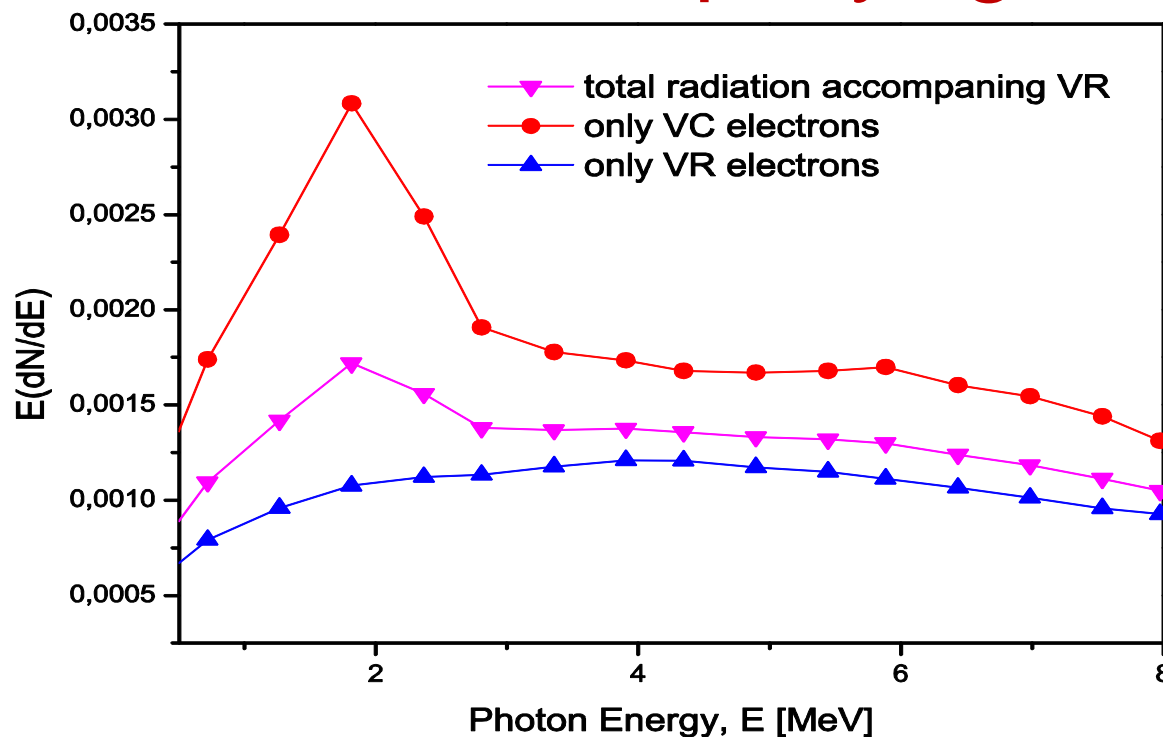


# Simulation of the contribution to radiation of incoherent scattering



**Simulation of the contribution of the scattering with nuclei and electrons to radiation spectral intensity  $[E(dN/dE)]$ .**

# Incoherent scattering contribution to radiation accompanying VR



**32% of volume-captured** and **68% of pure volume-reflected** electrons.  
The contribution of VC particles maintains the electromagnetic radiation accompanying VR close in intensity to that for CR over the whole angular acceptance

