

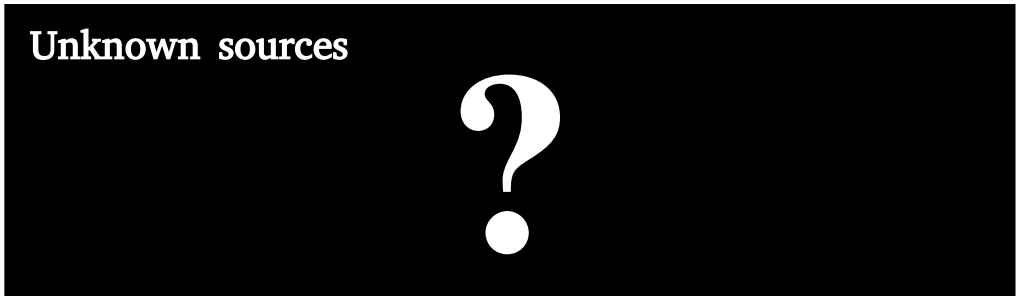
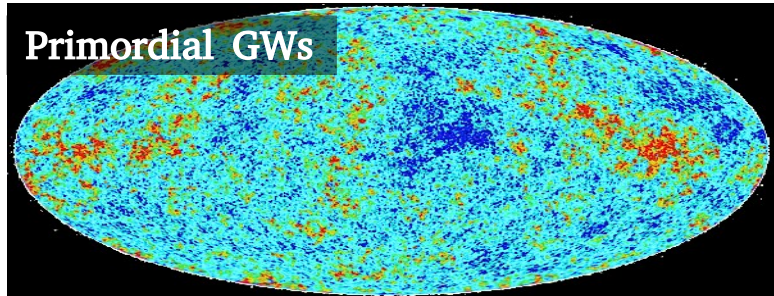
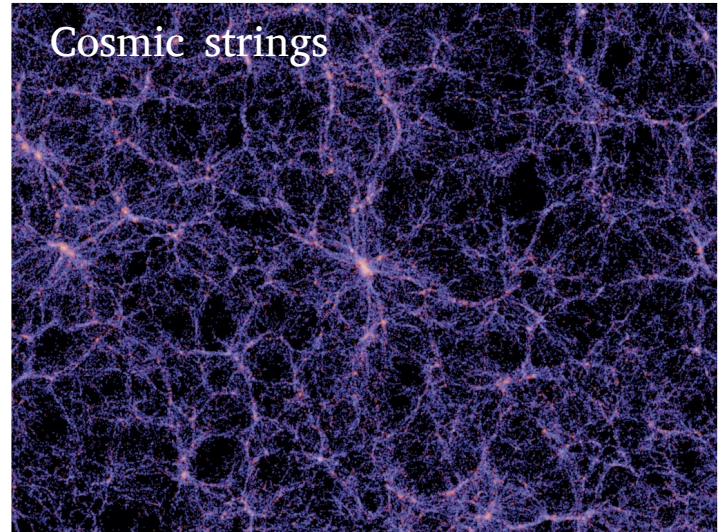
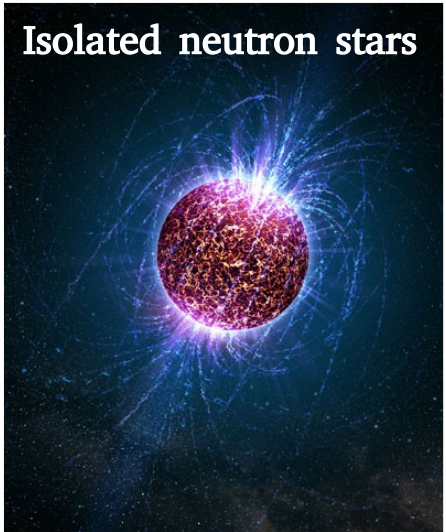
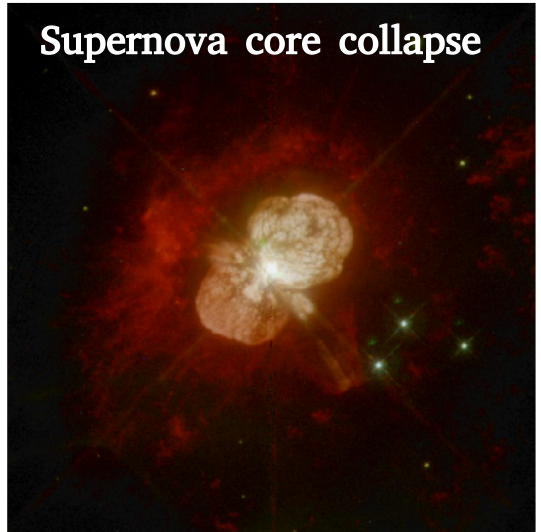
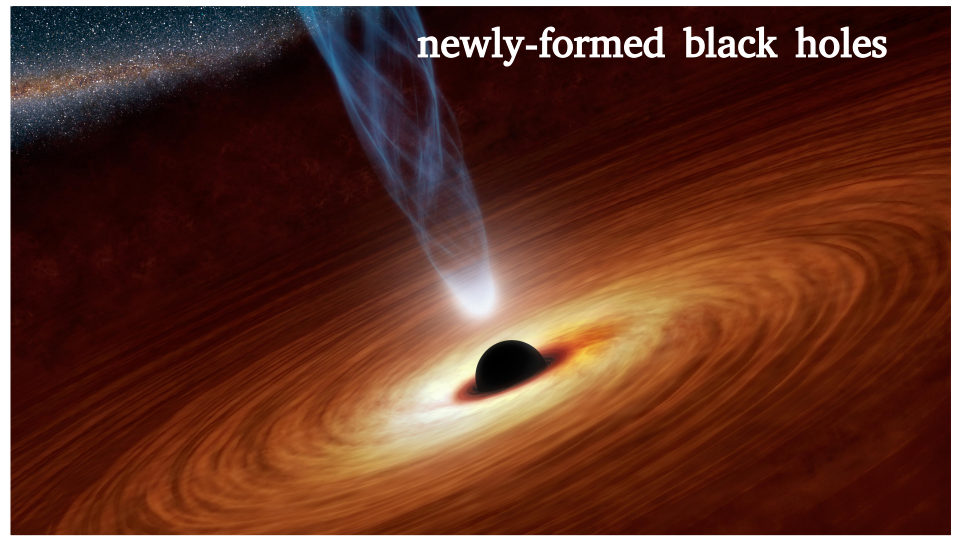
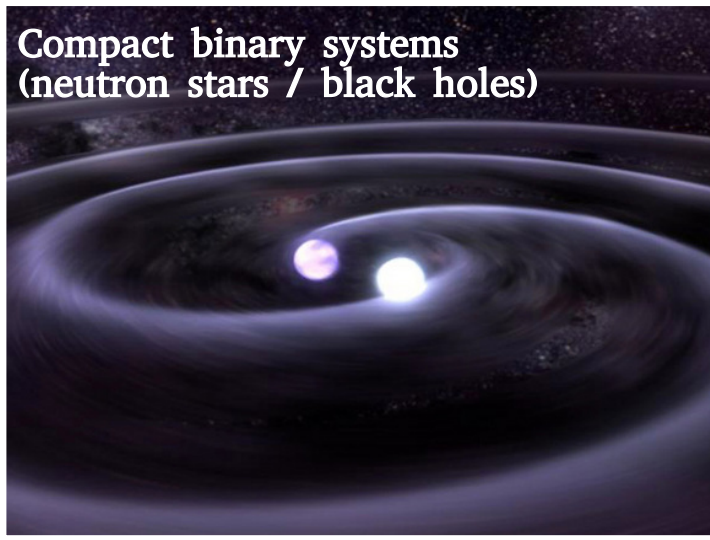


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Searches for gravitational waves with ground-based interferometers

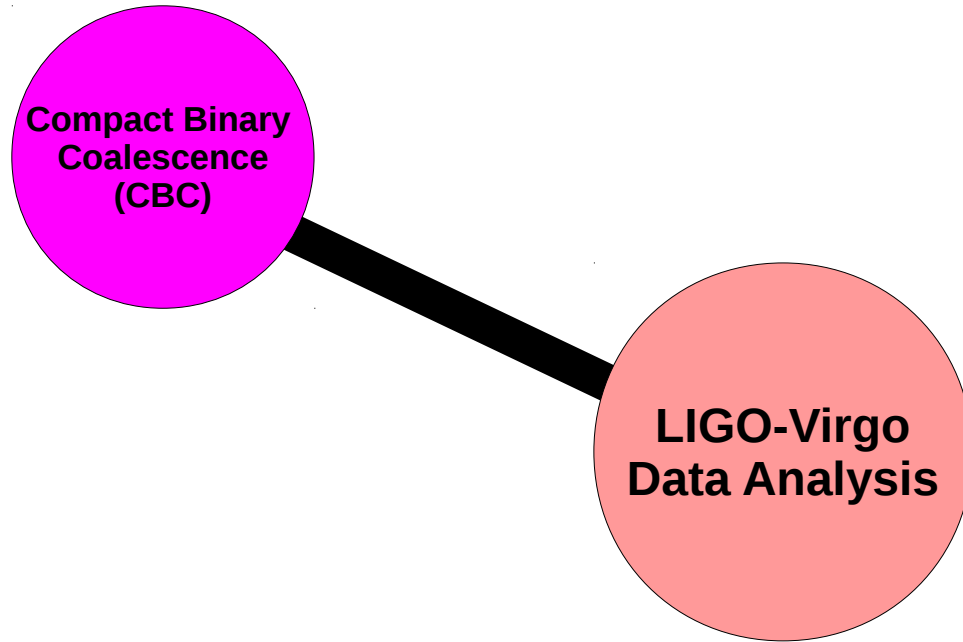
Searches for gravitational waves with ground-based interferometers

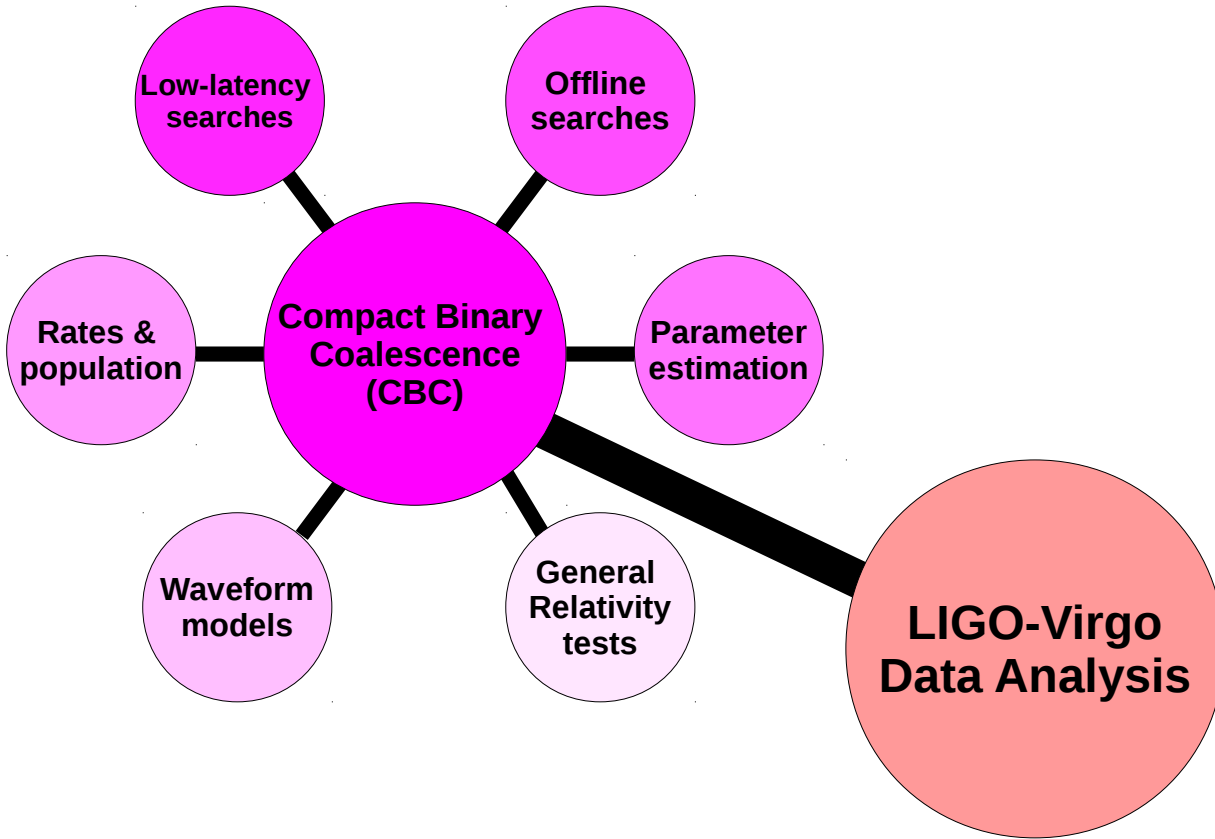
- I. LIGO-Virgo analysis groups
- II. Analysis methods for transient searches
- III. Recent gravitational-wave detections
- IV. Other search results
- V. O2 and conclusions

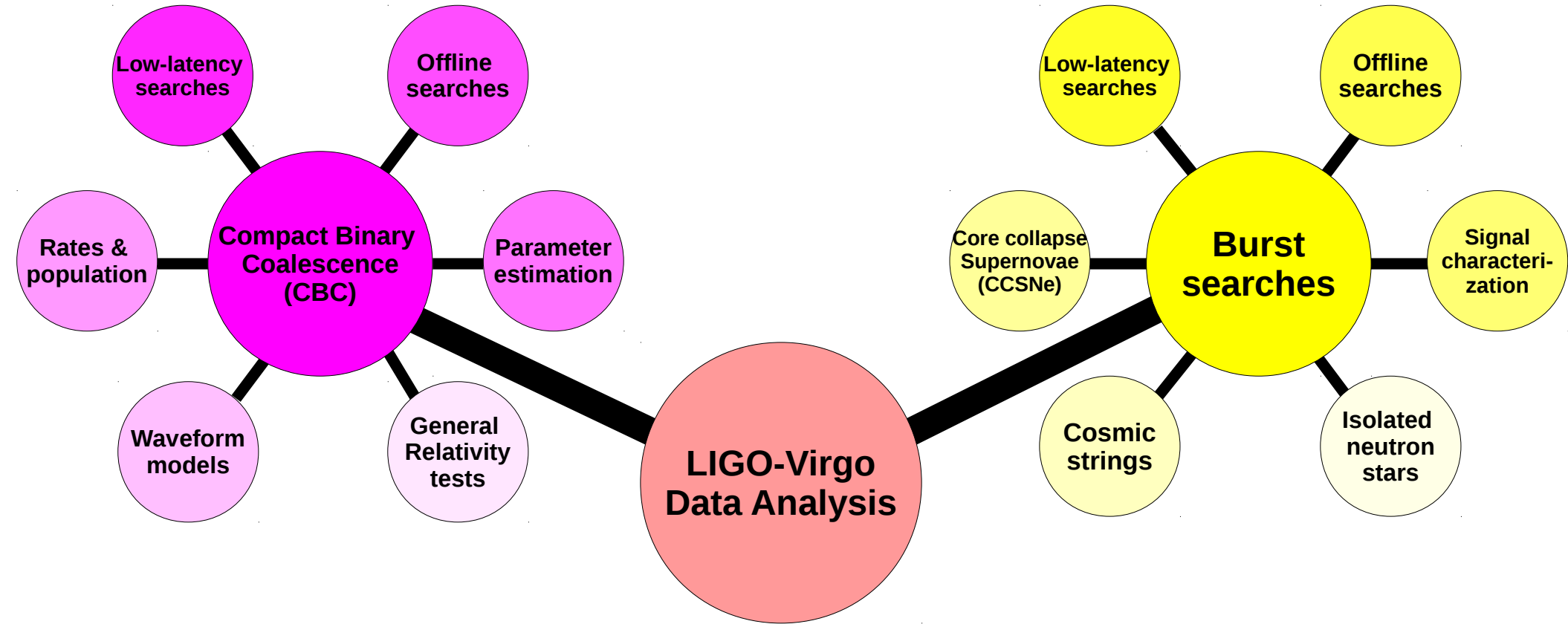


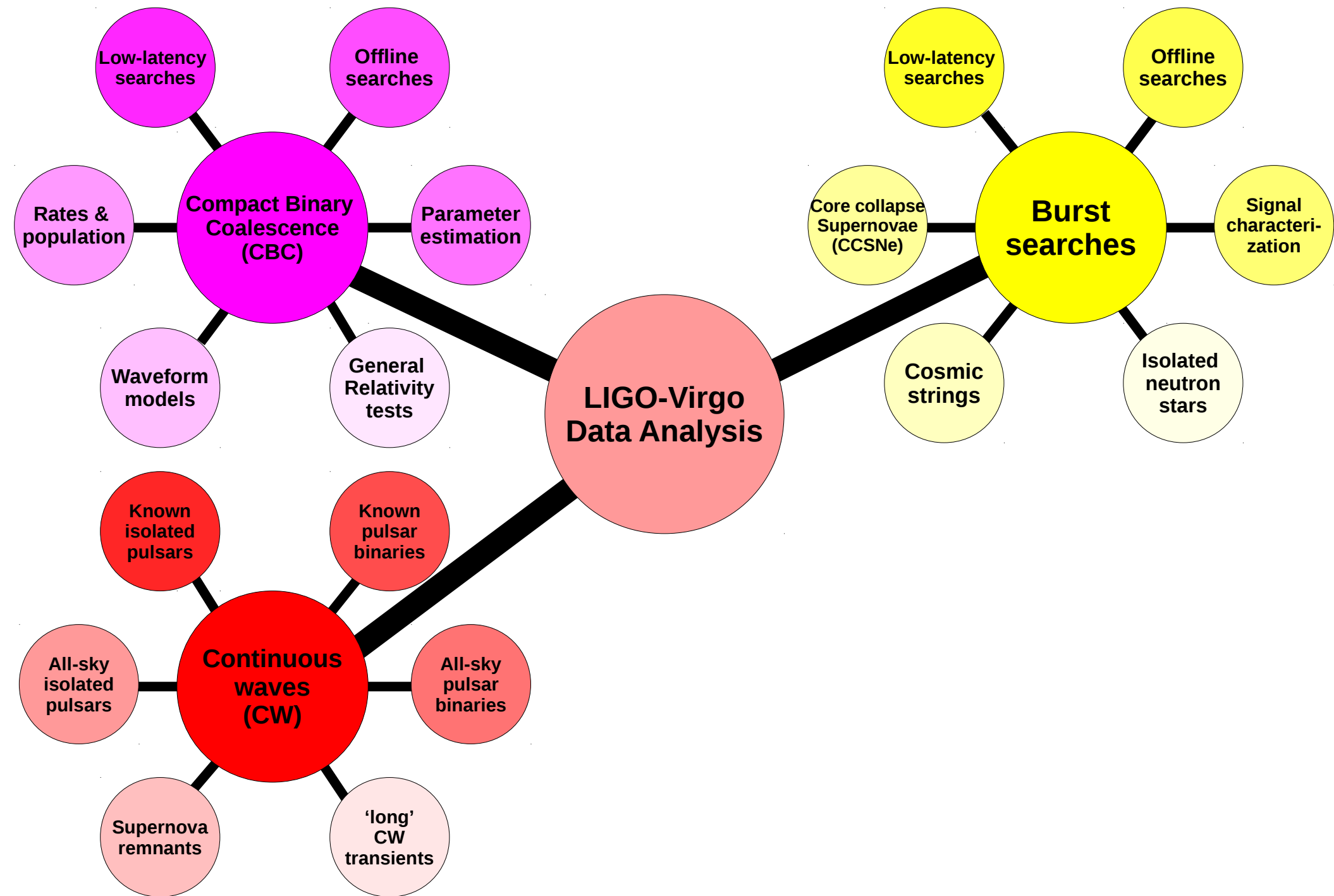


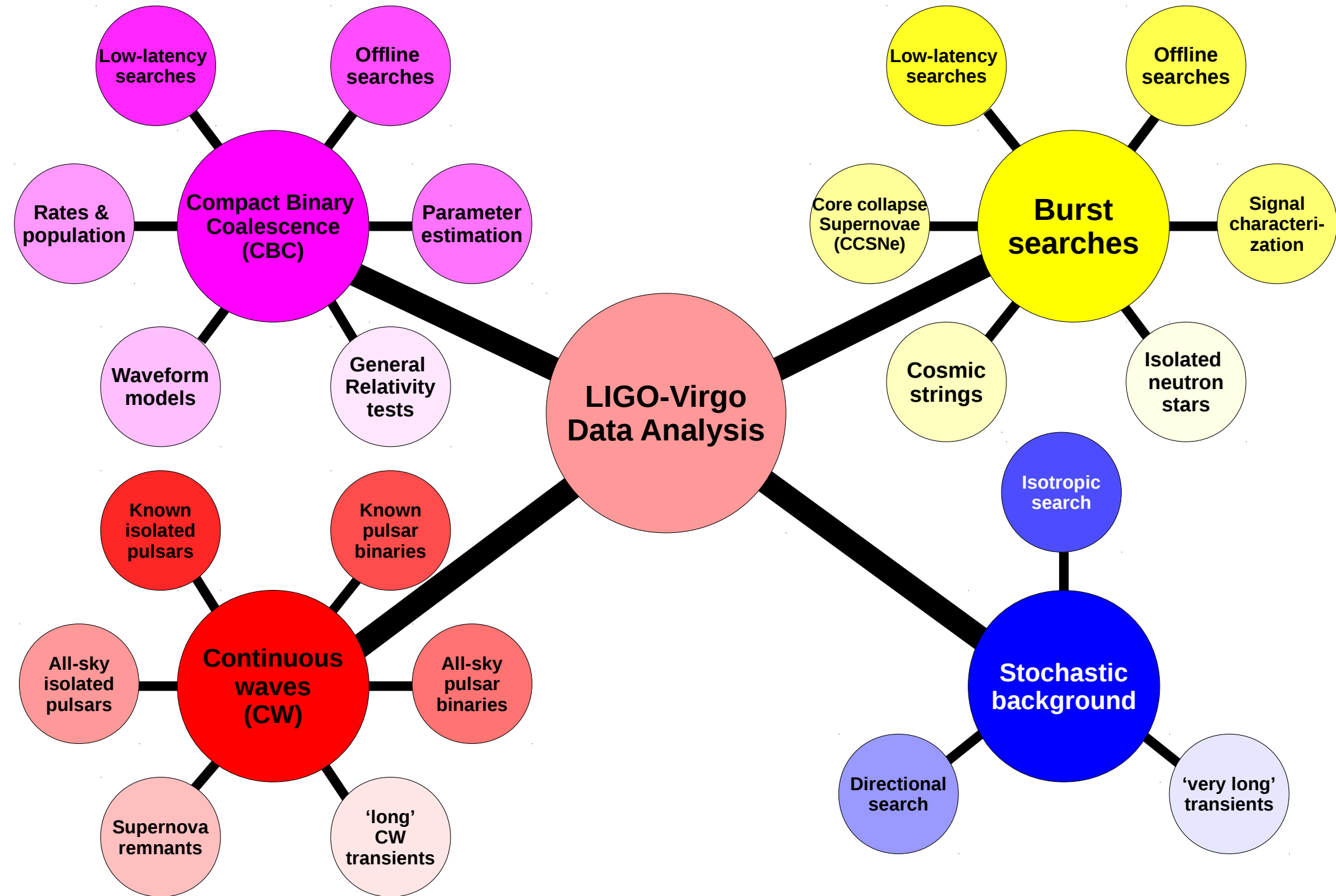
**LIGO-Virgo
Data Analysis**

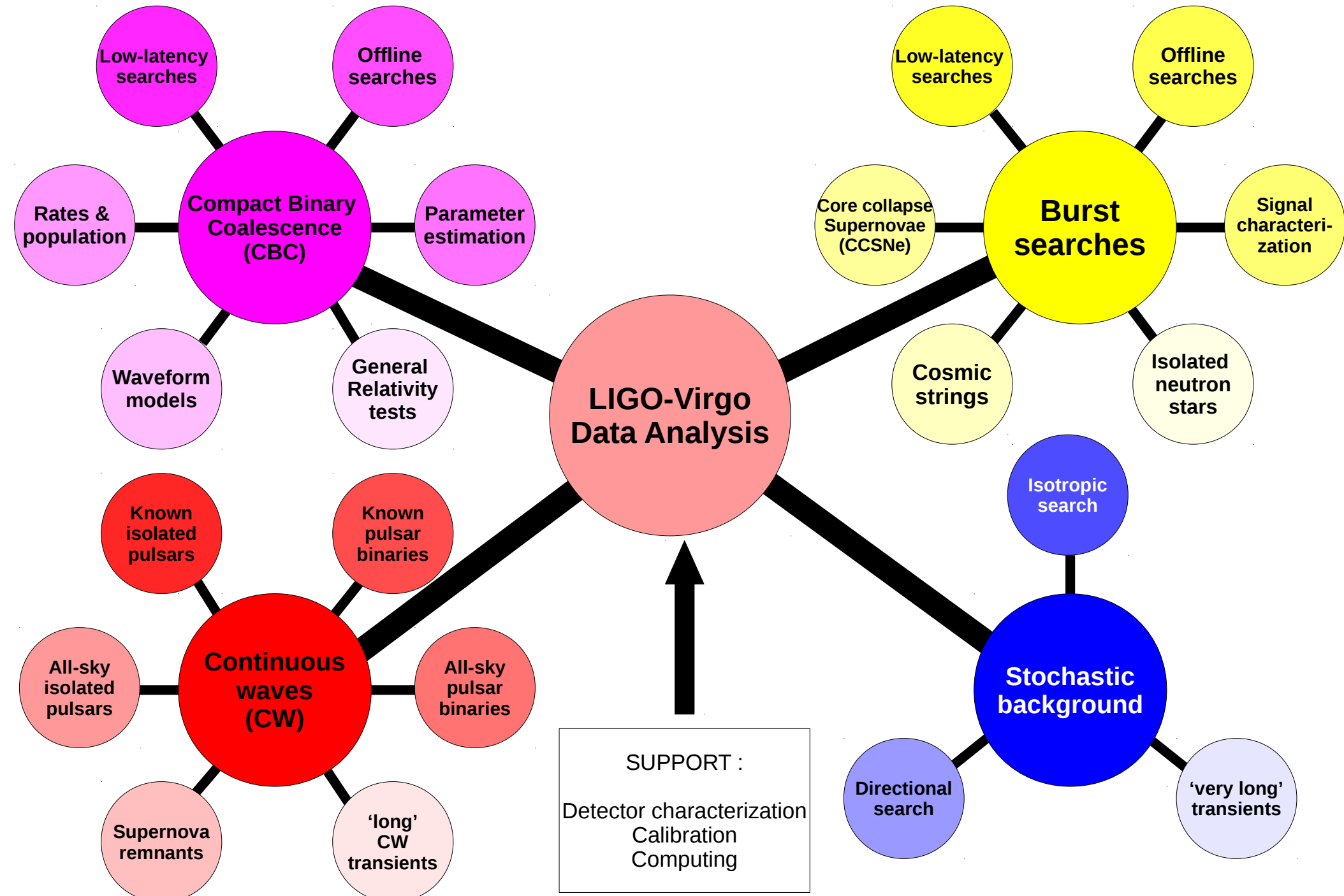


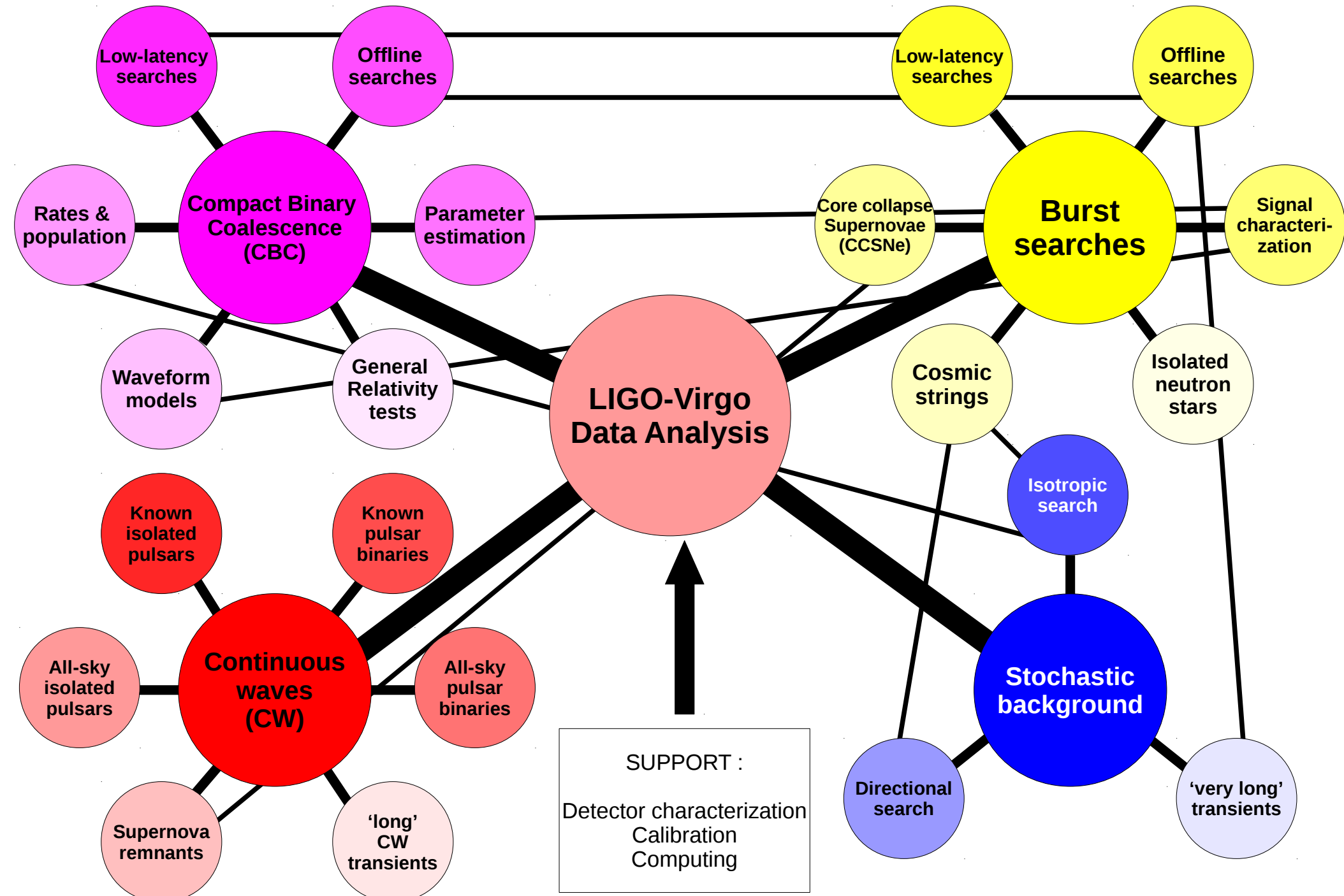


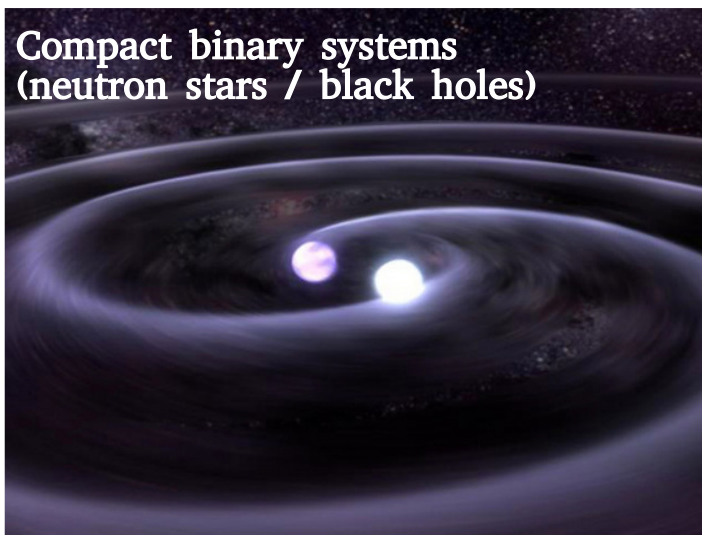




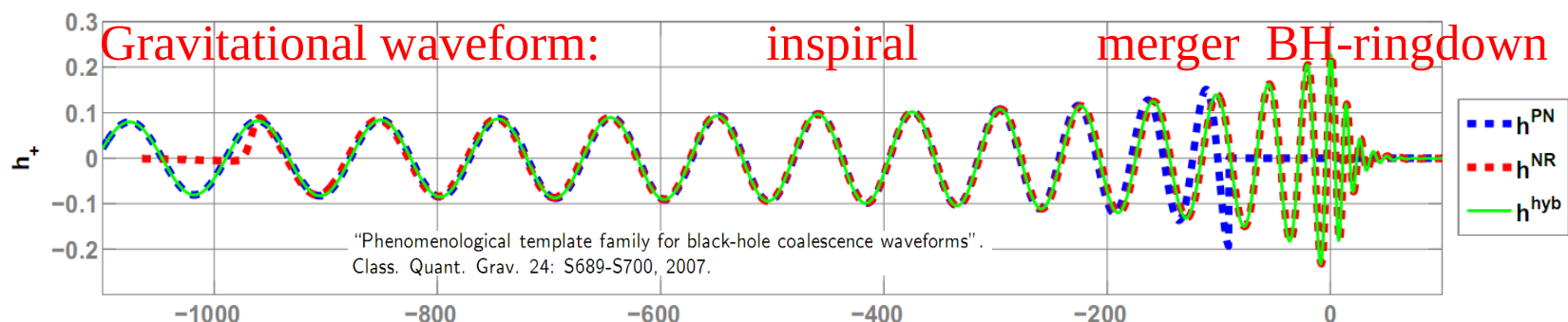




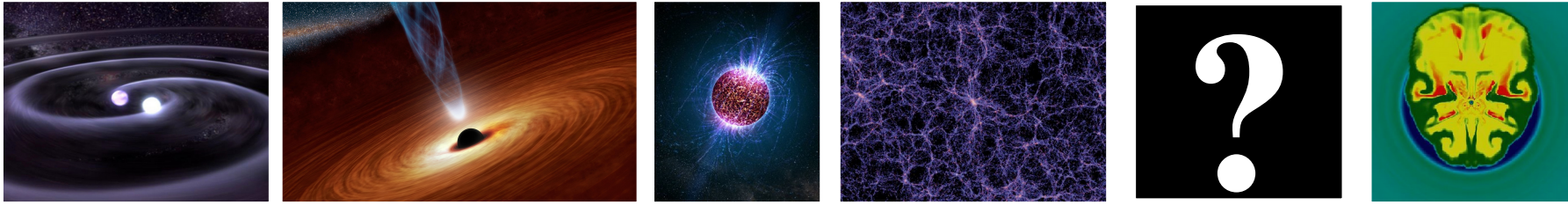




- Compact binary objects: Two neutron stars and/or black holes.
- Inspiral toward each other. Emit gravitational waves as they inspiral.
- Amplitude and frequency of the waves increases over time, until the merger.
- Waveform relatively well understood, → matched template searches.
- Unique way to study string field gravity and the structure of the nuclear matter in the most extreme conditions

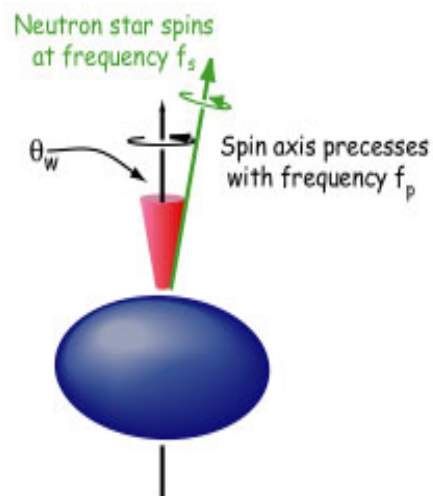
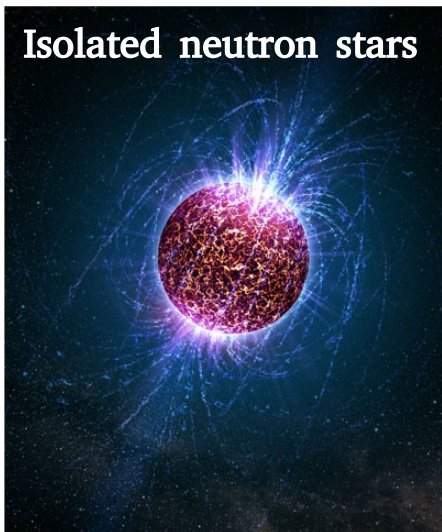


Waveform carries lots of information about binary masses, orbit, merger, spins, ...



- Many transient sources:
 - CBC
 - Supernovae: probe the explosion mechanisms.
 - Gamma Ray Bursts: collapse of rapidly rotating massive stars or neutron star mergers.
 - Pulsar glitches.
 - Cosmic strings cusps and kinks.
- Models are ok, but not essential:
 - Search for power excess in the data.
 - Search for any short signal with measurable strain signal.

