

THE ArDM PROJECT

2nd Symposium On

Neutrinos and Dark Matter in Nuclear Physics (NDMo6)

Paris September 3-9 2006



Sergio Navas
Univ. Granada, Spain

The ArDM working group (6 institutes, 31 members)

A. Badertscher, A. Baeztner, R. Chandrasekharan, L. Kaufmann,
L. Knecht, M. Laffranchi, M. Messina, G. Natterer, P. Otiougova,
A. Rubbia, J. Ulbricht
ETH Zurich, Switzerland

C. Amsler, V. Boccone, A. Buechler-Germann, C. Regenfus
Zurich University, Switzerland

A. Bueno, M.C. Carmona-Benitez, J. Lozano, S. Navas
University of Granada, Spain

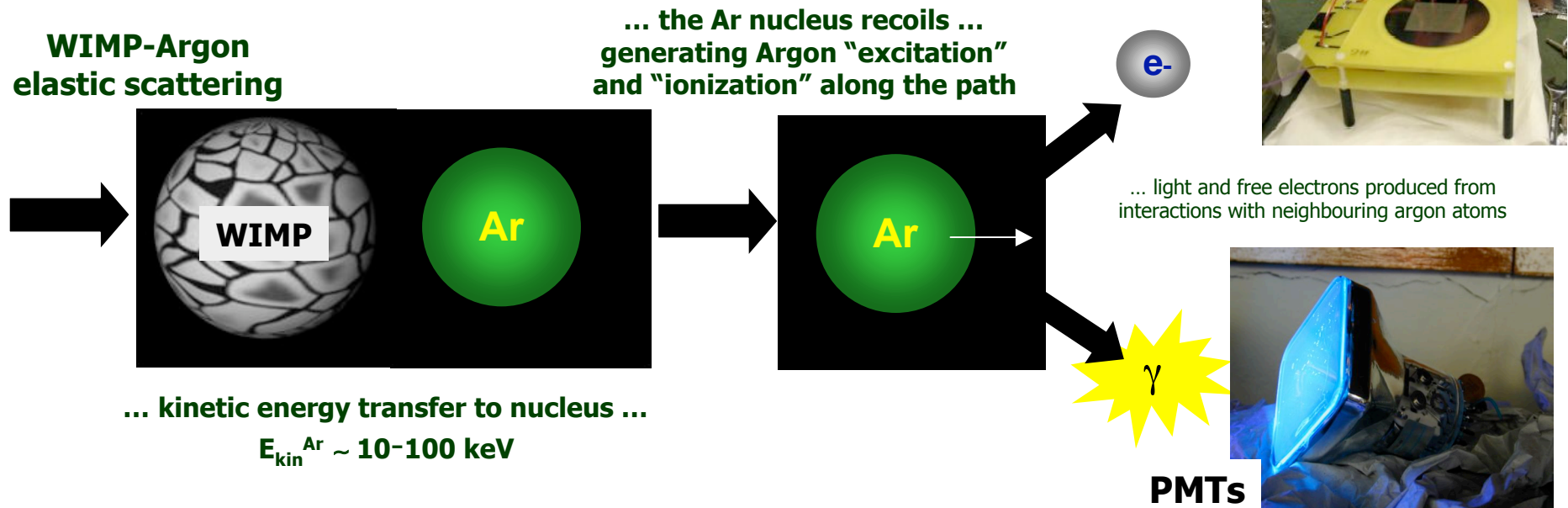
M. Daniel, P. Ladron de Guevara, L. Romero
CIEMAT, Spain

P. Mijakowski, P. Przewlocki, E. Rondio
Soltan Institute Warszawa, Poland

H. Chagani, E. Daw, V. Kudryavtsev, P. Lightfoot,
P. Majewski, N. Spooner
University of Sheffield, England

We acknowledge informal contribution from **LNF, Italy**. Interest from JINR, Russia

- We are constructing a **1 ton argon prototype** at CERN ... the goal (1st phase) is **to demonstrate the validity of the design**
- It has been shown that liquid Xenon or **Argon** can act as a **target for WIMP** detection (NIM A 327 (1993) 205 & NIM A 449 (2000) 147)
- Our aim is to detect the **ionization charge** and **scintillation light** independently:



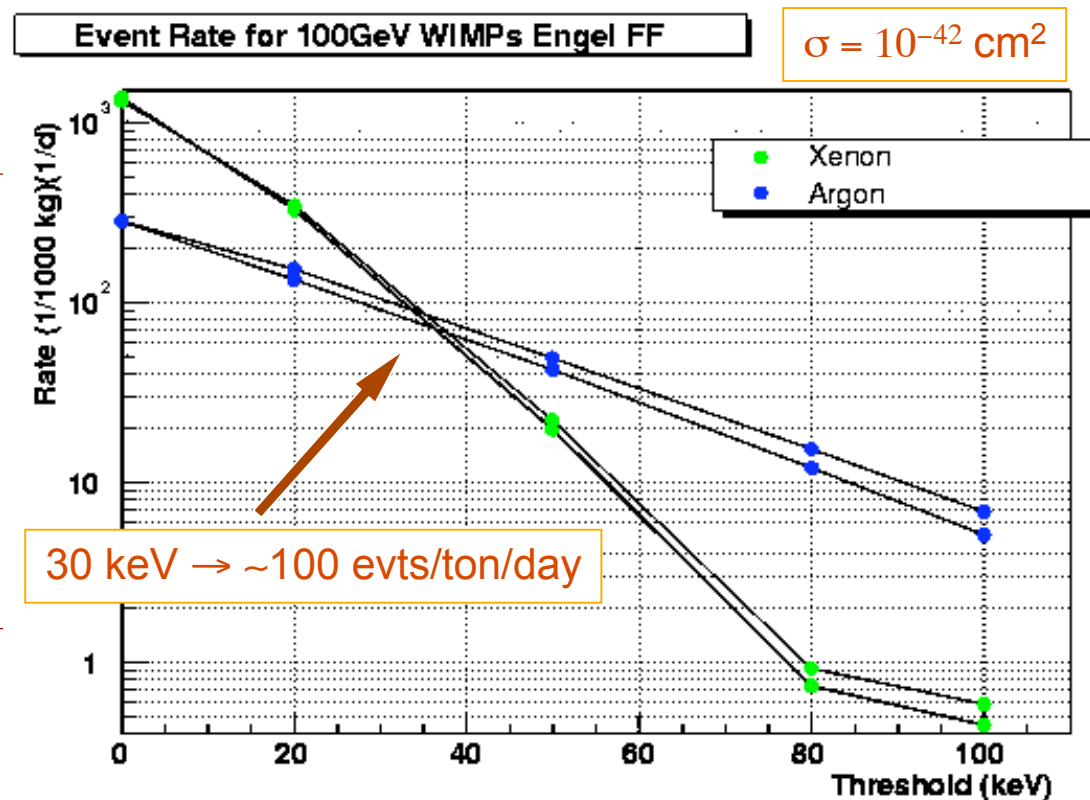
- **Event vertices** will be precisely localized (in space), important for γ -ray and neutron background rejection from surrounding elements
- **Event topology** will be imaged and used for background rejection
- The **Ratio of ionization to scintillation** used to reject β/γ -rays against nuclear recoils
- The **time distribution of the scintillation light** will be used to further discriminate between heavy recoils and other backgrounds