# COMPARISON WITH EXPERIMENTS

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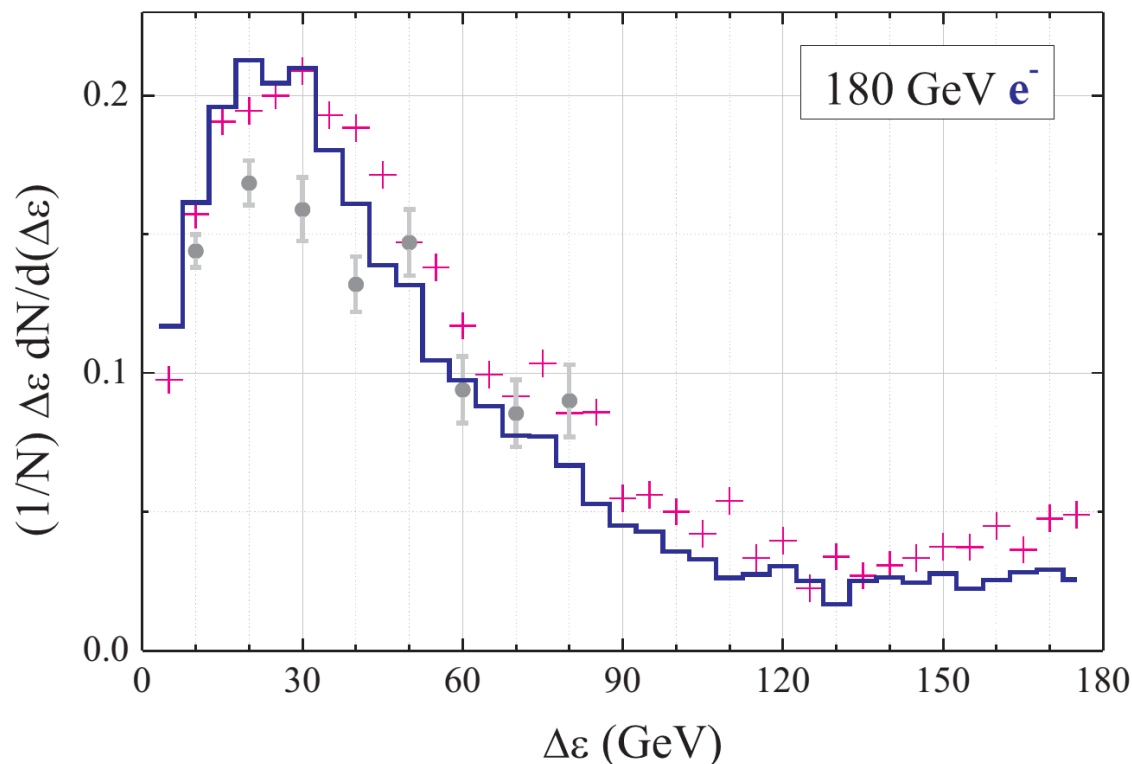
100 GeV energy range:

- Quantum correction;
- Multiple photon emission.

# Multi photon emission

- In principle, the BK formula should be integrated along the whole particle trajectory.
- At very-high energy, the total probability of radiation may exceed unity -> **multiple photon emission!**
- **Separation of particle trajectory in intermediate lengths**  $>$  coherence length and  $\ll$  typical distance between two sequential photon emission points. **Total probability of radiation on such trajectory part does not exceed 0.1.**
- The trajectory-part ends are neglected as the interference between them.

# Simulation of PLANAR volume reflection



## Si Crystal parameters:

Length = 0.84 mm;  
Bending radius = 12 m;  
Plane (111)

## Beam divergence:

$\sigma_x = 25 \mu\text{rad}$  and  $\sigma_y = 46 \mu\text{rad}$

## Incidence angle:

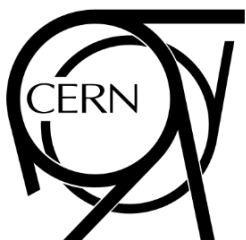
$\Theta_{x0} = 40 \mu\text{rad}$

Photon energy  $\geq 1$  GeV has  
been selected.

**Energy loss spectral intensities:  $(dn/dE)*E$   
of 180 GeV/c volume reflected electrons**

Experiment: W. Scandale, et al., Phys.Rev.A79, 012903 (2009).

