



**Reactor θ_{13} measurement and
Double Chooz Multi Detector Results**
(Nu2-WP2, Nu03)

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Contents of this talk

- * **FJPPL and Double Chooz experiment**
- * **Neutrino Oscillation and Reactor θ_{13} measurement**
- * **Double Chooz θ_{13} Result**
- * **Measured θ_{13} and CPV measurement**
- * **Summary**

FJPPL and Double Chooz

Application for AIL Project (2006)

| | | | | | | |
|--------------------------|---|--------------------|--------------------|-----------------------|---------------|------------------------|
| ID: Project Title | Nu_2-WP2: R&D of detectors for future high statistics, high precision experiment (R&D for reactor anti-neutrino experiments) | | | | | |
| Member List | French Group | | | Japanese Group | | |
| | Name | Title | Affiliation | Name | Title | Affiliation |
| | R&D for reactor anti-neutrino experiments | | | | | |
| | Leaders | | | Leader | | |
| | H. de Kerret | Dr. | CNRS/IN2P3/APC | F. Suekane | Assoc. | Tohoku U. |
| | Th. Lasserre | | | | sf. | Niigata U. |
| | | | | | sf. | Tokyo Metropolitan U. |
| | M. Cribier | | | M. Kuze | Associ. Prof. | Tokyo Inst. Technology |
| A. Tonazzo | | Université Paris 7 | | | | |
| D. Krym | | CNRS/IN2P3/APC | | | | |

In 2006, KASKA group (planned reactor θ_{13} exp. in Japan) joined Double Chooz group and applied to FJPPL.

DCJapan group has been funded for ~\$5M for DC thanks to JSPS grants in aid.

FJPPL Report 2016-2017
Fiscal year April 1st 2016 – March 31st 2017
Please replace the red examples by the appropriate data in black

| | |
|--------------------------------------|--|
| Summary of 2016-17 Activities | <p>The year of 2016 was a historical year for Double Chooz. Accumulating Near Detector data since the beginning of 2015, we have first released the Near+Far analysis result in Moriond (March 2016). Since then, we developed a new analysis strategy to enlarge the active neutrino interaction volume by a factor of three, by using not only the Gd-capture events in the Target region, but also the H-capture events in the Gamma-Catcher region. This also mitigated the effect of a tiny leak of Gd-scintillator from Target to Gamma-Catcher region, since Gd and H events are analyzed altogether in one large volume. The cost for the volume gain is the uncertainty in the proton number in the Gamma Catcher region, which was originally not planned to be used as a neutrino target. The evaluation of the systemati</p> <p style="text-align: center;">Precise measurement of neutrino oscillation angle θ_{13} using reactor neutrinos</p> <p>An int value of 0.119 ± 0 interestin measurer experime Currently</p> <p>Another experiments, Daya Bay, Double Chooz and RENO. Initiated by Double Chooz's proposal during London conference, it was realized in October in Seoul, Korea. Limited number of core experts gathered and spent concentrated three days of discussion. This workshop helped very much the understanding between the experiments on the details of analyses by each. The next gathering was promised to happen in 2017 in Paris. The ultimate goal would naturally be the combined value of the mixing angle by the three experiments.</p> <p>This year concludes the four-year Nu-03 program of TYL/FJPPL successfully. The support and annual workshops were very useful and are much appreciated for the French-Japanese collaboration in Double Chooz, which of course will continue even beyond 2017.</p> |
|--------------------------------------|--|

In 2017, the FJPPL-DC is successfully finishing, after the measurements of θ_{13} .

FJPPL and Double Chooz

Nu2-WP2
H.Kerret
F.Suekane

| | | | |
|------|------|--|---|
| June | 2006 | Double Chooz proposal | arXiv:0606025[hep-ex] |
| May | 2008 | Started FD construction | |
| Apr. | 2011 | Started FD data taking | |
| Nov. | 2011 | 1st θ_{13} result (Gd) | Reported in LowNu2011 Phys. Rev. Lett. 108 (2012) 131801 |
| June | 2012 | Started ND construction | |
| Sep. | 2012 | 2nd θ_{13} result (Gd) | Phys. Rev. D 86 (2012) 052008 |
| June | 2013 | 1st θ_{13} result (H) | Phys. Lett. B 723 (2013) 66 |
| Oct. | 2014 | 3rd θ_{13} result (Gd) | JHEP 10 (2014) 086 |
| Jan. | 2015 | Started ND data taking | |
| Jan. | 2016 | 2nd θ_{13} result (H) | JHEP 01 (2016) 163 |
| Mar. | 2016 | θ_{13} result w/ two detectors (Gd) | Reported in Moriond 2016 |
| Sep. | 2016 | θ_{13} result w/ two detectors (Gd+H) | Reported in CERN seminar https://indico.cern.ch/event/548805/ |

Nu03
C.Anatael
M.Kuze

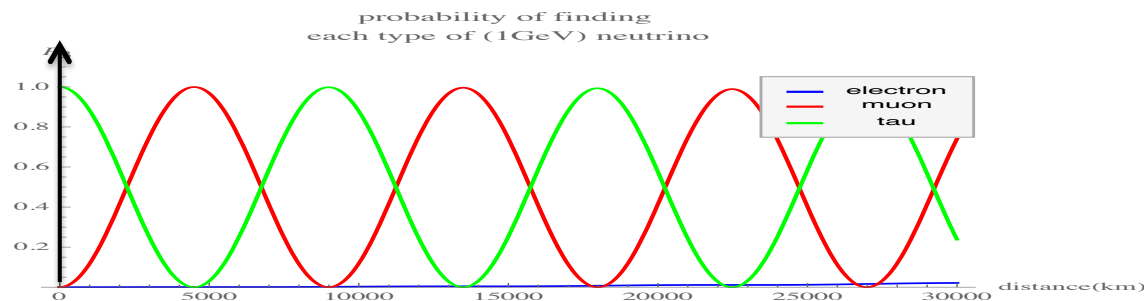
➔ This talk: summary of 10+ years of FJPPL/DC activities.

What is Neutrino Oscillation?

Electron stays as electron while it travels in space.



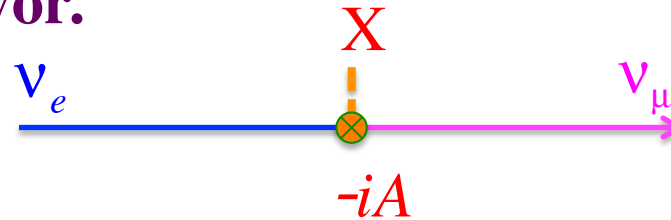
However, neutrinos change their flavors periodically.



This phenomenon is called neutrino oscillation

What causes the neutrino to oscillate?

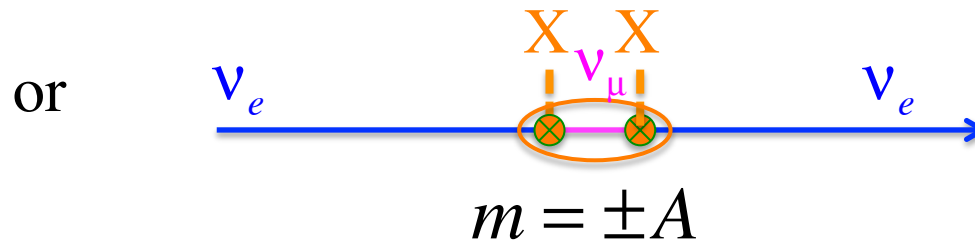
We do not know yet. But in order for N.O. to exist, something(**X**) has to change ν flavor.



"**A**" indicates the strength of the transition (amplitude).

In this case the equations of motion of ν are

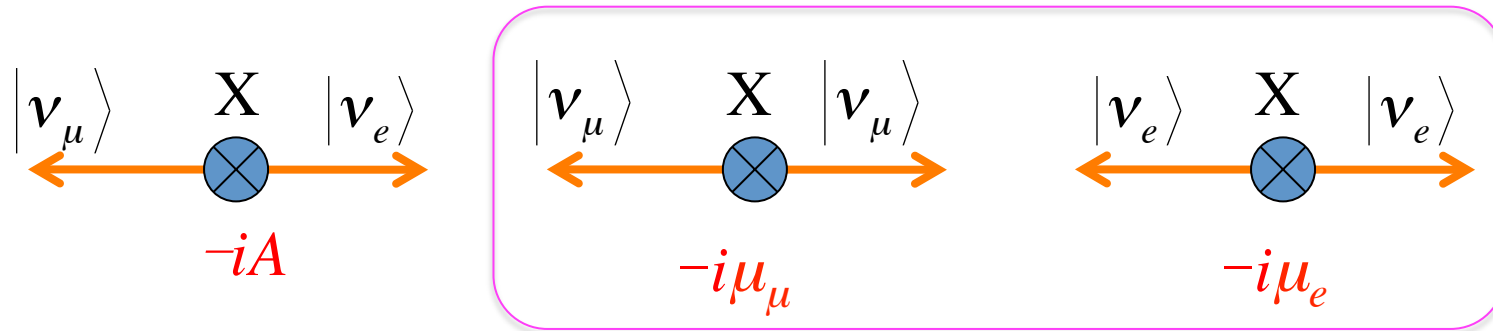
$$\frac{d}{dt} \nu_e = -iA \nu_\mu, \quad \frac{d}{dt} \nu_\mu = -iA \nu_e$$



$$\frac{d^2}{dt^2} \nu_e = -A^2 \nu_e$$

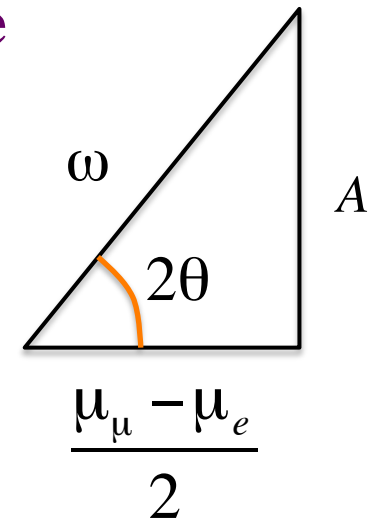
➔ If this transition exists, neutrinos obtain mass.