Isolated neutron stars



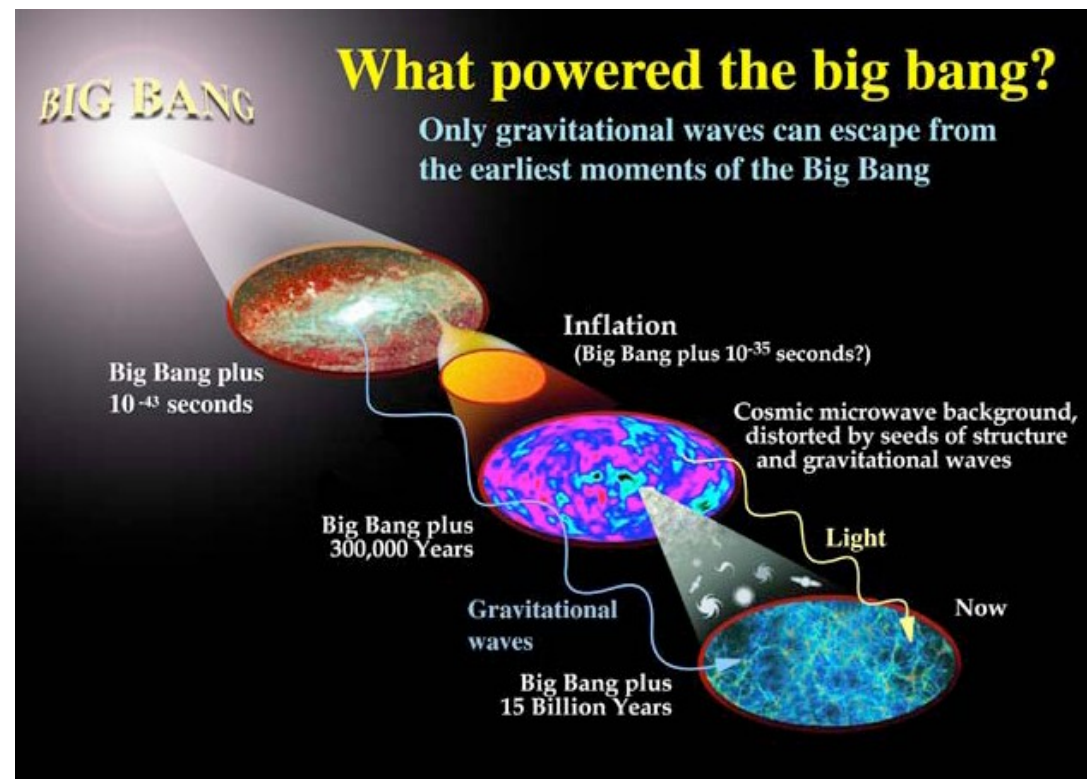
Persistent signals associated to sources with mass quadrupole moment varying in time in a nearly periodic way

- Pulsars with mass non-uniformity:
 - distortion due to elastic stresses or magnetic field
 - distortion due to matter accretion
 - free precession around rotation axis
 - excitation of long-lasting oscillations (e.g. r-modes)
- Produce gravitational-waves, often at twice the rotational frequency.
- Waveform well-understood:
 - Sinusoidal but Doppler modulated
- Continuous source

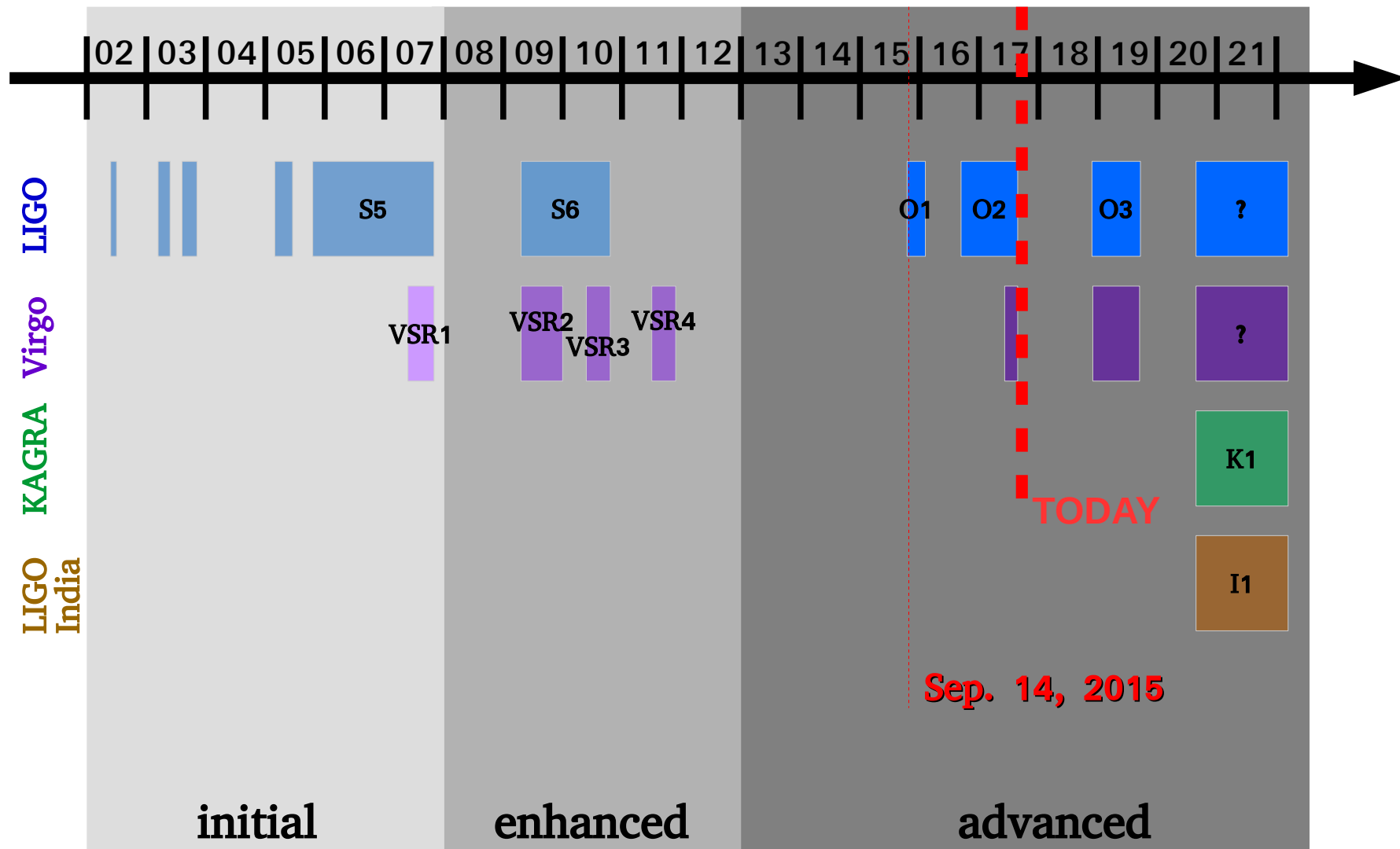
Signal amplitude:
$$h_0 \cong 10^{-27} \left(\frac{I_{zz}}{10^{38} \text{ kg} \cdot \text{m}^2} \right) \left(\frac{10 \text{ kpc}}{r} \right) \left(\frac{f}{100 \text{ Hz}} \right)^2 \left(\frac{\varepsilon}{10^{-6}} \right)$$

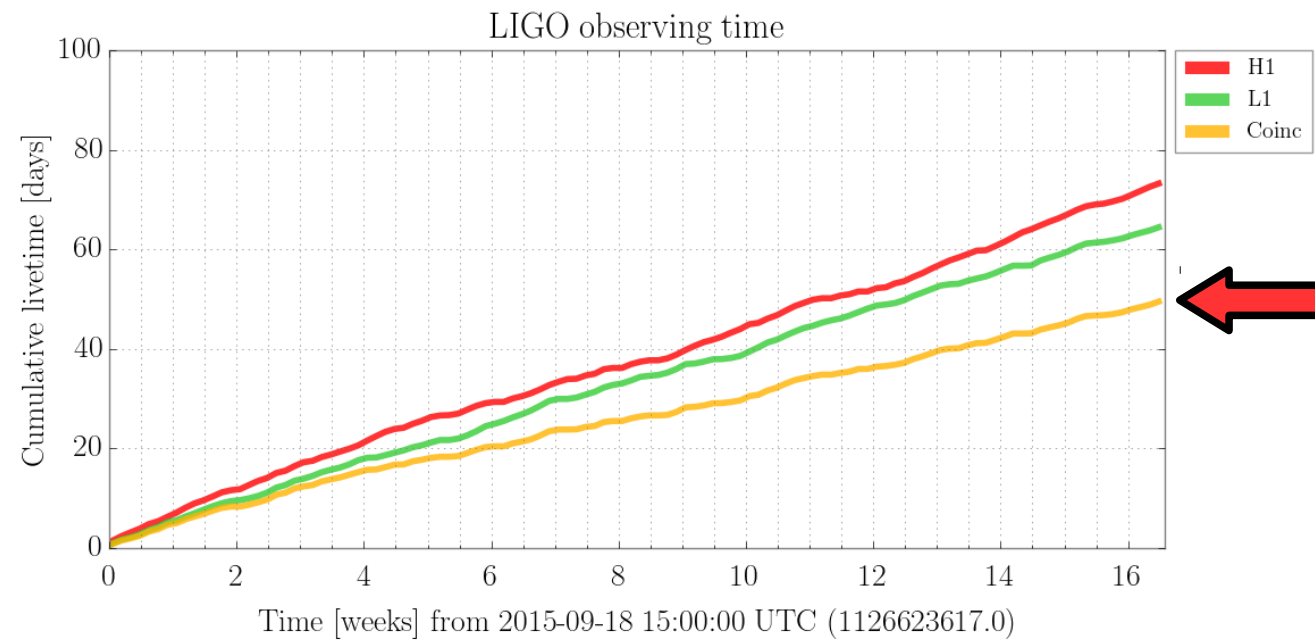
ε : ellipticity (adimensional number measuring the star's degree of asymmetry)
 f : signal frequency, proportional to star rotation frequency

- Incoherent superposition of many unresolved sources.
- Cosmological:
 - Inflationary epoch, preheating, reheating
 - Phase transitions
 - Cosmic strings
 - Alternative cosmologies
- Astrophysical:
 - Supernovae
 - Magnetars
 - Binary black holes
- Potentially could probe physics of the very-early Universe.

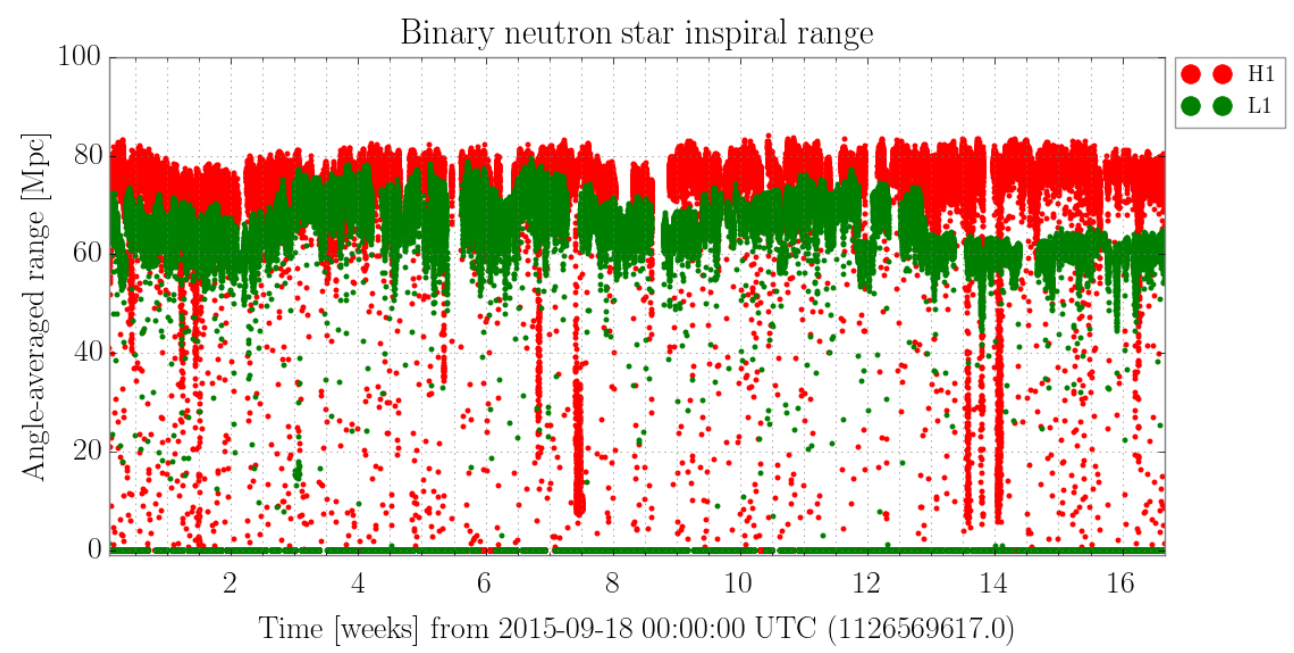


$$\Omega_{GW}(f) = \frac{f}{\rho_c} \frac{d\rho_{GW}}{df}$$





← ~49 days of double coincidence



LIGO-Hanford (H1)

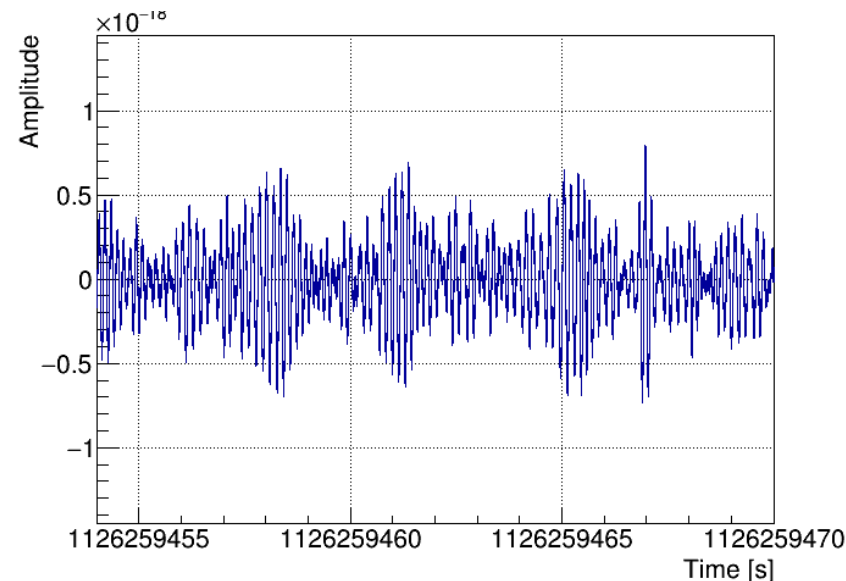


LIGO-Livingston (L1)

GW detectors' readout system provides at any instant an estimate of strain: a quantity that is sensitive to arms' length difference:

→ Digitized discrete time series: $raw(t)$ (sampled at 16384 Hz or 20000 Hz) and synchronized with GPS clocks.

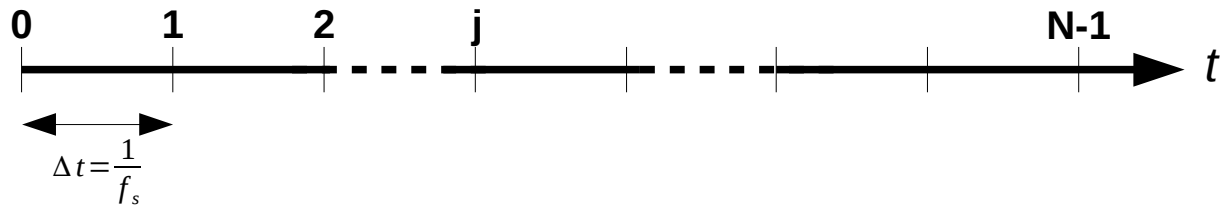
→ Calibration of $raw(t)$: apply a frequency dependent factor [in reality this is a bit more complicated ...]



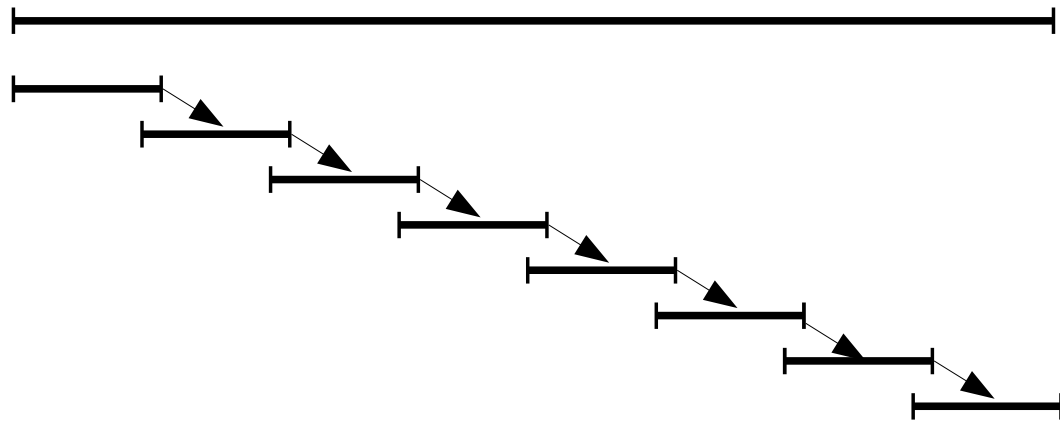
→ $h_{det}(t)$ time series that is detector noise plus all hypothetical GW signals

$$h_{det}(t) = n(t) + GW(t)$$

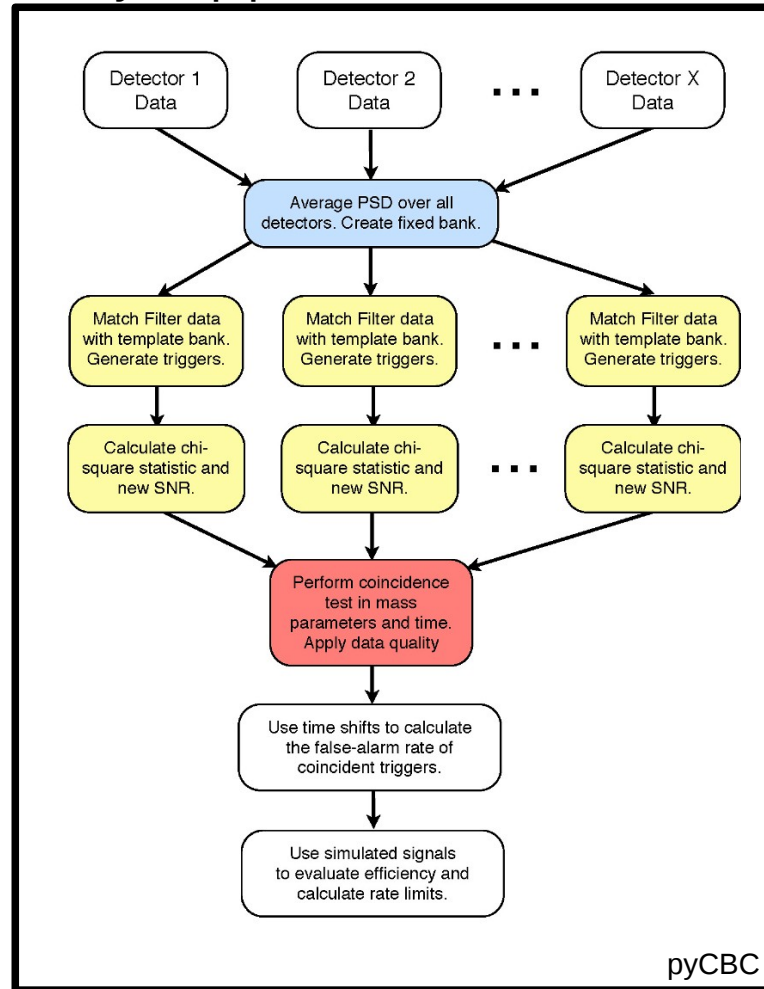
$$h_{det}(t) \rightarrow h_{det}[j]$$



GW search: load the data iteratively (analysis window)



Analysis pipeline



Fourier transform

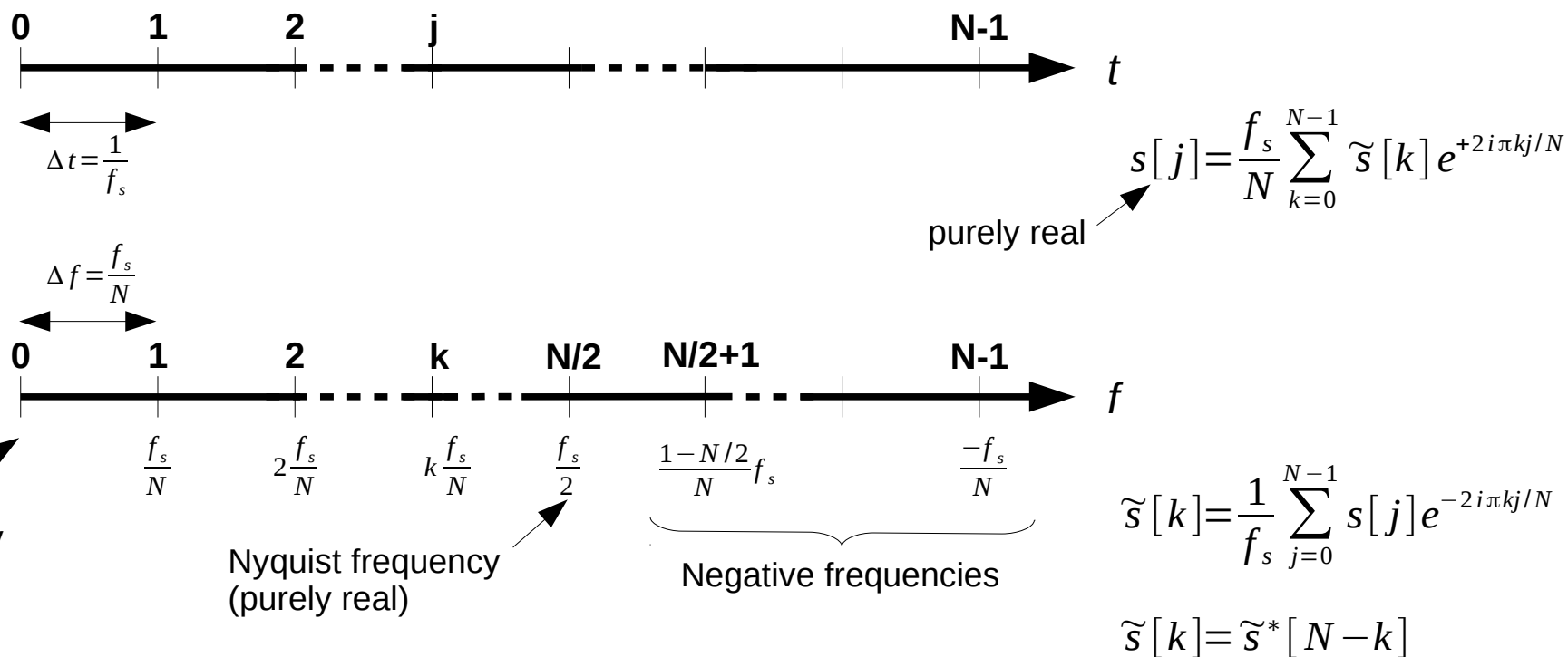
A time series $s(t)$ can be projected over a basis of sinusoidal functions:

$$\tilde{s}(f) = \int_{-\infty}^{\infty} s(t) e^{-2i\pi ft} dt \quad (\text{forward})$$

$$s(t) = \int_{-\infty}^{\infty} \tilde{s}(f) e^{2i\pi ft} df \quad (\text{backward})$$

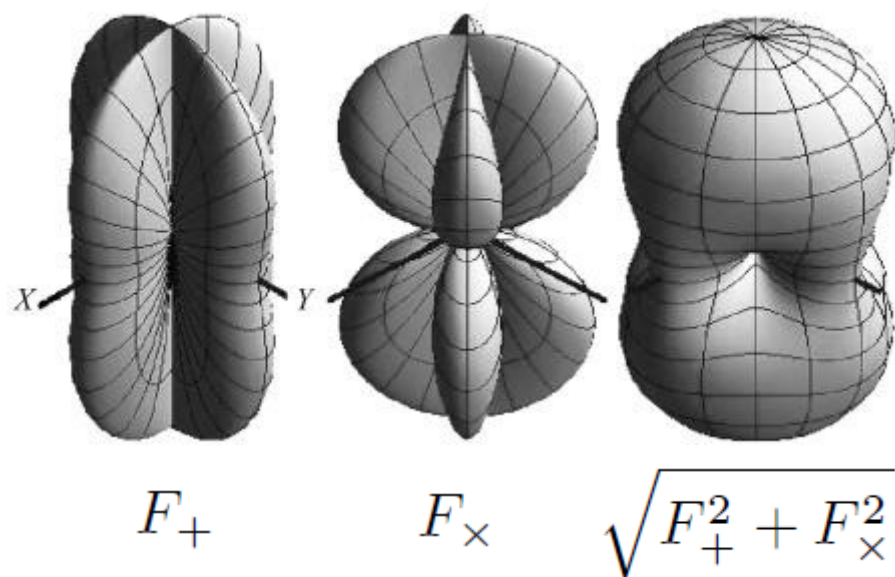
The signal is decomposed in characteristic frequencies

Discrete Fourier transform



$$h_{det}(t) = n(t) + \mathbf{GW}(t)$$

The detector's sensitivity over the sky is not uniform



$$GW(t) = F_+(t, \underline{ra}, \underline{dec}, \underline{\Psi}) \times h_+(t) + F_x(t, \underline{ra}, \underline{dec}, \underline{\Psi}) \times h_x(t)$$

Source position Source polarization angle

Power spectral density:
(PSD) $\lim_{T \rightarrow \infty} \frac{1}{T} |\tilde{x}_T(f)|^2$

Power spectral density estimator for finite data set: Periodogram = $\frac{1}{T} |\tilde{x}_T(f)|^2$

Improved estimator:

- average multiple periodograms (M) to reduce the variance
- noise is non-stationary: T should not be too long (a few minutes)
- use windowed data to limit spectral leakage
- Welch approach: average of periodograms computed over overlapping windowed data segments

Sensitivity measured using the noise power spectral density :

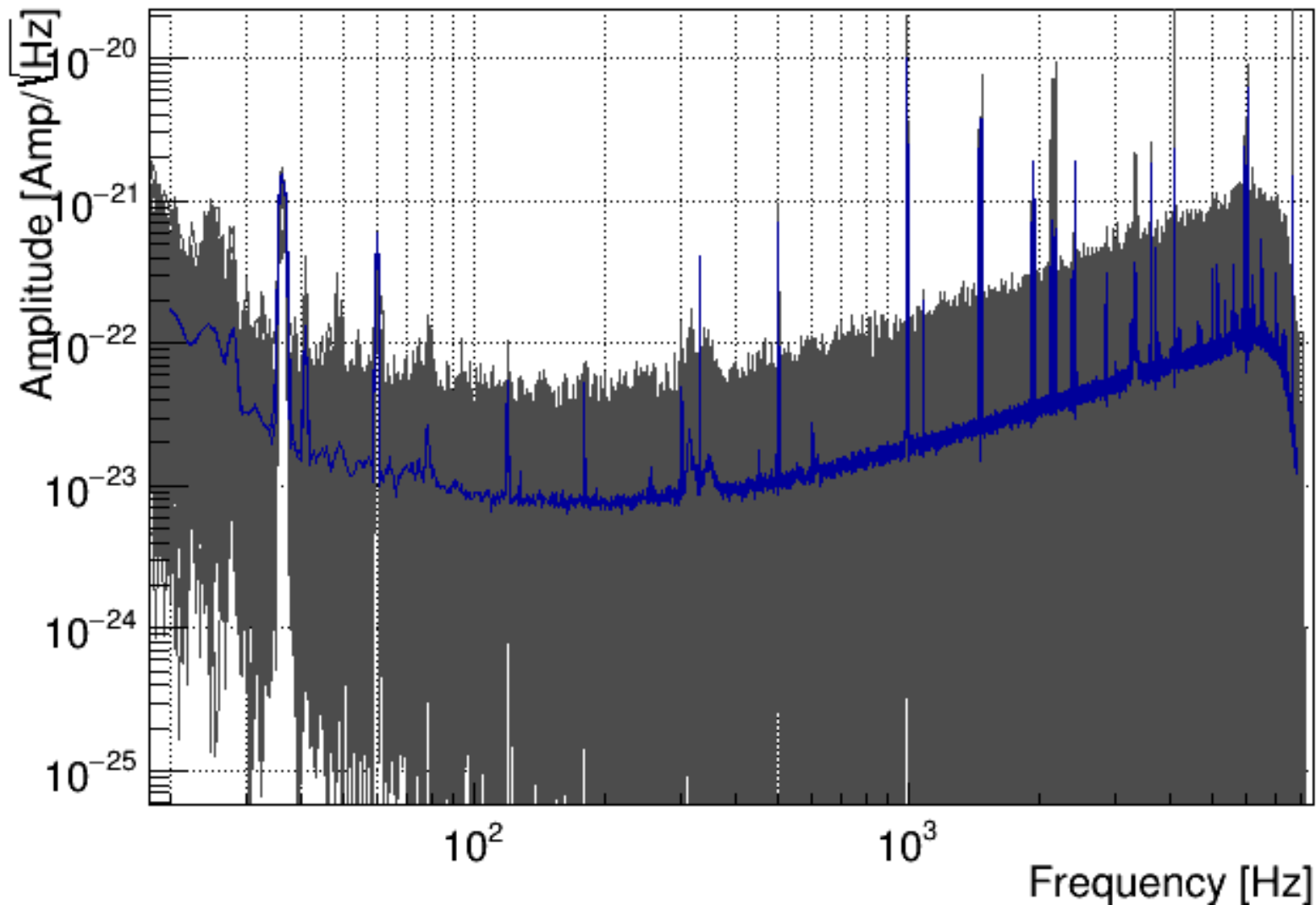
$$S_n(k) = \text{Median}_{0 \leq m < M} \left\{ \frac{1}{N f_s} \left| \sum_{j=0}^{N-1} x_m[j] w[j] e^{-2i\pi jk/N} \right|^2 \right\}$$

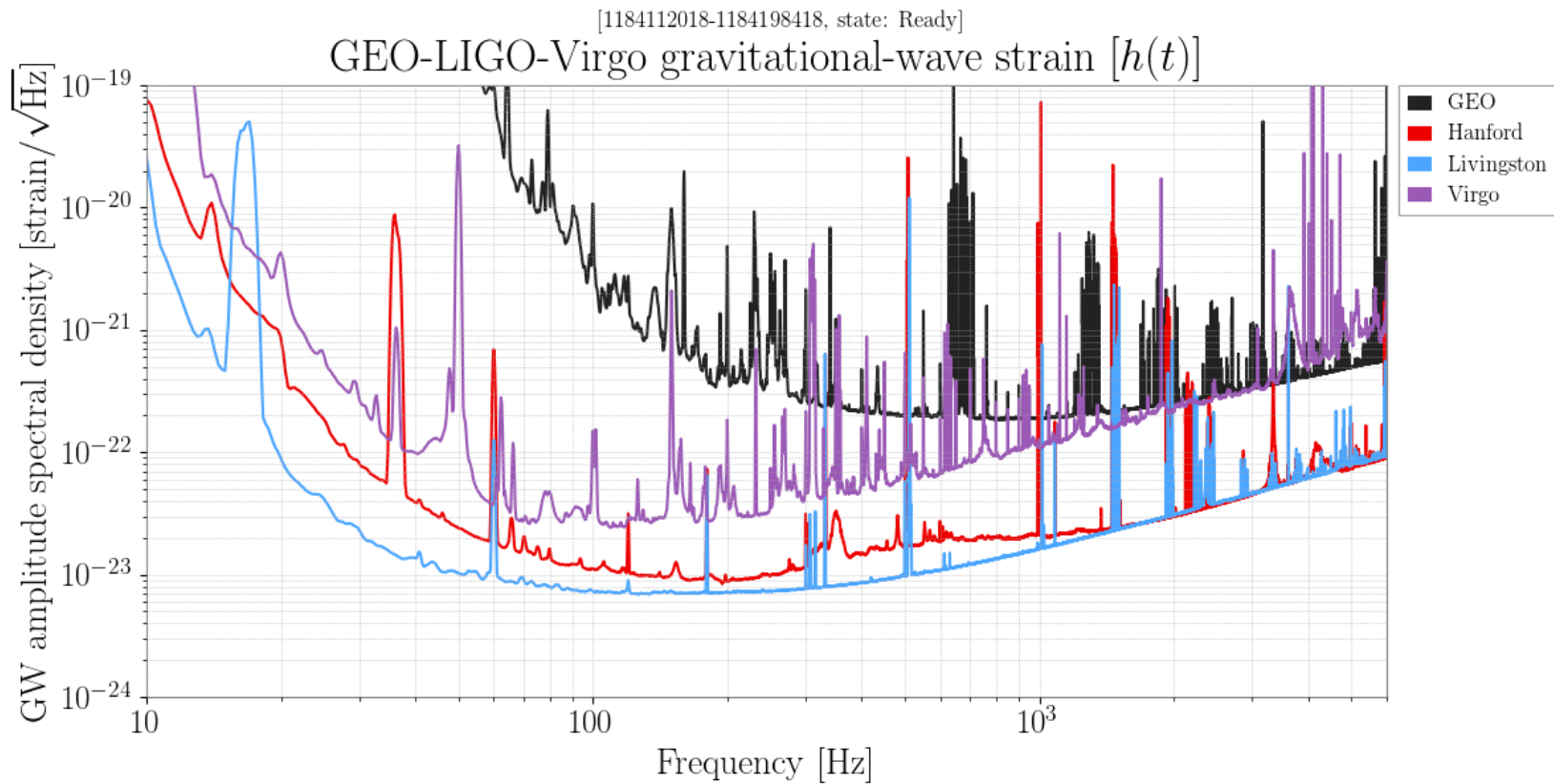
+ median-to-mean correction

One-sided / Two sided PSDs

Amplitude power spectral density: $\sqrt{S_n(k)}$

Noise amplitude spectral density of $h_{det}(t) = n(t)$

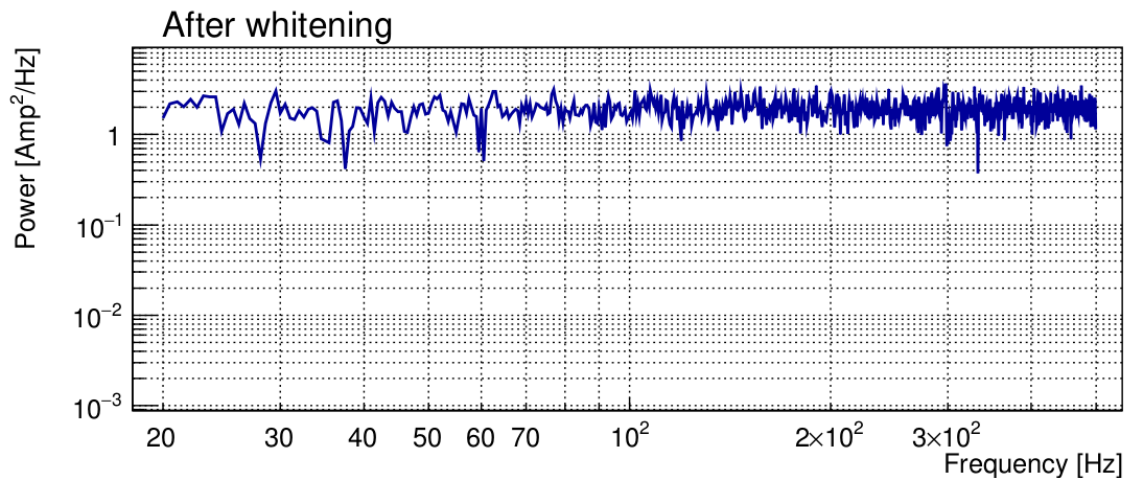
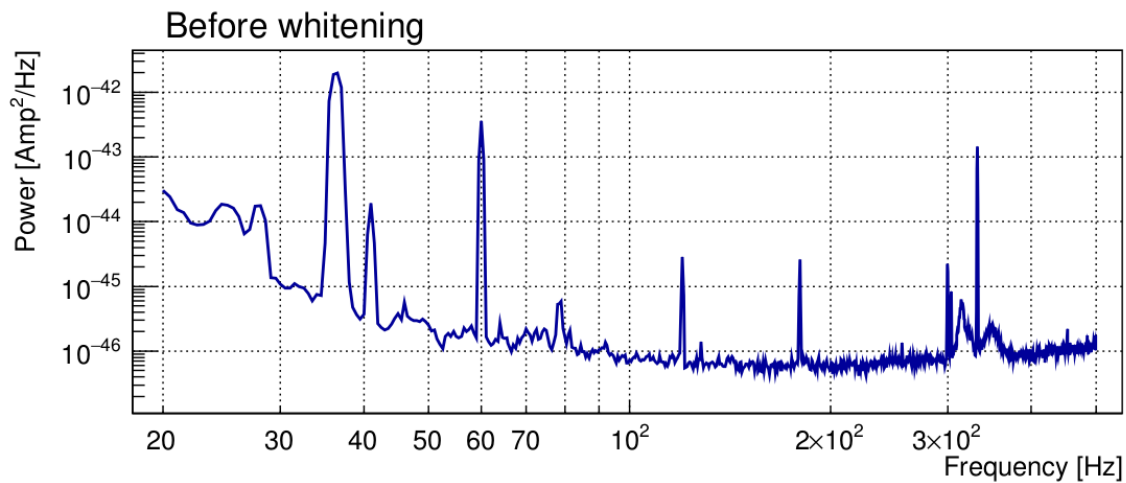




O2 Best noise spectrum achieved by **LIGO Hanford**, **LIGO Livingston** and **Virgo**

GW data is whitened: $\tilde{h}_{det}(f) \rightarrow \tilde{h}_{det}^w(f) = \frac{\tilde{h}_{det}(f)}{S_n(f)}$

→ white noise is mandatory for statistical interpretation of the data



Example : Q-transform

$$X(\tau, \phi, Q) = \int_{-\infty}^{+\infty} h_{det}(t) w(t - \tau, \phi, Q) e^{-2i\pi\phi\tau} dt$$

→ window width $\sim 1/\phi$

~ short Fourier transform with a Gaussian window

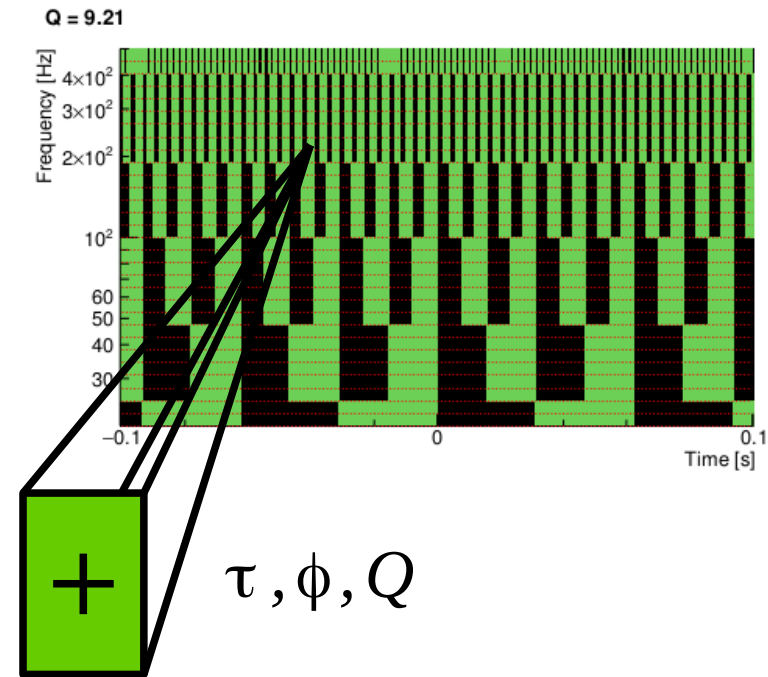
→ Goal : cover a parameter space as large as possible

$$\text{Noise only: } \langle |N(\tau, \phi, Q)|^2 \rangle = \int_{-\infty}^{+\infty} |\tilde{w}(\phi - f, \phi, Q)|^2 S_n(f) df$$

$$\text{Whitened noise + window normalization: } \langle |N^w(\tau, \phi, Q)|^2 \rangle = 1$$

→ Signal-to-noise ratio estimator

$$\hat{\rho}^2(\tau, \phi, Q) = |X^w(\tau, \phi, Q)|^2 - \langle |N^w(\tau, \phi, Q)|^2 \rangle = |X^w(\tau, \phi, Q)|^2 - 1$$