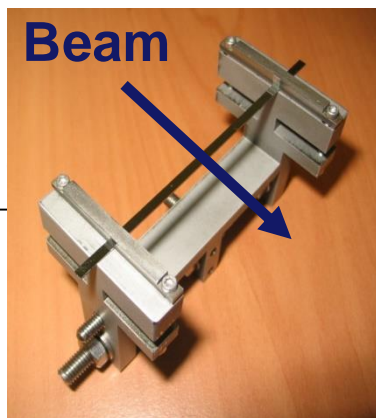
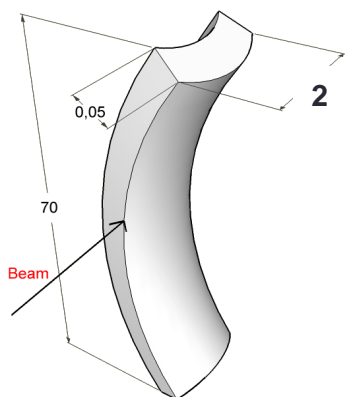
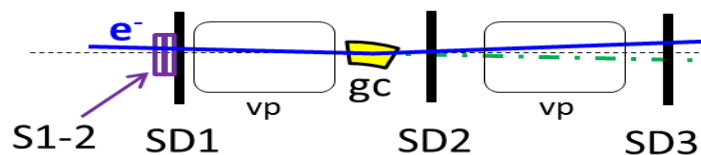
Simulation: V. Guidi, L. Bandiera, [V. Tikhomirov](#), Phys. Rev. A 86 (2012) 042903



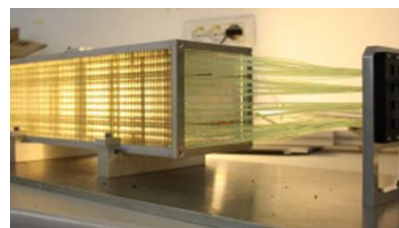
# Interaction of 120 GeV/c electrons with a 2 mm long bent Si crystal

## H4 extracted line from CERN SPS

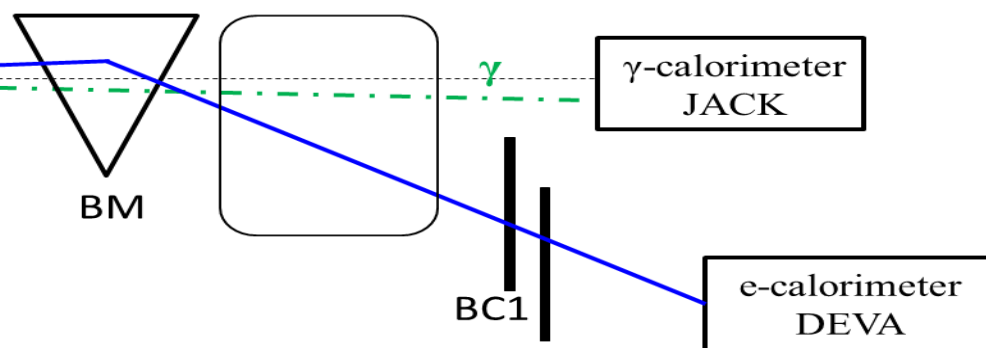
Double sided silicon detectors SDi (300 $\mu$ m)  
[5- $\mu$ m spatial resolution]



(110) Bent Si planes



$\gamma$ -beam *shashlik*  
calorimeter for a total of  
19 X0 to measure the  
emitted photons



A bending magnet (BM) and a pair of two single sided  
silicon beam chambers, BCI, form the  
**SPECTROMETER** to measure the particle momentum