ν oscillation



If there are **self-transitions**, the mass eigenstate become the superposition of flavor eigenstate;

$$\begin{pmatrix} \nu_1 \\ \nu_2 \end{pmatrix} = \begin{pmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{pmatrix} \begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix}$$



The neutrino masses are

$$\begin{cases} m_1 = \bar{\mu} - \omega \\ m_2 = \bar{\mu} + \omega \end{cases}, \quad \bar{\mu} = \frac{\mu_\mu + \mu_e}{2}, \quad \omega = \frac{A}{\sin 2\theta}$$

Oscillation probability

$$P[\nu_e \rightarrow \nu_\mu] = \left| \begin{array}{c} |\nu_\mu\rangle \\ \uparrow \cos\theta \\ e^{-iE_2 t} |\nu_2\rangle \\ |\nu_2\rangle \\ \downarrow \sin\theta \\ |\nu_e\rangle \end{array} + \begin{array}{c} |\nu_\mu\rangle \\ \uparrow -\sin\theta \\ e^{-iE_1 t} |\nu_1\rangle \\ |\nu_1\rangle \\ \downarrow \cos\theta \\ |\nu_e\rangle \end{array} \right|^2$$

$$P[\nu_e \rightarrow \nu_\mu] = \left| \sin\theta \cos\theta e^{-iE_1 t} - \sin\theta \cos\theta e^{-iE_2 t} \right|^2$$

$$\rightarrow \sin^2 2\theta \sin^2 \frac{m_2^2 - m_1^2}{4E} L$$

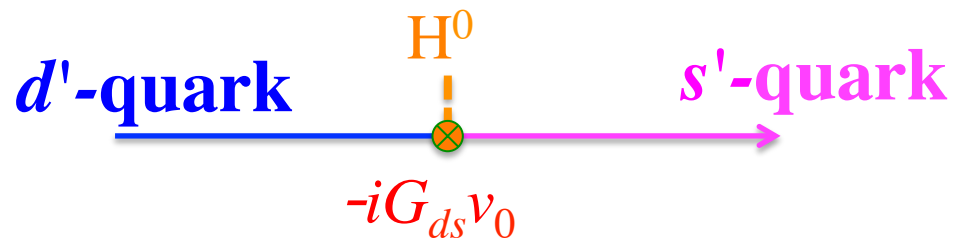
Oscillations are playing important roles

- * $K^0 \Leftrightarrow \bar{K}^0$, $B^0 \Leftrightarrow \bar{B}^0$ Oscillation. \rightarrow CP violation/
- * spin precession by B ($=|\uparrow\rangle \Leftrightarrow |\downarrow\rangle$ oscillation) \rightarrow Formation of Q.M.
- * $|u\bar{u}\rangle \Leftrightarrow |d\bar{d}\rangle$ oscillation in π^0 \rightarrow Hadron mass pattern
- * $d \Leftrightarrow s$ oscillation \rightarrow Cabibbo angle, quark mass.
- * $B \Leftrightarrow W_3$ oscillation \rightarrow Weinberg angle,
 \rightarrow We have learned a lot from these "Oscillations"

We should be able to learn more from ν oscillations, also.

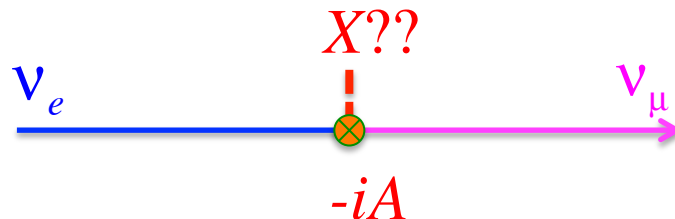
Comparison with quark case

For quark case, there are similar transitions which causes the quark masses, Cabibbo angle and CP violation. The transition is caused by the Yukawa Coupling to the Higgs field



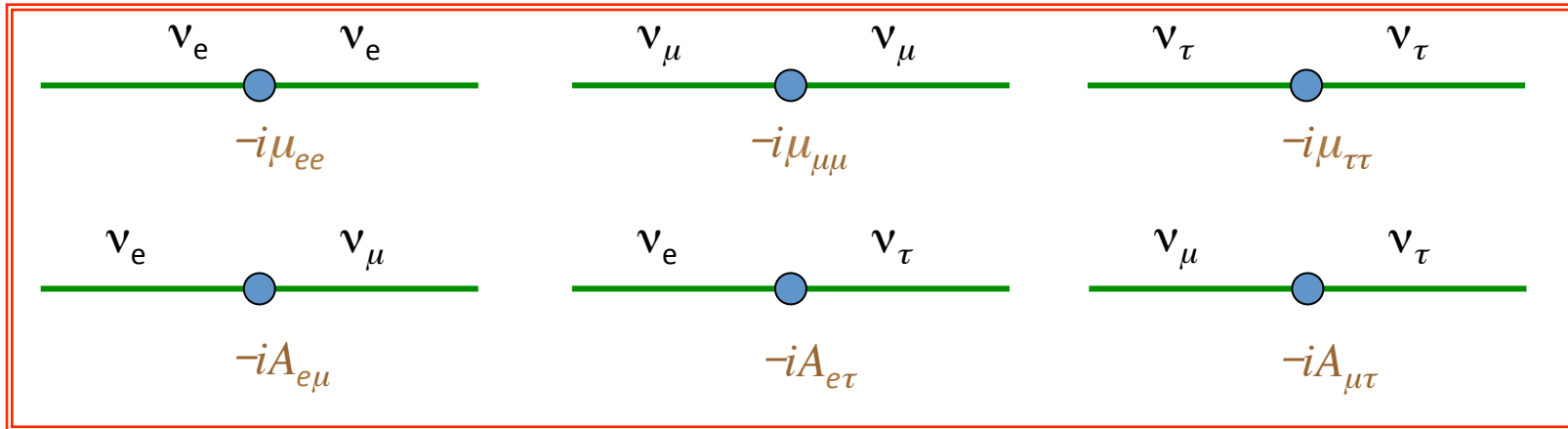
$$\tan 2\theta_C = \frac{2G_{ds}}{G_{ss} - G_{dd}}, \quad m_d = \bar{G}_{ds} - \frac{G_{ds}}{\sin 2\theta_C}, \quad m_s = \bar{G}_{ds} + \frac{G_{ds}}{\sin 2\theta_C}$$

For neutrino case, we do not know what X is.



Study of N.O. = Study of X

3 Flavor Neutrino Case



$A_{\alpha\beta}$ can be complex number

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

There are 6 oscillation parameters;
 $\theta_{12}, \theta_{23}, \theta_{13}, \delta_{CP}, \Delta m^2_{12}, \Delta m^2_{23}$

$$= \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

Prehistory (<2006) of Double Chooz

Atmospheric (Super Kamioka, etc.)

Accelerator (K2K, MINOS)

$$P[\nu_\mu \rightarrow \nu_\mu : @ \Delta m_{23}^2] \sim 1 - \sin^2 2\theta_{23} \frac{\Delta m_{23}^2}{4E}$$

$$\sin^2 2\theta_{23} \sim 1, \quad |\Delta m_{23}^2| \sim 3 \times 10^{-3} \text{ eV}^2$$

Solar (SNO etc.)

Reactor (KamLAND)

$$P[\nu_e \rightarrow \nu_e : @ \Delta m_{12}^2] \sim 1 - \sin^2 2\theta_{12} \frac{\Delta m_{12}^2 L}{4E}$$

$$\sin^2 2\theta_{12} \sim 0.85, \quad \Delta m_{12}^2 \sim 8 \times 10^{-5} \text{ eV}^2$$

Reactor (Chooz, Palo Verde)

$$P[\nu_e \rightarrow \nu_e : @ \Delta m_{23}^2] \sim 1 - \sin^2 2\theta_{13} \frac{\Delta m_{23}^2 L}{4E}$$

$$\sin^2 2\theta_{13} < 0.1$$

$$\theta_{12}, \theta_{23}, \theta_{13}, \delta_{CP}, \Delta m^2_{12}, \Delta m^2_{23}$$

→ θ_{13} & δ_{CP} measurements were next important step.

(1) Reactor Measurement of θ_{13} :

$$P[\bar{\nu}_e \rightarrow \bar{\nu}_e; @ \Delta m^2_{23}] \sim 1 - \sin^2 2\theta_{13}$$

→ Pure measurement of θ_{13} .

← Double Chooz

(2) Accelerator Measurement :

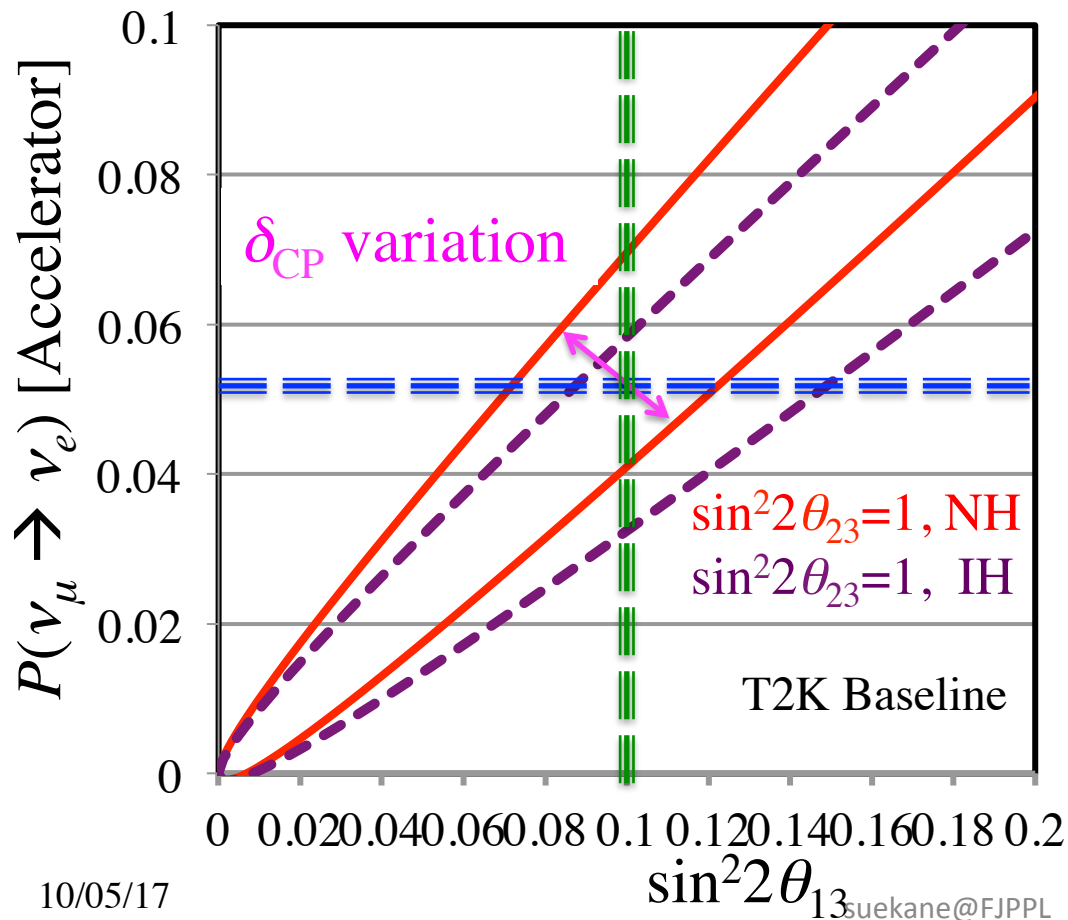
$$P[\nu_\mu \rightarrow \nu_e; @ \Delta m^2_{23}] \sim 0.5 \sin^2 2\theta_{13} - 0.04 \sin 2\theta_{13} \sin \delta_{CP}$$