WIMP – Argon elastic scattering

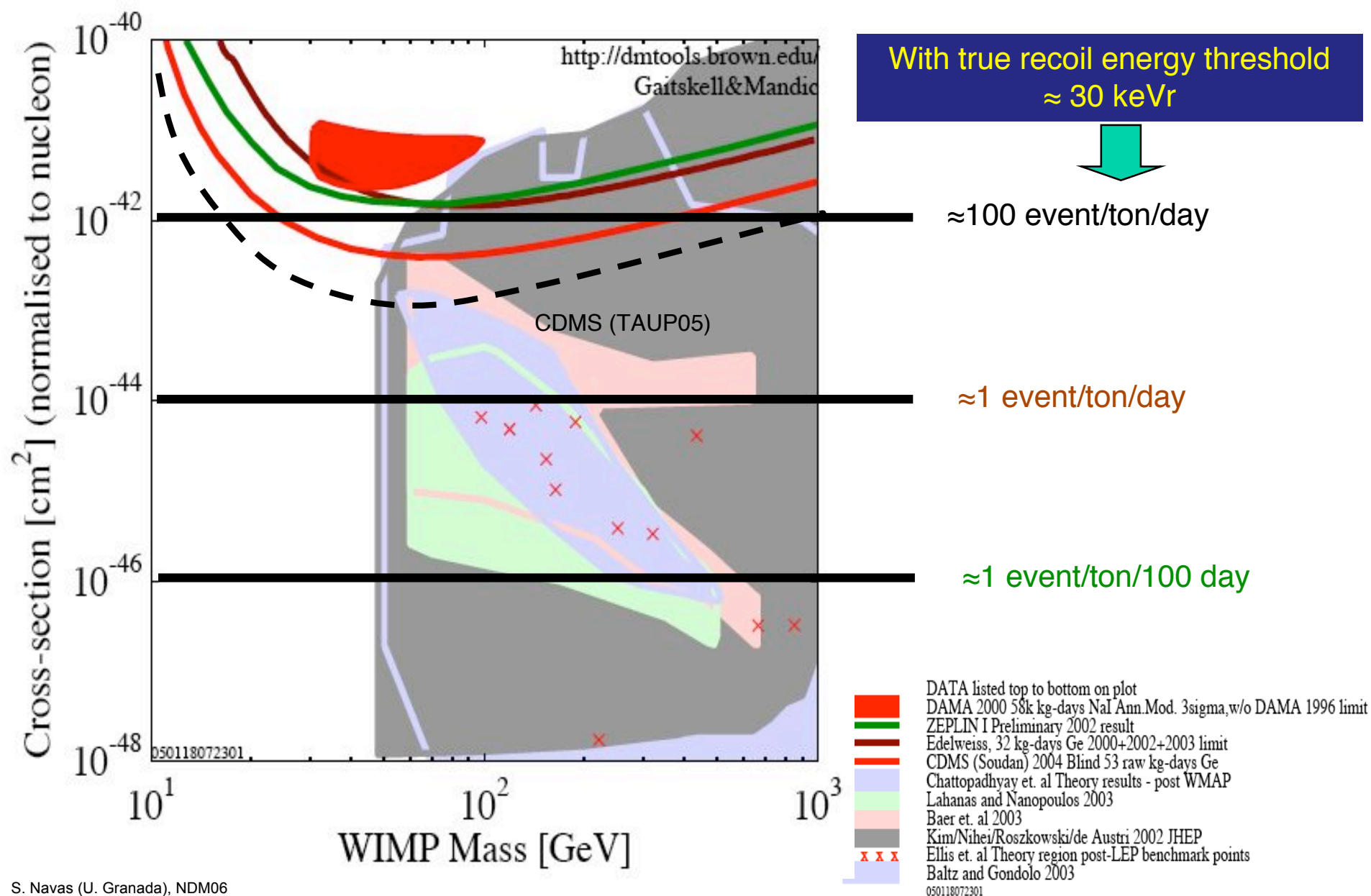
- Event rate in argon is less sensitive to threshold on recoil energy than for xenon (Form Factors)
- Recoil spectra in xenon and argon are different, providing an important cross-check in case of positive signal