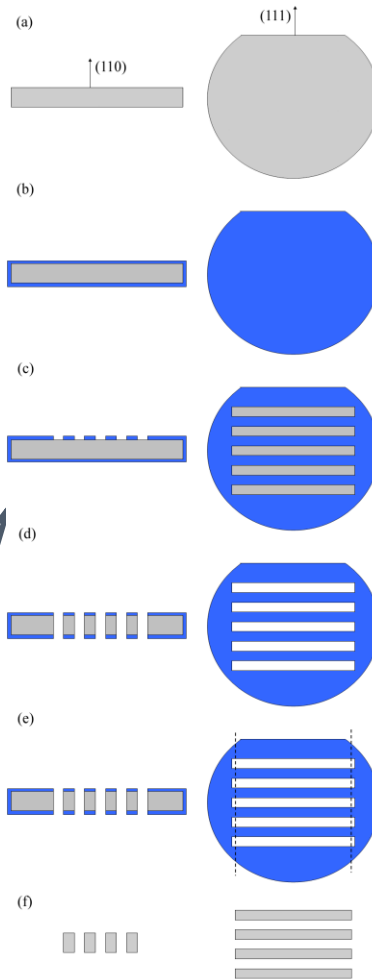
Sensor and  
Semiconductor  
Laboratory

# Silicon Crystal Fabrication

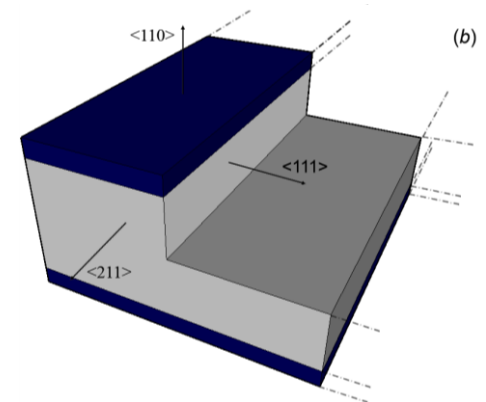
LPCVD deposition of  
silicon nitride thin layer



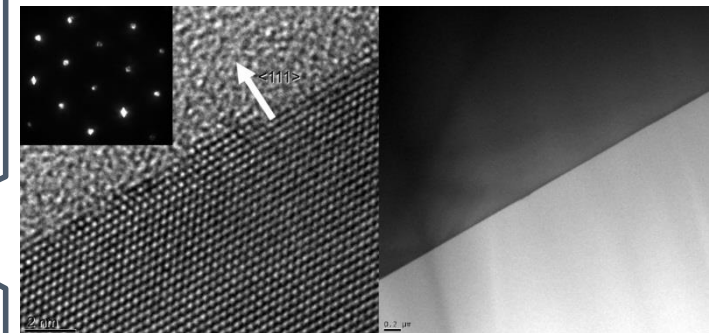
Silicon nitride  
patterning



Anisotropic etching:  
Etching rate on different silicon  
planes for KOH 20% at 40 °C