→ Depends on both θ_{13} & δ_{CP}

(1) + (2) → θ_{13} & δ_{CP}

Reactor Accelerator complementarity

$$P[\nu_\mu \rightarrow \nu_e; @ \Delta m_{23}^2] \sim \frac{0.5 \sin^2 2\theta_{13}}{(1 - a_m)^2} - 0.04 \frac{\sin 2\theta_{13}}{1 - a_m} \sin \delta_{CP}$$



(Matter effect exists.
 $\sim \pm 0.05$ for T2K.
 Sign depends on the
 mass hierarchy.)

If a_m is known, δ_{CP} can
 be determined by combining
 reactor and accelerator
 experiments.

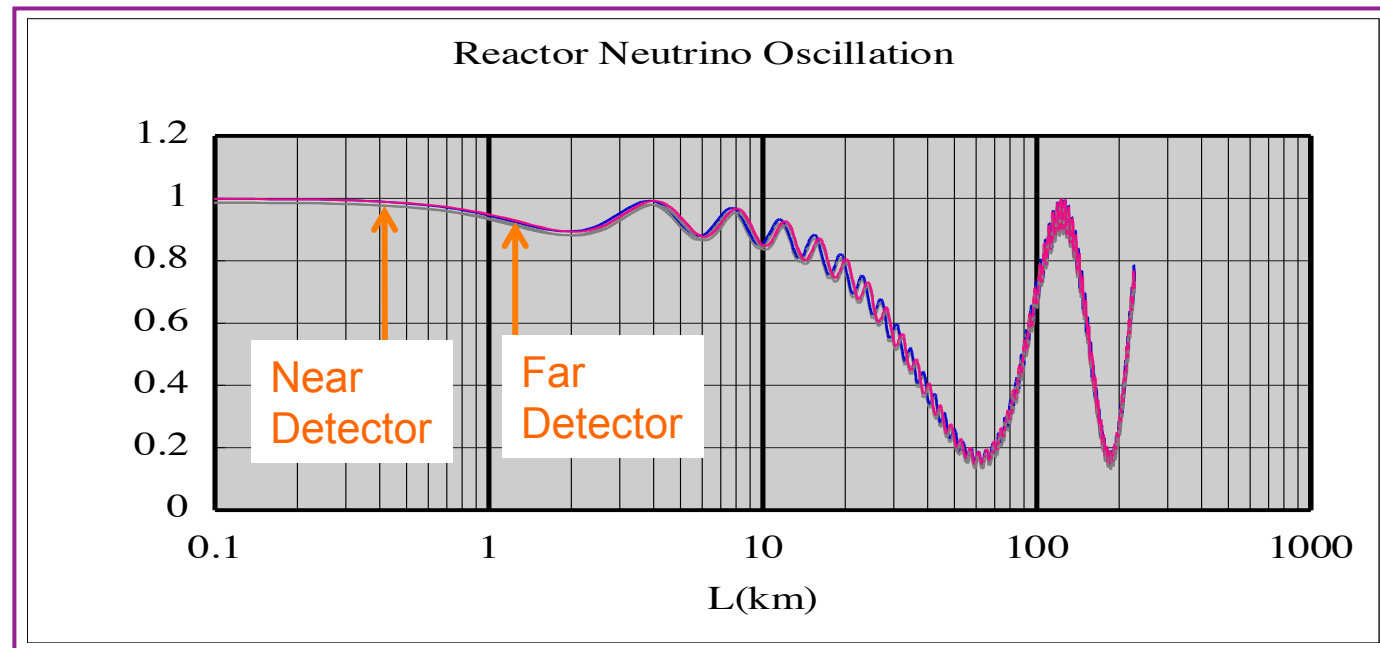
Reactor measurement of θ_{13}

$$P[\bar{\nu}_e \rightarrow \bar{\nu}_e; @ \Delta m_{23}^2] \sim 1 - \sin^2 2\theta_{13}$$

Difficulty: To achieve $\delta \sin^2 2\theta_{13} = 0.01$, we need to distinguish 99000 events from 100000 events.

However, usually we suffer from a few % of error on both ν flux and detection efficiency.

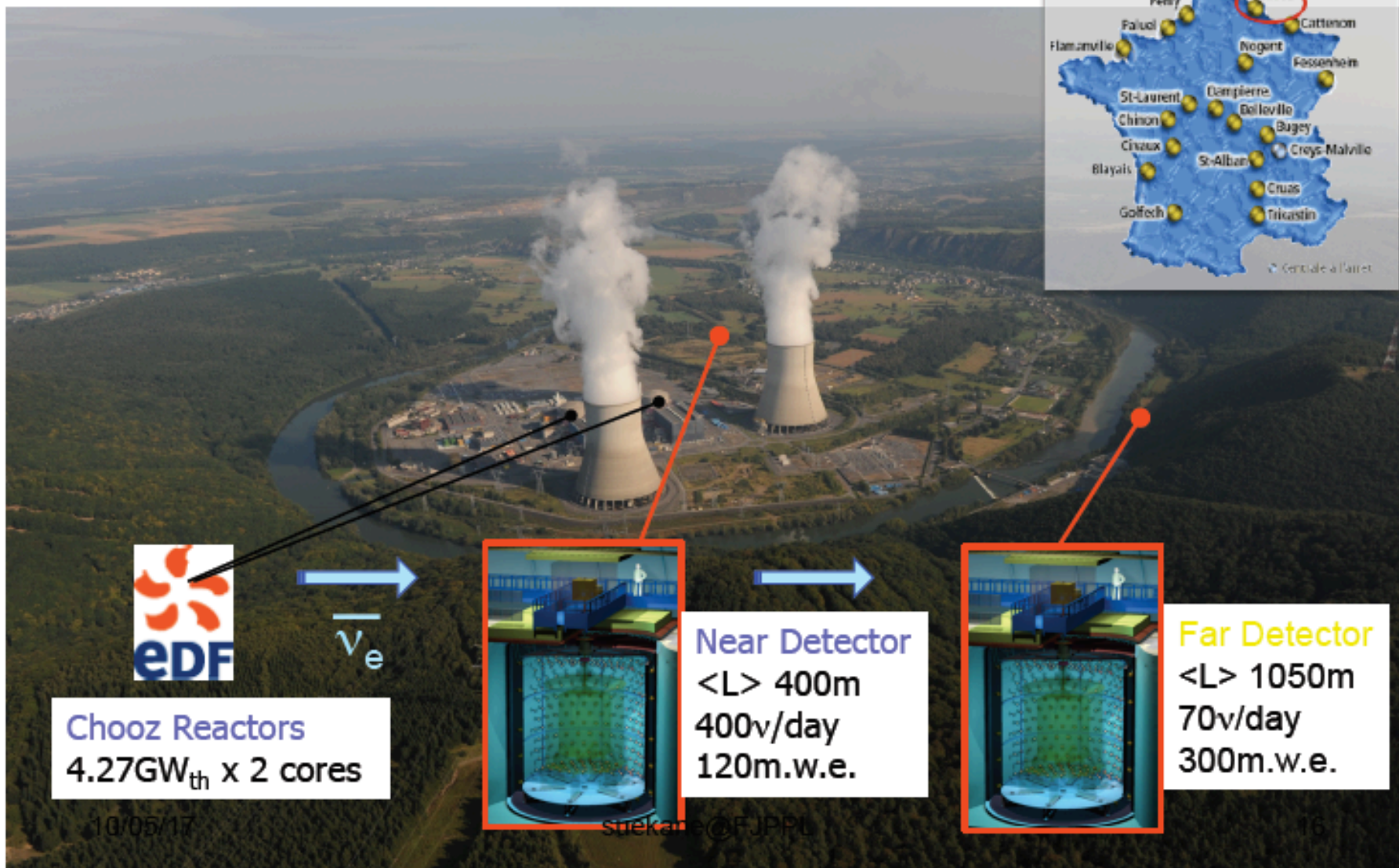
$$P(\bar{\nu}_e \rightarrow \bar{\nu}_e)$$



➔ **Near/Far multi detector scheme:**

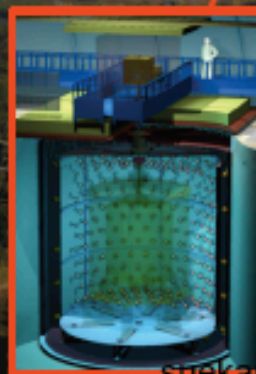
Use Near detector to cancel the systematic errors

Then we are performing Double Chooz experiment to measure Pure θ_{13}

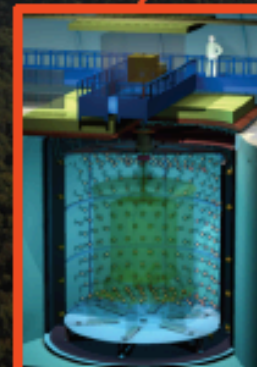


Chooz Reactors
4.27GW_{th} x 2 cores

$\bar{\nu}_e$



Near Detector
<L> 400m
400v/day
120m.w.e.



Far Detector
<L> 1050m
70v/day
300m.w.e.