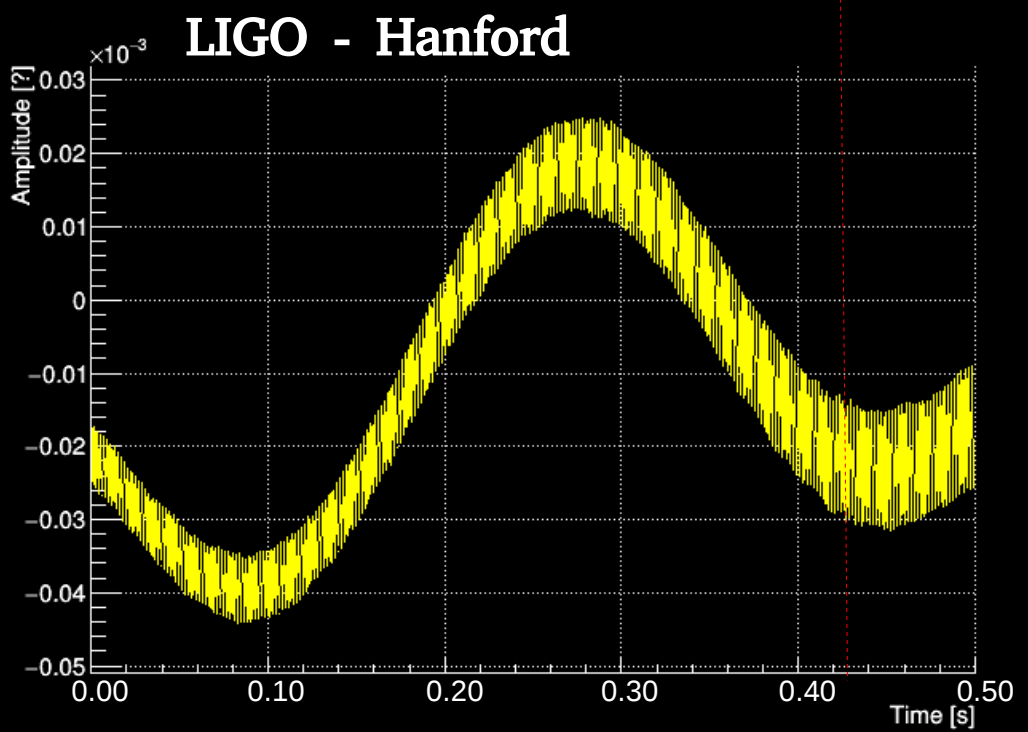
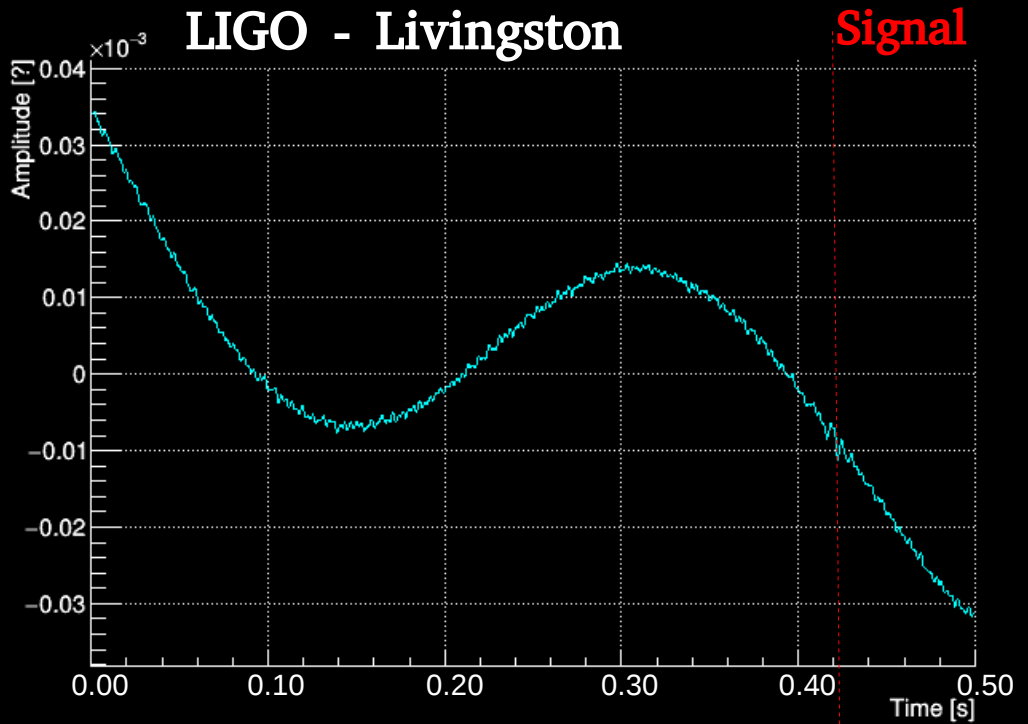
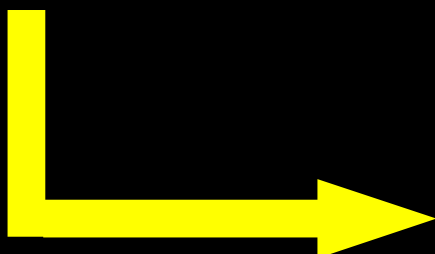
➡ Statistical interpretation: noise is Gaussian-distributed with unit variance

GW150914

Output power



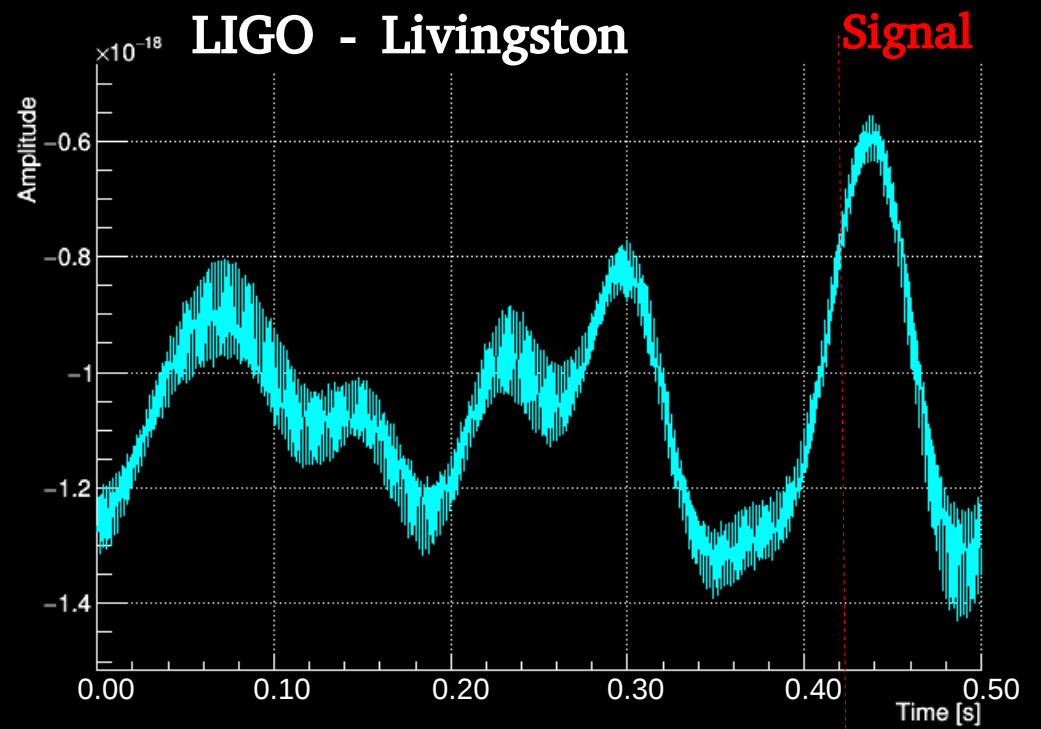
Output power



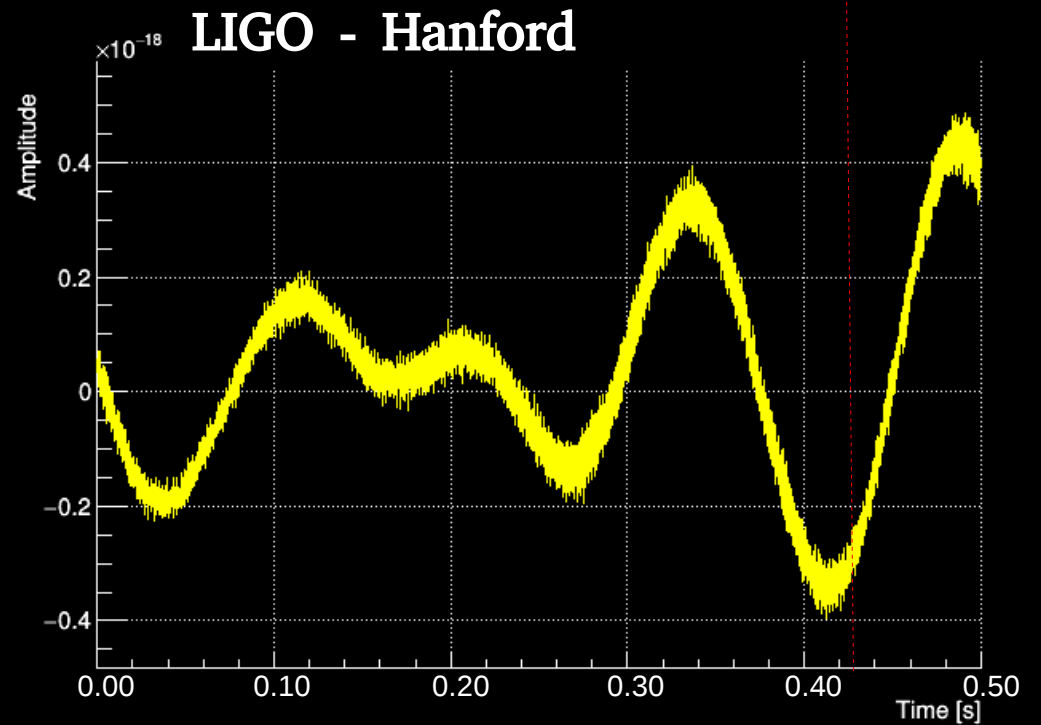
GW150914

$$h(t)$$

Data is calibrated
→ GW strain amplitude $h(t)$

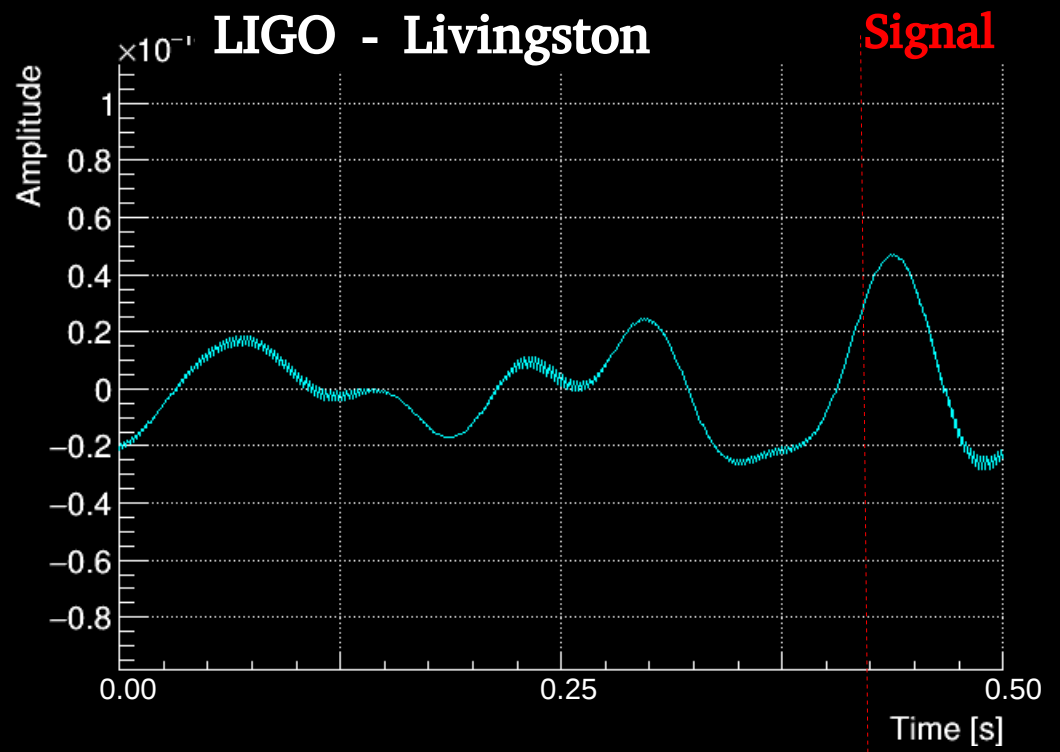


$$h(t)$$

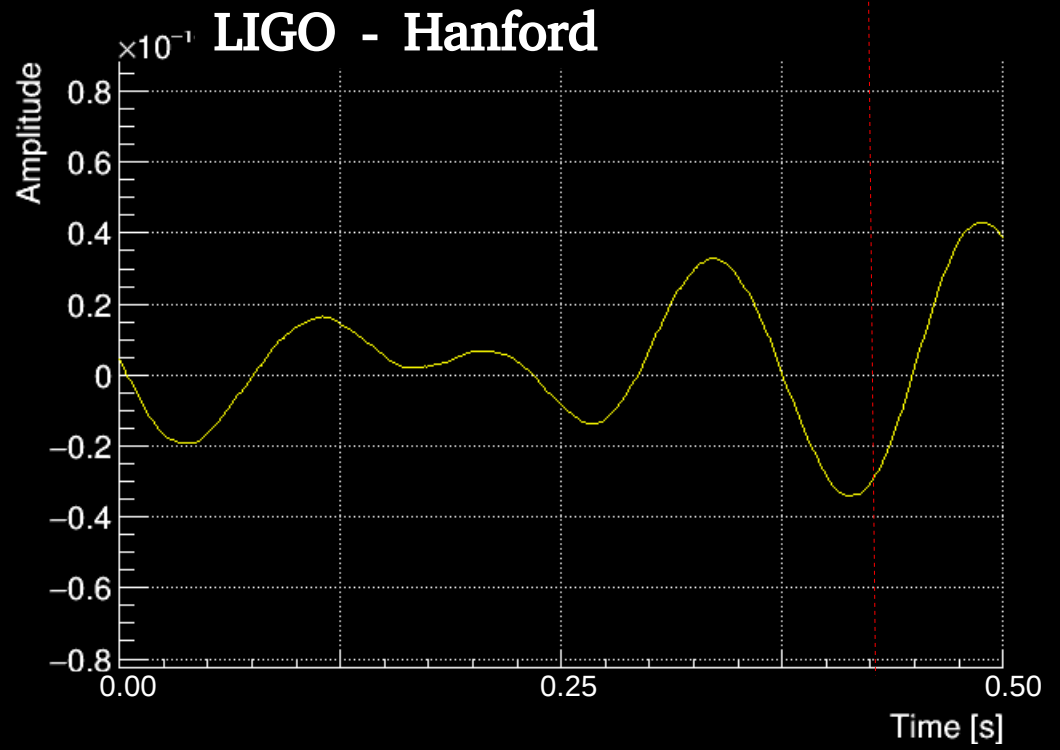


GW150914

Data are low-pass filtered
(here, < 500 Hz)

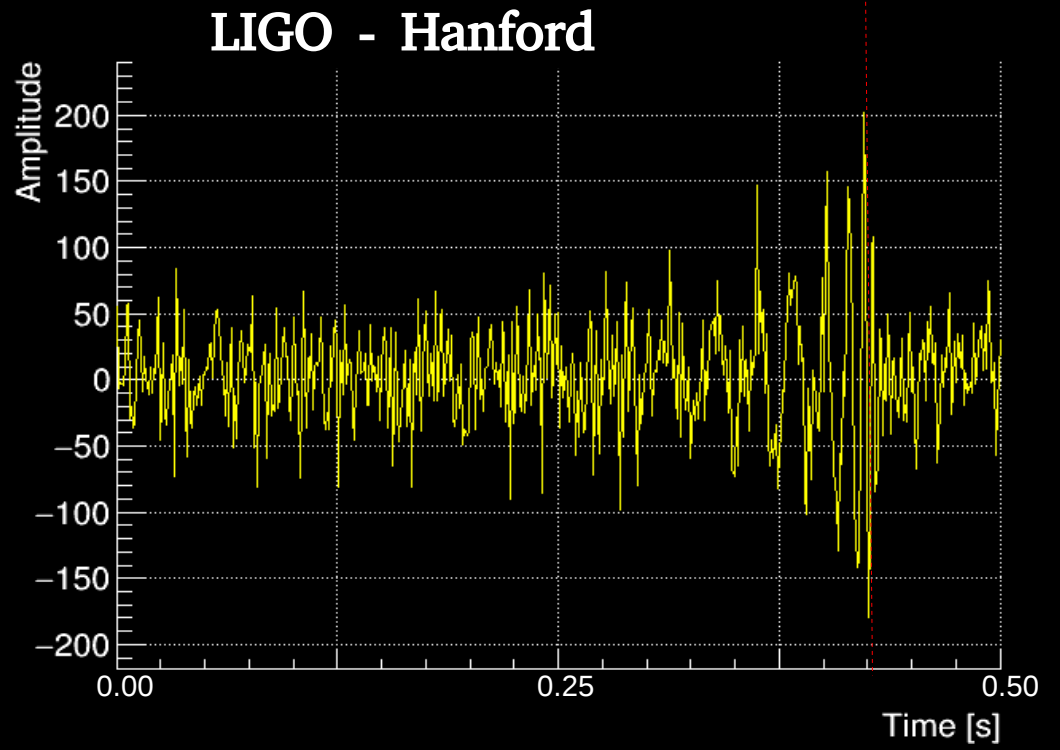
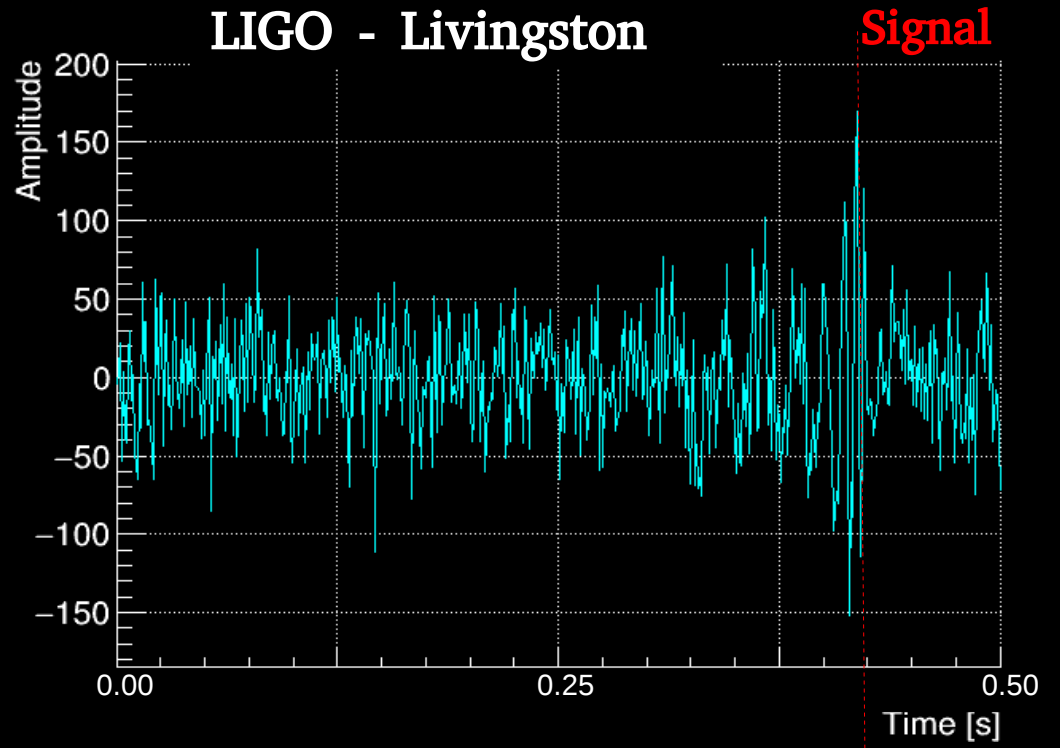


30

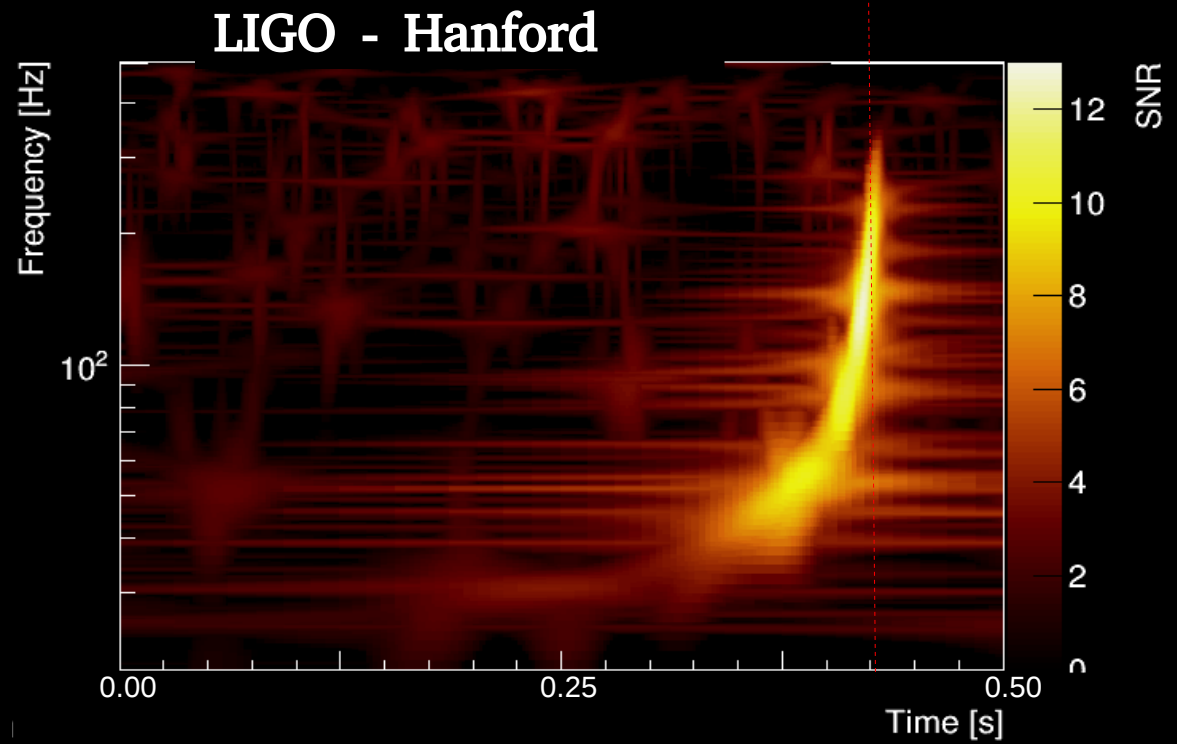
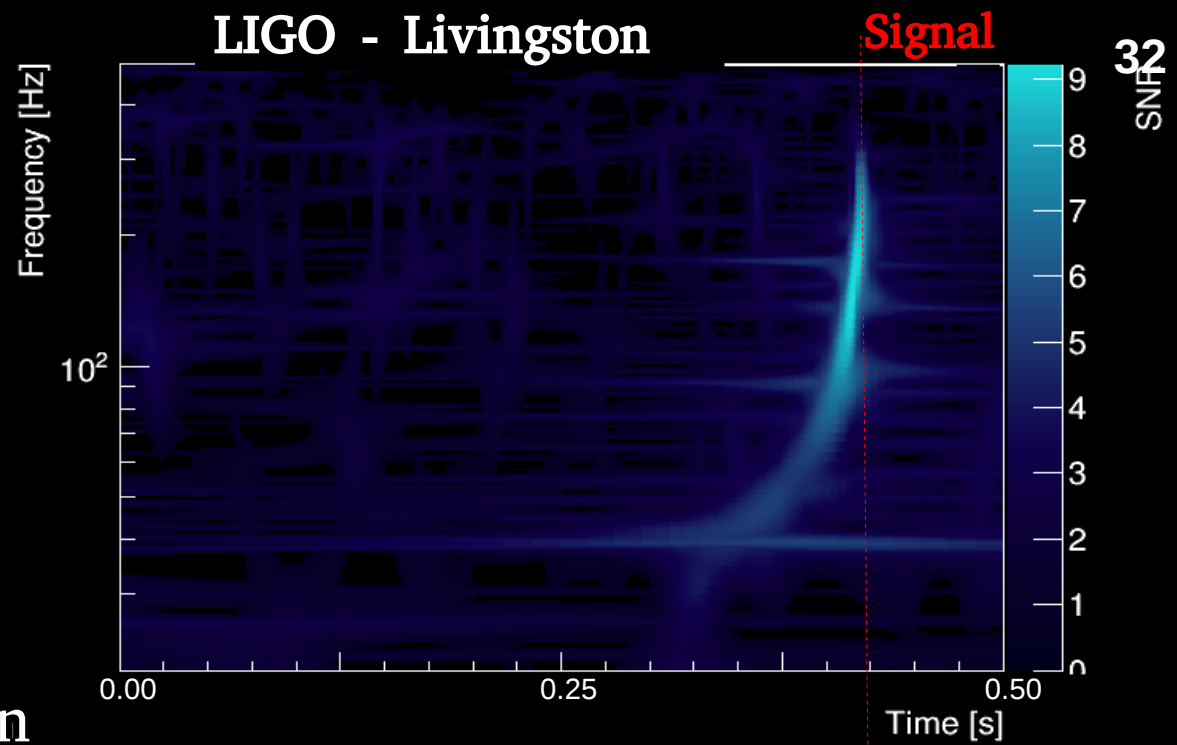


GW150914

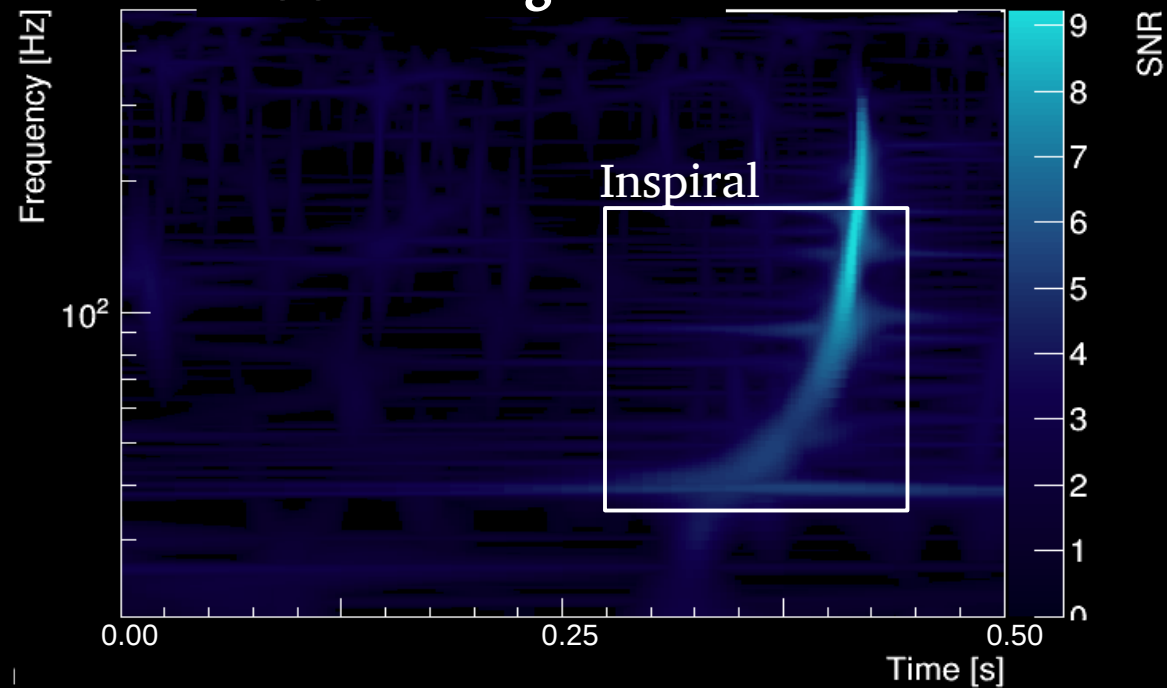
Data are whitened



Time-frequency decomposition
(Short Fourier transforms)



LIGO - Livingston



Phase evolution dictated by the chirp mass:

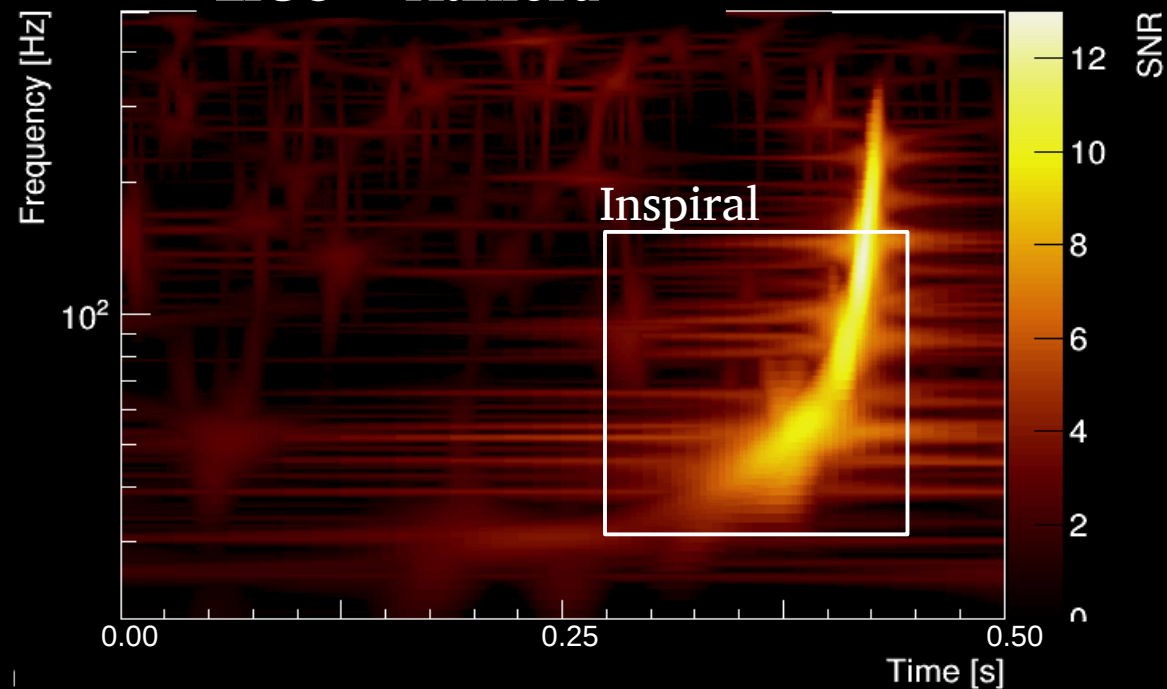
$$M_{chirp} = \frac{(m_1 m_2)^{3/5}}{(m_1 + m_2)^{1/5}} = \frac{c^3}{G} \left[\frac{5}{96} \pi^{-8/3} f^{-11/3} \dot{f} \right]^{3/5}$$

$$M_{chirp} \simeq 30 M_{sun}$$

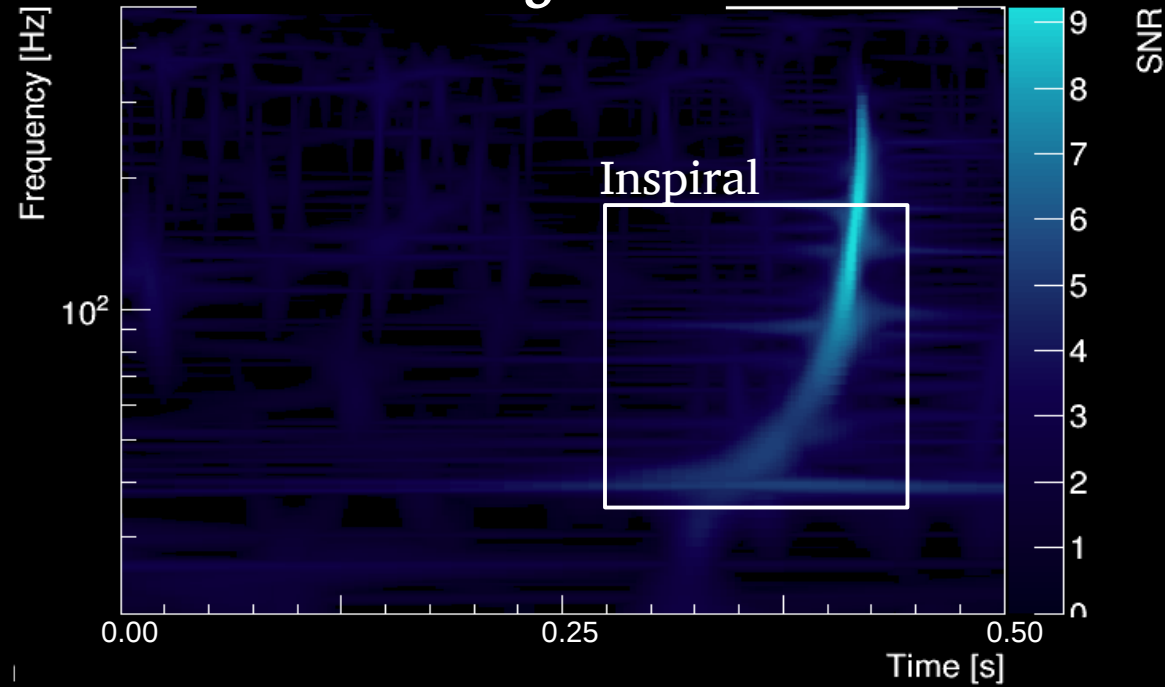
→ Total mass:

$$M_{tot} = m_1 + m_2 \geq 70 M_{sun}$$

LIGO - Hanford



LIGO - Livingston



Phase evolution dictated by the chirp mass:

$$M_{chirp} = \frac{(m_1 m_2)^{3/5}}{(m_1 + m_2)^{1/5}} = \frac{c^3}{G} \left[\frac{5}{96} \pi^{-8/3} f^{-11/3} \dot{f} \right]^{3/5}$$

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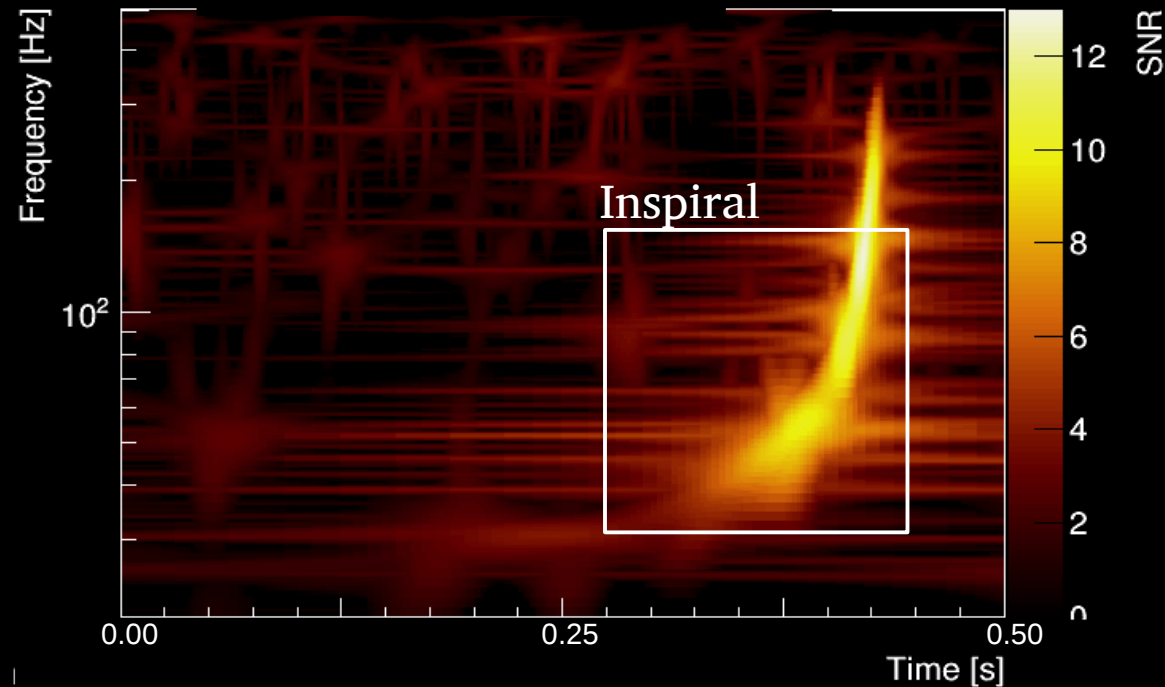
→ Total mass:

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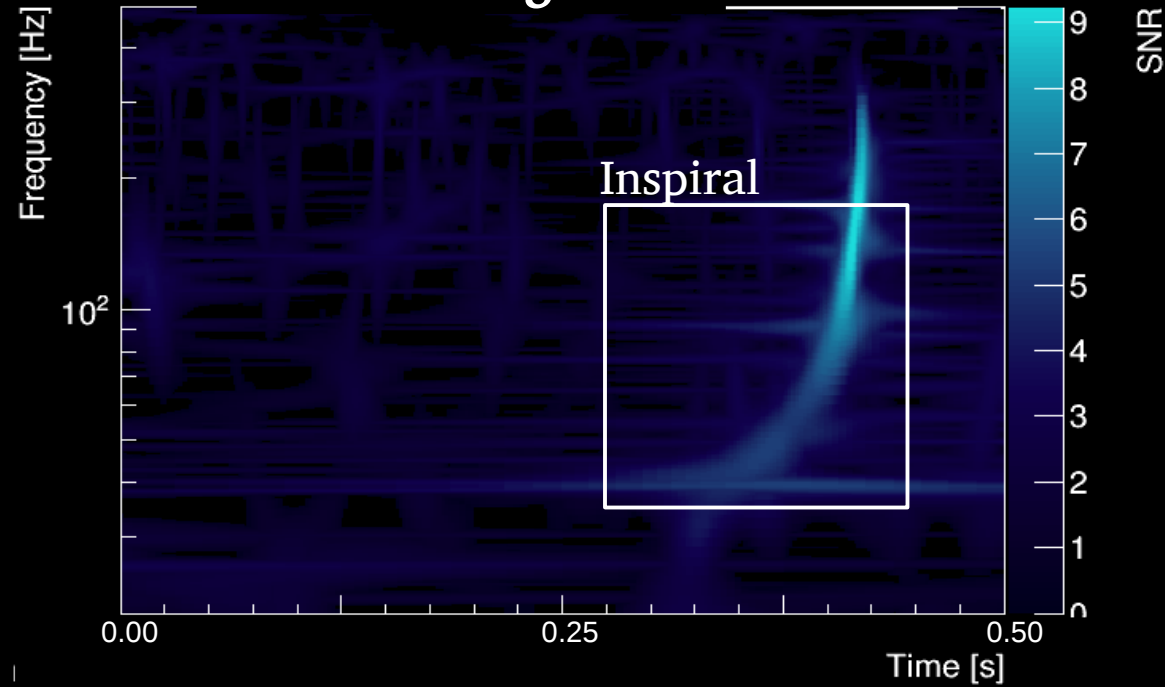
Schwarzschild radius:

$$\frac{2GM_{tot}}{c^2} \geq 210 \text{ km}$$

LIGO - Hanford



LIGO - Livingston



Phase evolution dictated by the chirp mass:

$$M_{chirp} = \frac{(m_1 m_2)^{3/5}}{(m_1 + m_2)^{1/5}} = \frac{c^3}{G} \left[\frac{5}{96} \pi^{-8/3} f^{-11/3} \dot{f} \right]^{3/5}$$

$$M_{chirp} \simeq 30 M_{sun}$$

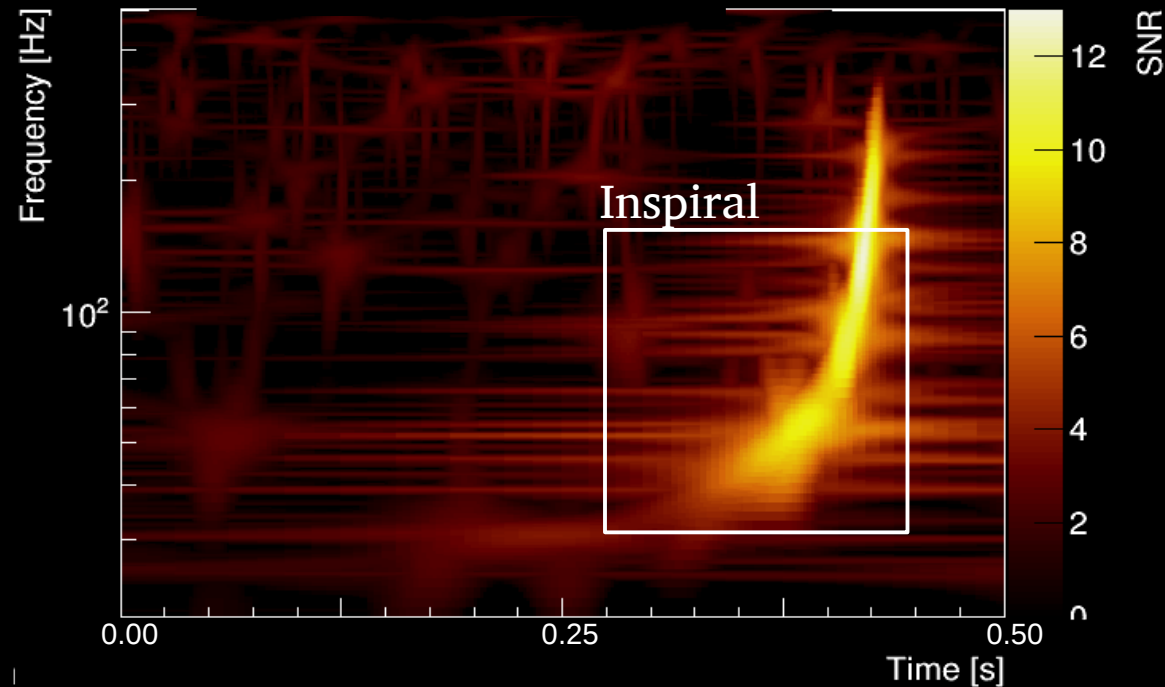
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Schwarzschild radius:

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LIGO - Hanford



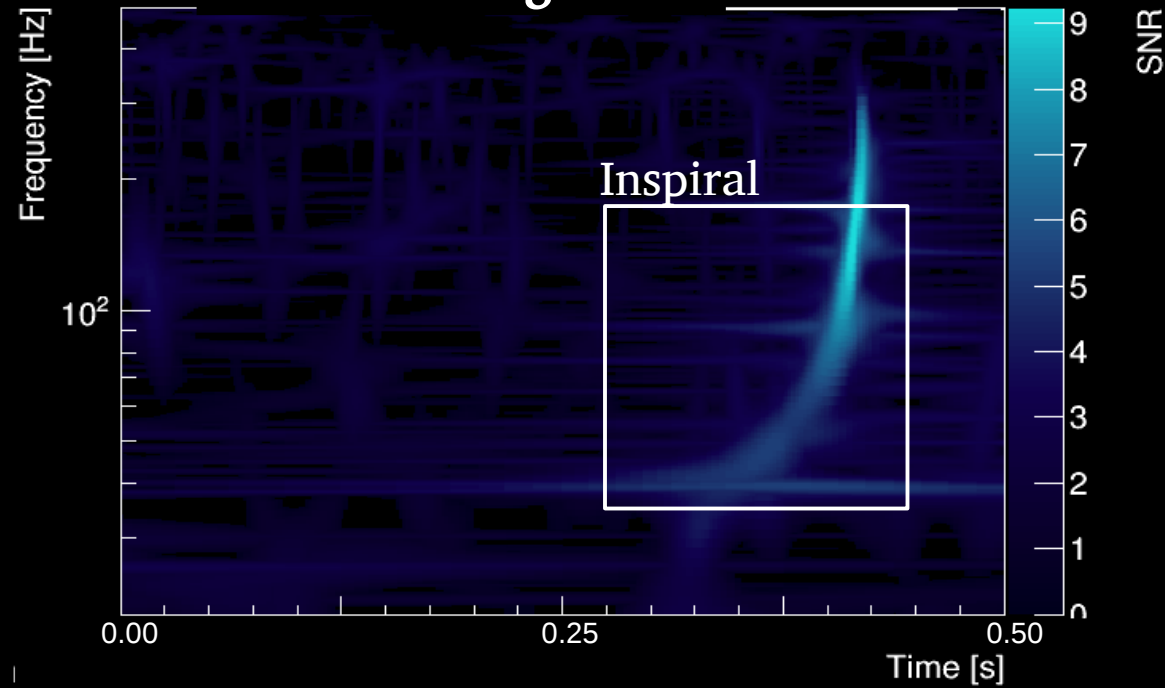
Orbital frequency:

$$f_{orbit} = f/2 \simeq 75 \text{ Hz}$$

Equal Newtonian point masses orbit:

$$d \simeq 350 \text{ km}$$

LIGO - Livingston



Phase evolution dictated by the chirp mass:

$$M_{\text{chirp}} = \frac{(m_1 m_2)^{3/5}}{(m_1 + m_2)^{1/5}} = \frac{c^3}{G} \left[\frac{5}{96} \pi^{-8/3} f^{-11/3} \dot{f} \right]^{3/5}$$

$$M_{\text{chirp}} \simeq 30 M_{\text{sun}}$$

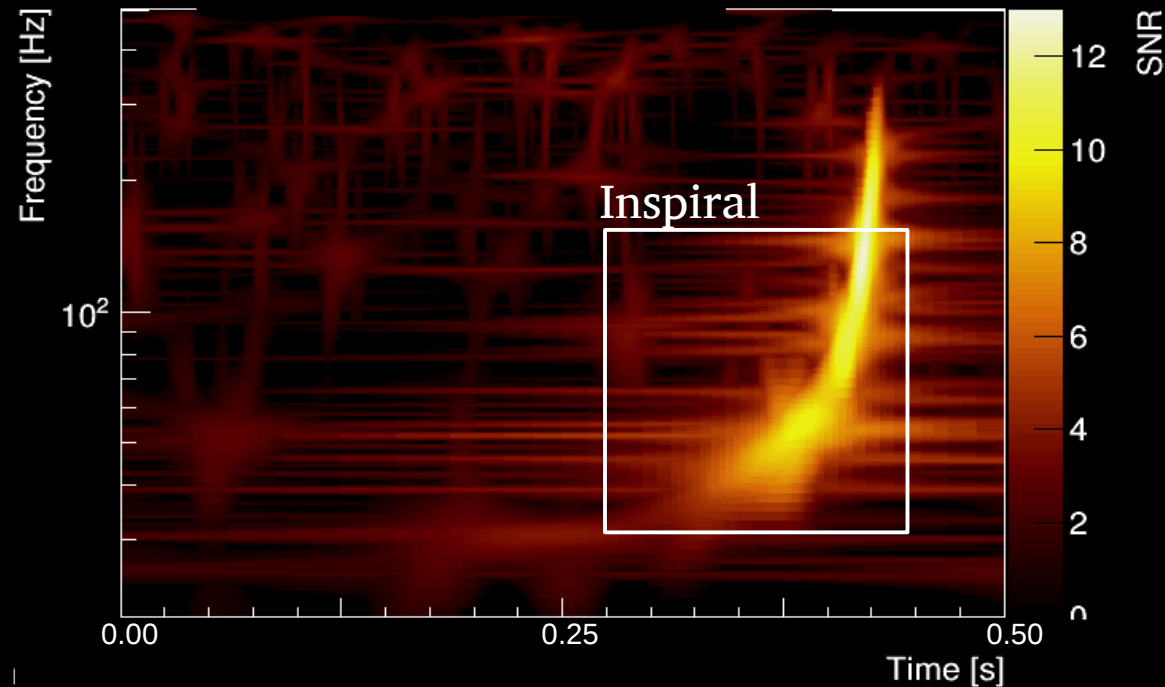
→ Total mass:

$$M_{\text{tot}} = m_1 + m_2 \geq 70 M_{\text{sun}}$$

Schwarzschild radius:

$$\frac{2GM_{\text{tot}}}{c^2} \geq 210 \text{ km}$$

LIGO - Hanford



Orbital frequency:

$$f_{\text{orbit}} = f/2 \simeq 75 \text{ Hz}$$

Equal Newtonian point masses orbit:

$$d \simeq 350 \text{ km}$$

→ Black hole binary

Known waveform → match-filtering technique

Simplest linear filter: correlation $C(t) = \int_{-\infty}^{+\infty} h_{det}(t')k(t-t')dt = \int_{-\infty}^{+\infty} \tilde{h}_{det}(f)\tilde{k}^*(f)e^{2i\pi ft}df$

$k(t)$ is the impulse response function of the filter : $h_{det}(t) = \delta(t) \Rightarrow C(t) = k(t)$

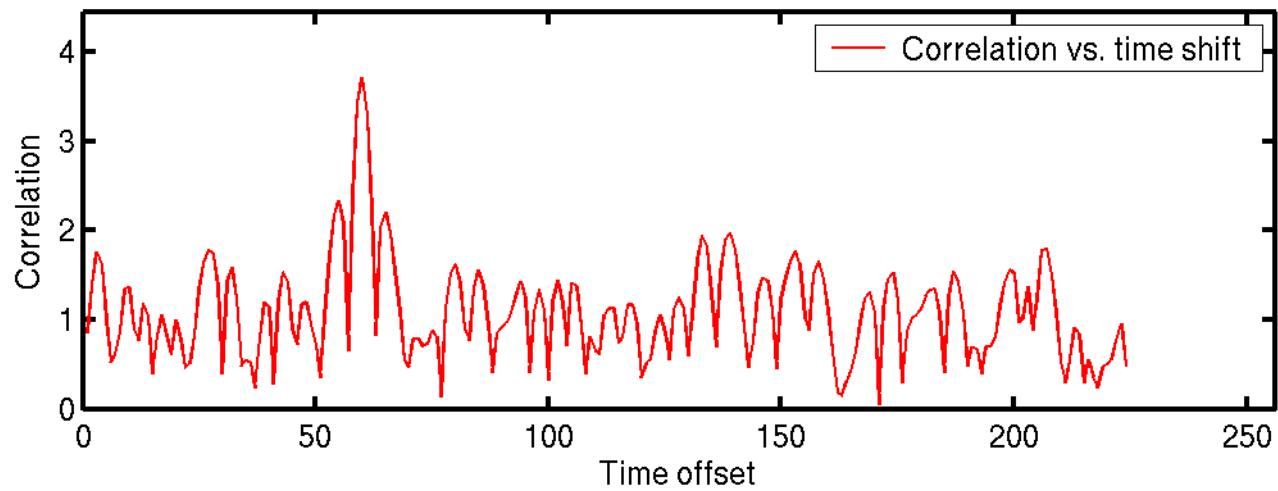
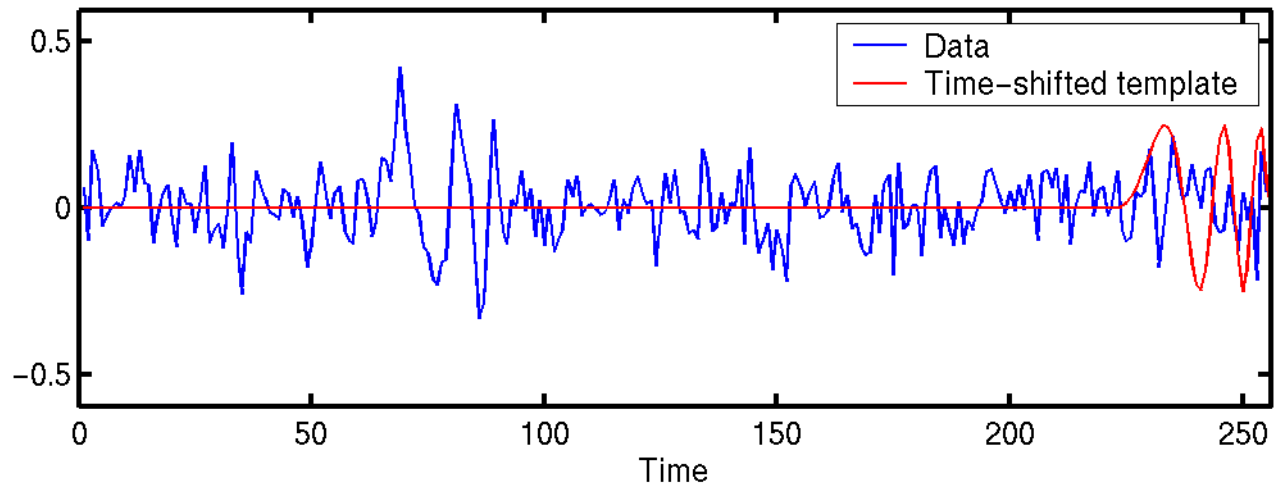
Match-filter: optimal filter maximizing the SNR in presence of additive noise

$$h_{det}(t) = n(t) + GW(t)$$

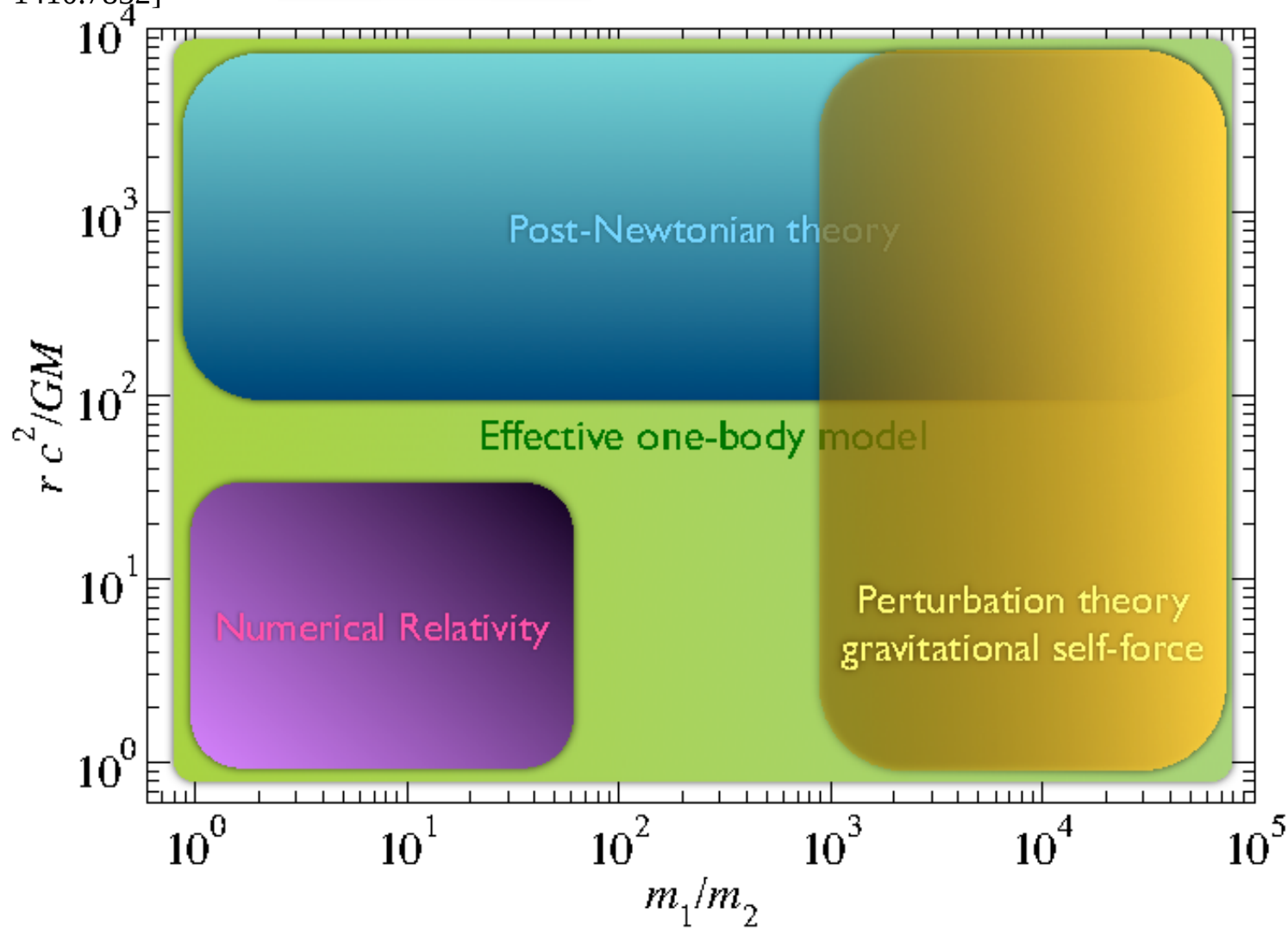
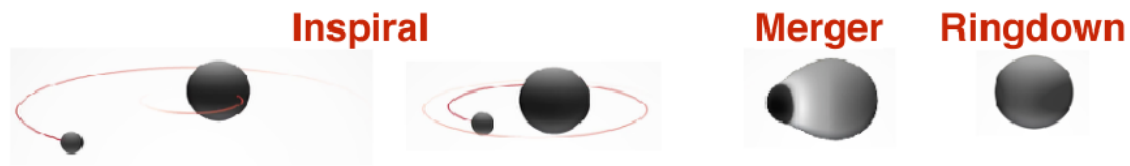
$$\rho(t) = \frac{C(t)}{\sqrt{\langle N^2(t) \rangle}} \quad \text{with} \quad \langle N^2(t) \rangle = \int_{-\infty}^{+\infty} |\tilde{k}(f)|^2 S_n(f) df$$

The SNR is maximized if $\tilde{k}(f) \propto \frac{G\tilde{W}^*(f)}{S_n(f)}$

$$\rho(t) = \int_{-\infty}^{+\infty} \frac{G\tilde{W}^*(f)\tilde{h}_{det}(f)}{S_n(f)} e^{2i\pi ft} df$$

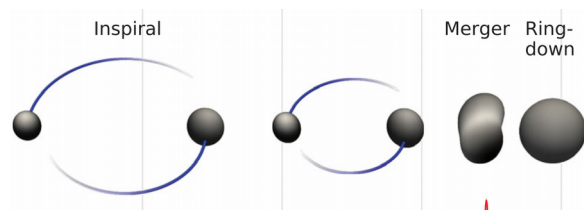


[arXiv 1410.7832]



Theoretical input:

- 90s: CBC PN waveforms (Blanchet, Iyer, Damour, Deruelle, Will, Wiseman, ...)
- 00s: CBC Effective One Body “EOB” (Damour, Buonanno)
- 06: BBH numerical simulation (Pretorius, Baker, Loustos, Campanelli)



The intrinsic waveform parameters:

- Masses (2 dofs):

$$M_{tot} = M_1 + M_2$$

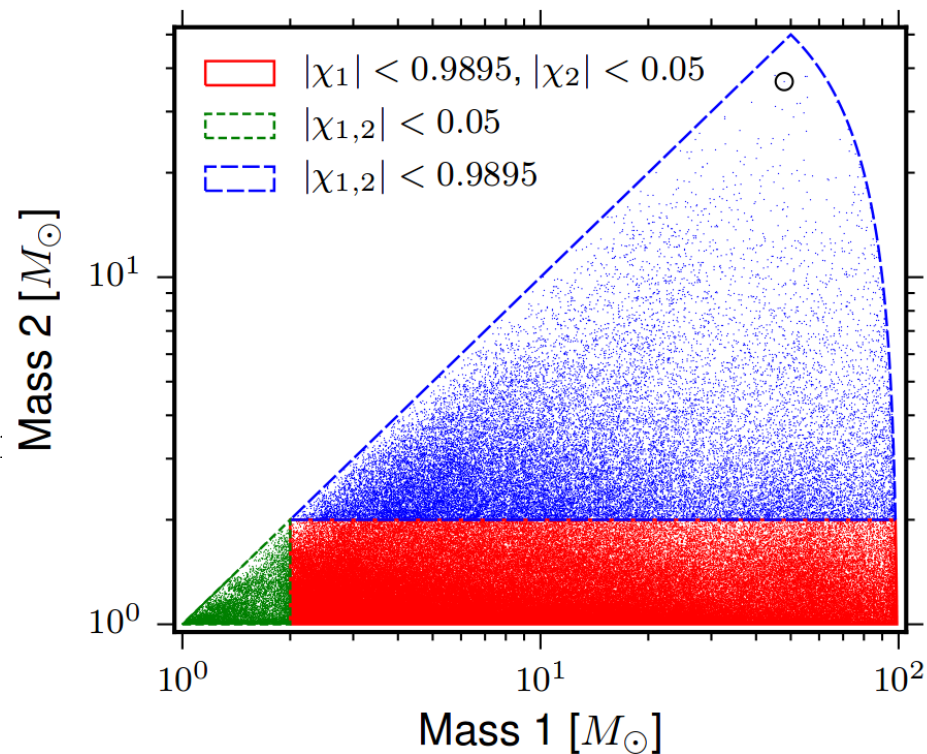
- Spins and orbital angular momentum (6 dofs):

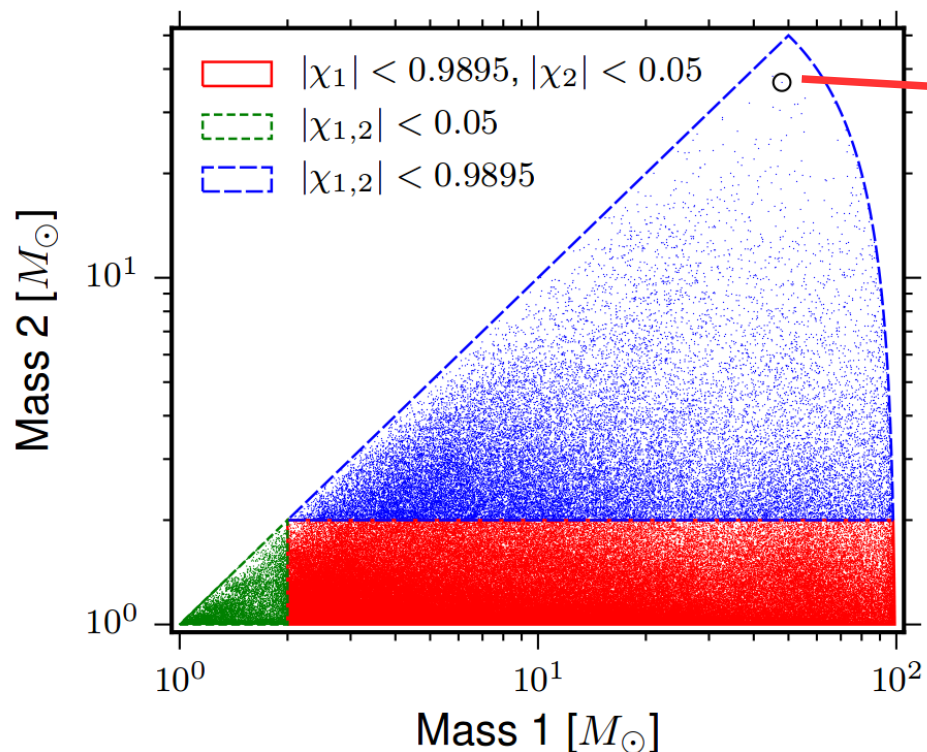
$$\vec{S}_{tot} = \vec{S}_1 + \vec{S}_2 + \vec{J}$$

The waveform models used for the search:

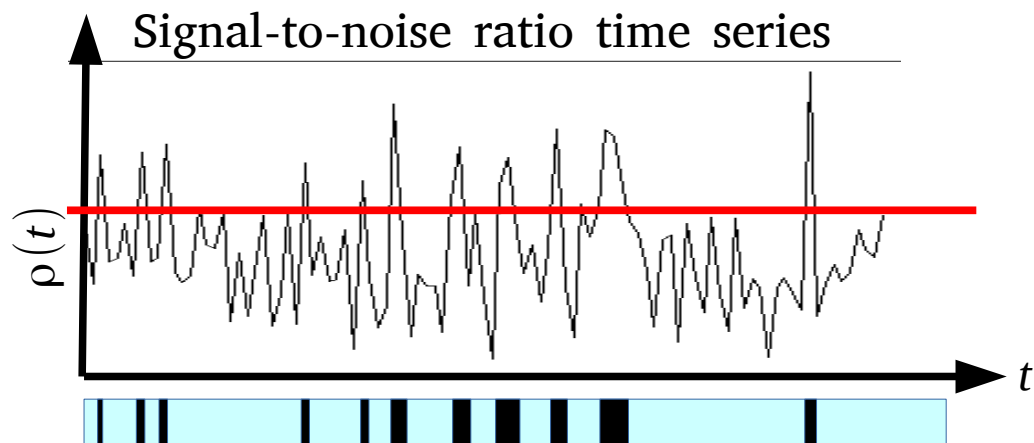
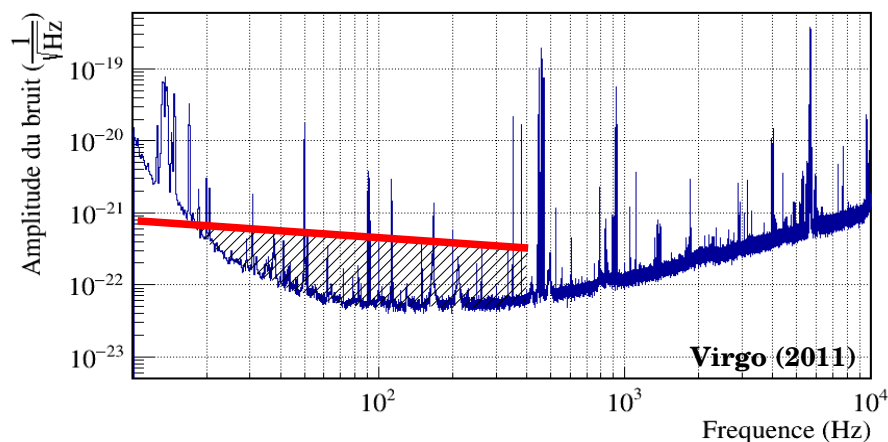
- Inspiral, PN3.5 for $M_{tot} < 4 M_{sun}$
- Inspiral/Merger/Ringdown EOB + numerical relativity for $M_{tot} > 4 M_{sun}$
- Spins and orbital angular momentum are aligned

Template bank → match-filtering technique



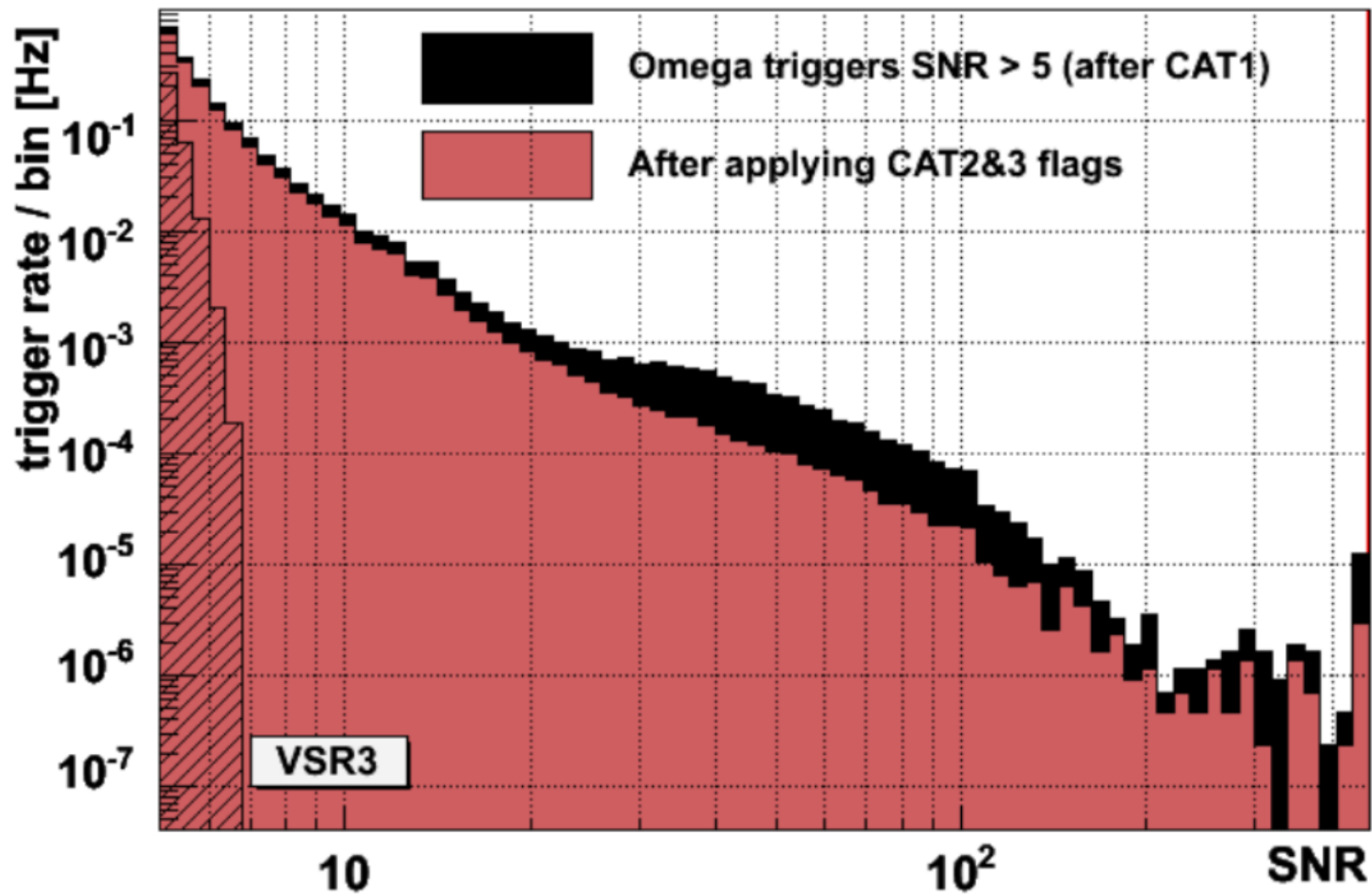


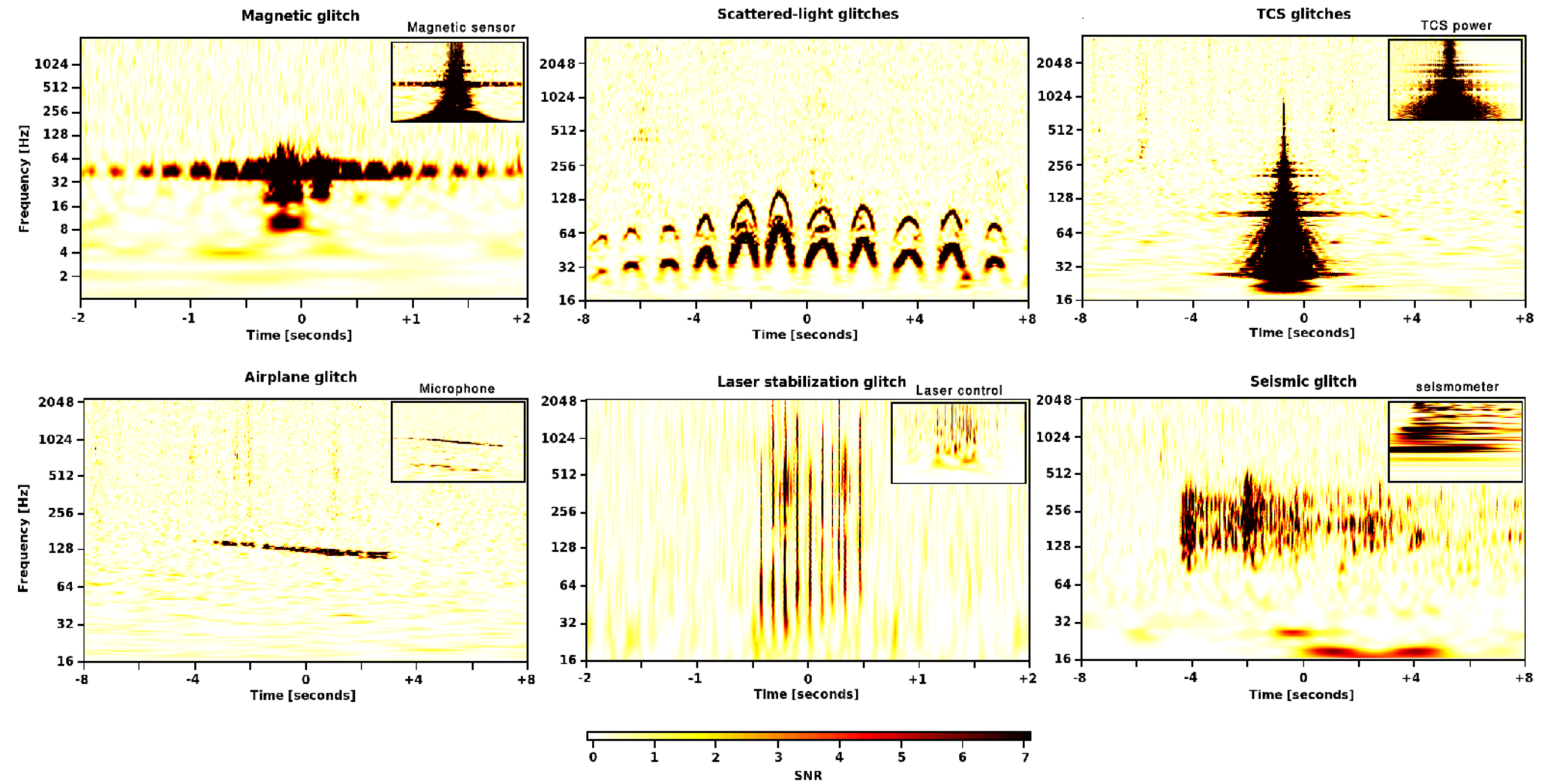
$$\rho(t) = \int_{-\infty}^{+\infty} \frac{G\tilde{W}^*(f)\tilde{h}_{det}(f)}{S_n(f)} e^{2i\pi ft} df$$



- A list of events is produced:
- start/end/peak times
 - SNR
 - template parameters (masses, spins)

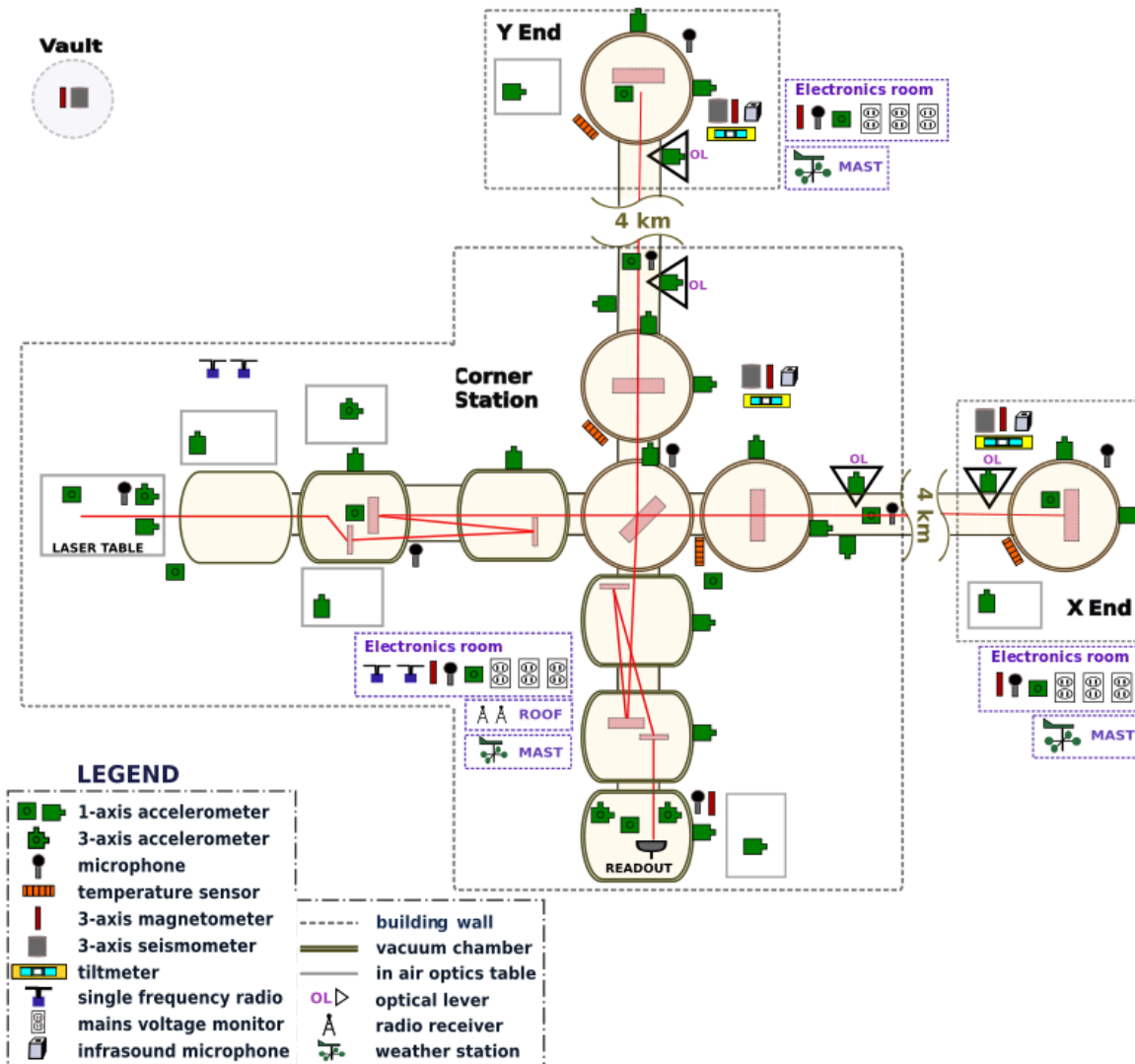
Now, the challenge is to reject noise events to better isolate true signals





Thousands of auxiliary channels are used to monitor the instruments

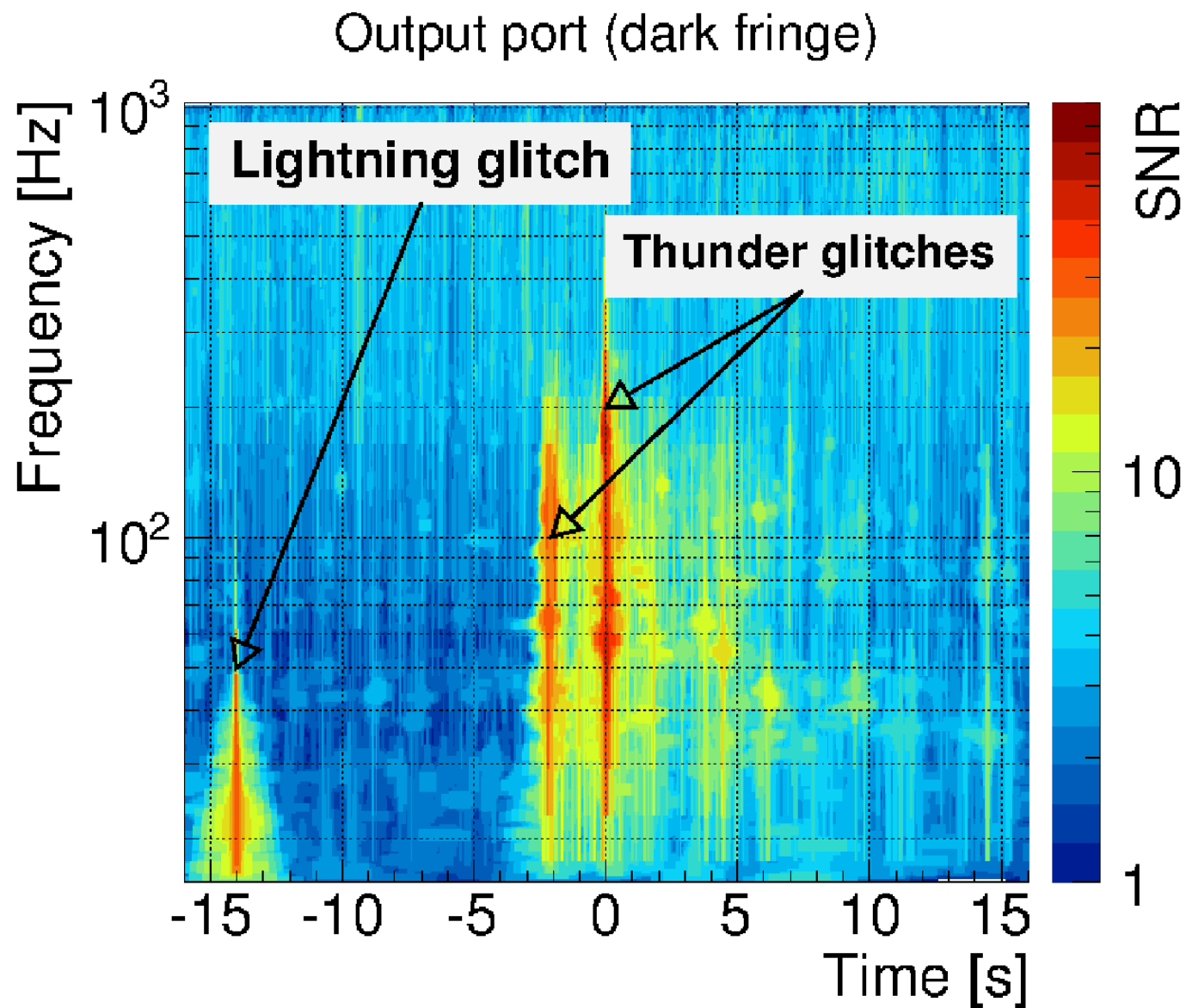
- environmental sensors
- detector sub-systems
- detector control

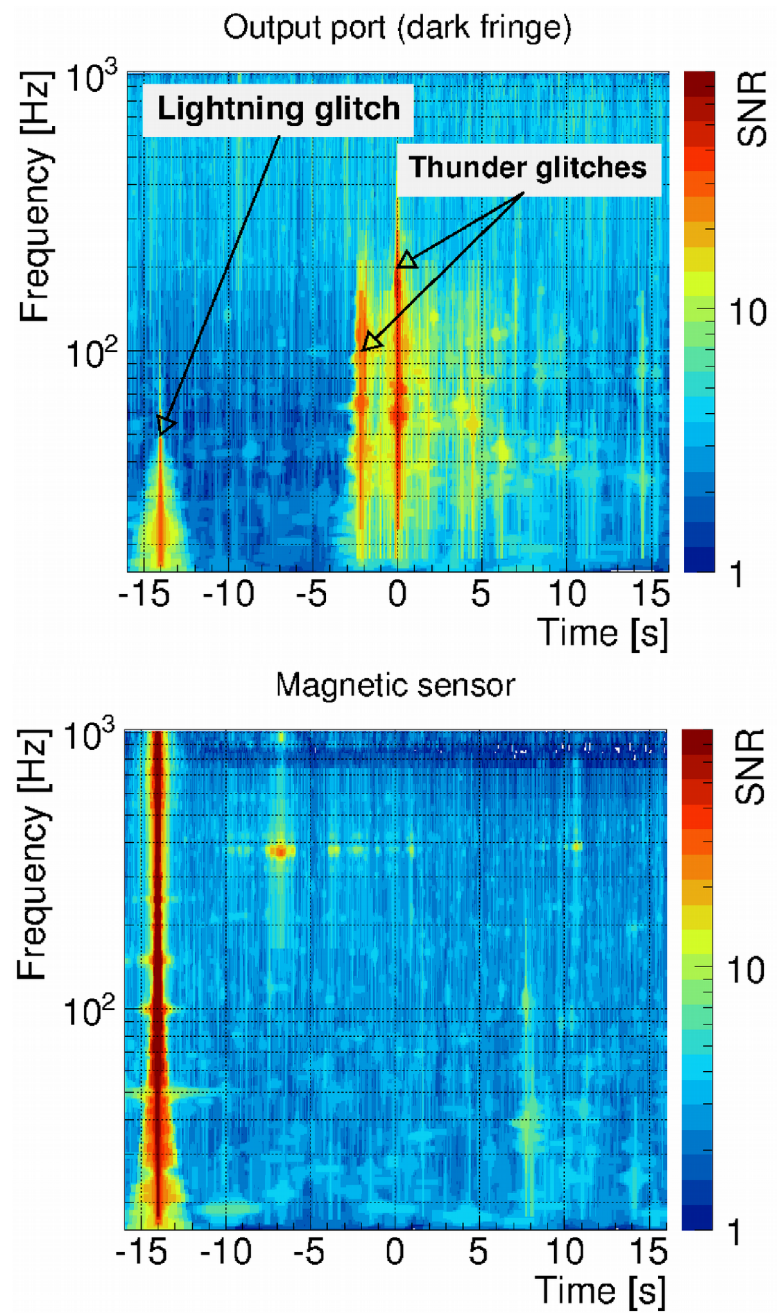


Noise injection campaigns are conducted to identify the detector's response to different noise stimulation

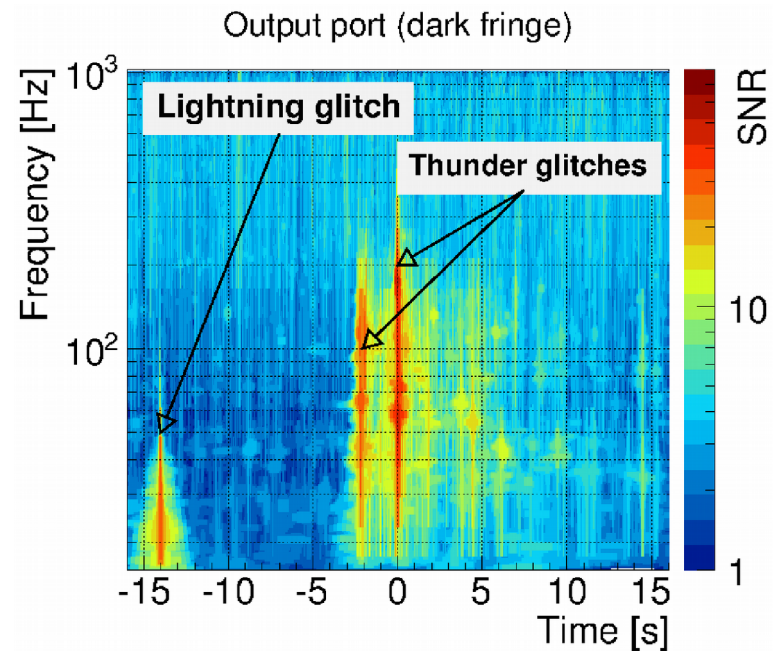
- Multiple transient noises were identified during the run
- Anthropogenic noise
 - Earthquakes
 - Radio-frequency modulation
 - ...