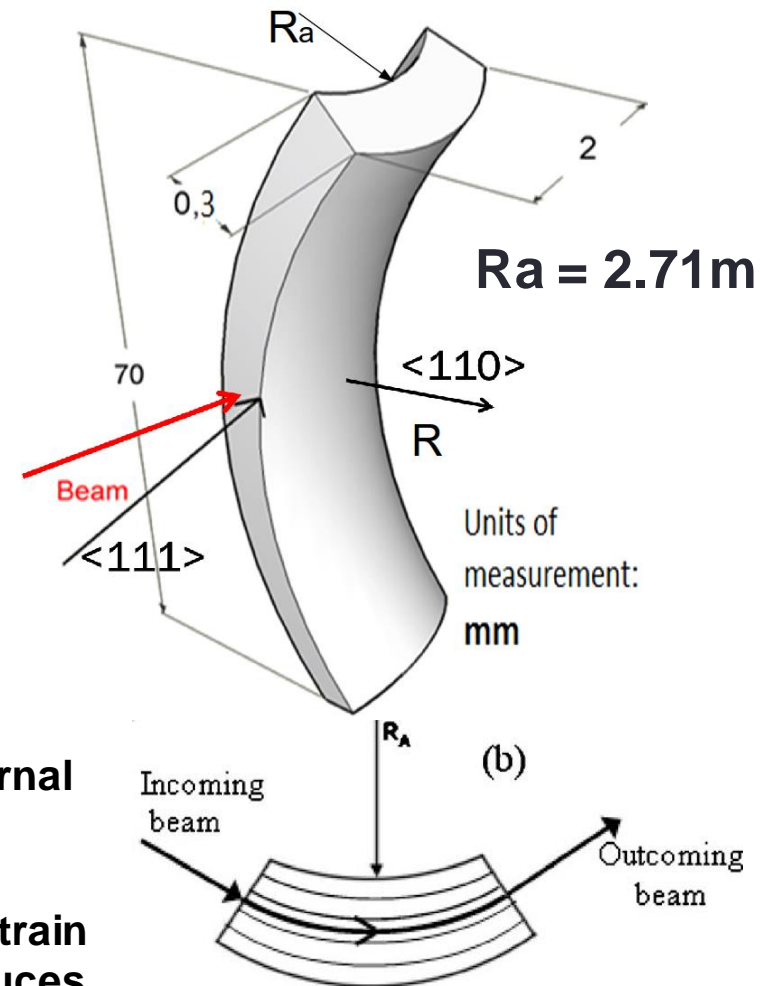
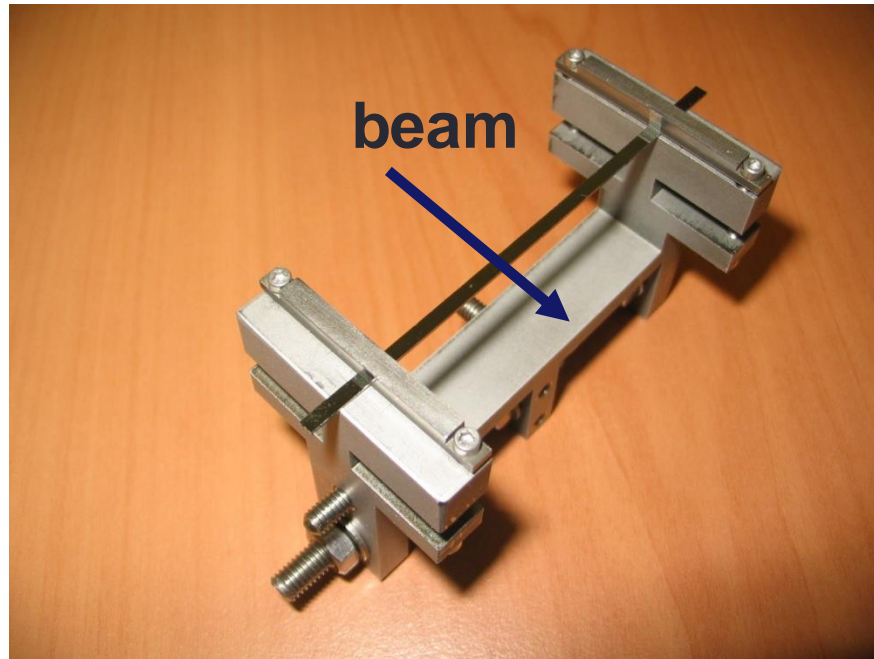
Crystalline surfaces





Sensor and  
Semiconductor  
Laboratory

# Bent Si crystal



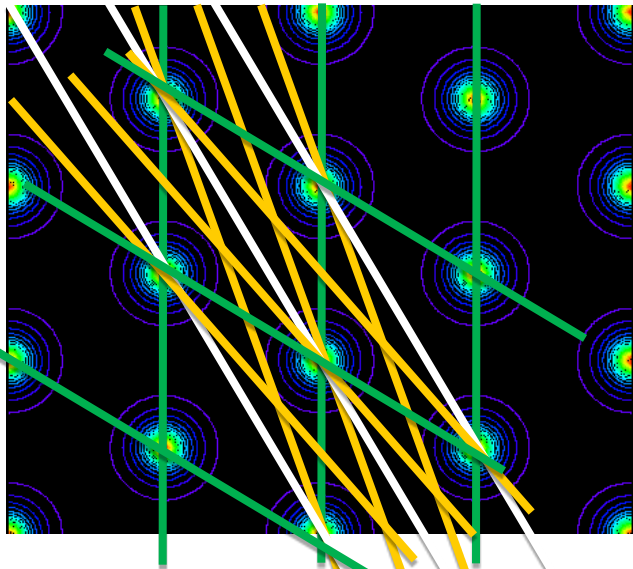
- A primary curvature is imparted by mechanical external forces, which result in a secondary (anticlastic) curvature;
- The mechanical holder used to impart the primary strain can set far apart from the particles and hence it reduces any wanted interaction with the beam.