Double Chooz collaboration



Brazil

CBPF
UNICAMP
UFABC



France

APC
CEA/DSM/
IRFU:
SPP, SPhN
SEDI, SIS
SENAC
CNRS/IN2P3:
Subatech
IPHC



Germany

EKU
Tübingen
MPIK
Heidelberg
RWTH
Aachen
TU München



Japan

Tohoku U.
Tokyo Inst. Tech.
Tokyo Metro. U.
Kitasato U.
Kobe U.
Tohoku Gakuin U.
Hiroshima Inst.
Tech.



Russia

INR RAS
IPC RAS
RRC
Kurchatov



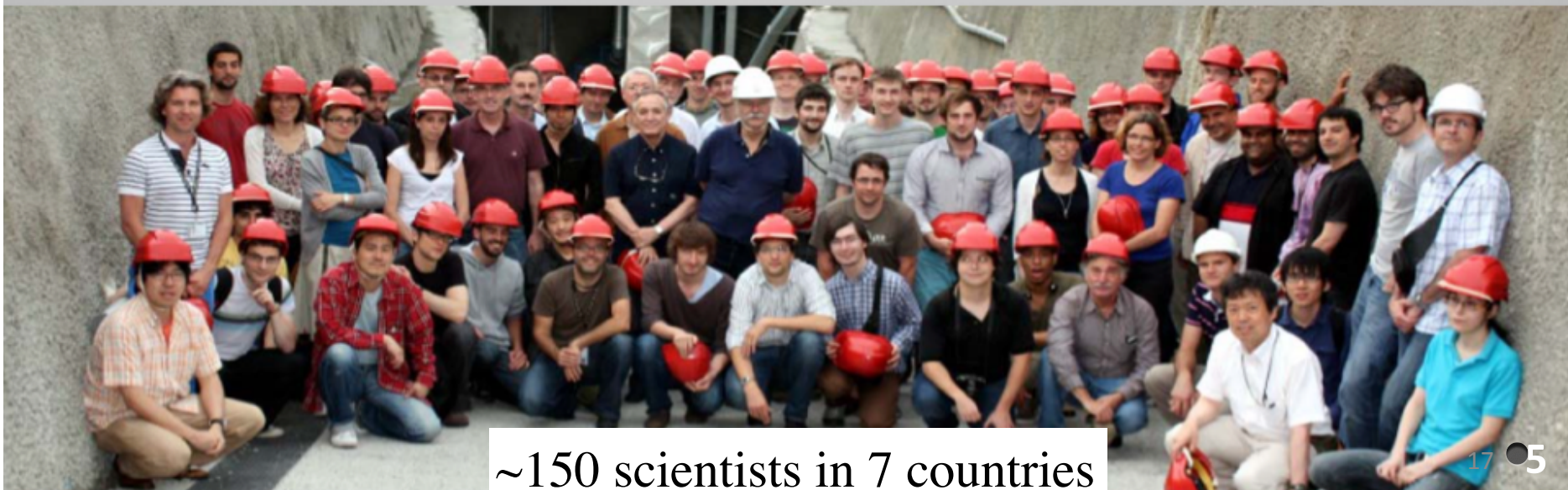
Spain

CIEMAT-
Madrid



USA

U. Alabama
ANL, U. Chicago
Columbia U.
UC Davis
Drexel U.
IIT, KSU, MIT,
U. Notre Dame
U. Tennessee
Virginia Tech



~150 scientists in 7 countries

Main Components of DC Detector

Target ν :
10m³ Gd loaded Liquid Scintillator
8mm Acrylic Tank

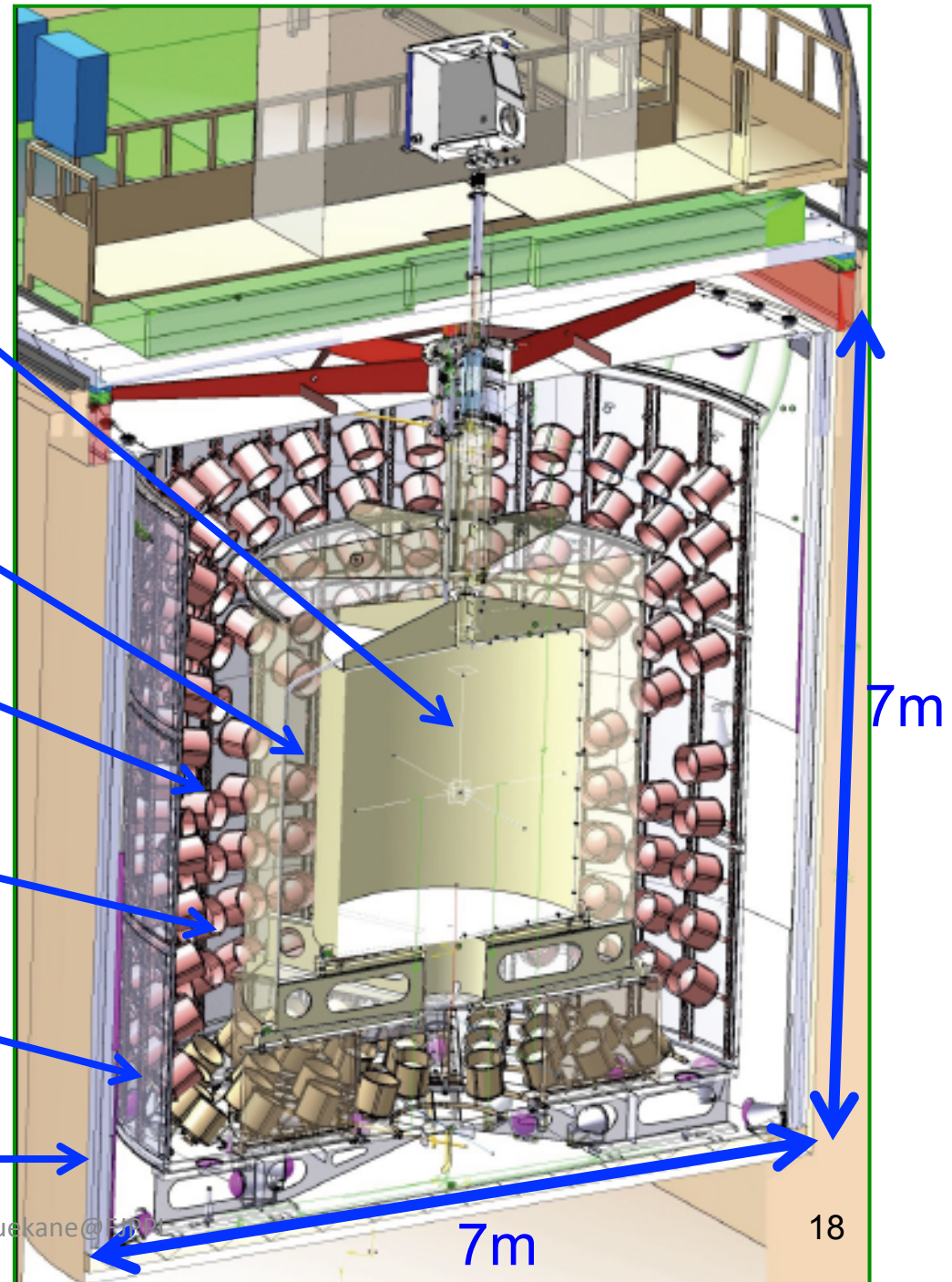
γ Catcher :
22m³ Liquid Scintillator
12mm Acrylic Tank

Light Detection:
390 Low BKG 10" PMTs

Buffer oil :
110m³ Paraffine Oil
3mm Stainless Steel Tank

Inner Muon Veto :
90m³ LS + 78 8" PMTs

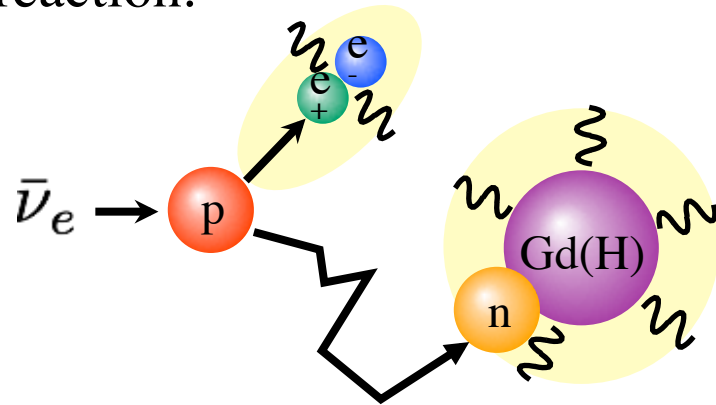
shield:
15cm Iron (Far), water (Near)



Detection principle

Inverse Beta Decay

IBD reaction:
 $\bar{\nu}_e + p \rightarrow e^+ + n$



$$E_\nu = E_{\text{vis}} + 0.78 \text{ MeV}$$

→ θ_{13} oscillation analysis w/ spectral shape gives further constraint

Delayed coincidence:

- Prompt signal

e^+ ionization & annihilation:

$$E_{\text{prompt}} = 1 \sim 8 \text{ MeV}$$

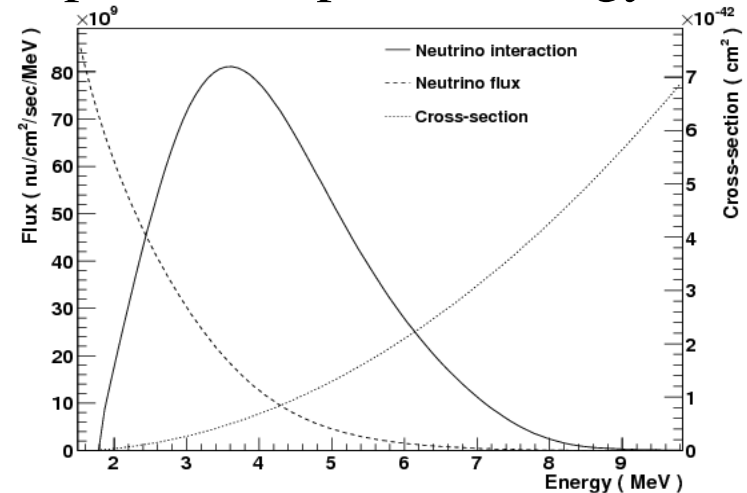
- Delayed signal

n capture on Gd (H):