- Option #1: fix the detector
- Option #2: remove transient events in the data

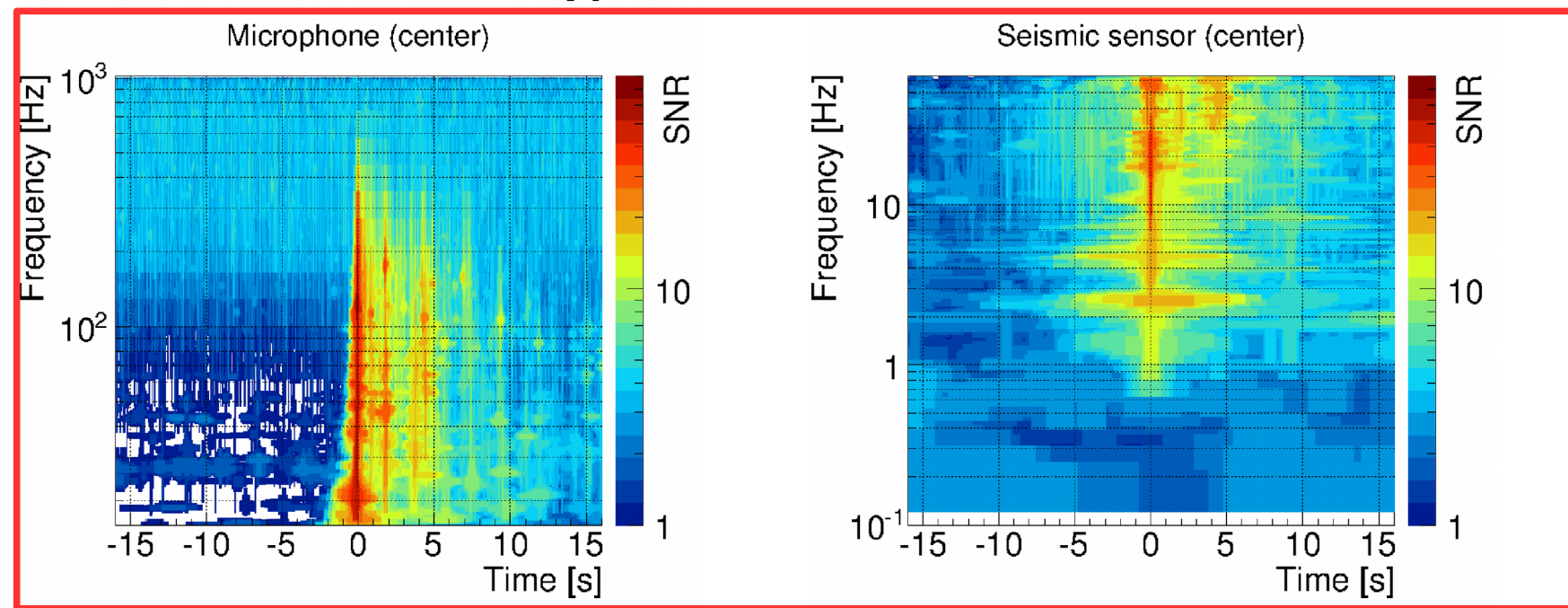




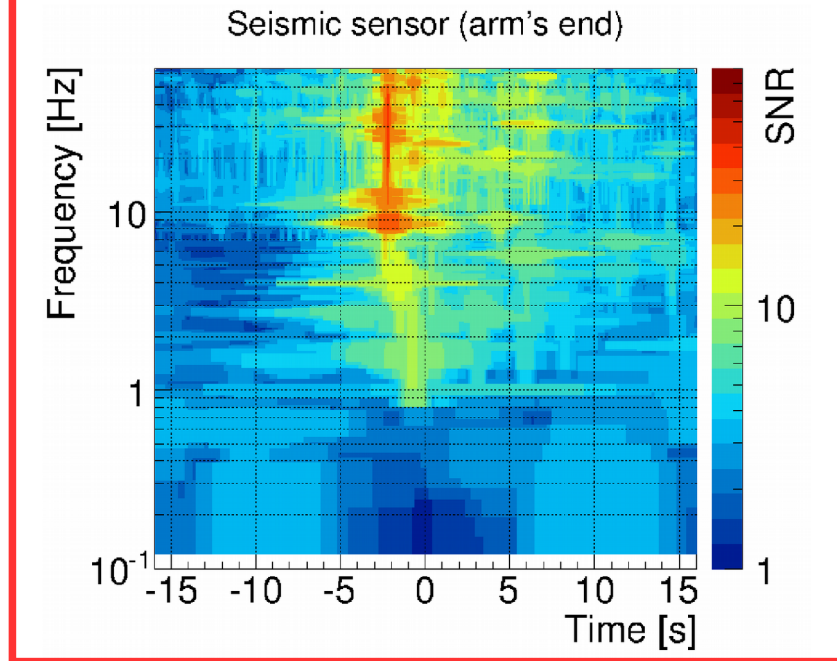
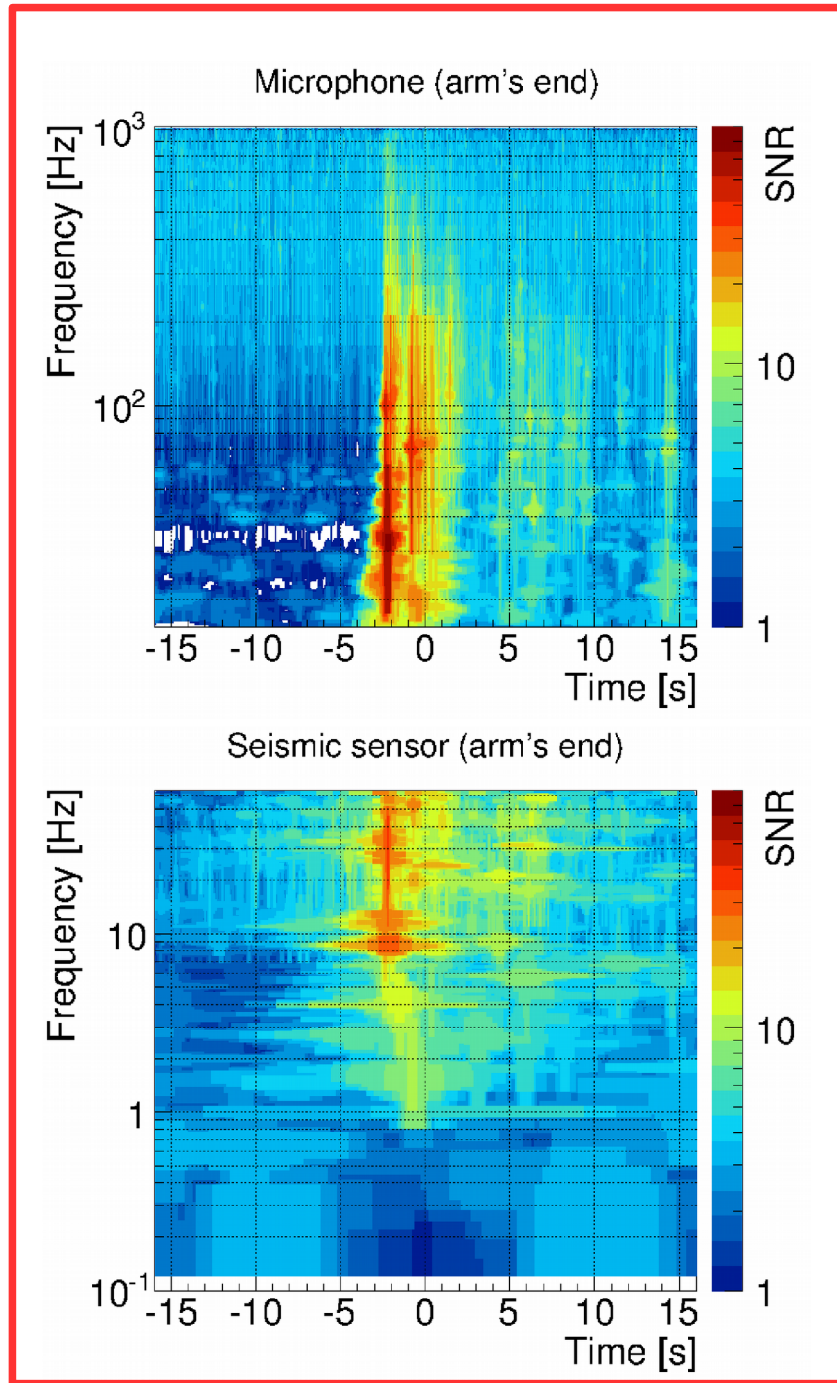
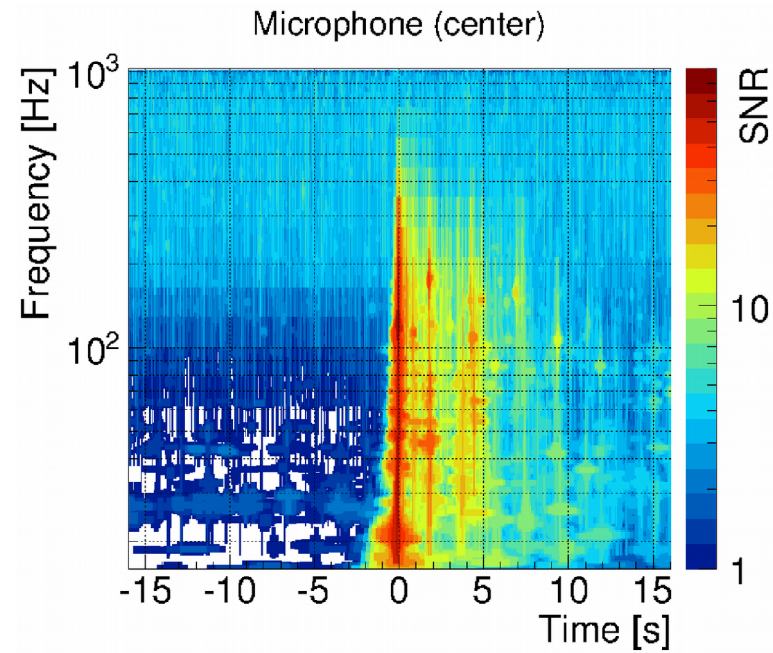
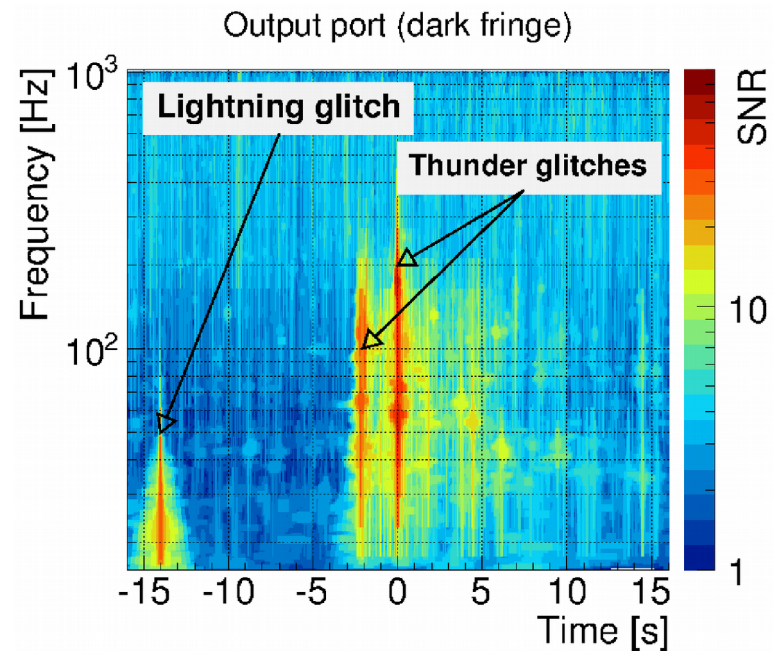
The lightnings are detected in magnetic sensors



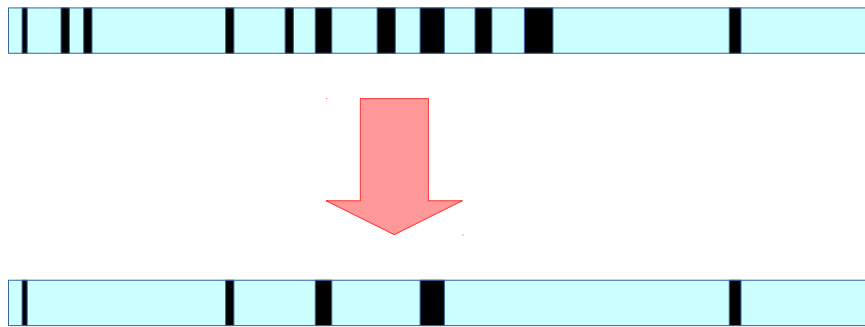
The thunder is detected in seismic sensors and microphones



Central building



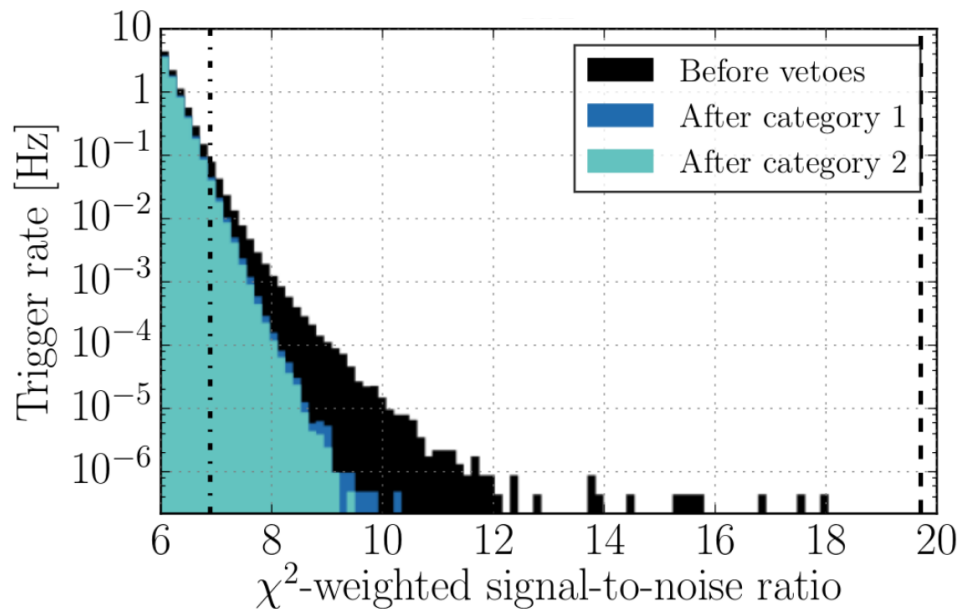
Arm's end



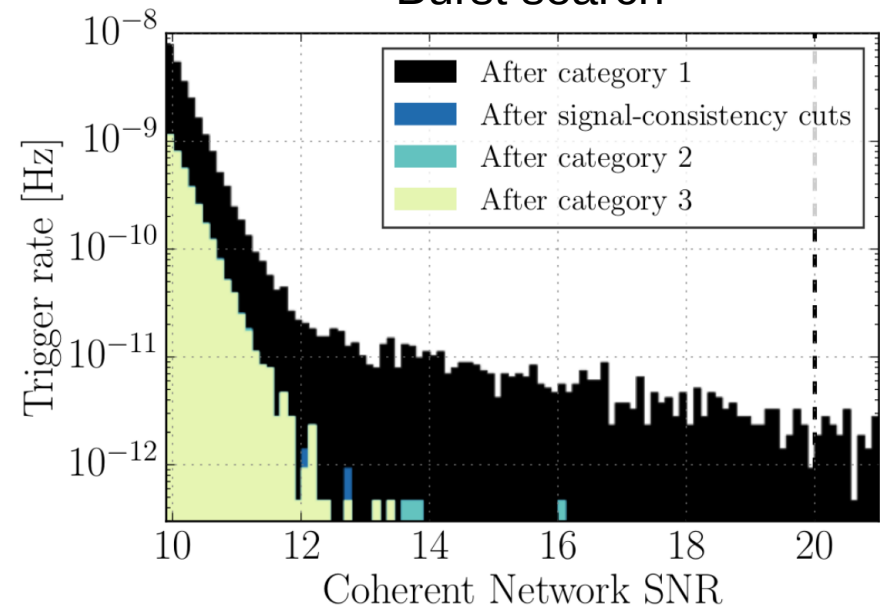
Analysis vetoes are created to remove transient noise events from known origins

- environment
- detector glitches (laser, control loops...)
- scattered light
- ...

CBC search

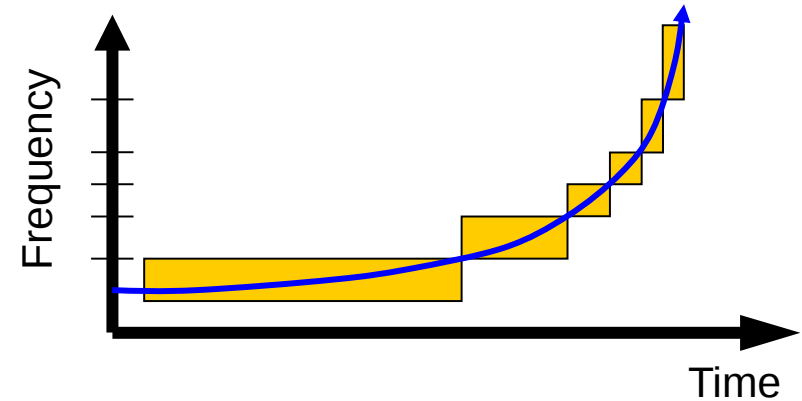


Burst search



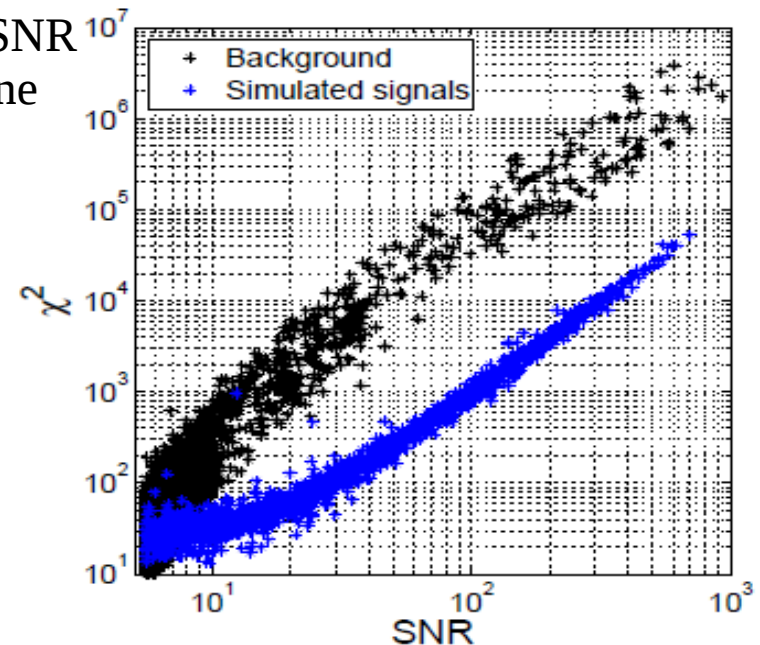
- Divide the “selected” template into p parts
- The frequency intervals are chosen so that for a true signal, the SNR is uniformly shared among the frequency bands.

$$\chi^2(t) = p \sum_{j=1}^p \left| \rho_j - \frac{\rho}{p} \right|^2$$



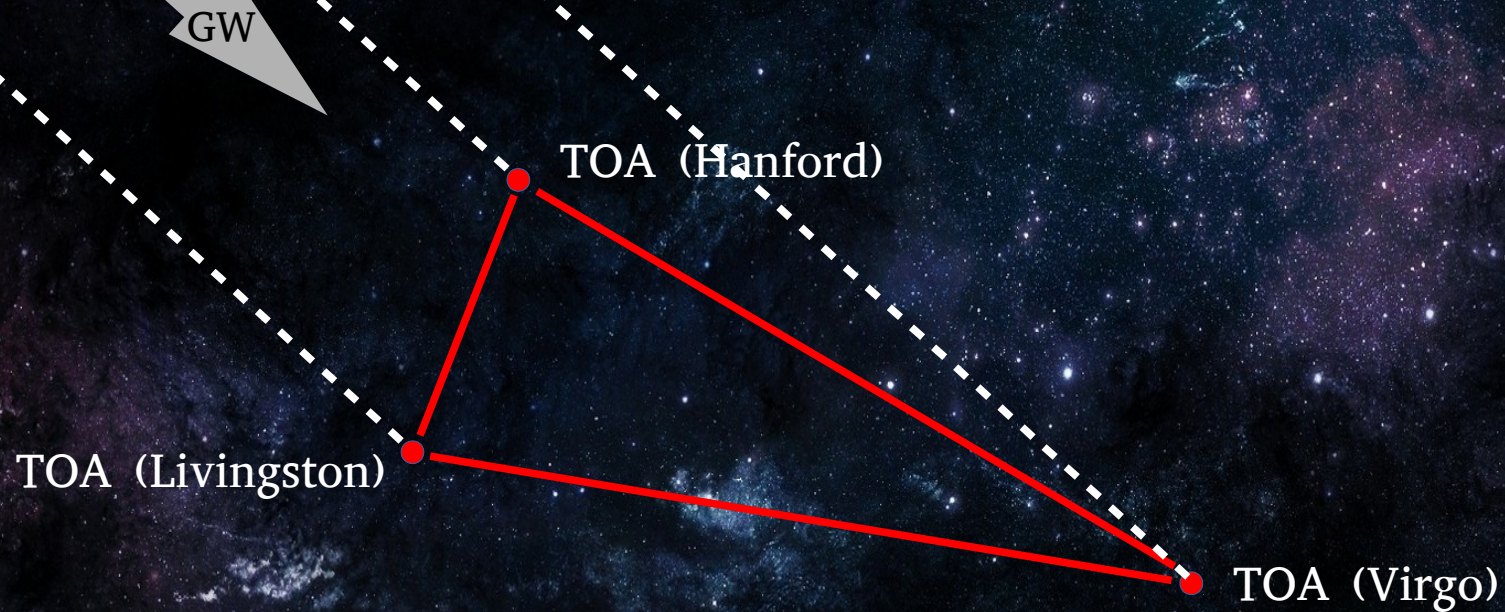
- For a stationary and Gaussian noise χ^2 has an expectation value: $\langle \chi^2 \rangle = p - 1$
- In practise χ^2 values are larger than expected for large SNR (discrete template banks effect) \rightarrow cut in (SNR, χ^2) plane
- Weighted SNR

$$\rho_{\text{new}} = \begin{cases} \rho, & \chi^2 \leq n_{\text{dof}} \\ \frac{\rho}{\left[\left(1 + \frac{\chi^2}{n_{\text{dof}}} \right)^{4/3} / 2 \right]^{1/4}}, & \chi^2 > n_{\text{dof}} \end{cases}$$



Coincidence between detectors

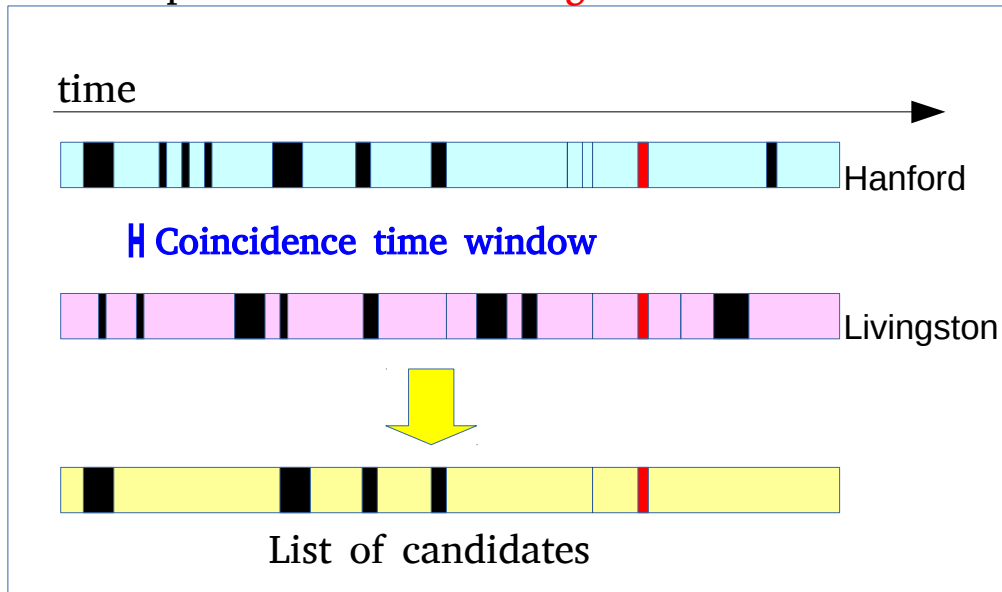




- The GW signal must be detected almost simultaneously across the network
 $\delta t \leq \text{Max light travel time}$
- The noise in the detectors is uncorrelated

A gravitational-wave signal is detected by multiple detectors almost simultaneously

True experiment = noise + **signal**



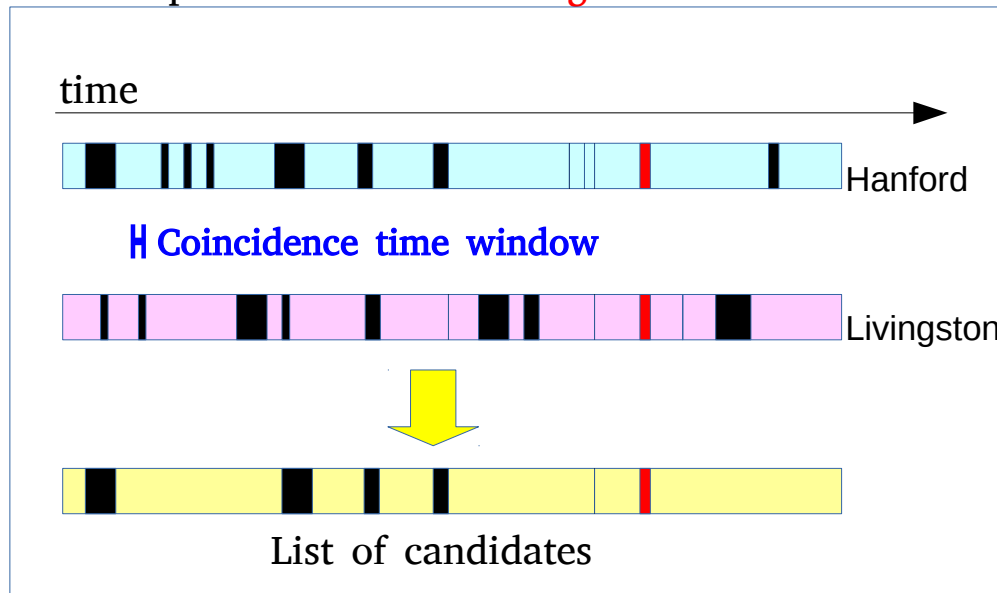
Coincidence rate:

$$R_{\text{coinc}} \sim R_H R_L \Delta t_{\text{win}}$$

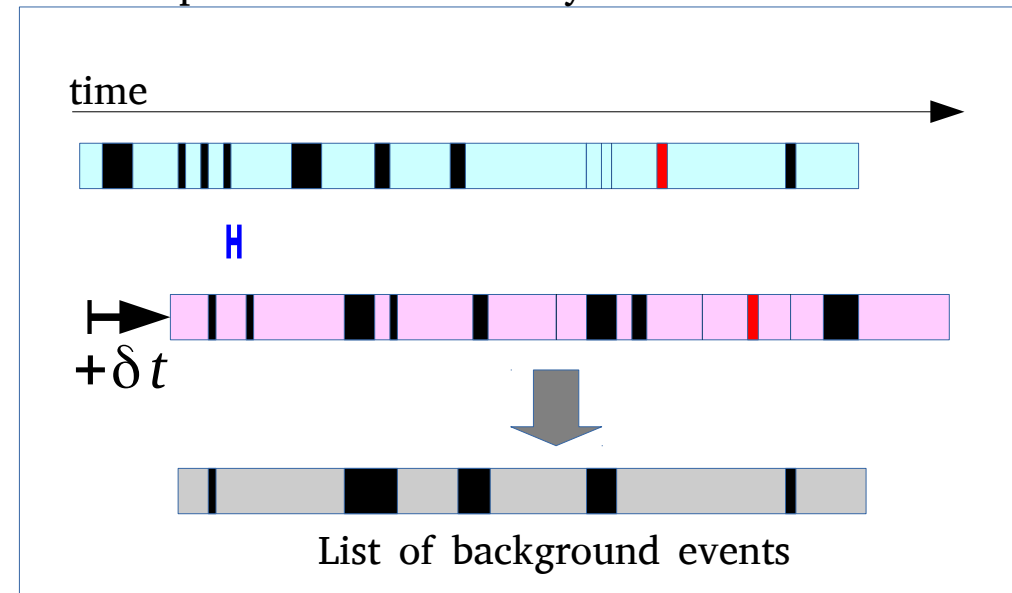
$$\sim (1 \text{ Hz}) \times (1 \text{ Hz}) \times (10^{-2} \text{ s}) = 10^{-2} \text{ Hz}$$

The background of a gravitational-wave search is estimated using the time-slide technique
 Assumption = uncorrelated noise between detectors

True experiment = noise + **signal**



Fake experiment = noise only

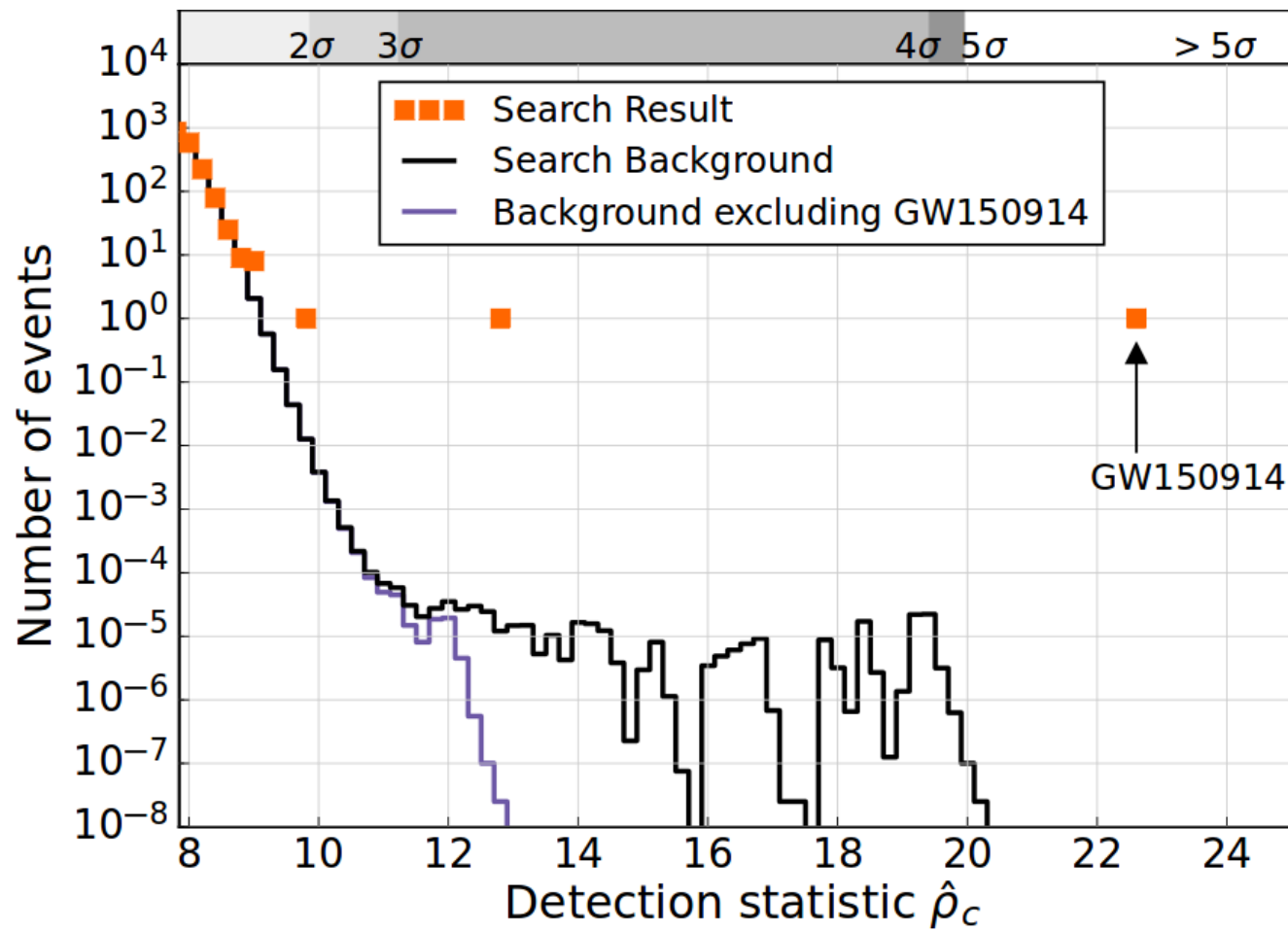


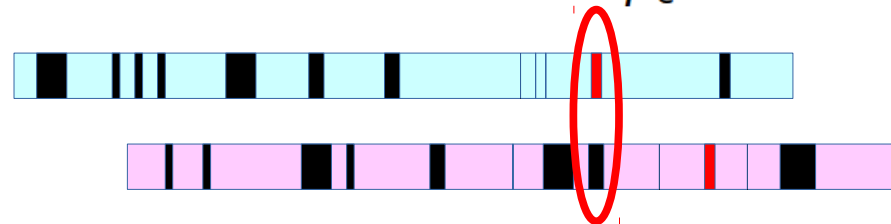
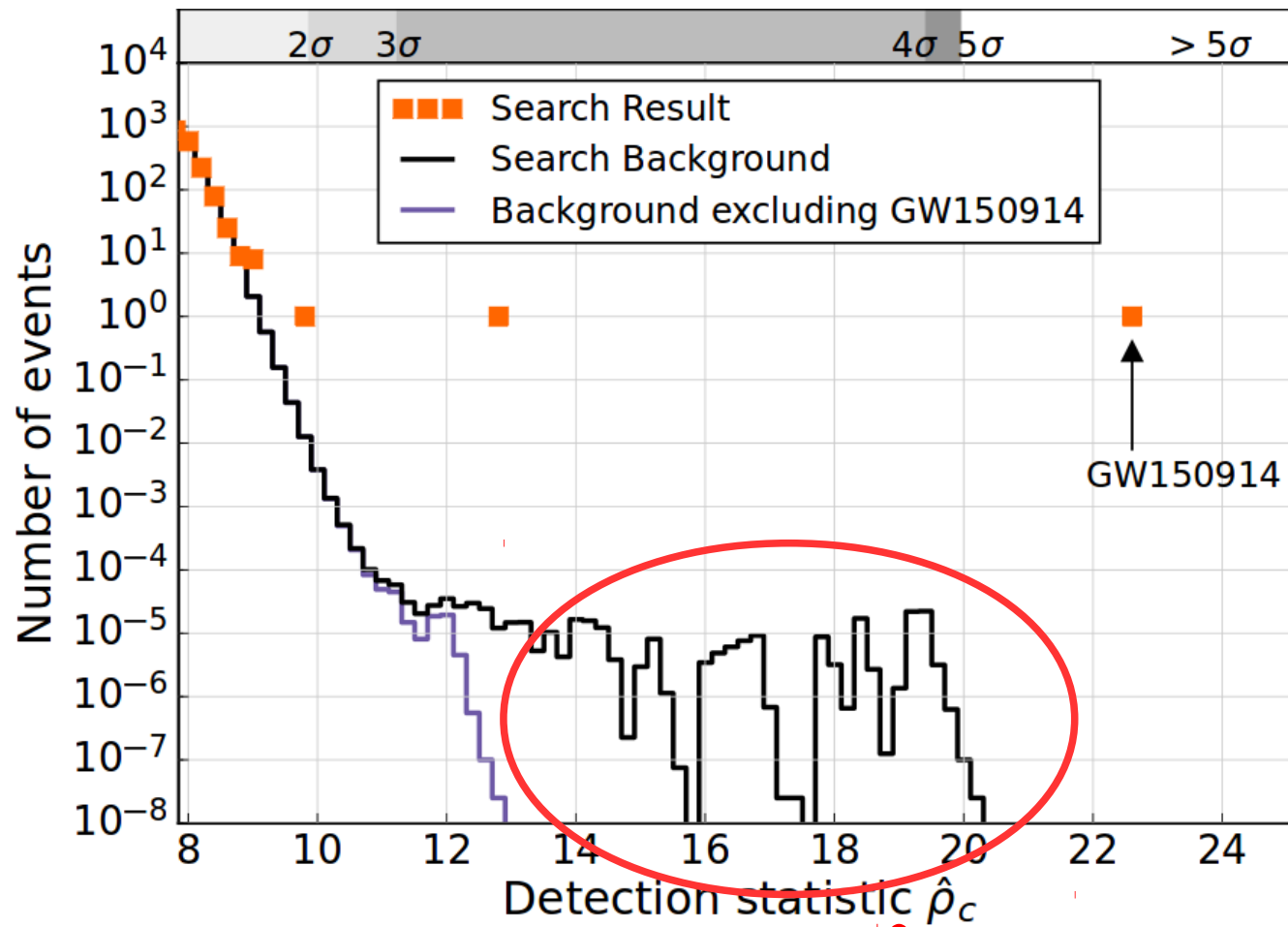
A very large number of fake experiments can be simulated using multiple offsets