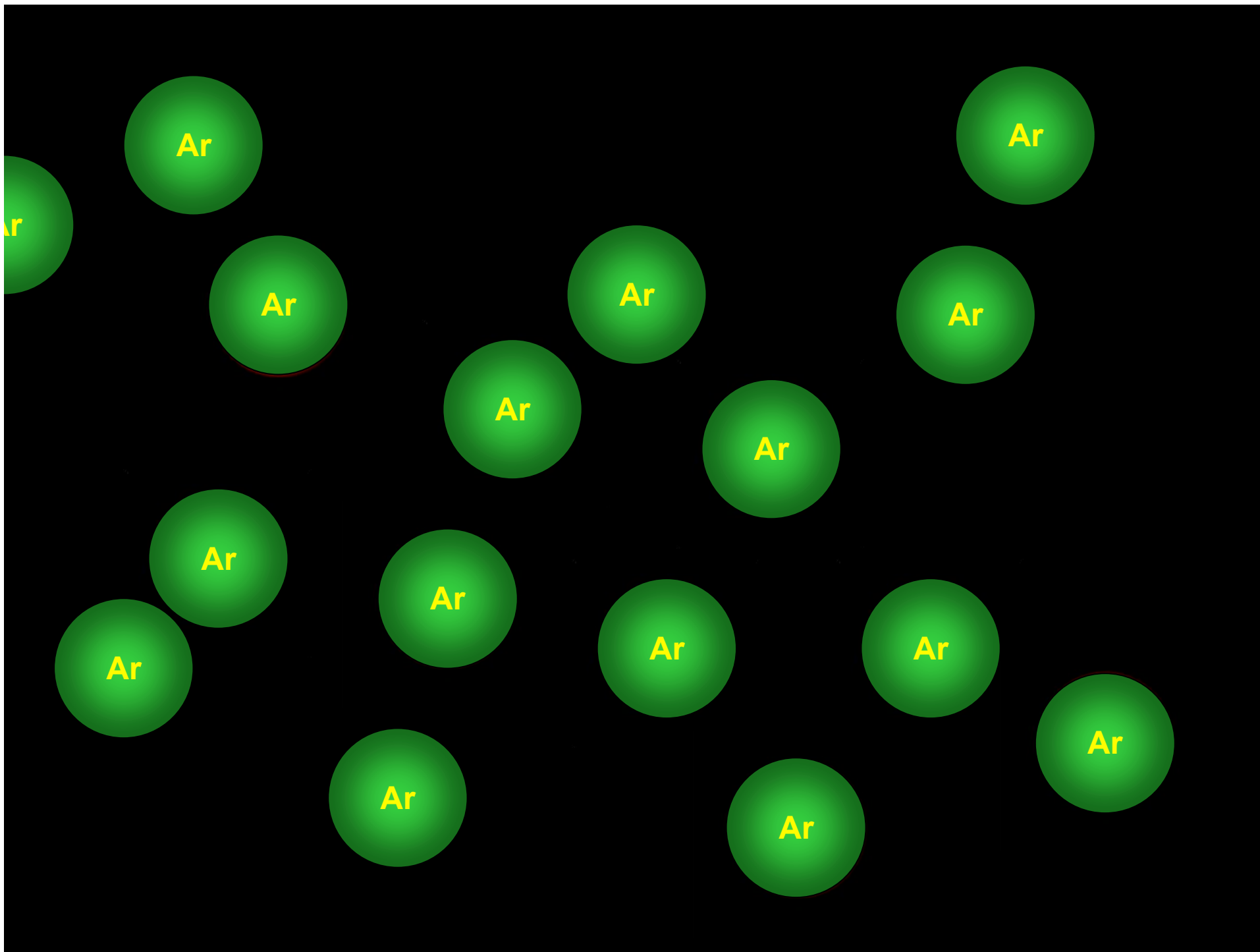
Assumptions for simulation:

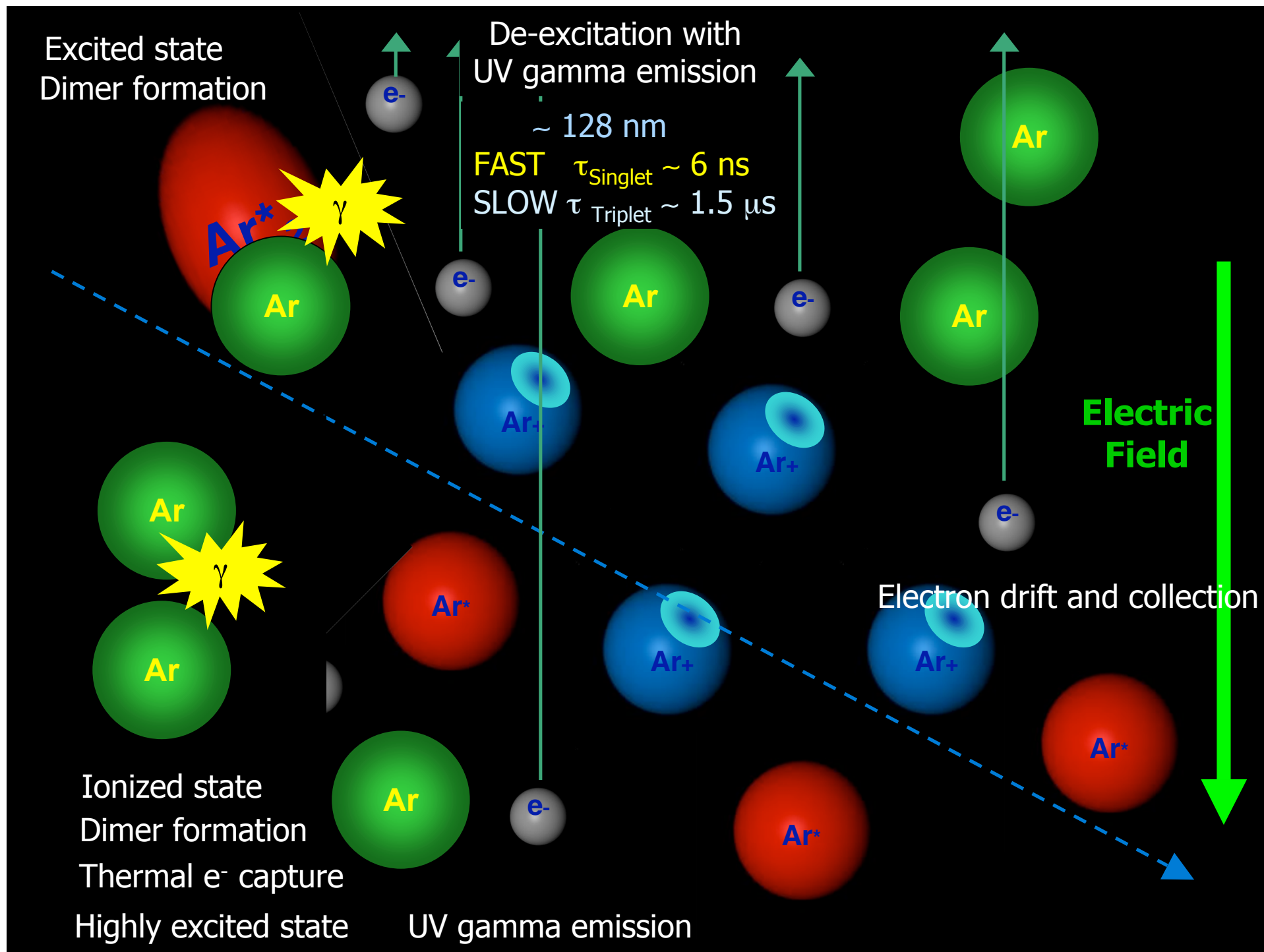
- Cross-section normalized to nucleon
 - $\sigma = 10^{-42} \text{ cm}^2 = 10^{-6} \text{ pb}$
 - $M_{\text{WIMP}} = 100 \text{ GeV}$
- Halo Model
 - WIMP Density = 0.5 GeV/cm^3
 - $v_{\text{esc}} = 600 \text{ km/s}$
- Interaction
 - Spin independent
 - Engel Form factor