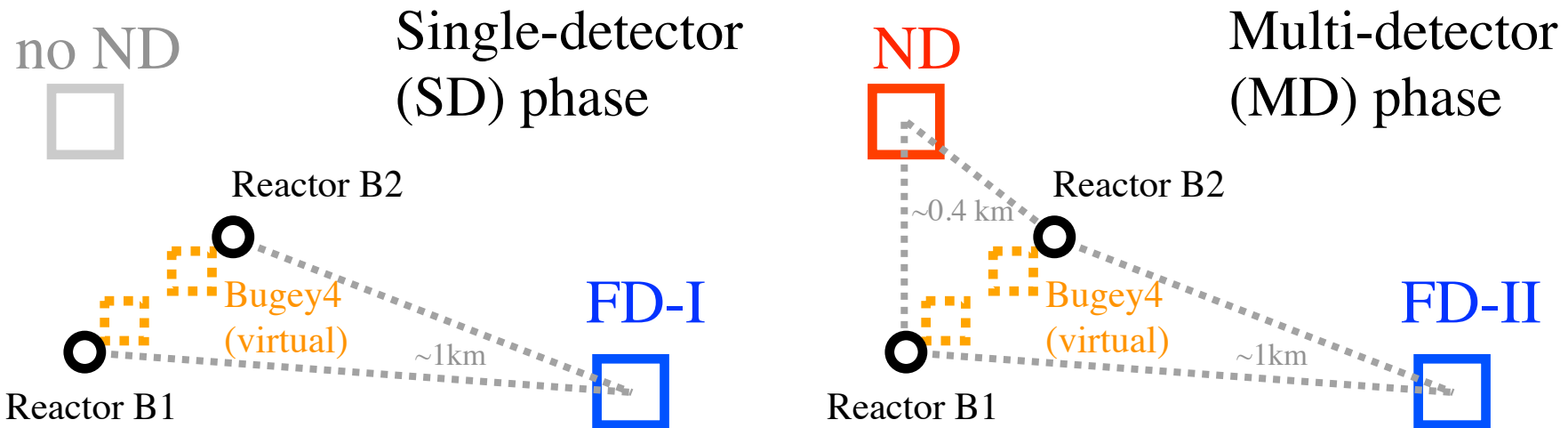
$$E_{\text{delayed}} = \sim 8 (\sim 2.2) \text{ MeV}$$

- Time coincidence of those

Spectral shape of ν energy



Error suppression by two detectors



(*) FD-I and FD-II data from same detector

SD phase (2R1D setup : **FD-I** & **Reactor-off** data)

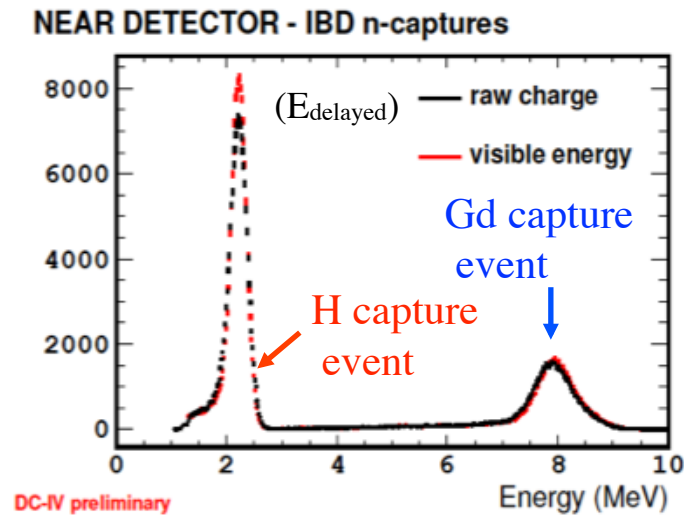
- Bugey4 is used as an anchor of reactor ν flux ($\sim 1.7\%$ of total flux precision)
- Reactor-off data (~ 7 days) is used to constrain BG

MD phase (2R2D setup : **FD-II** & **ND** data)

- Nearly iso-flux setup can suppress ν flux error ($\sim 0.1\%$ of total flux precision)
- Identical detector cancels correlated errors like detection efficiency

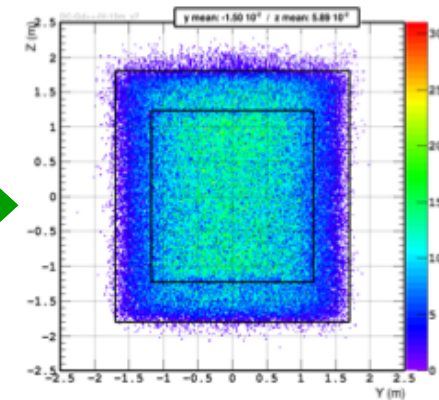
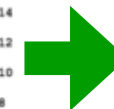
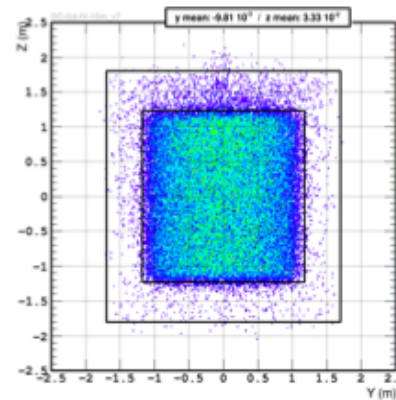
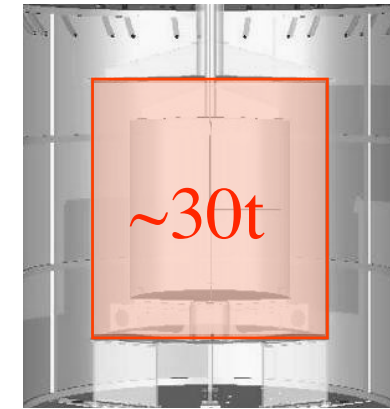
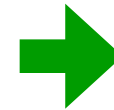
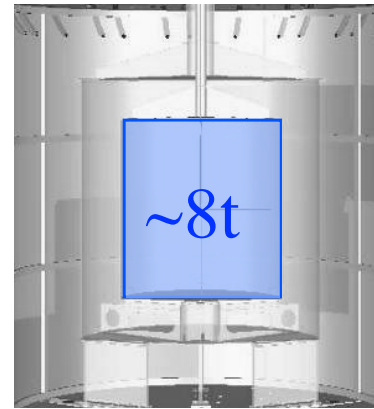
Boosted statistics by Gd+H analysis

Enlargement of ν target volume



| IBD rate | Gd analysis | Gd+H analysis |
|----------|---------------------------|---------------------------|
| FD | $\sim 40 \text{ d}^{-1}$ | $\sim 100 \text{ d}^{-1}$ |
| ND | $\sim 300 \text{ d}^{-1}$ | $\sim 800 \text{ d}^{-1}$ |

~ 2.5 times



N_ν ~ 2.5 times of original design by Gd+H analysis, trading off a part of cancellation of detection efficiency systematics.

BG veto & leak @ ND

Backgrounds

- Accidental coincidence: e.g.) environmental γ + spallation n
- Fast n / Stopping μ : $n + p \rightarrow \text{recoil } p + n$ $\mu \rightarrow e + \nu + \nu$
- Spallation product: e.g.) ${}^9\text{Li} \rightarrow {}^8\text{Be} + e + \nu + n$

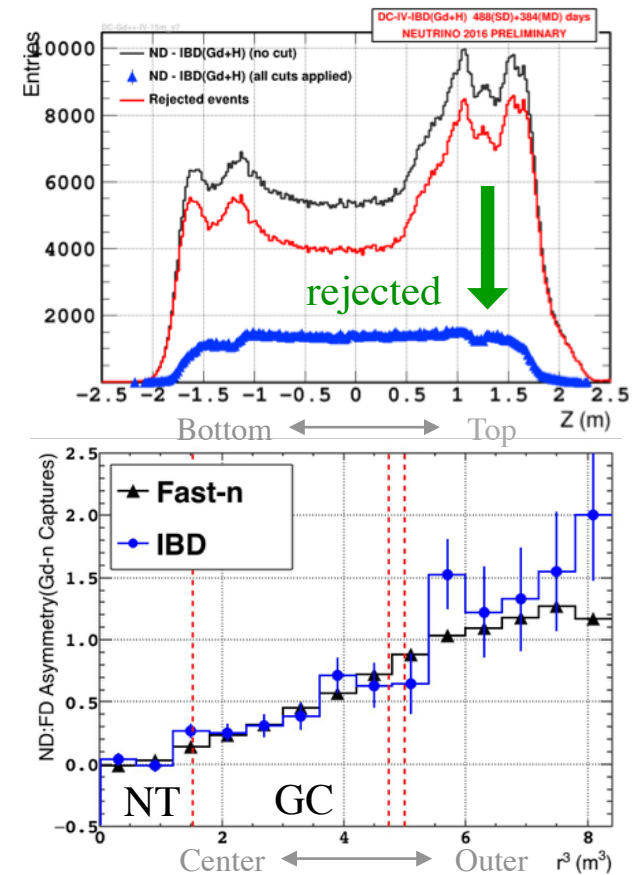
→ Vetoed by dedicated cuts like ANN

LS on Buffer @ ND

- Increased Stop- μ BG. Rejected by BG veto
→ No effect in our analysis (ND:FD consistent)
- Cause is not evident (Filling or Running?)
→ Monitoring stability