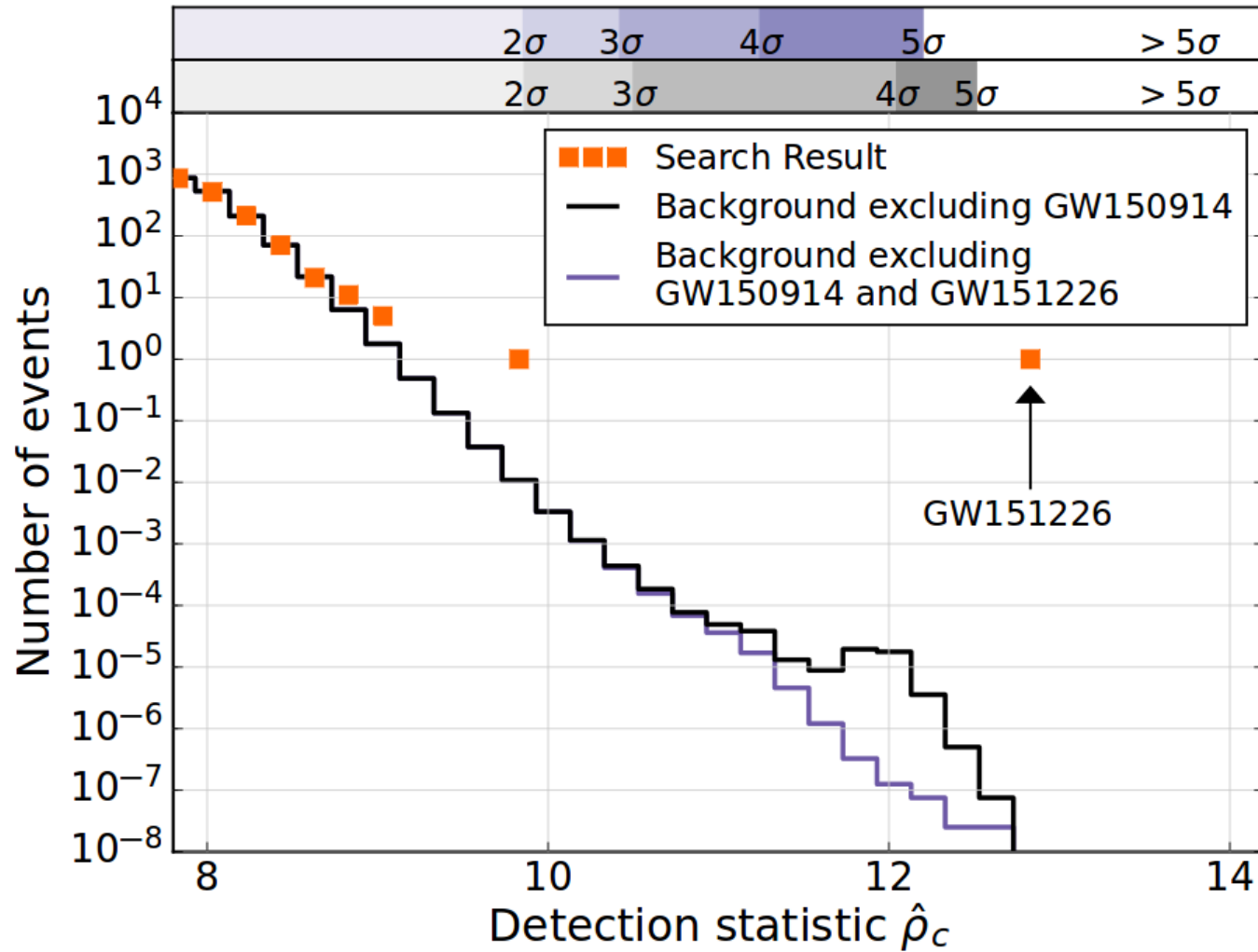
LIGO O1 analysis:

– $O(10^6)$ time offsets

→ background estimated using a fake experiment of $O(100,000)$ years







**Monday September 14, 2015
09:50:45 UTC**

GW



The LIGO detectors are both operational
and stable (O1)

Virgo is off (upgrade in progress)

It is daytime in Europe,
the middle of the night in the US



GW150914 was detected within 3 minutes by a burst search

GraceDB — Gravitational Wave Candidate Event Database

HOME	SEARCH	CREATE	REPORTS	RSS	LATEST	OPTIONS	DOCUMENTATION	AUTHENTICATED AS: FLORENT ROBINET	
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Basic Info

UID	Labels	Group	Pipeline	Search	Instruments	UTC Event Time	FAR (Hz)	Links	UTC Submitted
G184098	H1OK L1OK	Burst	CWB	AllSky	H1,L1	2015-09-14 09:50:45 UTC	1.178e-08	Data	2015-09-14 09:53:51 UTC

Analysis-Specific Attributes

start_time	1126259461	central_freq	123.8285	false_alarm_rate	
start_time_ns	750000000	bandwidth	51.8386	ligo_axis_ra	130.9219
duration	2.477e-02	amplitude	1.410e+01	ligo_axis_dec	4.4808
peak_time	None	snr	23.4521	ligo_angle	None
peak_time_ns	None	confidence		ligo_angle_sig	None

From Marco Drago
 Subject: [detchar] Very Interesting event on ER8
 Reply to: detchar LIGO, Marco Drago
 To: lvc-burst@sympa.ligo.org
 Cc: cbc@ligo.org BinariesGroup, daswg@ligo.org, Calibration <calibration@ligo.org>, dac@sympa.ligo.org, <burst@ligo.org>, 3 more

Hi all,
 cWB has put on gracedb a very interesting event in the last hour.
<https://gracedb.ligo.org/events/view/G184098>

This is the CED:
https://ldas-jobs.ligo.caltech.edu/~waveburst/online/ER8_LH_ONLINE/JOB/112625/1126259540-1126259600/OUTPUT_CED/ced_1126259420_180_1126259540-1126259600_slag0_lag0_1_job1/L1H1_1126259461.750_1126259461.750/

Qscan made by Andy:
https://ldas-jobs.ligo.caltech.edu/~lundgren/wdq/L1_1126259462.3910/
https://ldas-jobs.ligo.caltech.edu/~lundgren/wdq/H1_1126259462.3910/

It is not flag as an hardware injection, as we understand after some fast investigation. Someone can confirm that is not an hardware injection?

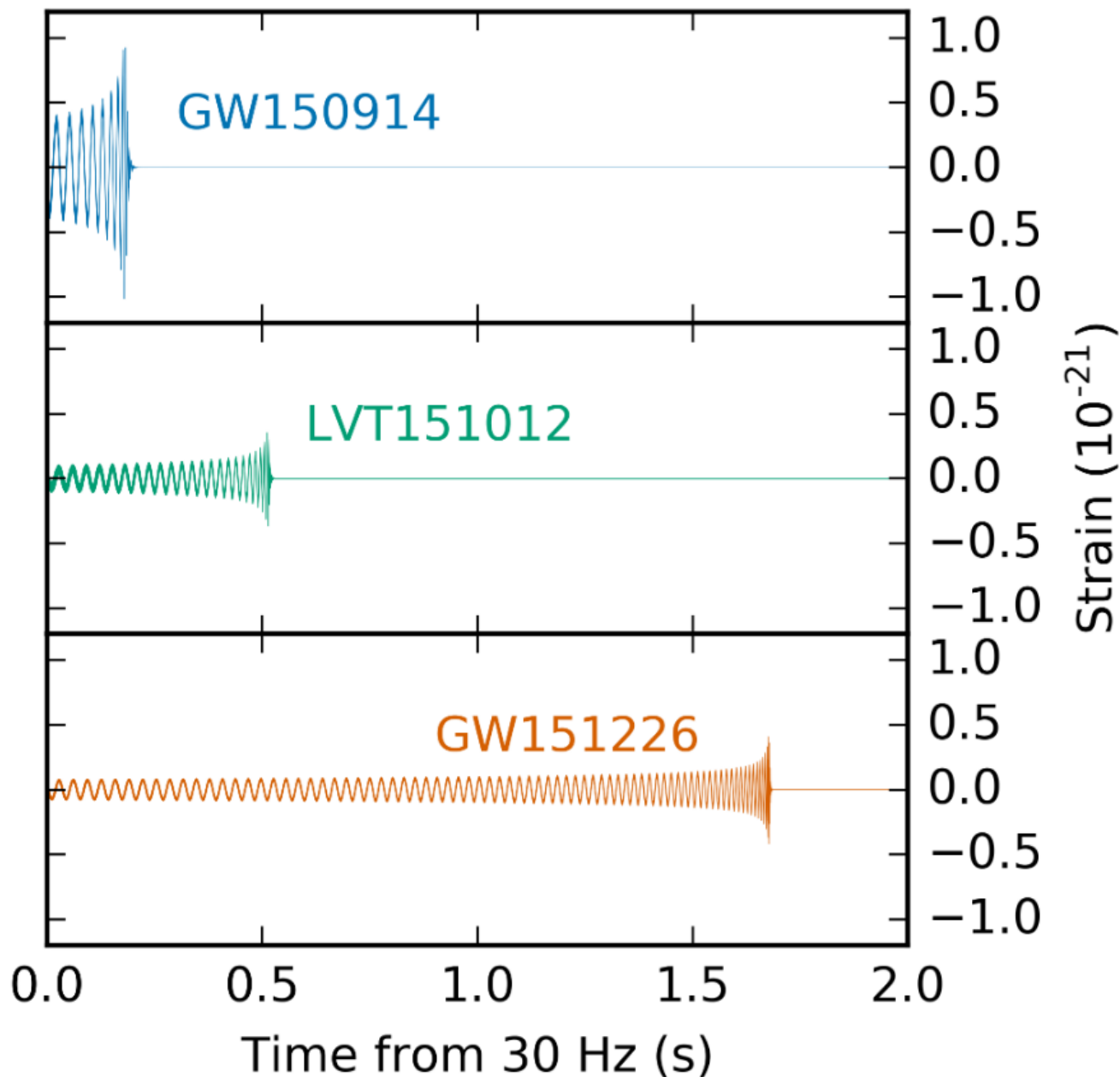
Marco

2+1 events associated to the coalescence of a black hole binary system

Short-duration and
high-amplitude event

Event with low
significance

Long-duration and
low-amplitude event



Full analysis of the data surrounding the event

- only input from searches: time of the event
- fully explore the parameter space
- include calibration uncertainty

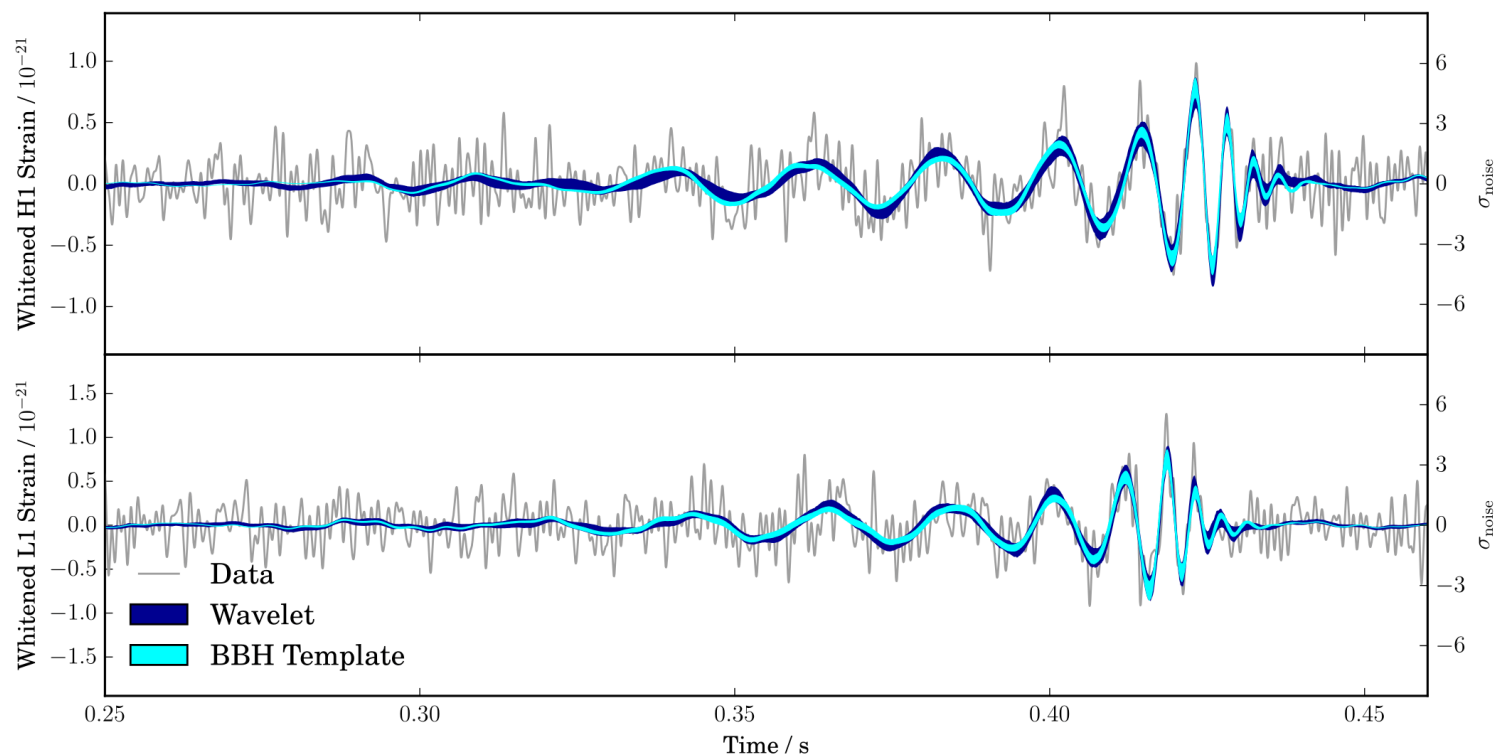
8 intrinsic parameters (masses and spins)

9 extrinsic parameters (distance, position, orientation, coalescence time and phase)

Orbital ellipticity is neglected

Dimensionless spin: $a = \frac{c|\vec{S}|}{Gm^2} \leq 1$

Frequency is redshifted → masses must be rescaled by a factor $(1+z)$

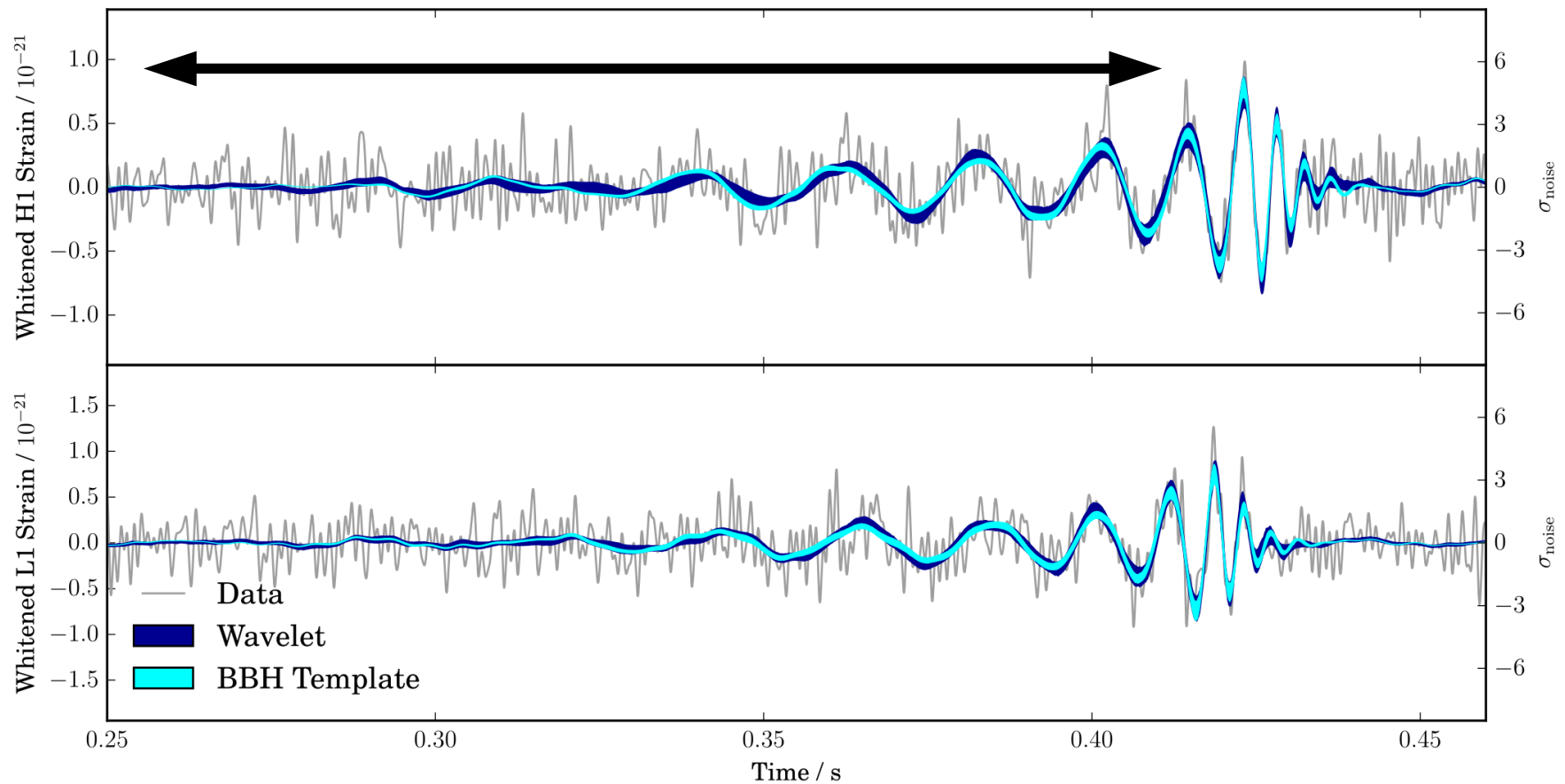


Inspiral phase: PN perturbative expansion (v/c)

Leading order \rightarrow phase evolution driven by the chirp mass
(tight constraints)

Next order \rightarrow m_2/m_1 and spins $\parallel \mathbf{L}$

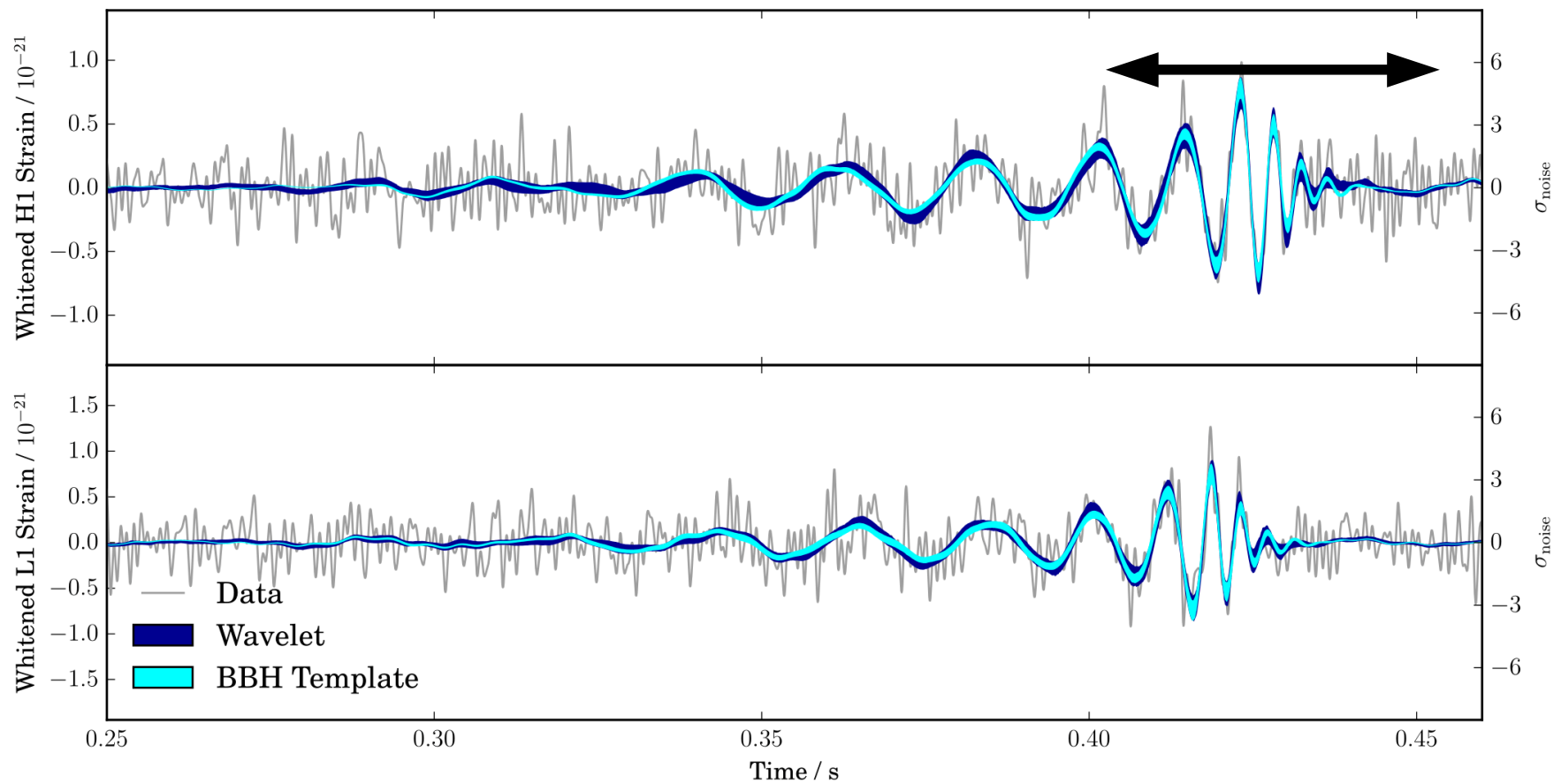
Next orders \rightarrow full spins



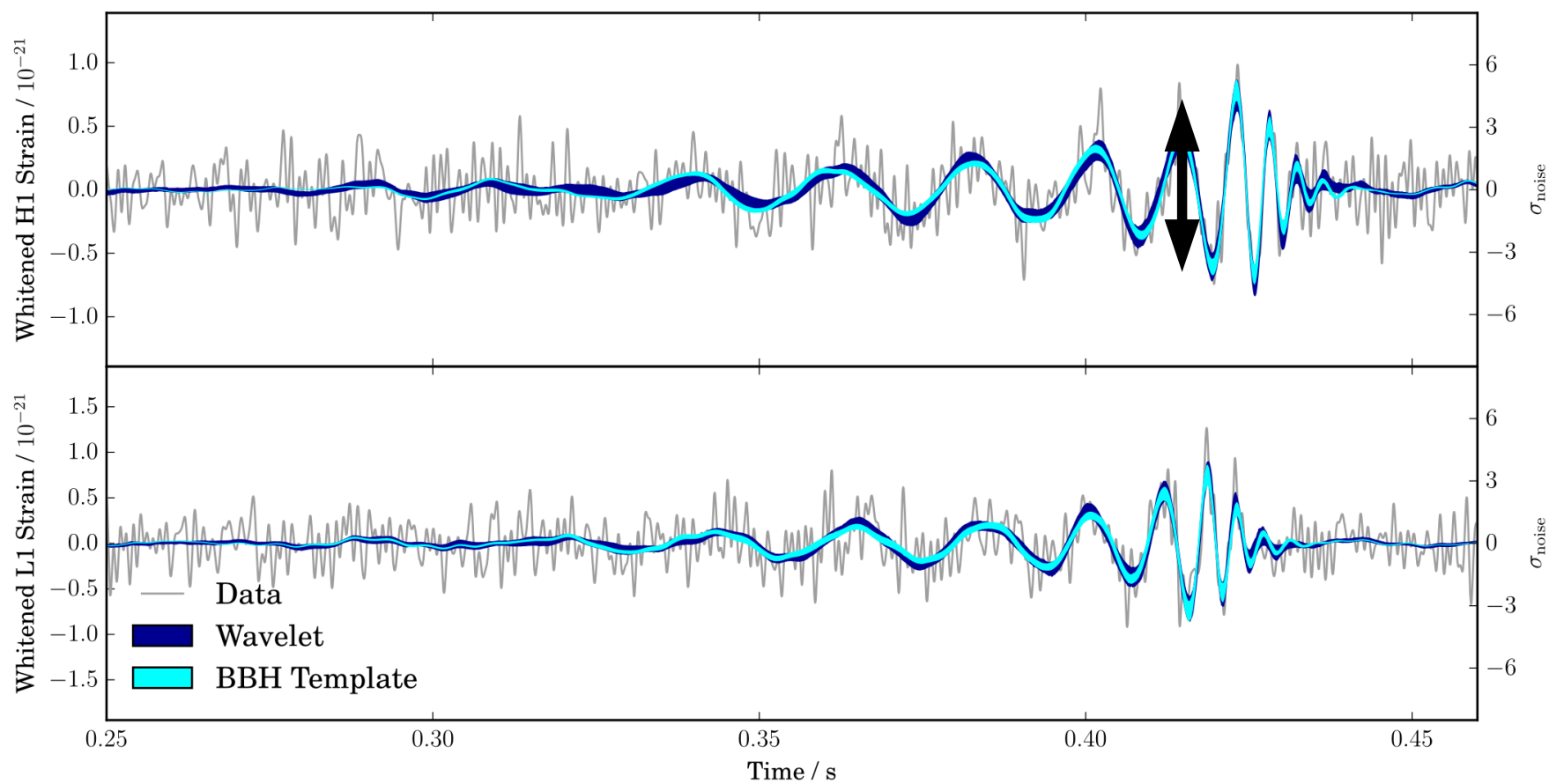
Late inspiral – merger – ringdown: numerical relativity waveforms

Late inspiral → total mass (+chirp mass + m_1/m_2) → individual masses

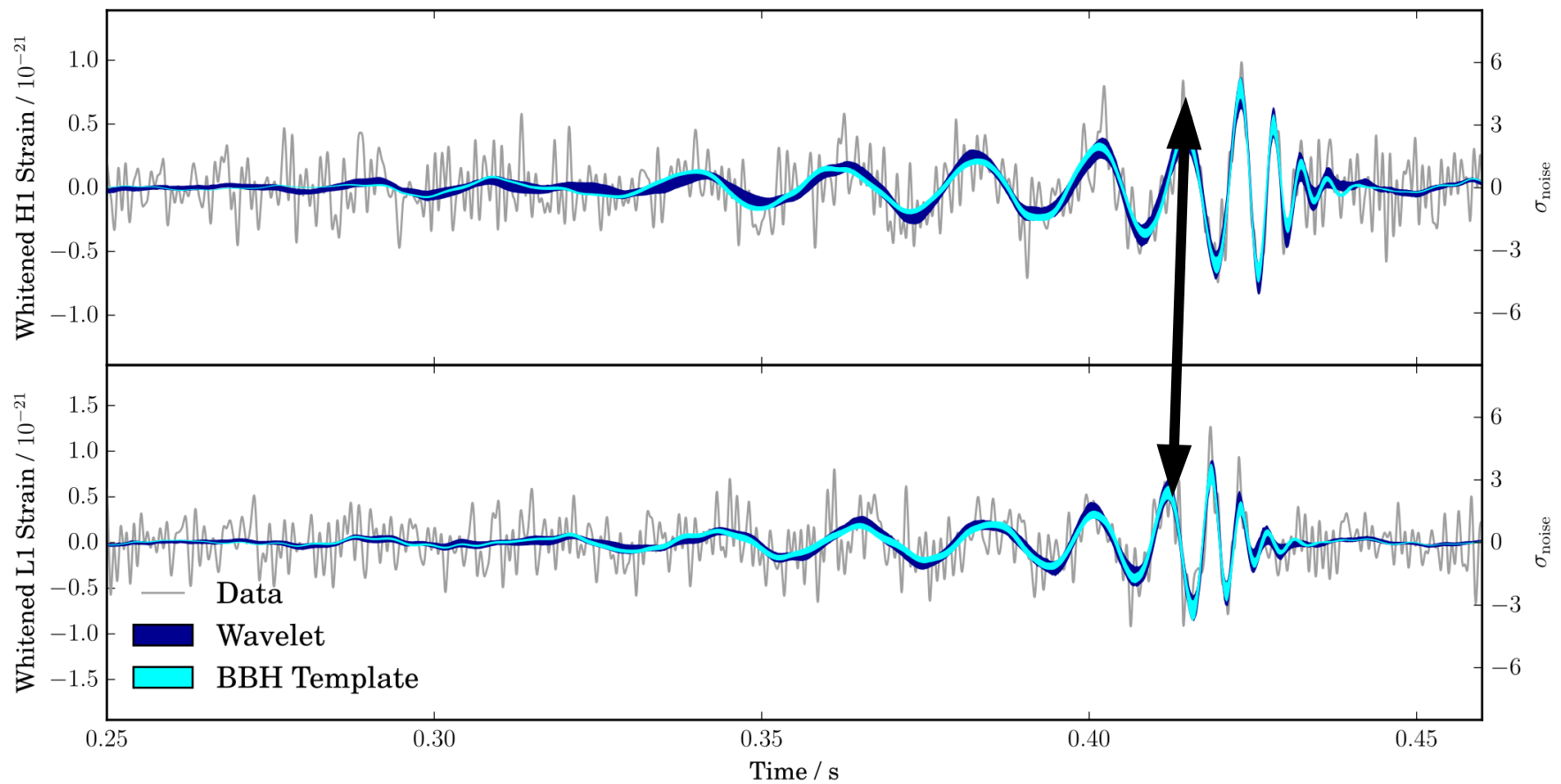
Ringdown → final BH mass and spin

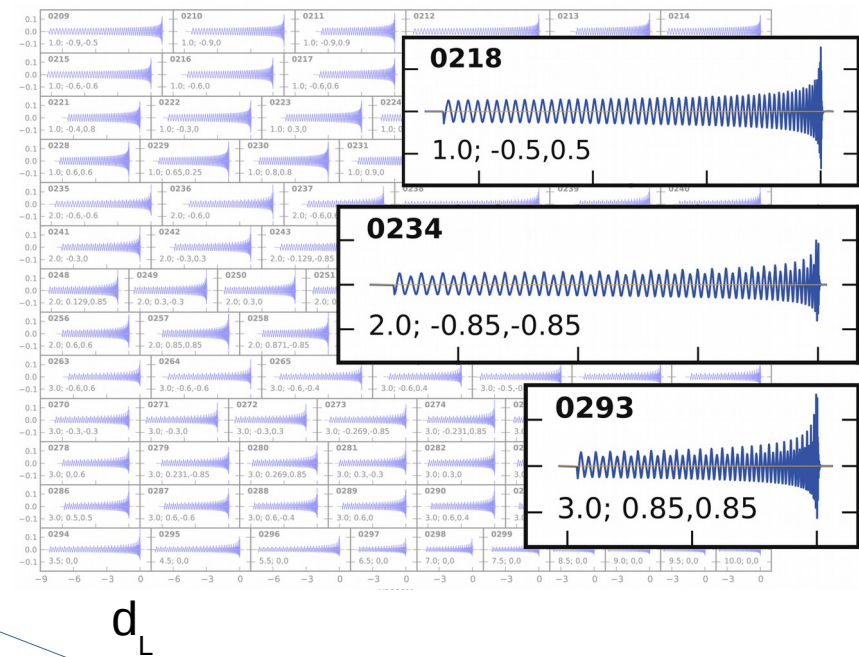
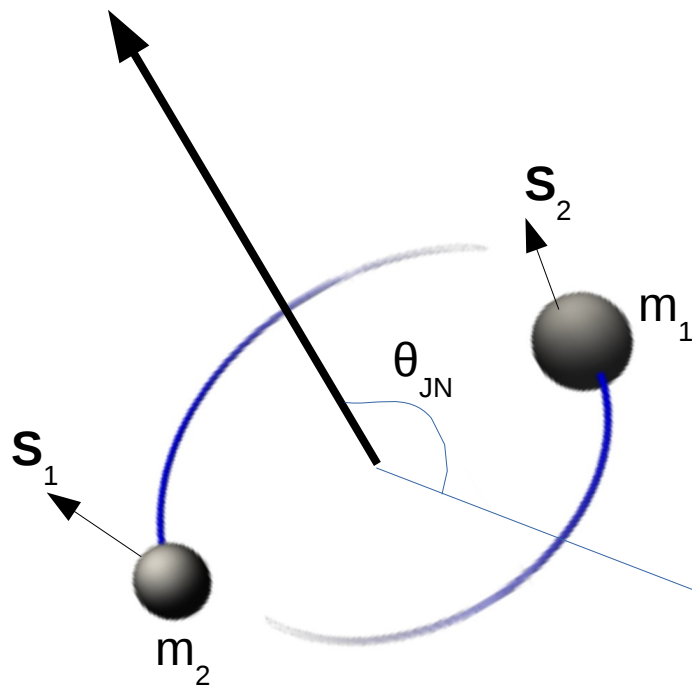


Amplitude: inversely proportional to the distance



Amplitude and phase difference between sites \rightarrow sky location
+ Amplitude and phase consistency





GWs (?)

noise

data from multiple instruments

$$\tilde{d}^{(D)}(f) = \tilde{h}^{(D)}(f) + \tilde{n}^{(D)}(f)$$

$$\vec{d} = \{d^{(H)}, d^{(L)}, \dots\}$$

Bayes theorem:

$$p(\vec{\theta}|\vec{d}, \mathcal{H}, I) = \frac{p(\vec{\theta}|\mathcal{H}, I)p(\vec{d}|\vec{\theta}, \mathcal{H}, I)}{p(\vec{d}|\mathcal{H}, I)}$$

posterior

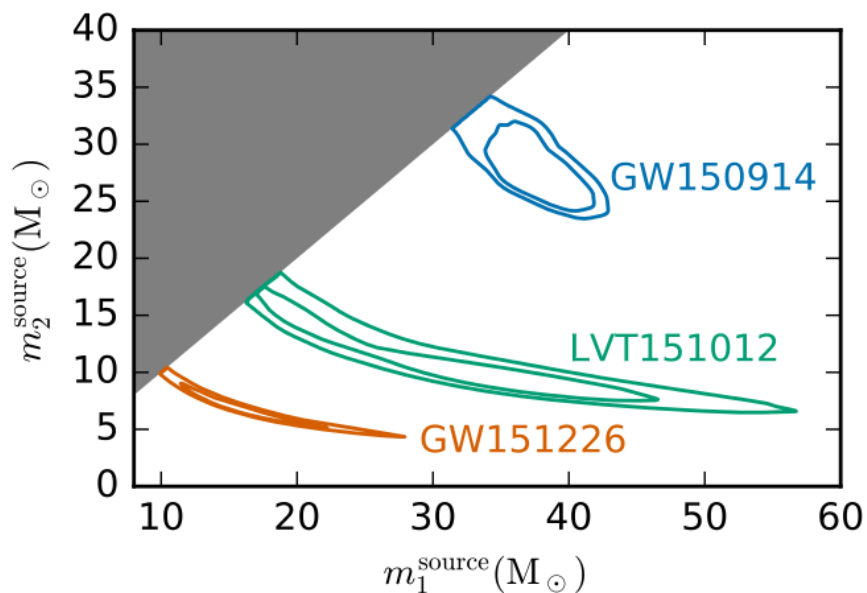
prior

likelihood

Marginalized PDF:

$$p(\vec{\theta}_A|\vec{d}, \mathcal{H}, I) = \int_{\Theta_B} p(\vec{\theta}|\vec{d}, \mathcal{H}, I)d\vec{\theta}_B$$