Estimated event rates on Argon

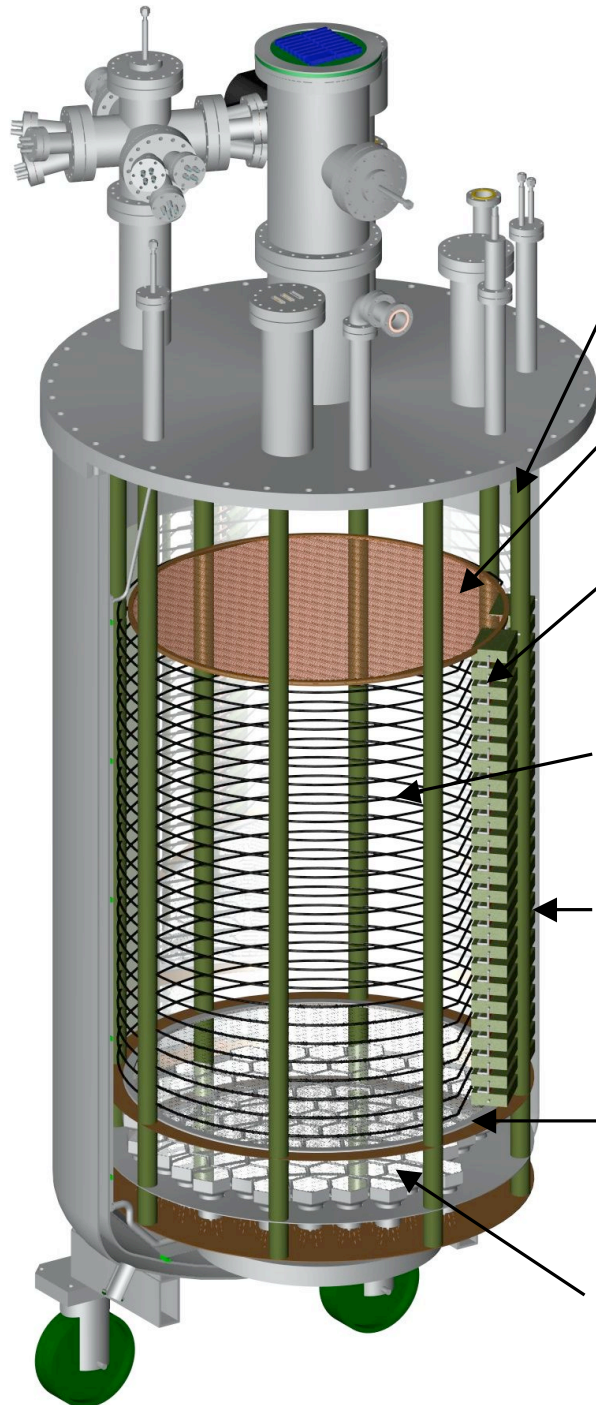






Prototype layout

- Cylindrical volume, drift length ≈ 120 cm
- 850 kg LAr target



10 Polyethylene pillars as mechanical support.

2x LEM for the electron multiplication and readout ($\text{Gain} \approx 10^3 - 10^4$)

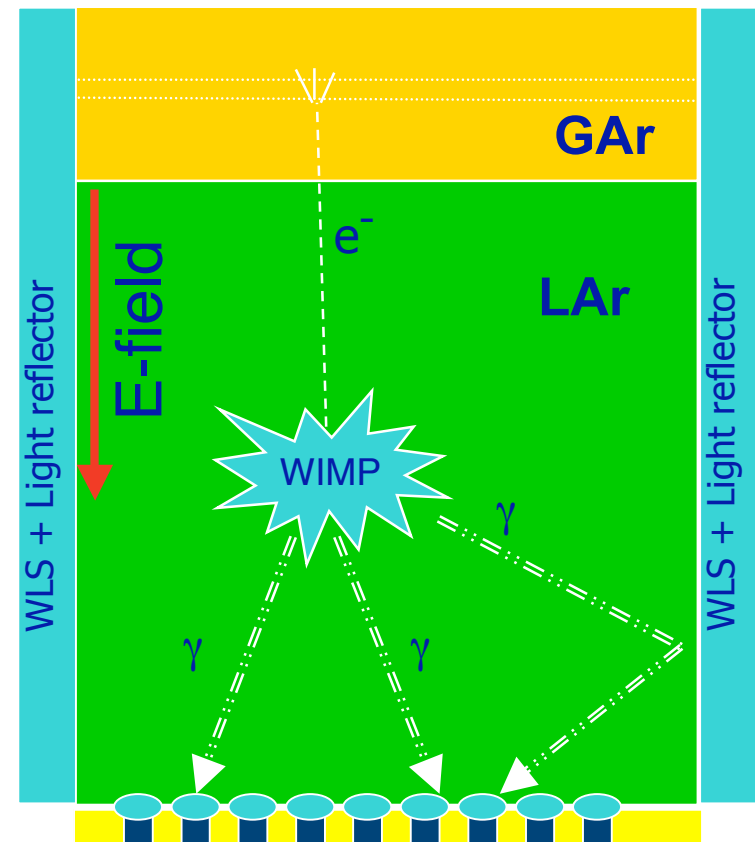
Greinacher chain: supplies the right voltages to the field shapers rings and the cathode up to 500kV $\rightarrow \approx 4$ kV/cm