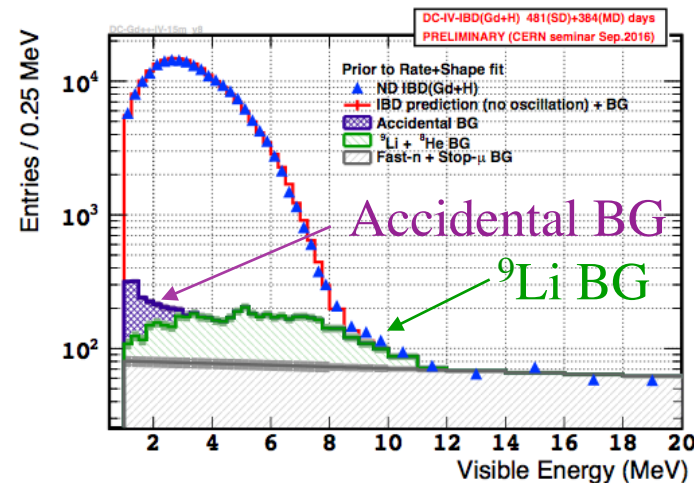
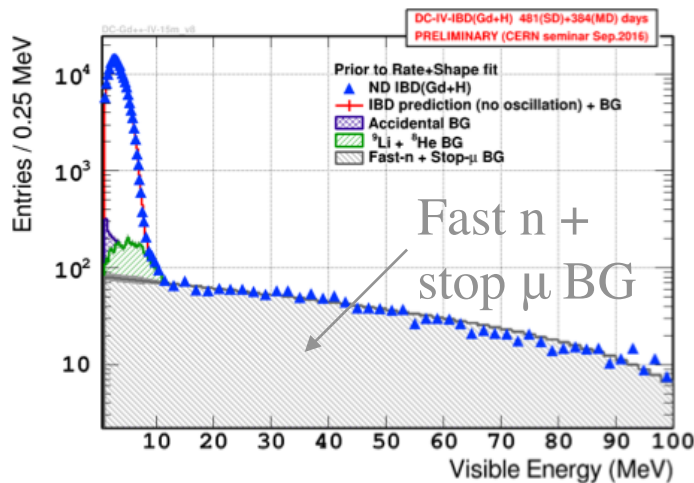
Gd concentration in GC @ ND

- Found in comparison with ND and FD
→ No effect in Gd+H analysis (w/ both volumes)
- Estimating effect to Gd analysis (x-ch)



Remaining BG estimation

- All backgrounds are measured from data
 - -Accidental BG: Off-time coincidence (Rate & Shape)
 - -Fast n + Stopping μ BG : High energy window (Rate)
IV/OV tagging (Shape)
 - - ${}^9\text{Li}$ BG : ${}^9\text{Li}$ enriched data (Shape)

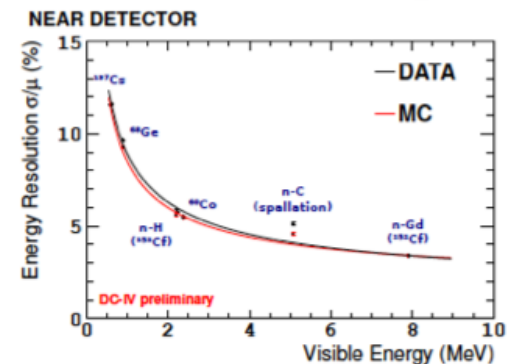
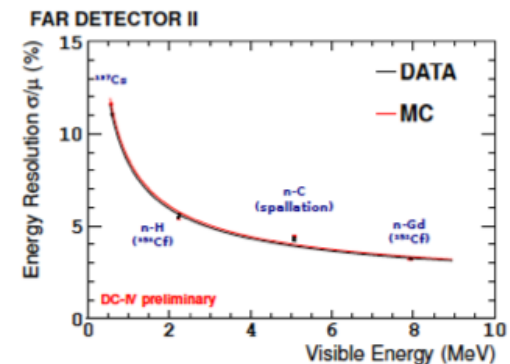
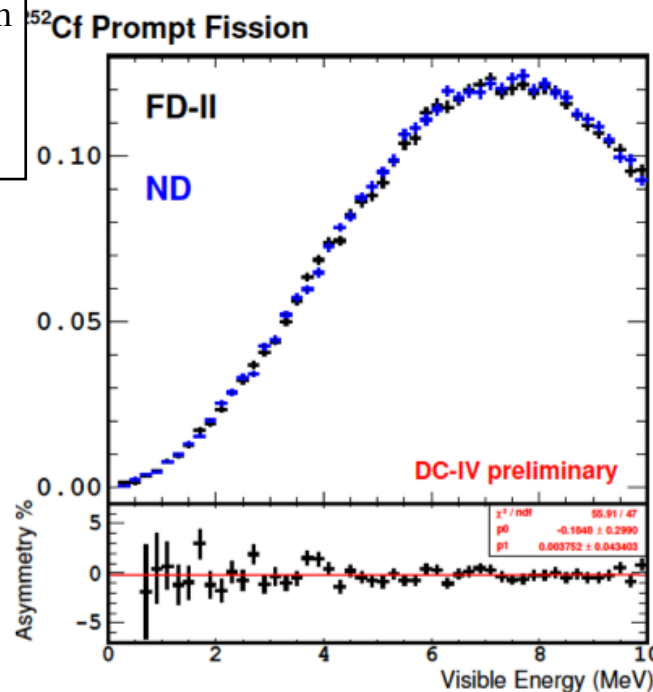
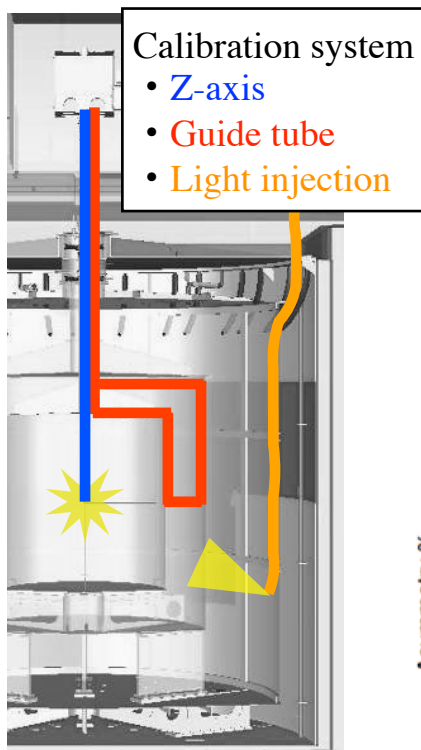


- All backgrounds have characteristic spectrum
- Both “Rate & Shape” are used in oscillation analysis except for ${}^9\text{Li}$ rate
→ ${}^9\text{Li}$ BG rate is constrained by the shape in the fit

Detector response

Important to understand detector response in ND and FD

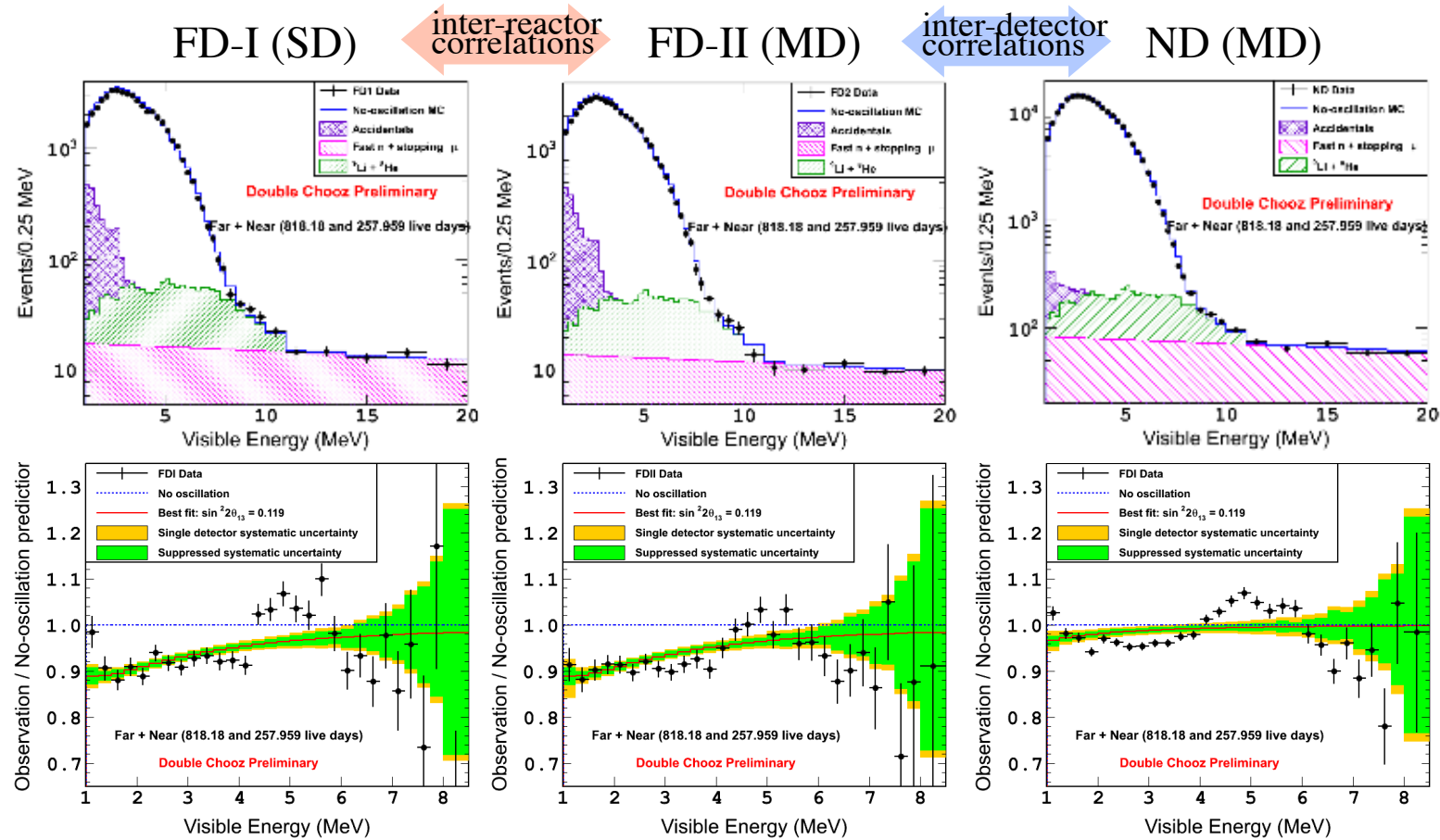
- Electronics calibration by the Light injection system
- Energy calibration by deployment and natural sources