JAP 107 (2010) 113534



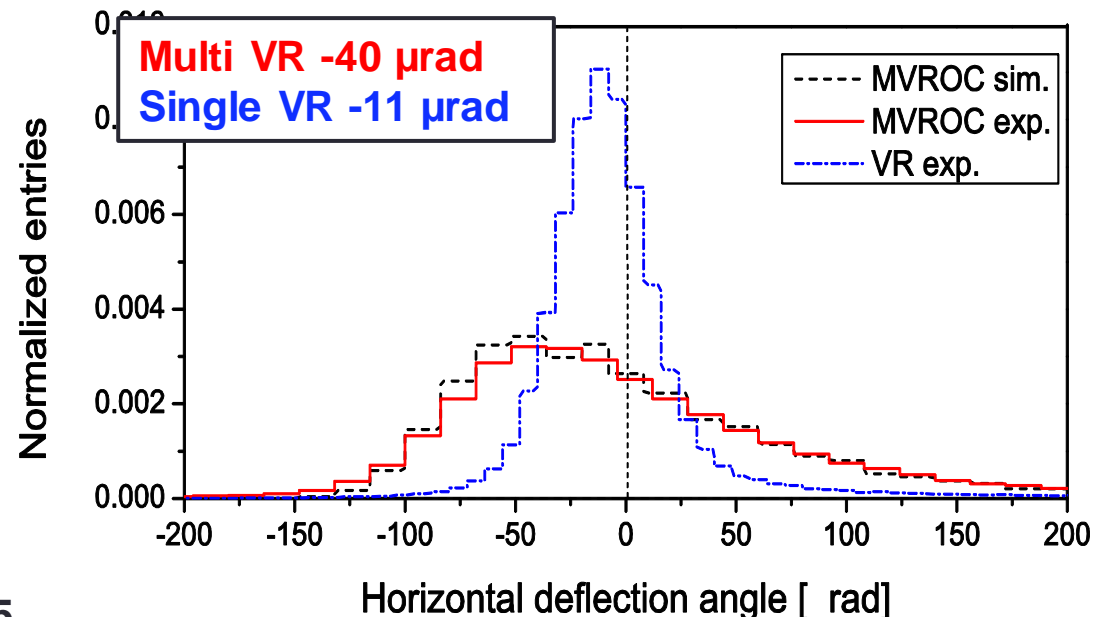
# Single vs Multiple Volume Reflection in One bent Crystal

- ✓ Single VR results in a small deflection angle, which decreases while increasing the particle energy.
- ✓ Possible solution to increase the VR deflection angle: Multi VR from different planes of a bent crystal becomes possible when particles move at a small angle with respect to a crystal axis.