The field shapers are needed to make an homogeneous \vec{E}

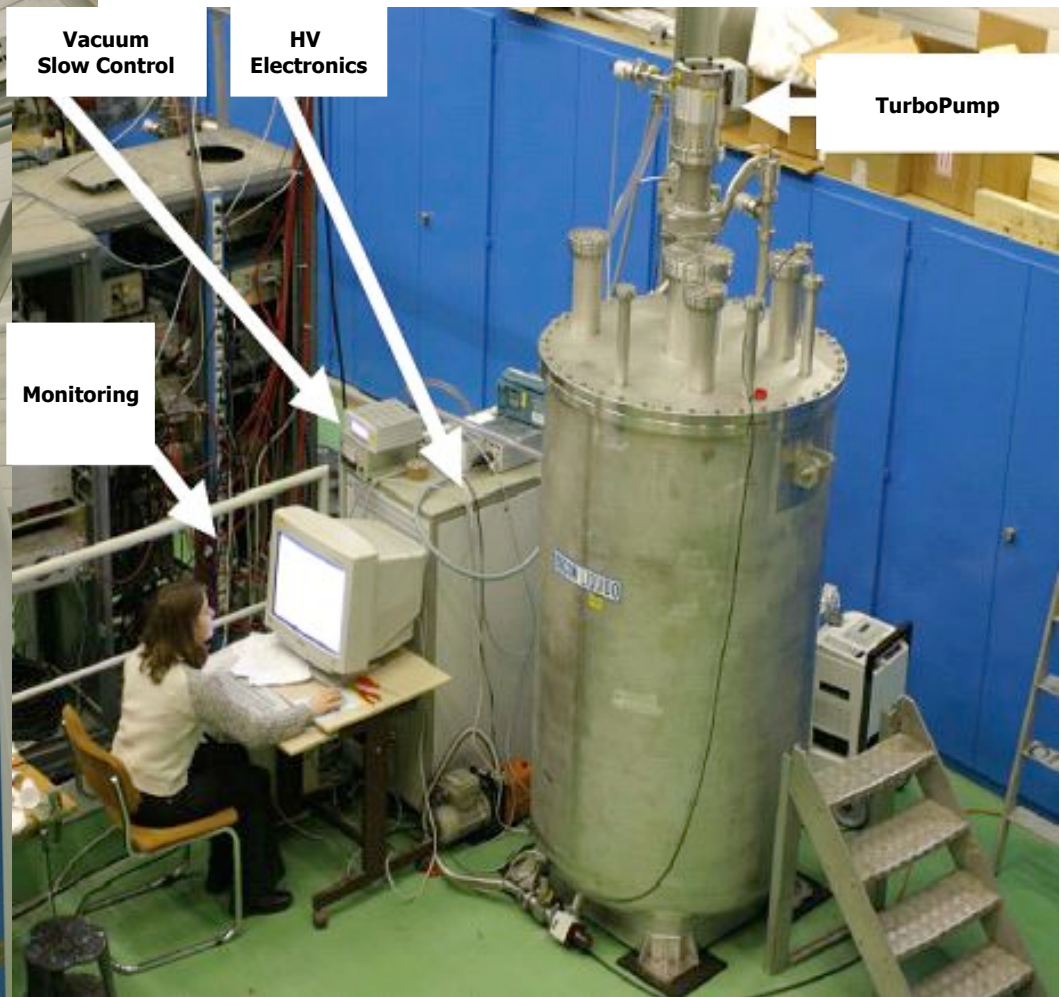
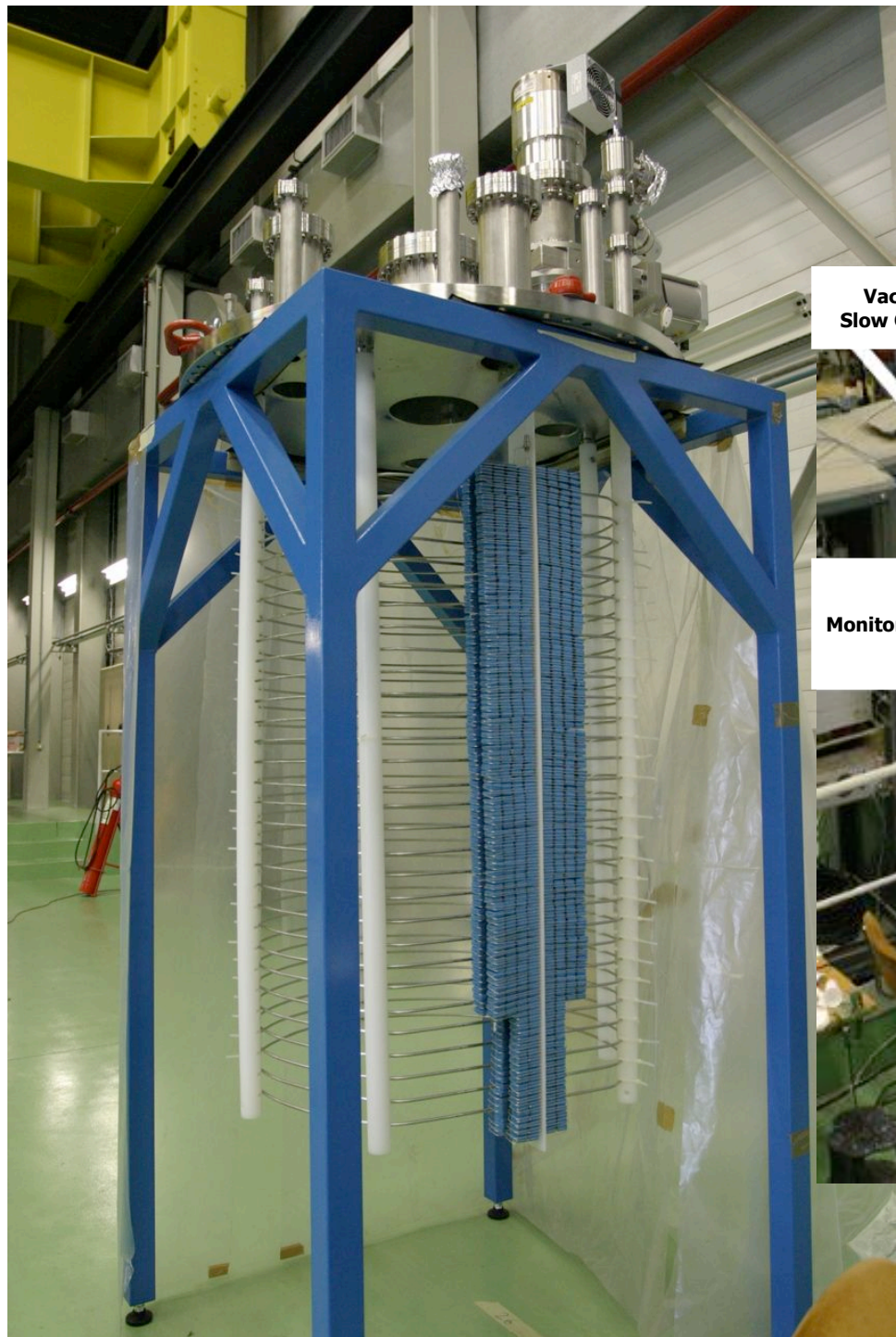
The aluminized Mylar reflects the scintillation light ($>90\%$)