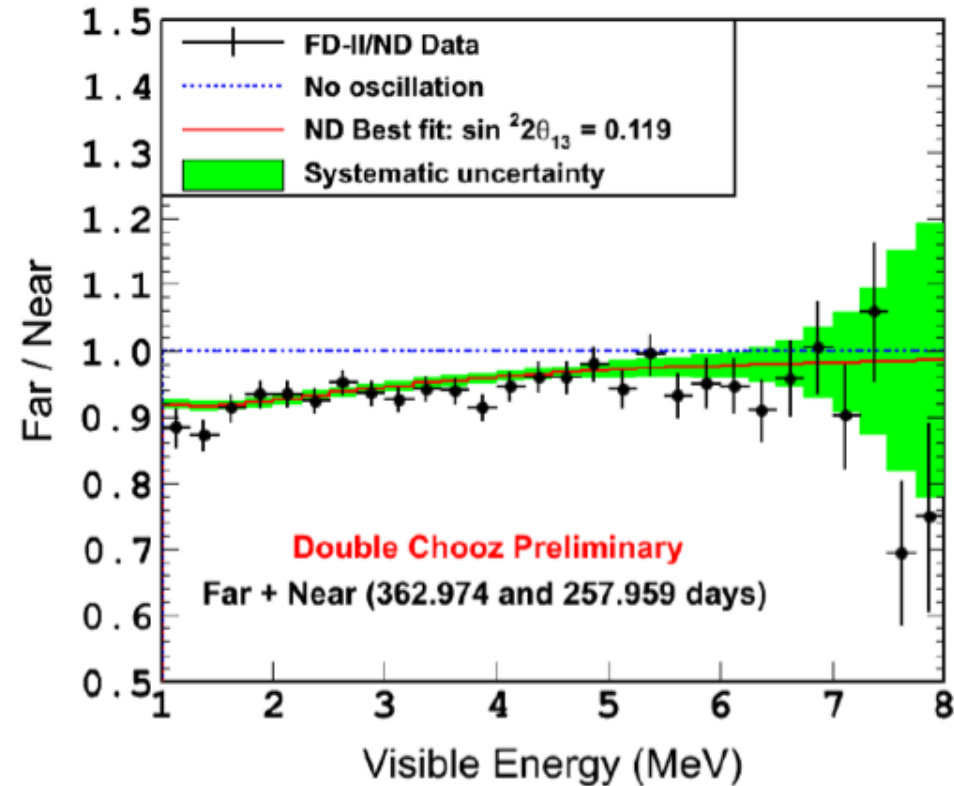
Detector performances are validated. Confirmed well tuned MC

Oscillation fit result

Simultaneous χ^2 fit with Data-to-MC comparison for each data set



FD-II/ND ratio



Common deviation is cancelled in FD-II/ND ratio

→ The deviation comes from flux prediction (under investigation)

$$\sin^2 2\theta_{13} = 0.119 \pm 0.016 \text{ (preliminary)}$$
$$\text{with } \chi^2/\text{ndf} = 236.2/114$$

Current θ_{13} in the world

Double Chooz
JHEP 1410, 086 (2014)

Preliminary
(CERN seminar 2016)

Daya Bay
PRL 115, 111802 (2015)

RENO
PRL 116 211801(2016)

T2K
PRD 91, 072010 (2015)

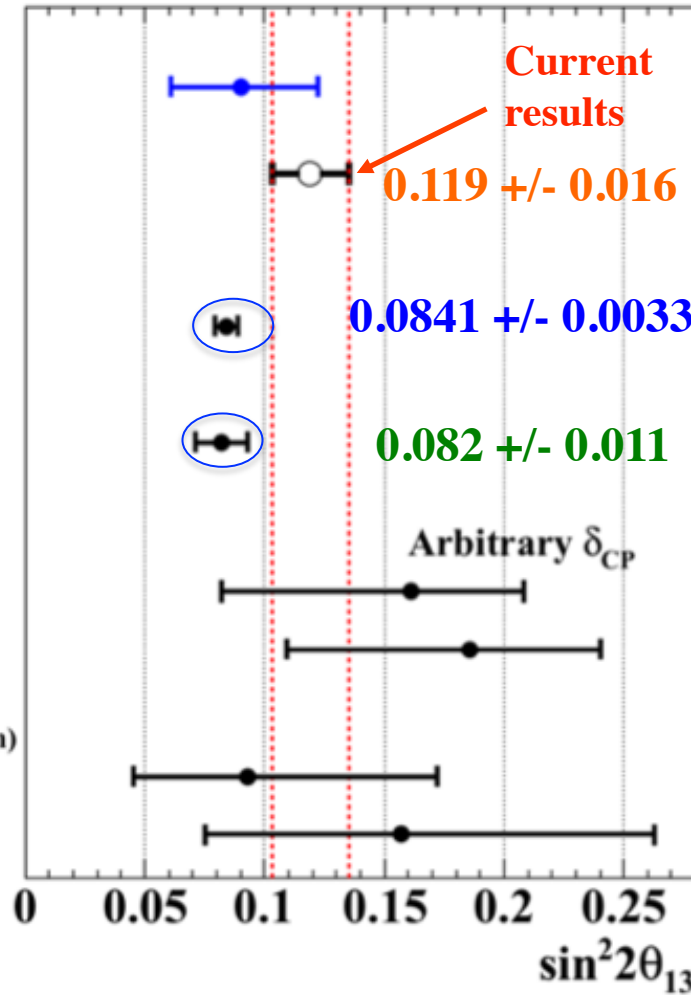
NOvA
Preliminary (private communication)