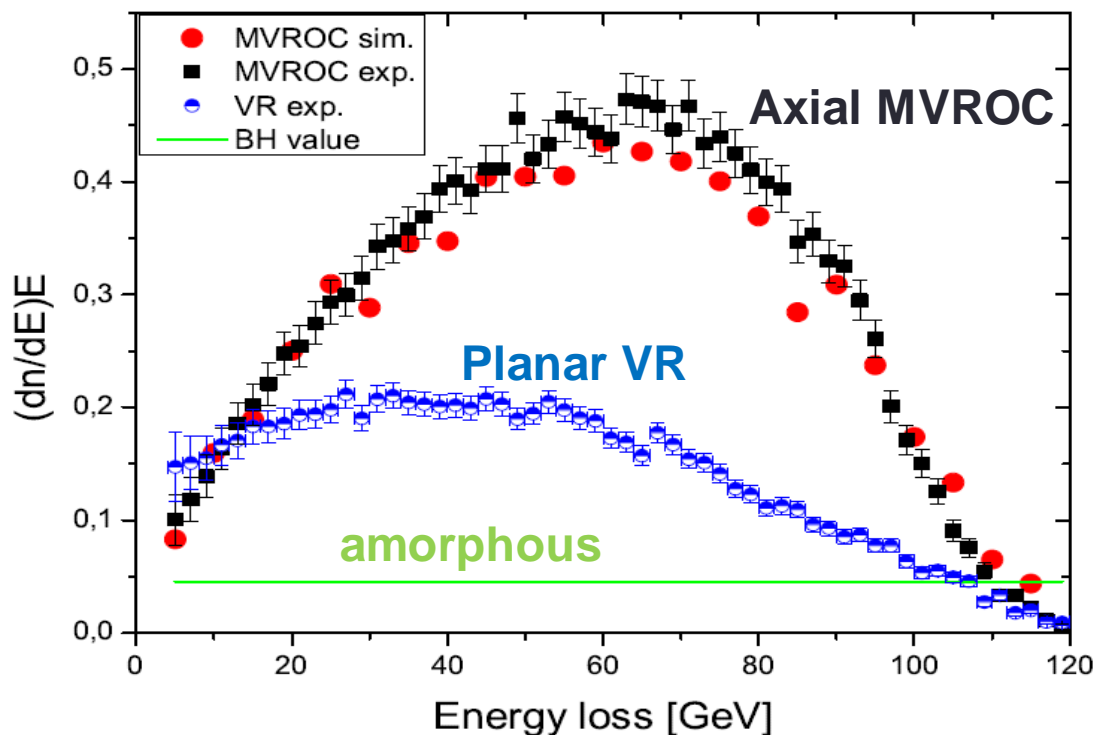


Potential along a Si crystal [111] direction

## Deflection distribution



# Simulation of AXIAL multi-volume reflection



## Si Crystal parameters:

Length = 2mm;

Bending radius = 2.7 m;

Plane (110)

Axis  $\langle 111 \rangle$

## Beam divergence:

$\sigma_x = 50 \mu\text{rad}$  and  $\sigma_y = 65 \mu\text{rad}$

## Incidence angle:

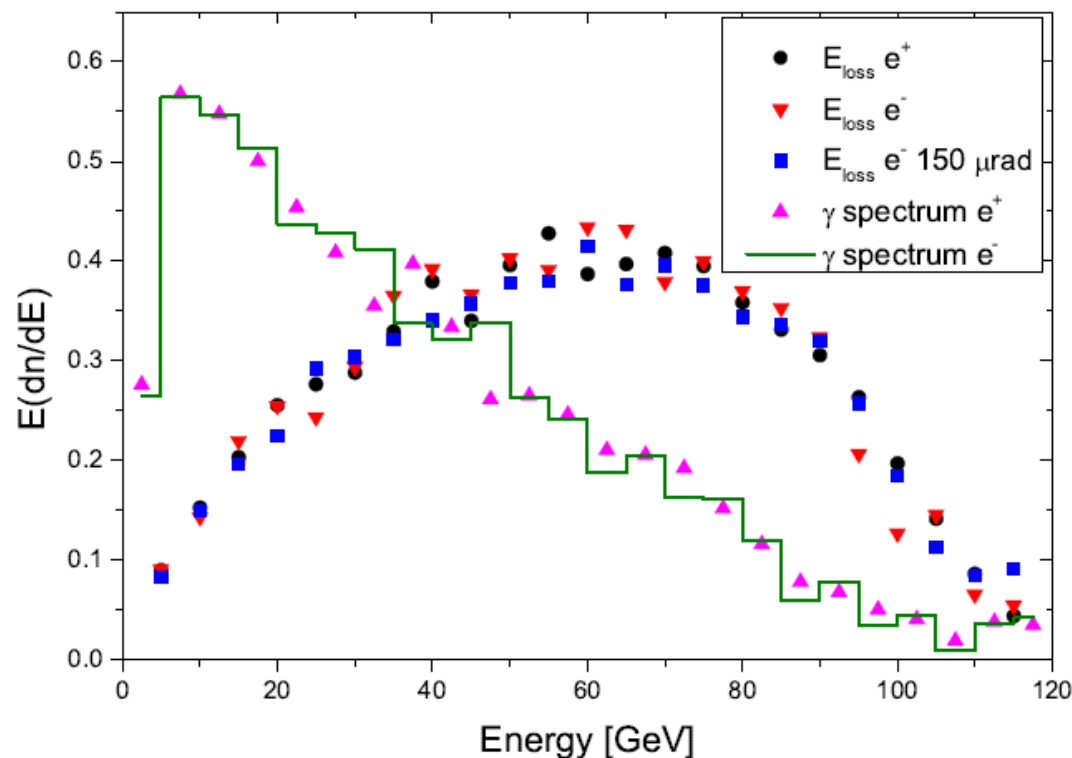
$\Theta_{X0} = 365 \mu\text{rad} \approx \text{half bending}$

$\Theta_{Y0}$  misaligned for planar VR  
 $= 205 \mu\text{rad}$  for axial MVROC

Photon energy  $\geq 1 \text{ GeV}$  has been selected.