Mostly sensitive to the chirp mass
 $\rightarrow m_1, m_2$ degeneracy

$$M_c = \frac{(m_1 m_2)^{3/5}}{(m_1 + m_2)^{1/5}}$$

GW150914

$$m_1 = 36.2^{+5.2}_{-3.2} M_{sun}$$

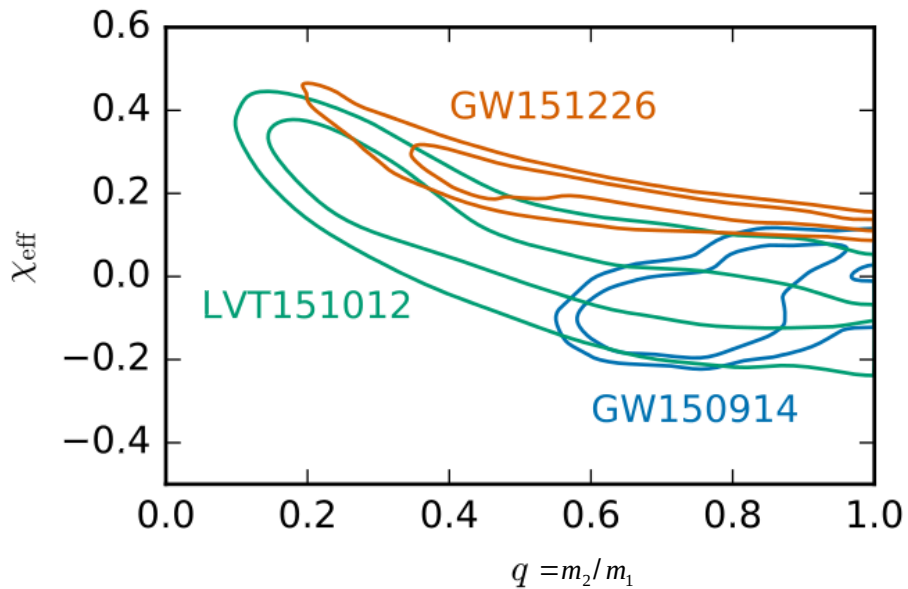
$$m_2 = 29.1^{+3.7}_{-4.4} M_{sun}$$

GW151226

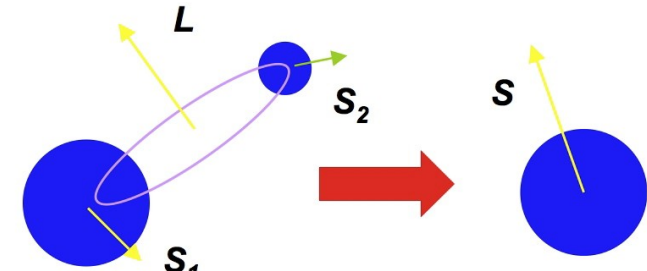
$$m_1 = 14.2^{+8.3}_{-3.7} M_{sun}$$

$$m_2 = 7.5^{+2.3}_{-2.3} M_{sun}$$

- \rightarrow All the components are black holes
- \rightarrow Very high masses for GW150914



$$\chi_{\text{eff}} = \frac{m_1 a_{1z} + m_2 a_{2z}}{m_1 + m_2}$$

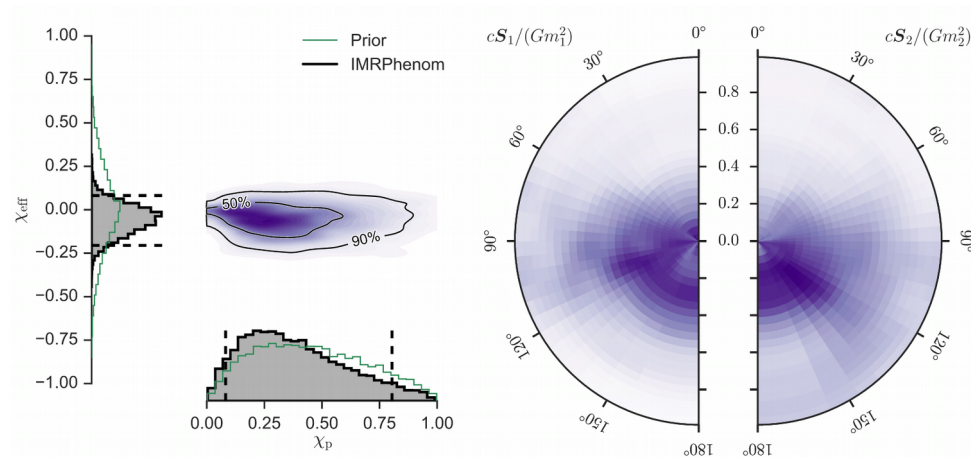


→ not well constrained

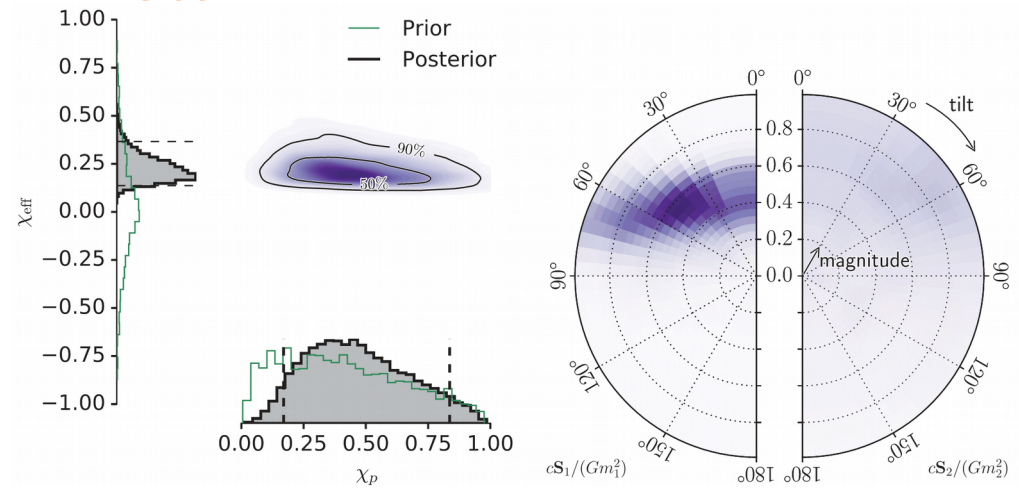
GW151226: at least one black hole is a Kerr black hole
spin > 0.2

Uninformative about precession

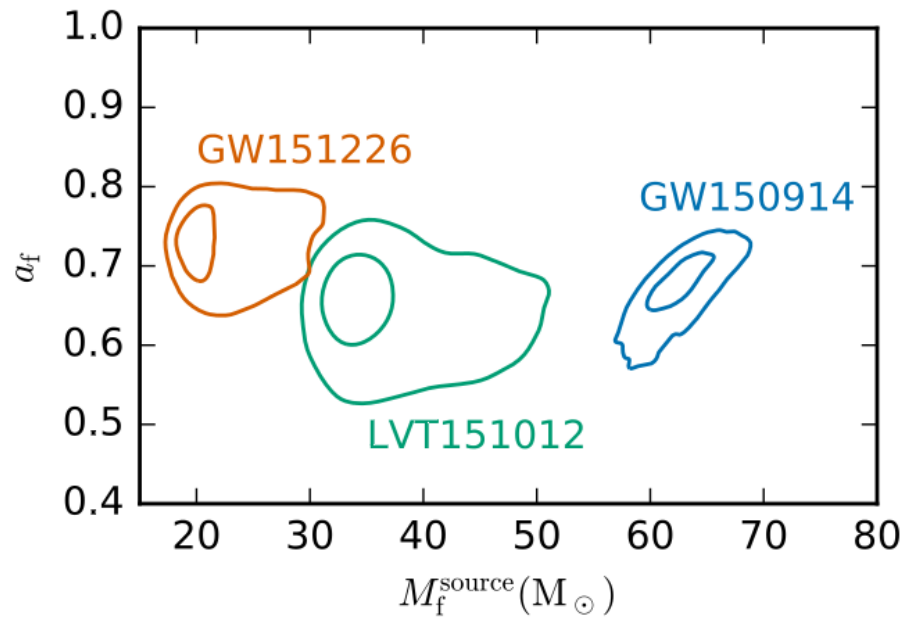
GW150914



GW151226



Final mass & spin

**GW150914**

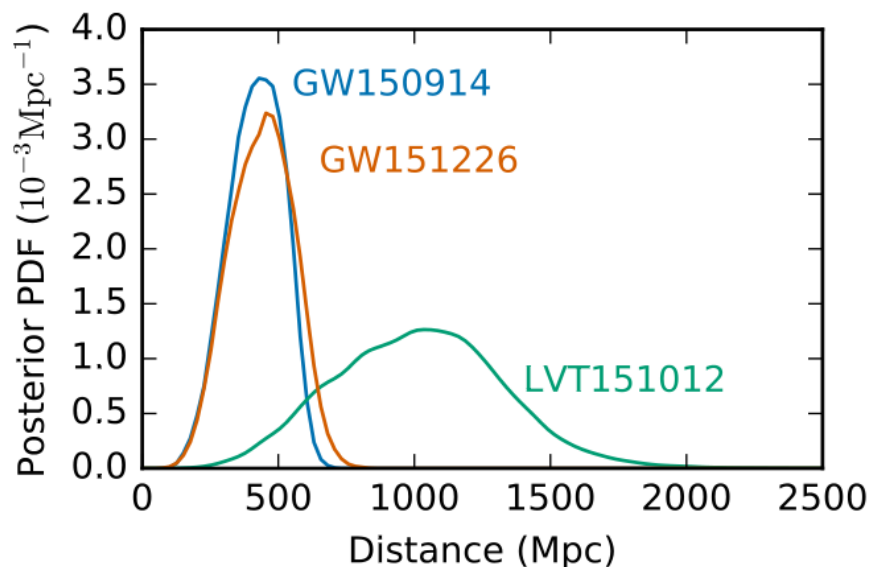
$$M_f = 62.3_{-3.1}^{+3.7} M_{\text{sun}}$$

$$a_f = 0.68_{-0.06}^{+0.05}$$

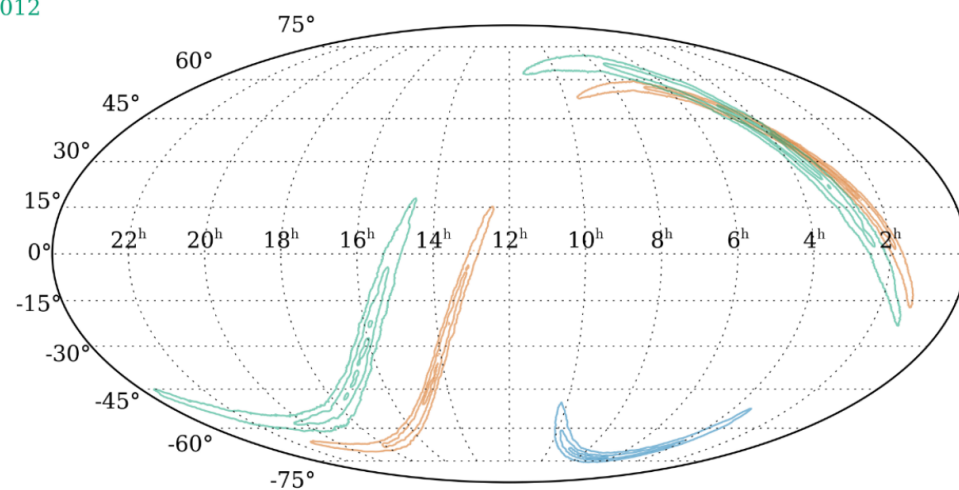
GW151226

$$M_f = 20.8_{-1.7}^{+6.1} M_{\text{sun}}$$

$$a_f = 0.74_{-0.06}^{+0.06}$$



GW150914
 GW151226
 LVT151012



90% credible region for sky location:

→ GW150914 = 230 deg²

→ GW151226 = 850 deg²

GW150914
 $D_L = 420_{-180}^{+150} \text{ Mpc} \quad z = 0.09_{-0.04}^{+0.03}$

GW151226
 $D_L = 440_{-190}^{+180} \text{ Mpc} \quad z = 0.09_{-0.04}^{+0.03}$

Limited accuracy with 2 detectors
 → will be improved with a 3rd detector (a few deg²)

(Lambda-CDM cosmology)

- Simulated signals with parameters drawn from astrophysical populations
- Noise distribution from GW searches

Mass distribution	$R/(\text{Gpc}^{-3}\text{yr}^{-1})$		
	PyCBC	GstLAL	Combined
Event based			
GW150914	$3.2^{+8.3}_{-2.7}$	$3.6^{+9.1}_{-3.0}$	$3.4^{+8.6}_{-2.8}$
LVT151012	$9.2^{+30.3}_{-8.5}$	$9.2^{+31.4}_{-8.5}$	$9.4^{+30.4}_{-8.7}$
GW151226	35^{+92}_{-29}	37^{+94}_{-31}	37^{+92}_{-31}
All	53^{+100}_{-40}	56^{+105}_{-42}	55^{+99}_{-41}
Astrophysical			
Flat	31^{+43}_{-21}	30^{+43}_{-21}	30^{+43}_{-21}
Power Law	100^{+136}_{-69}	95^{+138}_{-67}	99^{+138}_{-70}

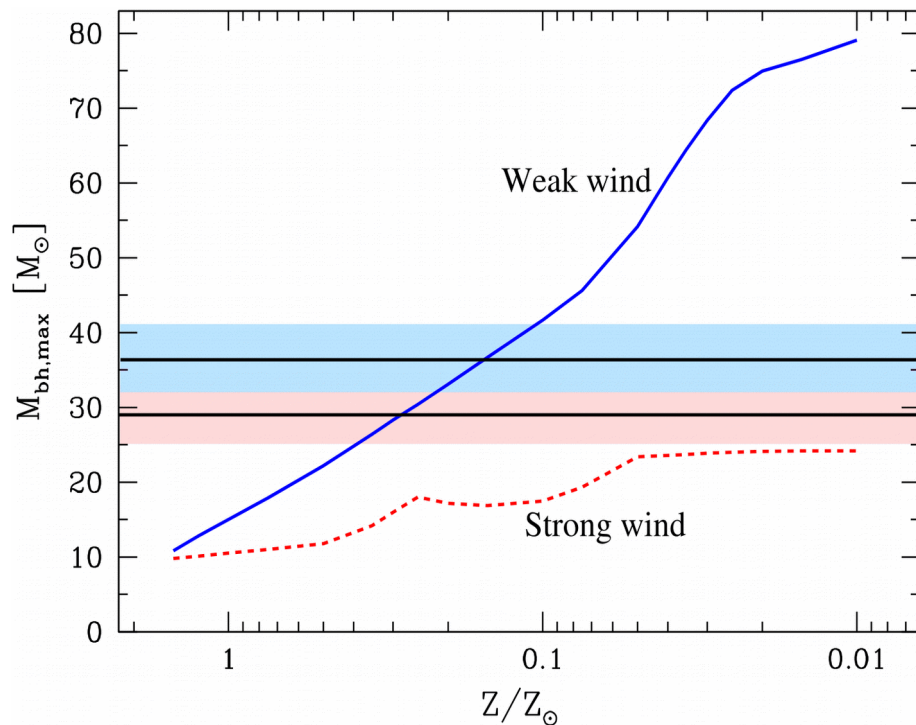
$p(m_1, m_2) \propto m_1^{-1} m_2^{-1}$
 ~underestimates the rate

$p(m_1) \propto m_1^{-2.35}$
 $p(m_2) = \text{cte}$
 ~overestimates the rate

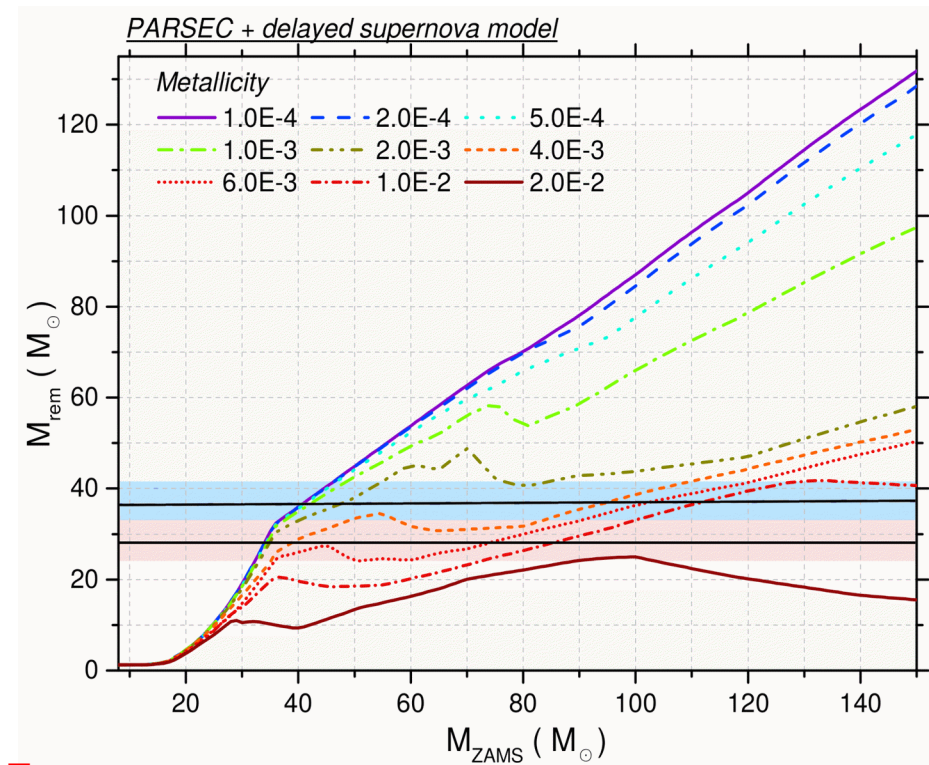
$$R = 9 - 240 \text{ Gpc}^{-3} \text{ yr}^{-1}$$

- Better understand BH and BBH formation scenarios.
- Testing GR in strong field.
- Core collapse supernovae explosion mechanisms.
- New standard sirens to measure the Universe expansion.
- Fundamental physics test:
 - EOS in neutron stars
 - Graviton mass, neutrino masses, ...
 - Primordial BH as dark matter
 - Primordial GWs and inflation

- First BBH system ever observed & heaviest stellar mass black holes (>25 Msun).



High mass stellar BH \rightarrow low metallicity $Z < \frac{1}{2} Z_{sun}$
 \rightarrow weak massive-star winds



- BBH formation: isolated binaries (low-Z to popIII) vs capture in dense clusters (globular clusters, galactic centers, ...): no way to discriminate between the 2 scenarios with 4 BBHs.

First opportunity to study GR in a strong-field regime

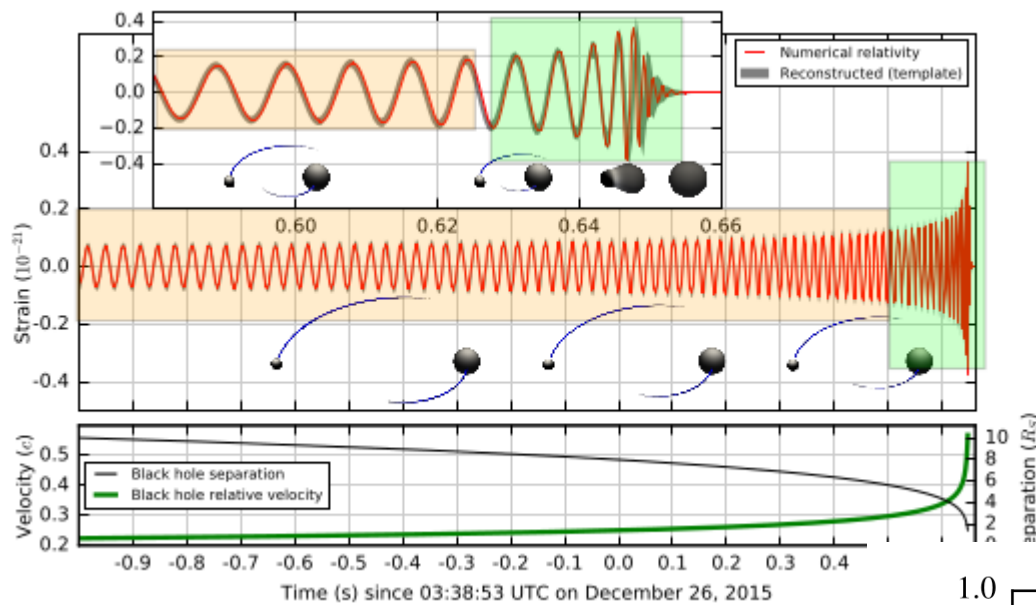
Test #1 → signal waveform/GR consistency: residual compatible with noise

Test #2 → BBH parameter consistency before/after merger: excellent

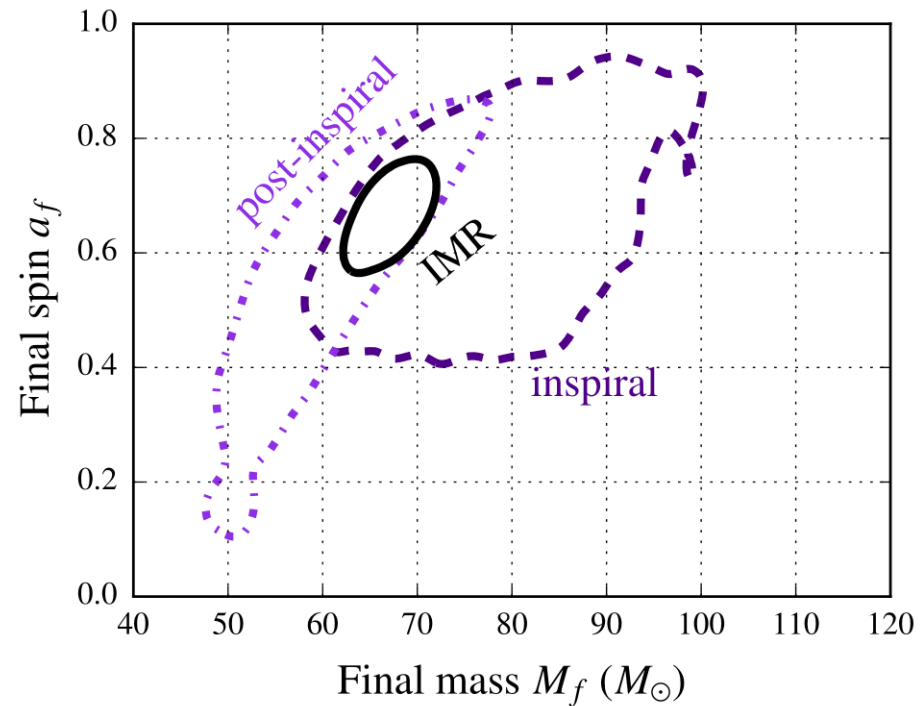
Test #3 → deviation from PN waveforms: constraints on PN coefficients

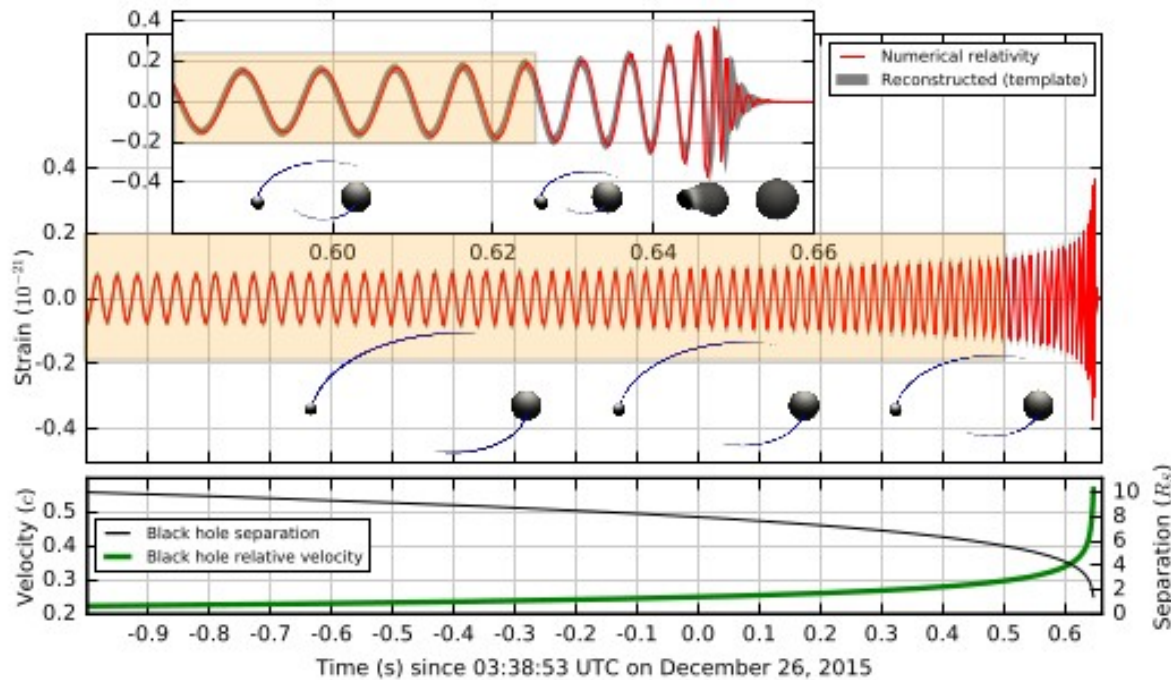
Test #4 → consistency with the least-damped quasi-normal-mode of the remnant black hole

Test #5 → theory with massive graviton: best constraints on the graviton mass



Verify self-consistency by comparing final mass and spin predicted from the “inspiral” and from the “post-inspiral”
 [Ghosh+, 2016]
 [LVC PRL(2016)]



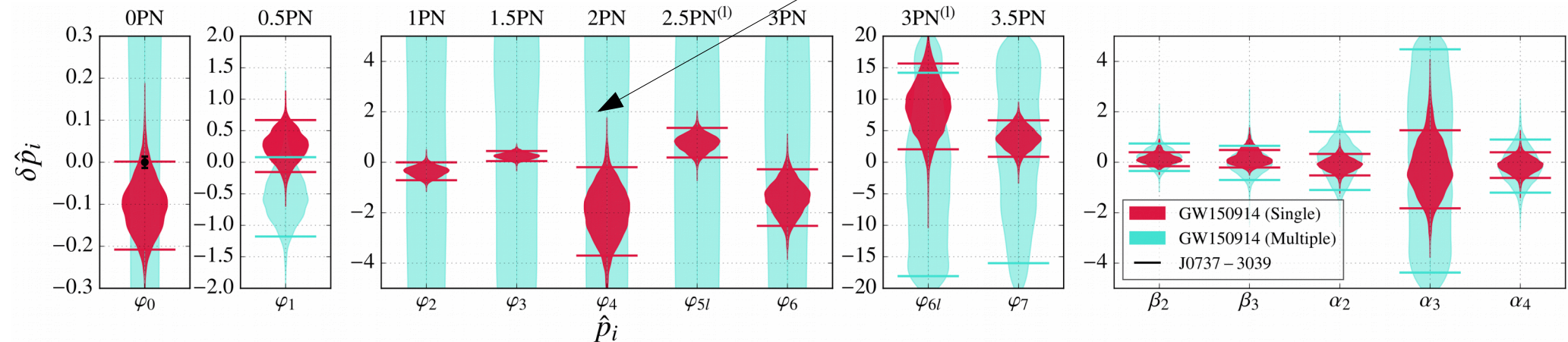


- **Waveform:**

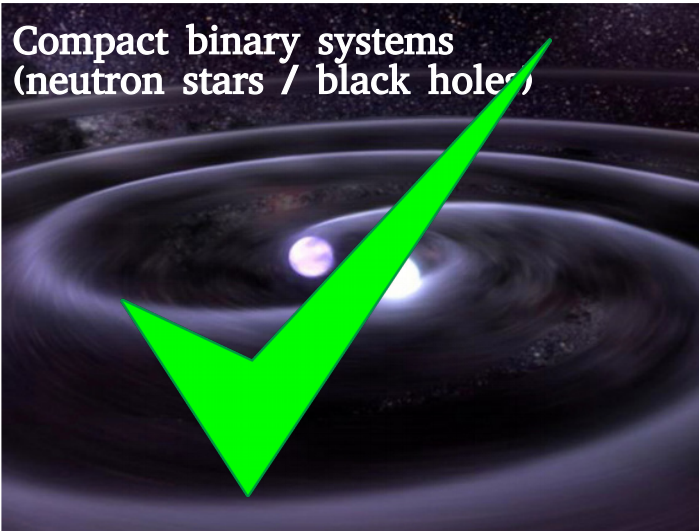
$$h(f, \theta) = A(f; \theta)e^{i\phi(f; \theta)},$$
- $$\phi = \phi_o + \sum \phi_k(\theta)(\pi M f)^{(k-5)}$$

$$\theta = \{m_1, m_2, s_1, s_2\}$$
- $$\phi_k = \phi_k^{GR}(1 + \delta\phi_k)$$

[LVC PRL(2016), PRX(2016), PRL(2017)]



Compact binary systems
(neutron stars / black holes)



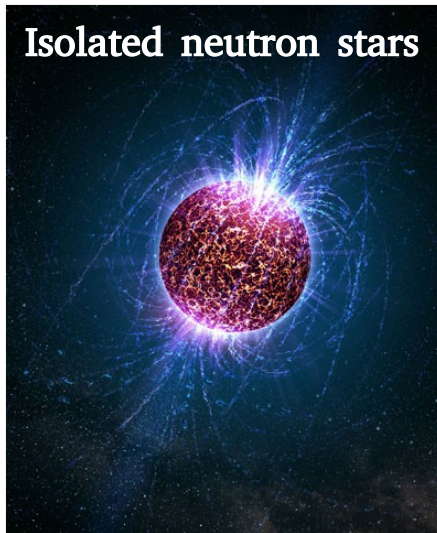
newly-formed black holes



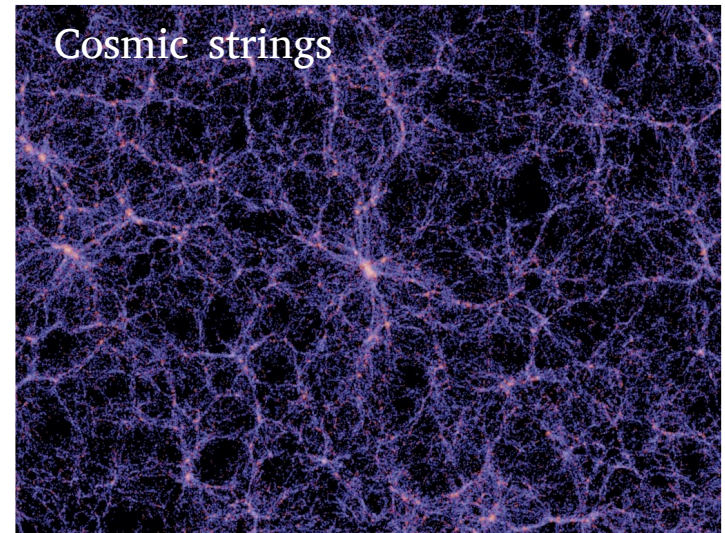
Supernova core collapse



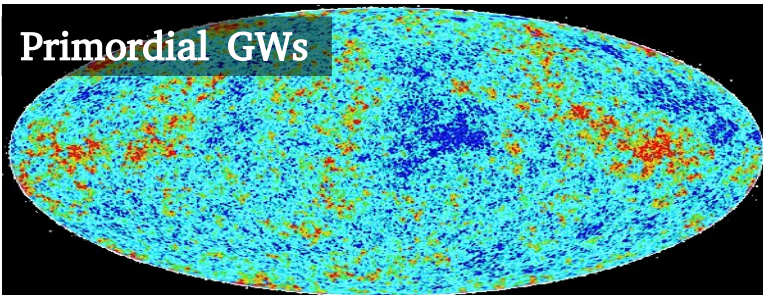
Isolated neutron stars



Cosmic strings



Primordial GWs

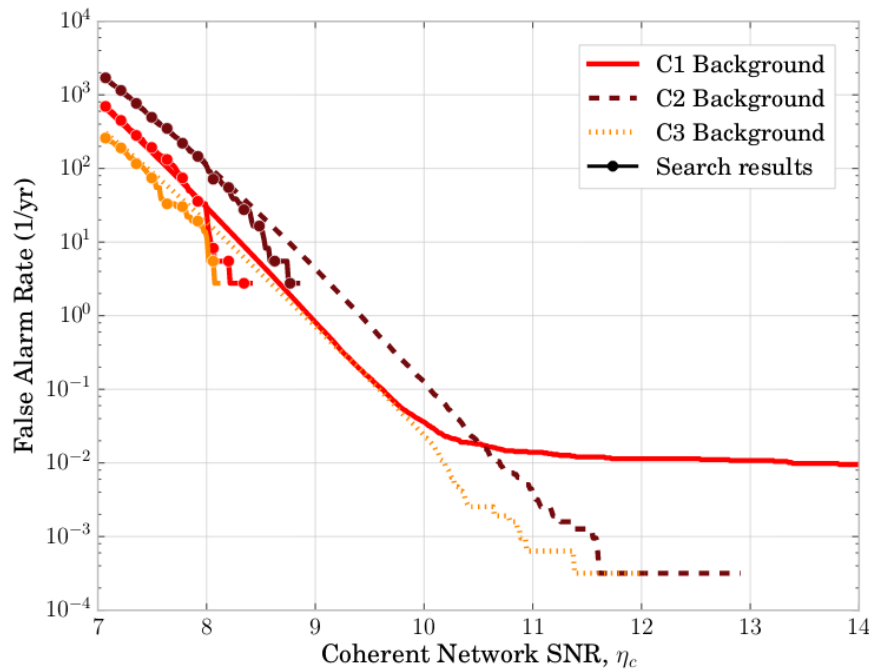


Unknown sources

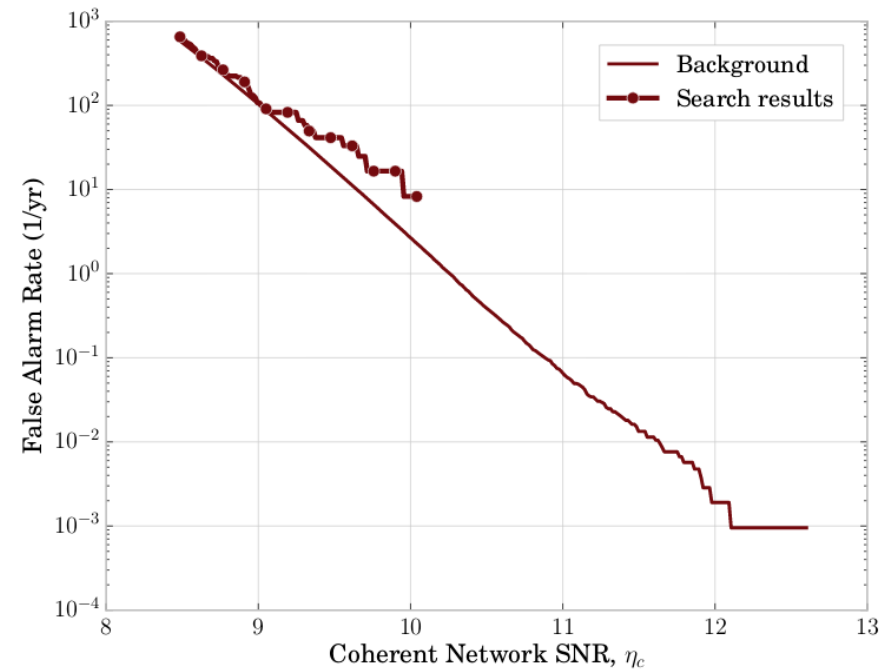


Target : transient signals with duration ranging from milliseconds to seconds over the frequency band of 32 to 4096 Hz

- compact binary mergers
- core-collapse supernovae
- neutron stars collapsing to form black holes
- pulsar glitches
- cosmic string cusps



(a) **cWB** 32-1024 Hz search classes: *C1* (red), *C2* (brown), *C3* (yellow).



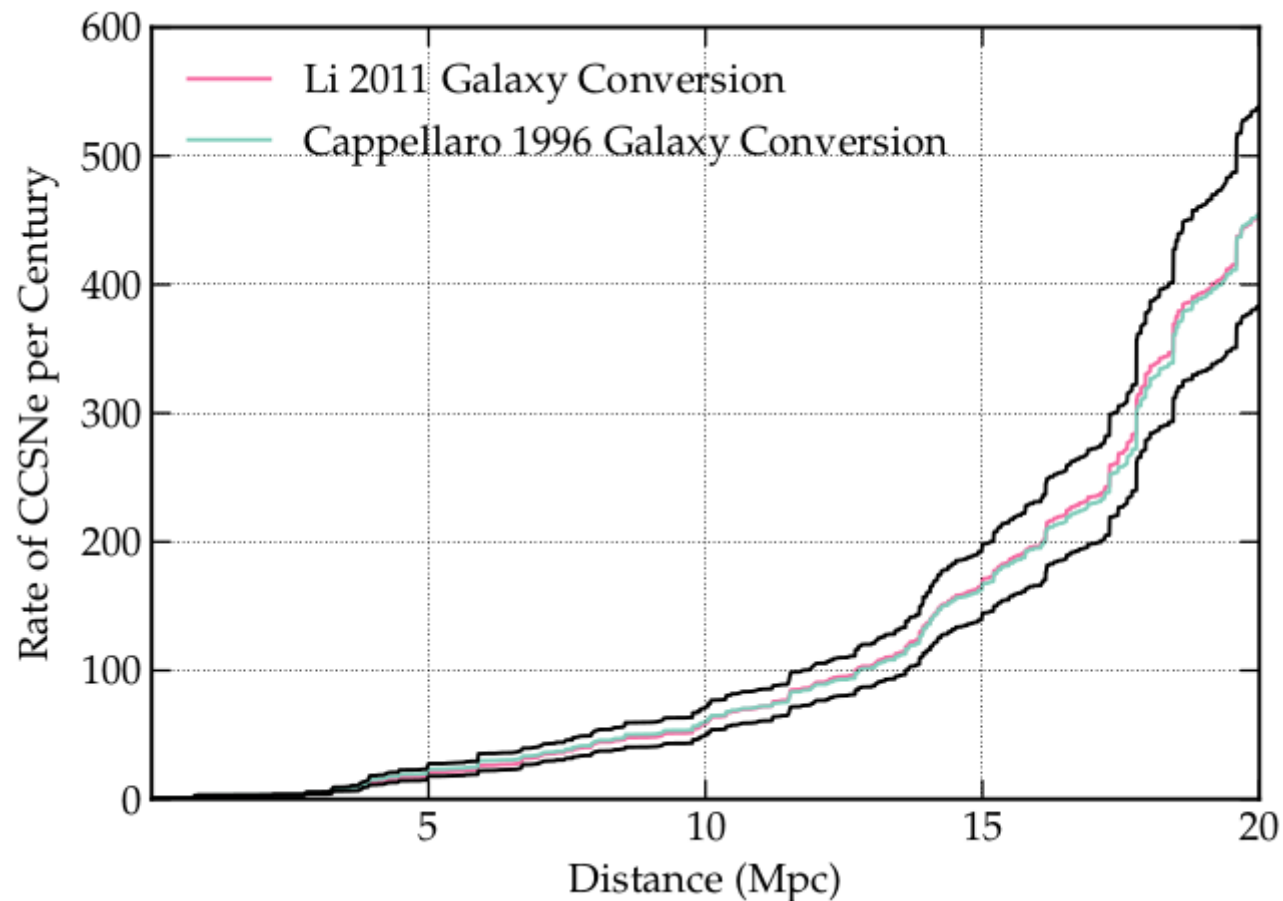
(b) **cWB** 1024-4096 Hz search class.