$\Delta m_{32}^2 > 0$

$\Delta m_{32}^2 < 0$

$\Delta m_{32}^2 > 0$

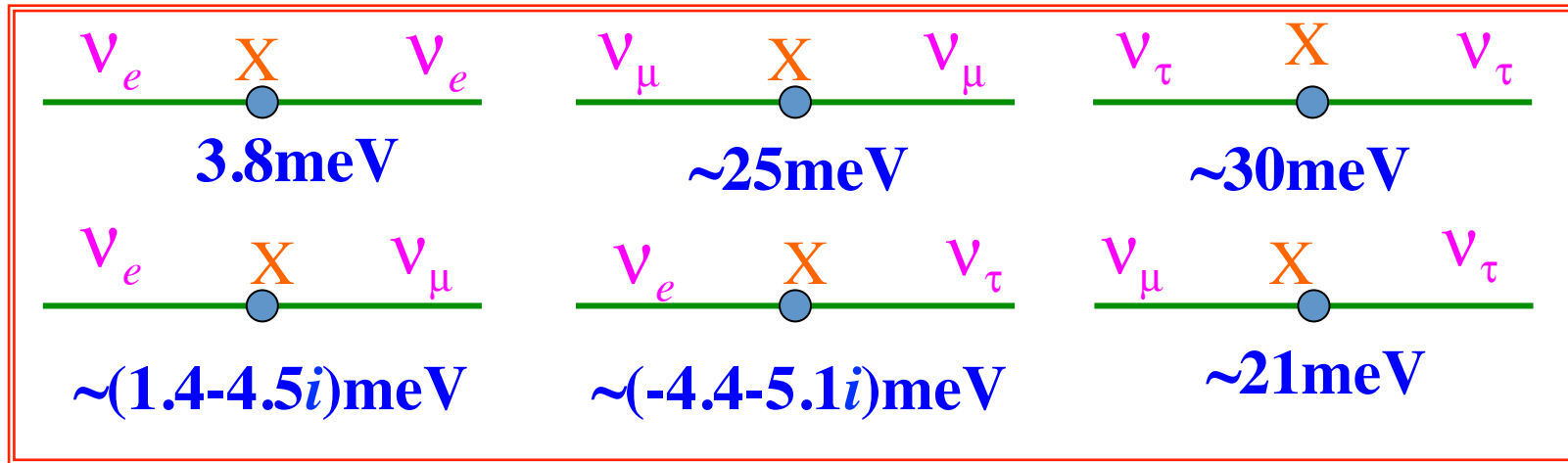
$\Delta m_{32}^2 < 0$



$\sim 2\sigma$ difference from
DayaBay/RENO

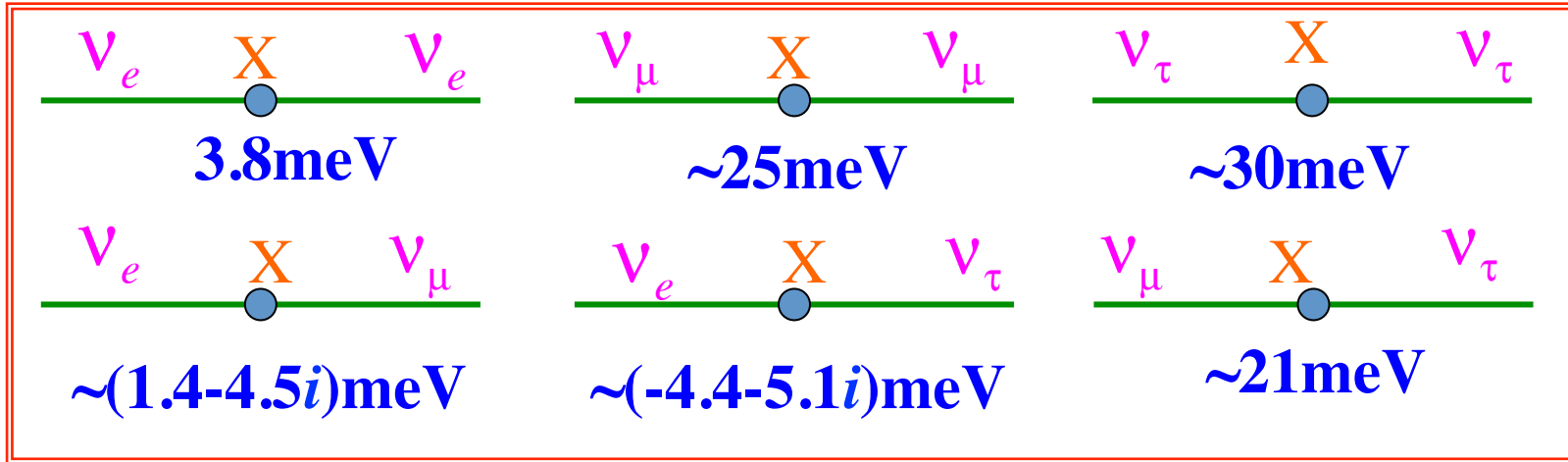
Neutrino transition amplitudes

Assumption: $m_3 > m_1 \sim 0$,

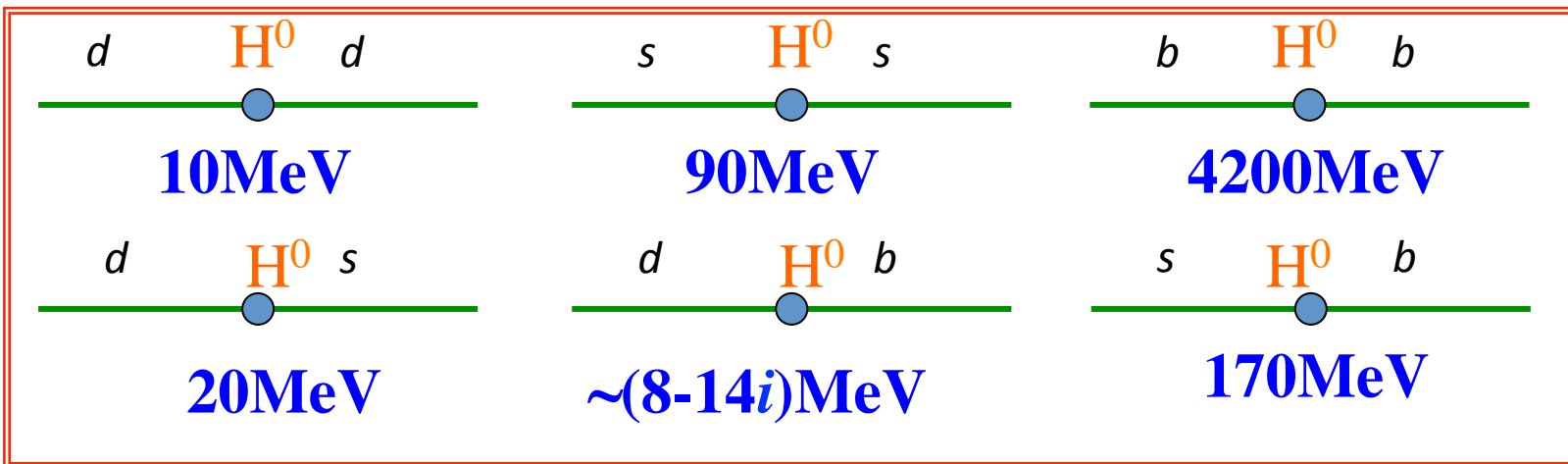


Neutrino transition amplitudes

Assumption: $m_3 > m_1 \sim 0$,

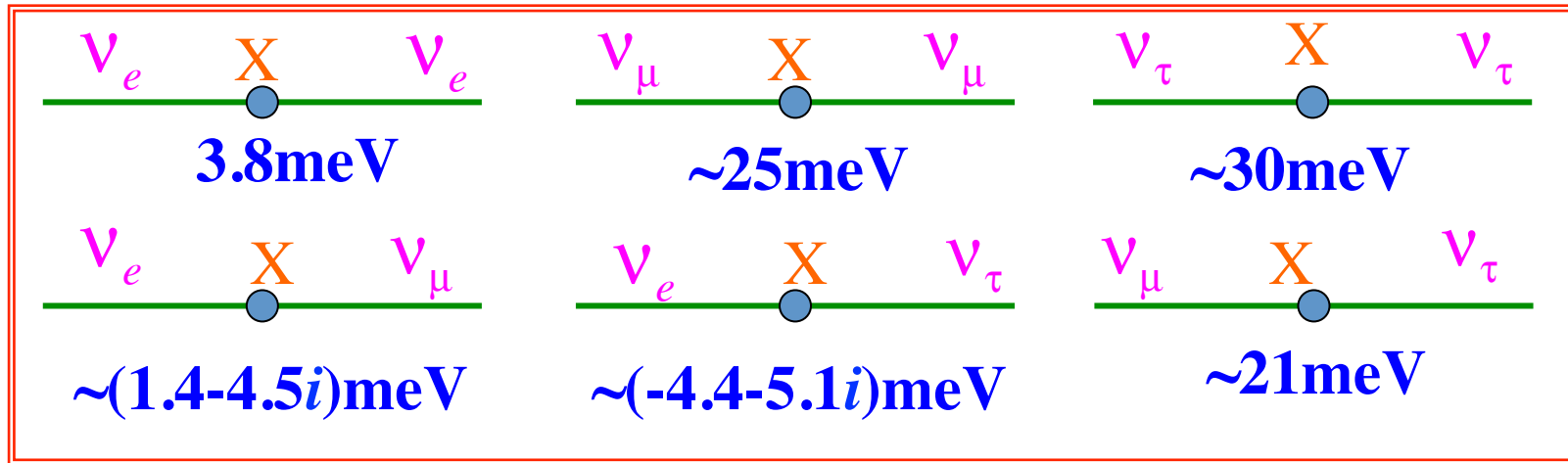


Comparison to the quark transition amplitudes



Neutrino transition amplitudes

Assumption: $m_3 > m_1 \sim 0$,



A is very small for **X**!! What is **X**??

θ_{13} and the CP Violation measurement

(matter effect ignored)

* CP asymmetry of N.O.:

$$A_{CP} = \frac{P[\nu_{\mu} \rightarrow \nu_e] - P[\bar{\nu}_{\mu} \rightarrow \bar{\nu}_e]}{P[\nu_{\mu} \rightarrow \nu_e] + P[\bar{\nu}_{\mu} \rightarrow \bar{\nu}_e]} \sim - \left| \frac{\Delta m_{12}^2}{\Delta m_{31}^2} \right| \frac{\pi \sin 2\theta_{12}}{\tan \theta_{23} \sin 2\theta_{13}} \sin \delta_{CP}$$

$$\sim - \frac{0.09}{\sin 2\theta_{13}} \sin \delta_{CP}$$

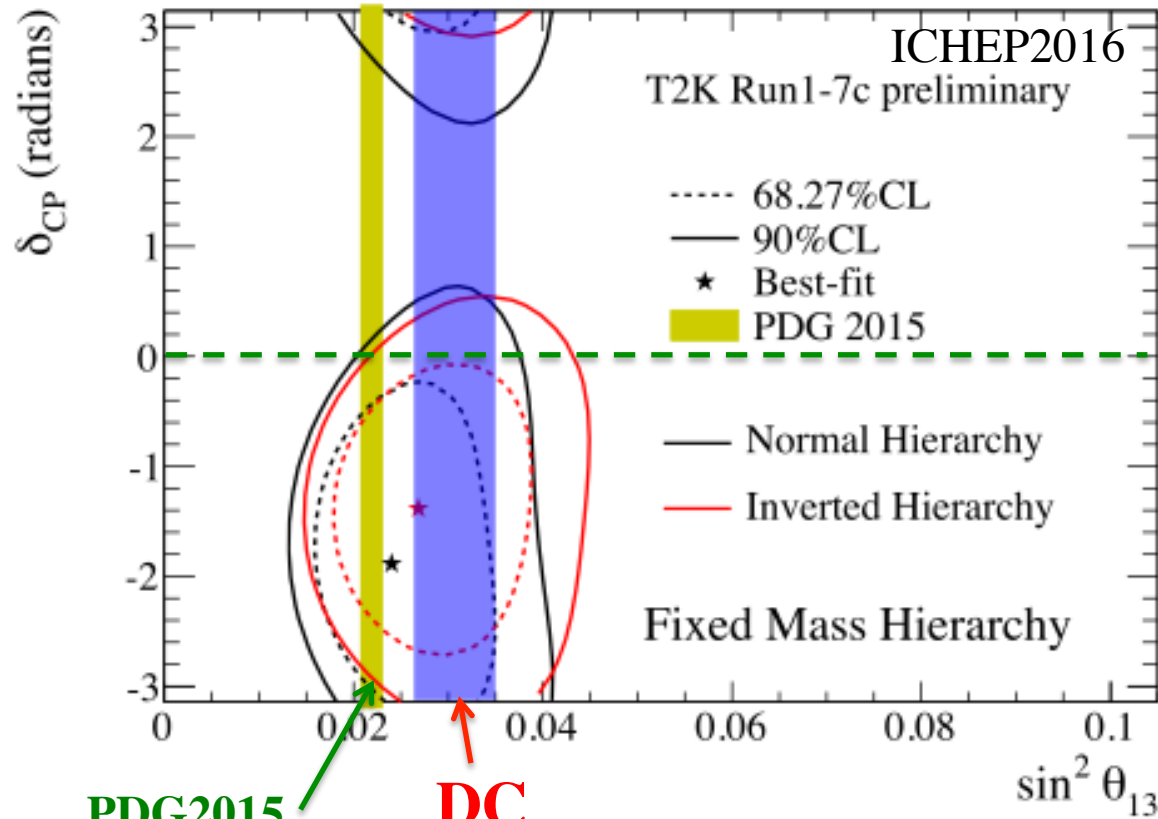
From measured A_{CP} , δ_{CP} is calculated as,

$$\sin \delta_{CP} \sim -11 \times \sin 2\theta_{13} \times A_{CP}$$

The error of δ_{CP} is,

$$\frac{\delta(\sin \delta_{CP})}{\sin \delta_{CP}} = \frac{\delta(\sin 2\theta_{13})}{\sin 2\theta_{13}} \oplus \frac{\delta A_{CP}}{A_{CP}} = \frac{1}{2} \frac{\delta(\sin^2 2\theta_{13})}{\sin^2 2\theta_{13}} \oplus \frac{\delta A_{CP}}{A_{CP}}$$

T2K result and Reactor θ_{13} .



← NO CPV

PDG2015 DC

- Broader range of δ_{CP} is allowed if large θ_{13}
- The discrepancy is not so significant now but further studies are necessary

Analysis Expert Workshop among DC, DB and RENO groups

The result of the current 3 reactor experiments will remain the most precise θ_{13} value in the foreseen future.

→ It is important to determine θ_{13} as precise and reliable as possible now.

The 1st analysis expert workshop among DC, DB and RENO groups was held in Korea in August 2016, initiated by the Double Chooz group.

The purpose of the workshop is to unify the 3 experiments at deep analysis level and obtain the single best value of θ_{13} .

The 2nd workshop is planned at APC in May 2017.

Summary

- * **Double Chooz has joined FJPPL since 2006.**
 - **In the 1st phase, it successfully reported the first reactor θ_{13} result in 2011. (Nu2-WP2)**
 - **In the 2nd phase, it is successfully reporting the precise θ_{13} result using 2 detectors. (Nu03)**
- * **The measured θ_{13} value at DC is, $\sin^2 2\theta_{13} = 0.119 \pm 0.016$.**
- * **There is some difference between DC and DB/RENO results.**
- * **Started discussion among the DC/DC/RENO analysis experts.**
- * **The precision of the θ_{13} measurement will improve and we can see the tension develops or shrinks.**