**Energy loss spectral intensities:  $(dn/dE)*E$   
of 120 GeV/c single and multi-reflected electrons**

# Radiation features: single photon spectra



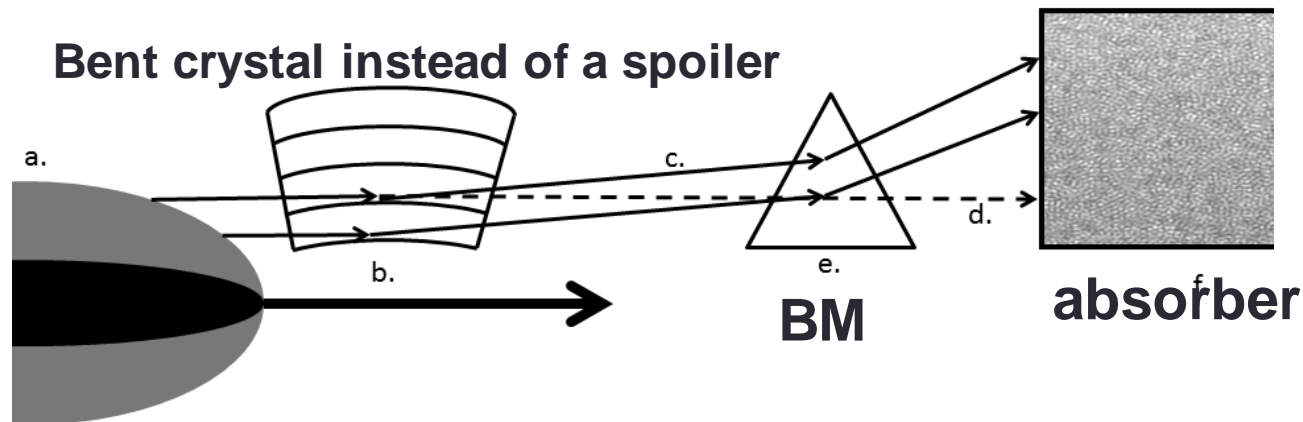
▲ Single photon spectrum for 120 GeV/c positrons

— Single photon spectrum for 120 GeV/c electrons

A mean number of photons emitted by each particles (multiplicity factor) equal to  $2.2/e^-$  for  $\hbar\omega > 1$  GeV.

# Possible application

**Some opportunities in the collimation system of future linear electron/positron colliders, such as ILC**



- The insertion of a short Si crystal ( $0.02-0.05 X_0$ ) instead of a long spoiler ( $0.5-1 X_0$ ) to clean the halo particles would diminish the perturbations on the beam.
- The deflection of particles under single and multiple VR and the increase in energy loss (more than 30% in  $0.02 X_0$ ) may increase the discrimination of halo particles, which are deflected more in the bending magnets placed downstream after the collimator.

**A. Seryi, Nucl. Instrum. Methods Phys. Res., Sect. A 623, 23 (2010)**

**L. Bandiera et al., et al., Journal of Physics: Conference Series 517 (2014) 012043**

# Summarizing

- An algorithm to compute of radiation emitted by relativistic  $e^\pm$  in crystals based on the Baier-Katkov method has been presented;
- Such algorithm has already been implemented in existing Monte Carlo codes for simulation of particle trajectories in crystals;
- Comparison with experiments show a very good agreement in a wide energy range (from 1 GeV to 100 GeV);
- Such a method can be inserted in the most general toolkits for the simulation of the passage of particles through matter, such as Geant 4, as an implementation of the radiation processes in oriented crystalline structures.