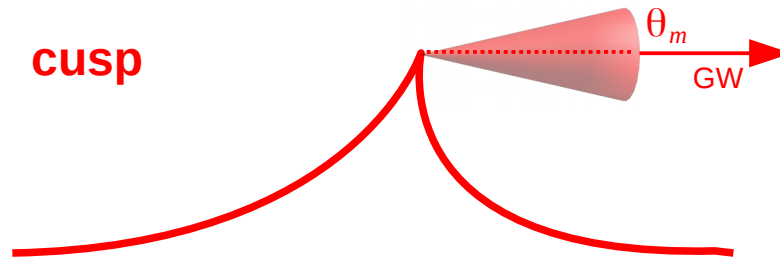
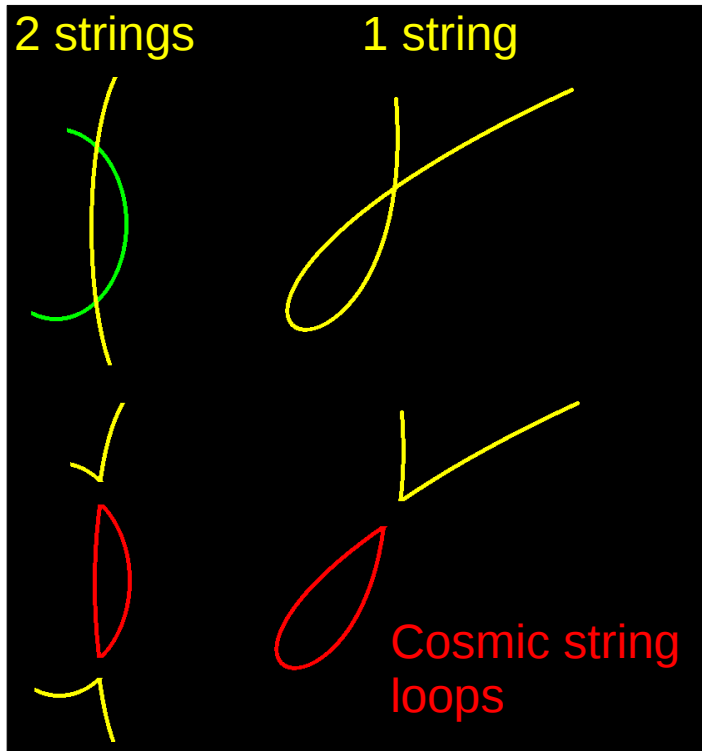
→ the search results are consistent with the expectations of accidental noise coincidences

With advanced LIGO and advanced Virgo:

Distance: between **100 kpc** (SASI and MHD) and **20 Mpc** (extreme model like disk fragmentation and bar mode) [Gossan et al arxiv:1511.02836]

Rate : [J. Gill et al in preparation]





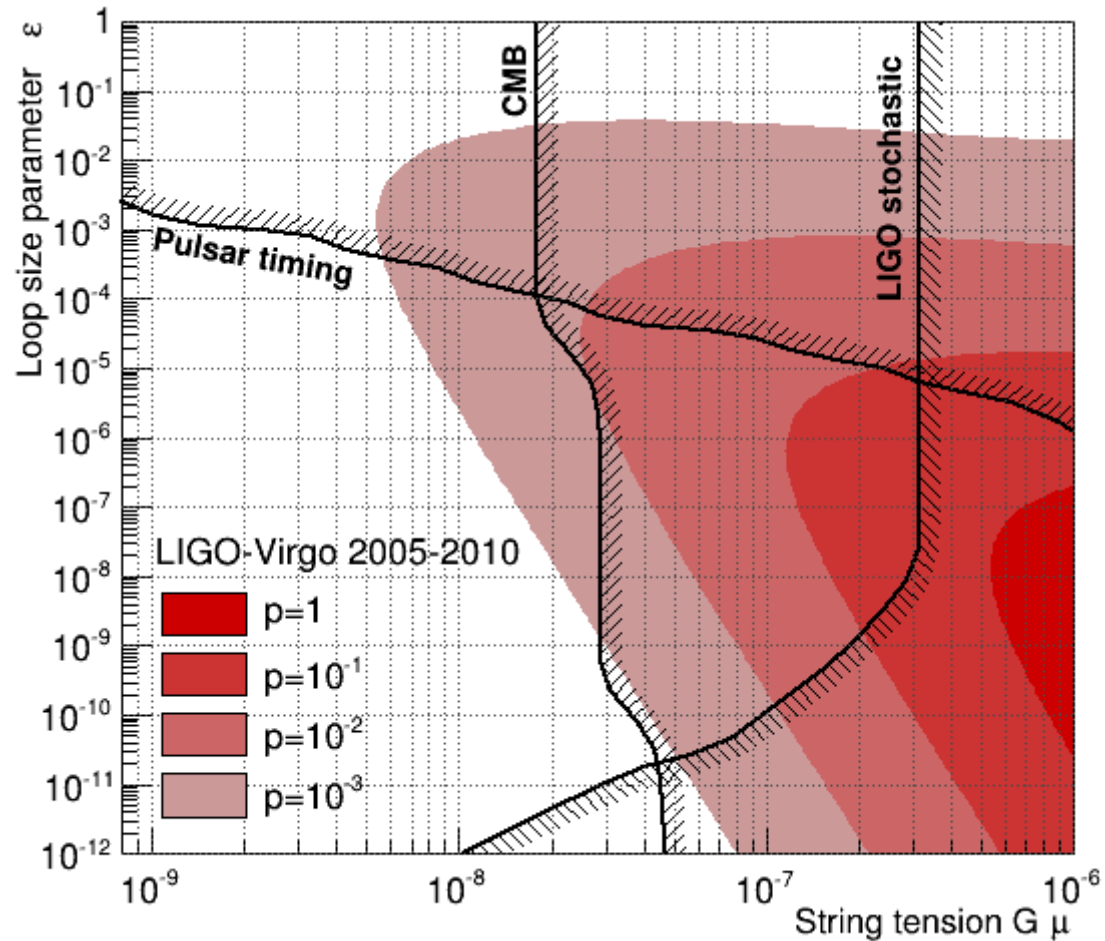
GW waveform:

$$h(\ell, z, f) = A_q(\ell, z) f^{-q} \Theta(f_h - f)$$

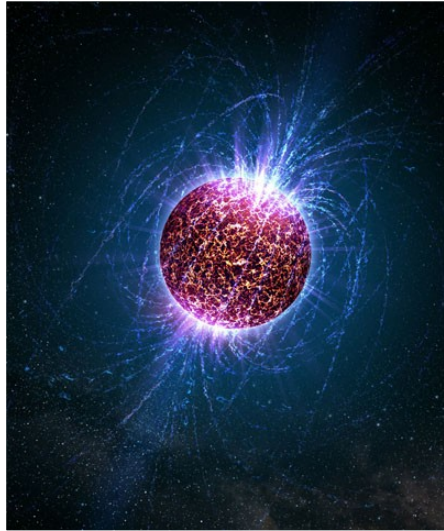
$$A_q(\ell, z) = g_1 \frac{G\mu\ell^{2-q}}{(1+z)^{q-1}r(z)}$$

$q = 4/3$ for cusps, $q = 5/3$ for kinks

+ model for loop distribution



[Enhanced LIGO/Virgo results]

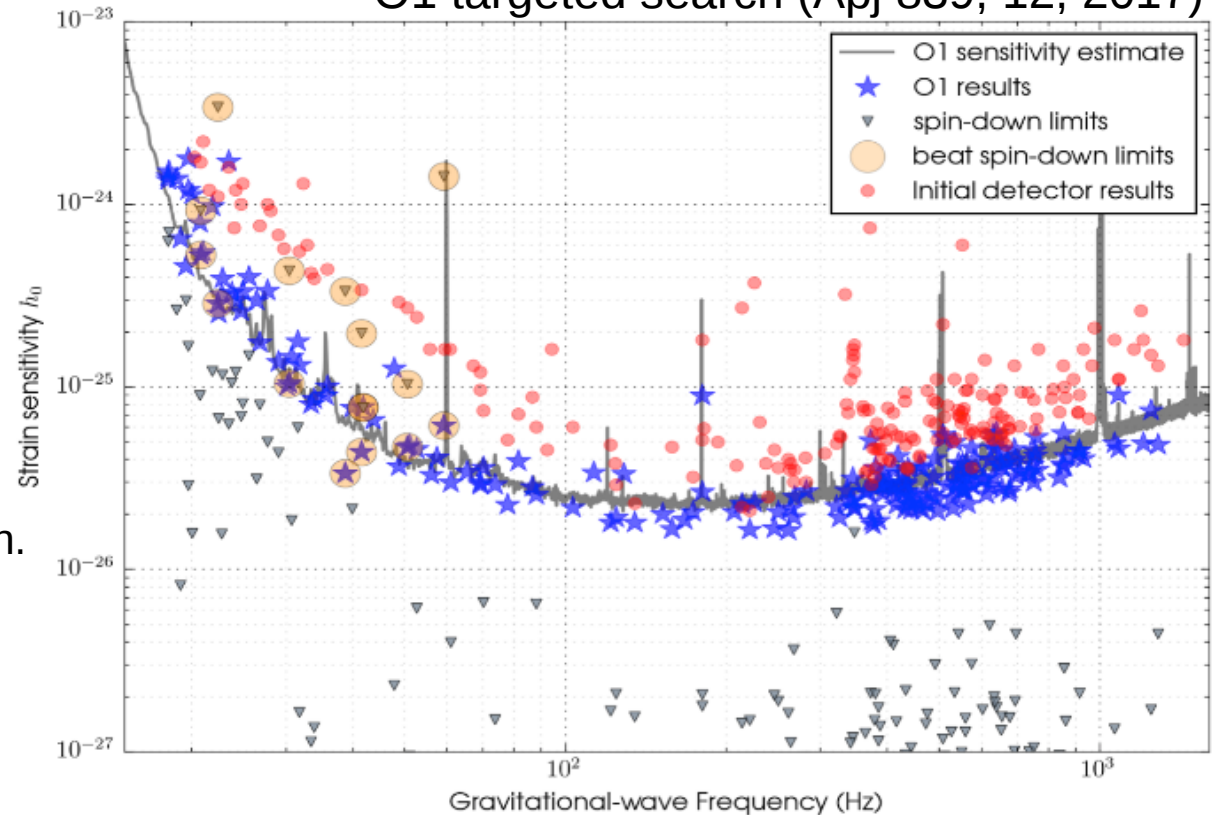


Signal amplitude:
$$h_0 \cong 10^{-27} \left(\frac{I_{zz}}{10^{38} \text{ kg} \cdot \text{m}^2} \right) \left(\frac{10 \text{ kpc}}{r} \right) \left(\frac{f}{100 \text{ Hz}} \right)^2 \left(\frac{\varepsilon}{10^{-6}} \right)$$

ε : ellipticity (adimensional number measuring the star's degree of asymmetry)
 f : signal frequency, proportional to star rotation frequency

- Persistent signals.
- Weak amplitudes (ellipticity unknown).
- Known sources in the galaxy.
- Multi-messenger analysis with radio telescope inputs.
- All-sky/targeted searches
- 8 known pulsars spin down limit beaten.

O1 targeted search (Apj 839, 12, 2017)



Assumption : stationary, unpolarized, and Gaussian stochastic background

→ Cross correlate the output of detector pairs to eliminate the noise

$$h_i = n_i + GW_i$$

$$\langle h_1, h_2 \rangle = \langle GW_1, GW_2 \rangle + \underbrace{\langle n_1, GW_2 \rangle}_0 + \underbrace{\langle GW_1, n_2 \rangle}_0 + \underbrace{\langle n_1, n_2 \rangle}_0$$

With $\langle x_1, x_2 \rangle = \int_{-\infty}^{+\infty} \tilde{x}_1^*(f) \tilde{Q}(f) \tilde{x}_2(f) df$

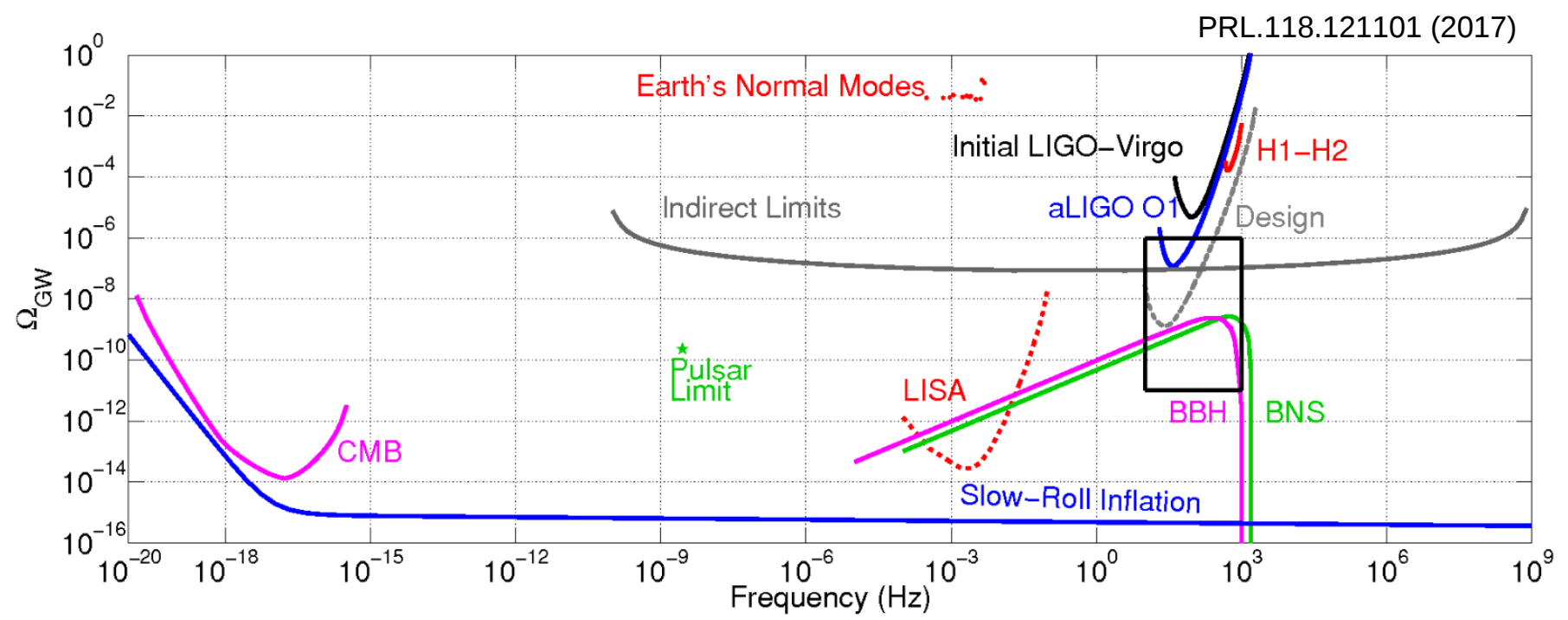
Optimal filter:

$$\tilde{Q}(f) \propto \frac{\gamma(f) \Omega_{GW}(f)}{f^3 S_{n,1}(f) S_{n,2}(f)}$$

← overlap of antenna pattern
 ← GW spectrum $\Omega_{GW}(f) = \Omega_\alpha f^\alpha$
 ↙ ↘
 Detector PSDs

O1 isotropic search, for $\alpha = 0$: $\Omega_{GW}(25 \text{ Hz}) < 1.7 \times 10^{-7}$

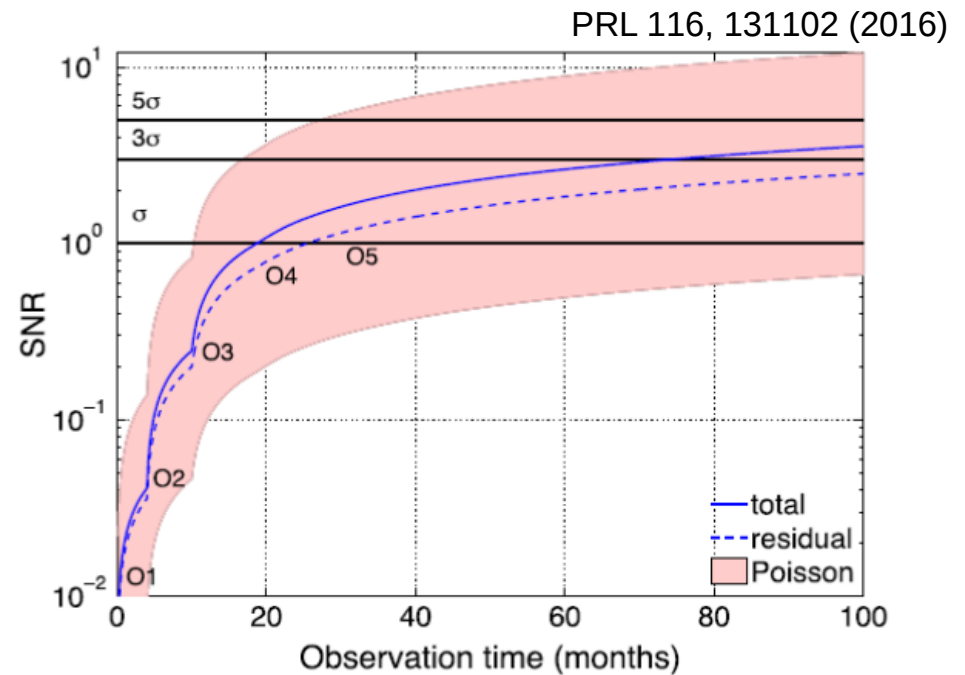
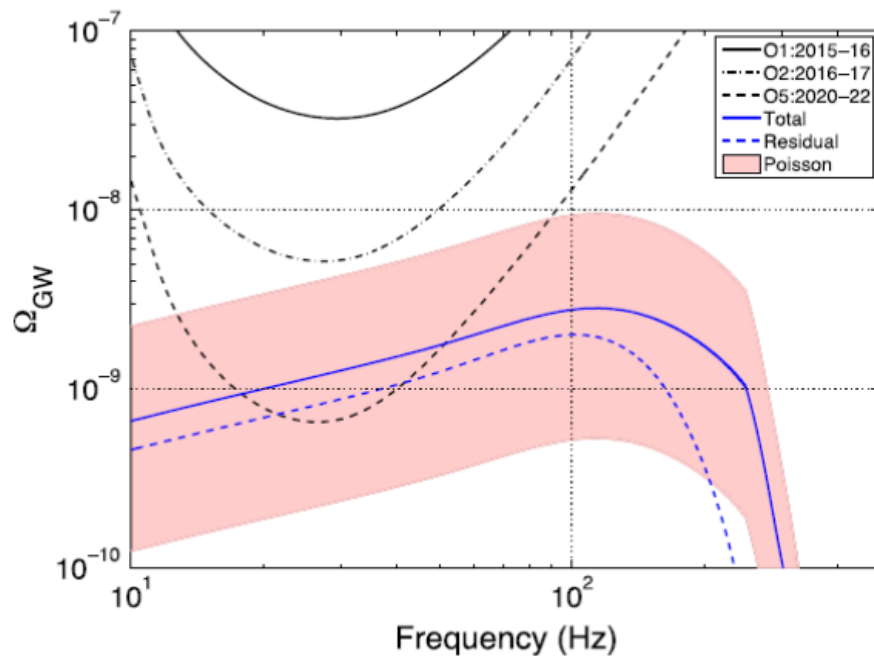
O1 isotropic search, for $\alpha=0$: $\Omega_{GW}(25\text{ Hz}) < 1.7 \times 10^{-7}$

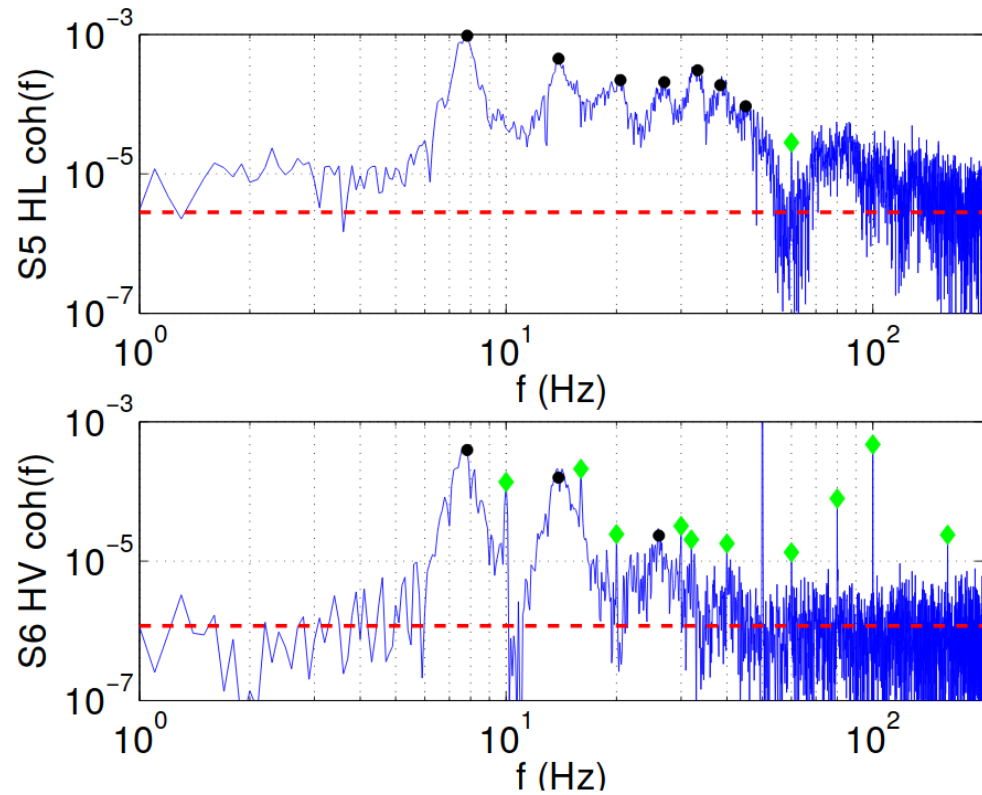


$$\Omega_{\text{GW}}(f; \theta_k) = \frac{f}{\rho_c H_0} \int_0^{z_{\text{max}}} dz \frac{R_m(z, \theta_k) \frac{dE_{\text{GW}}}{df_s}(f_s, \theta_k)}{(1+z)E(\Omega_M, \Omega_\Lambda, z)}$$

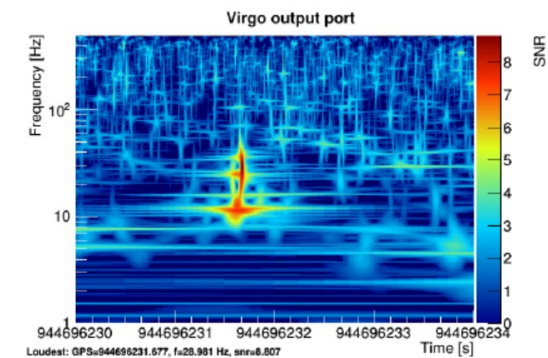
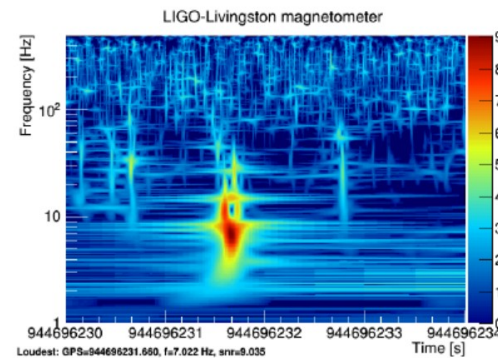
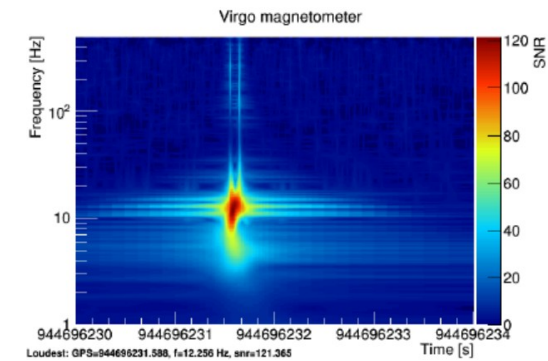
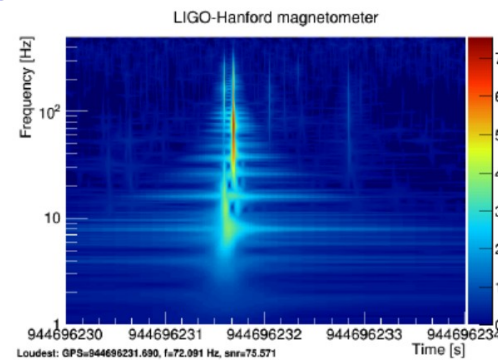
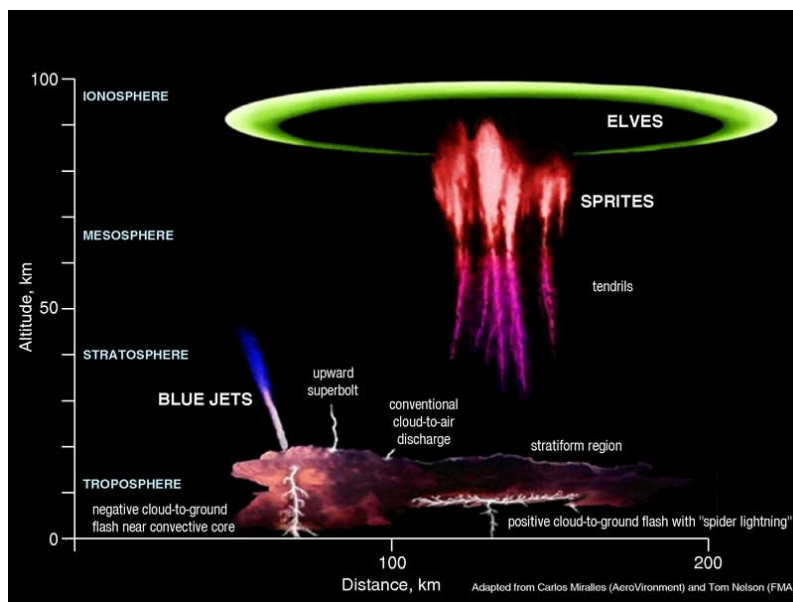
Merger Rate
Source Energy Spectrum

↑
Cosmology

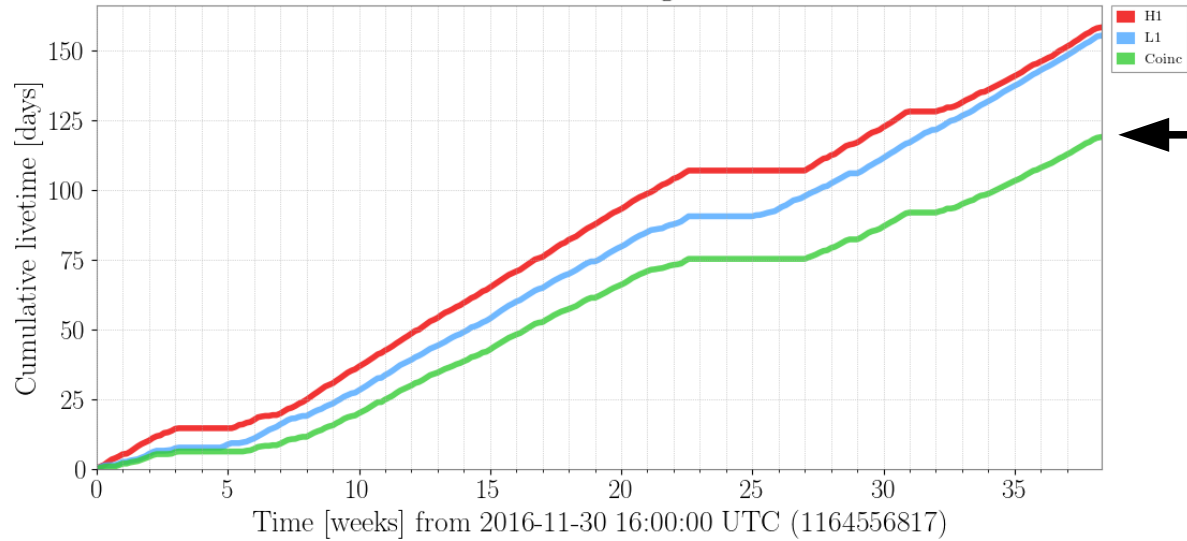




Schumann resonance

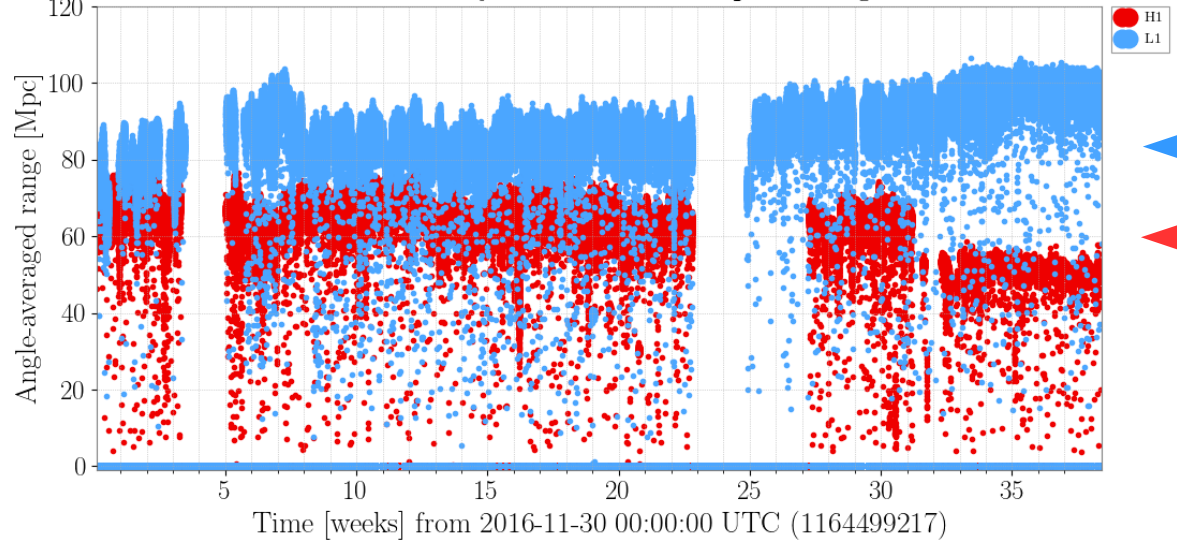


LIGO observing time



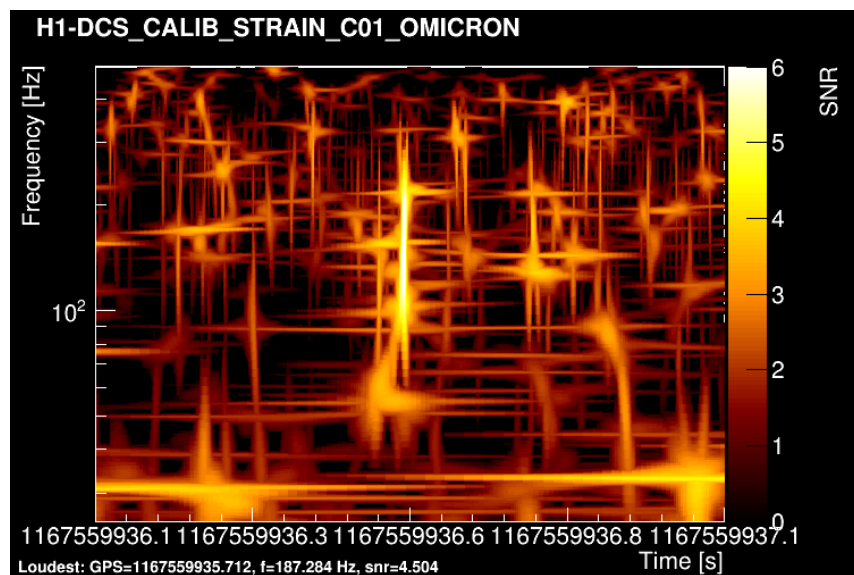
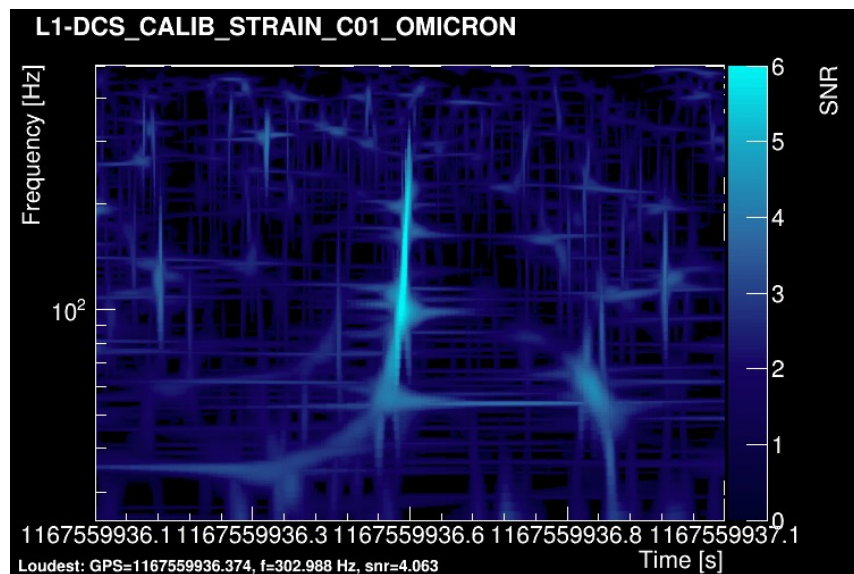
← Livetime x 2.5 /O1

LIGO binary neutron star inspiral range



← L1 sensitivity x 1.3 /O1

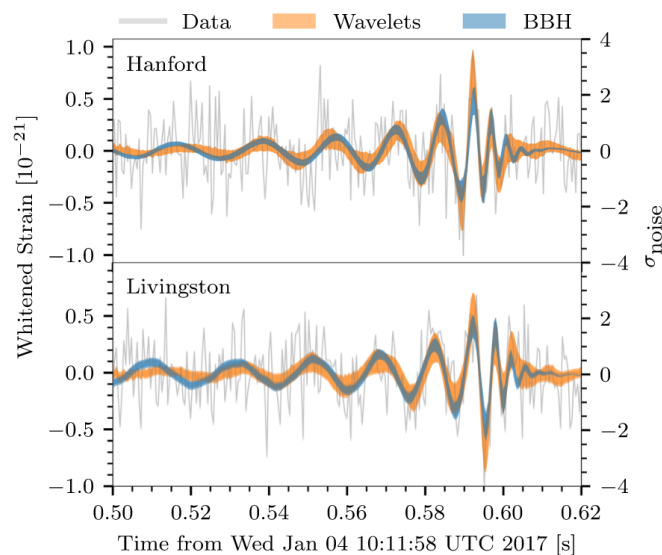
← H1 sensitivity x 0.8 /O1



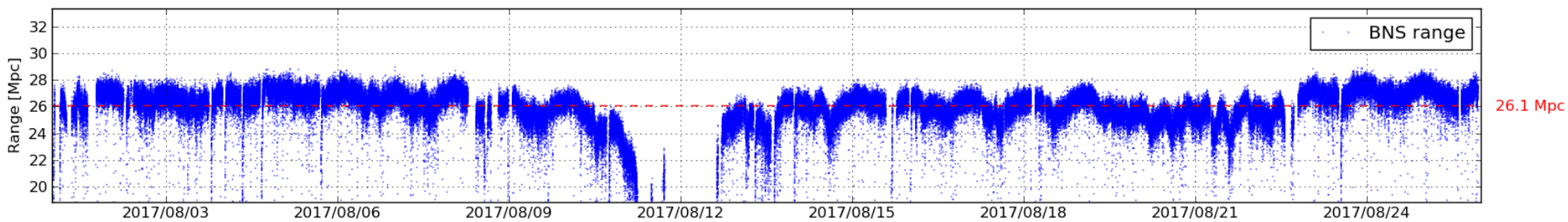
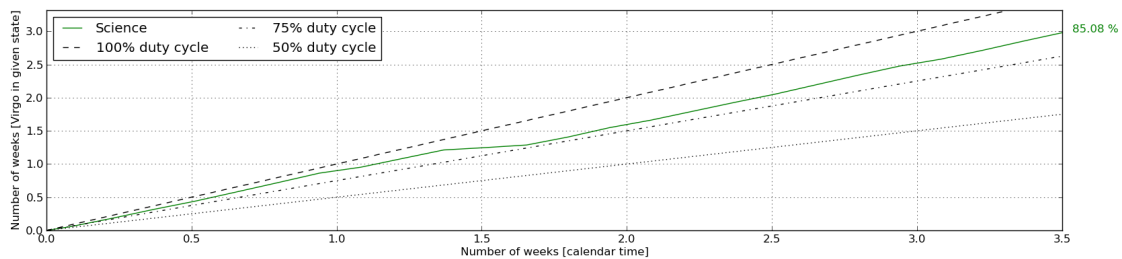
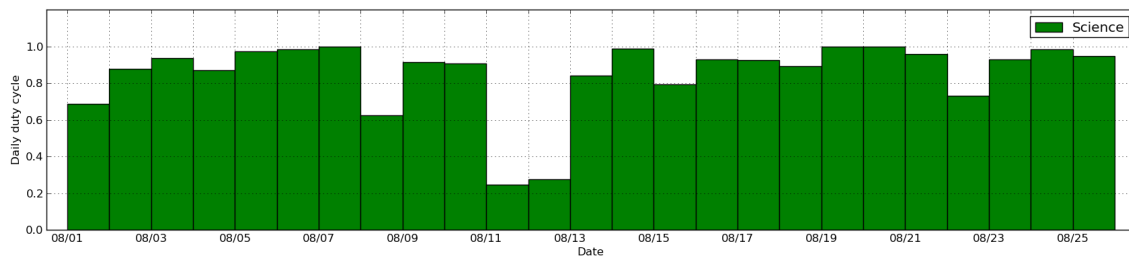
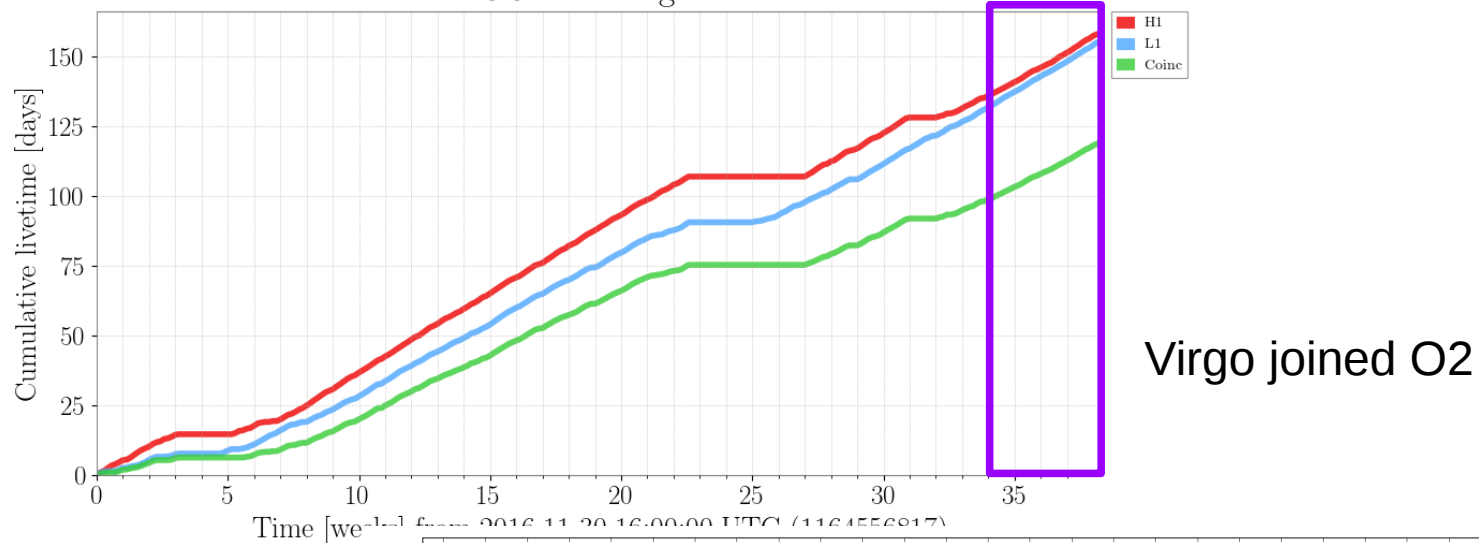
GW170104

Primary black hole mass m_1	$31.2^{+8.4}_{-6.0} M_{\odot}$
Secondary black hole mass m_2	$19.4^{+5.3}_{-5.9} M_{\odot}$
Chirp mass \mathcal{M}	$21.1^{+2.4}_{-2.7} M_{\odot}$
Total mass M	$50.7^{+5.9}_{-5.0} M_{\odot}$
Final black hole mass M_f	$48.7^{+5.7}_{-4.6} M_{\odot}$
Radiated energy E_{rad}	$2.0^{+0.6}_{-0.7} M_{\odot} c^2$
Peak luminosity ℓ_{peak}	$3.1^{+0.7}_{-1.3} \times 10^{56} \text{ erg s}^{-1}$
Effective inspiral spin parameter χ_{eff}	$-0.12^{+0.21}_{-0.30}$
Final black hole spin a_f	$0.64^{+0.09}_{-0.20}$
Luminosity distance D_L	$880^{+450}_{-390} \text{ Mpc}$
Source redshift z	$0.18^{+0.08}_{-0.07}$

PRL118,221101 (2017)



LIGO observing time



The Virgo and LIGO Scientific Collaborations have been observing since November 30, 2016 in the second Advanced Detector Observing Run 'O2', searching for gravitational-wave signals, first with the two LIGO detectors, then with both LIGO and Virgo instruments operating together since August 1, 2017. **Some promising gravitational-wave candidates have been identified in data from both LIGO and Virgo during our preliminary analysis, and we have shared what we currently know with astronomical observing partners.** We are working hard to assure that the candidates are valid gravitational-wave events, and it will require time to establish the level of confidence needed to bring any results to the scientific community and the greater public. We will let you know as soon we have information ready to share.