Cathode: semi-transparent in order to let the scintillation light pass trough ...

~ 80 PMTs below the cathode to detect the scintillation light.



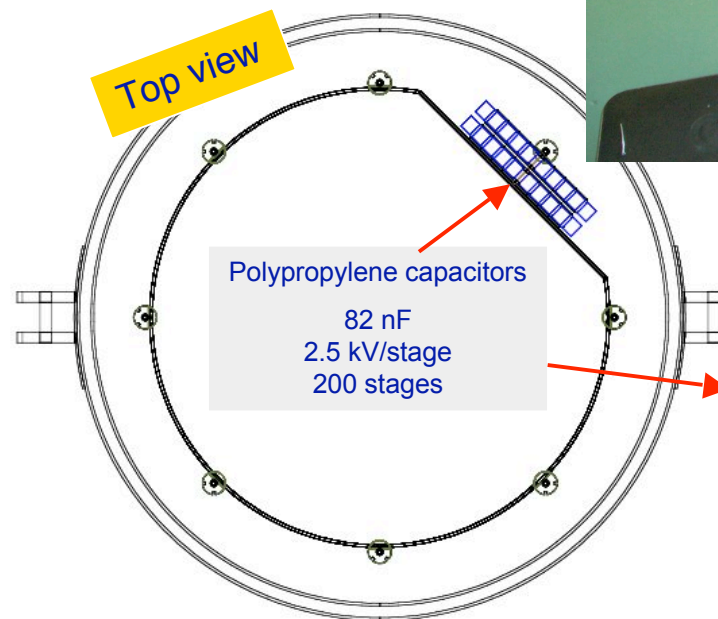
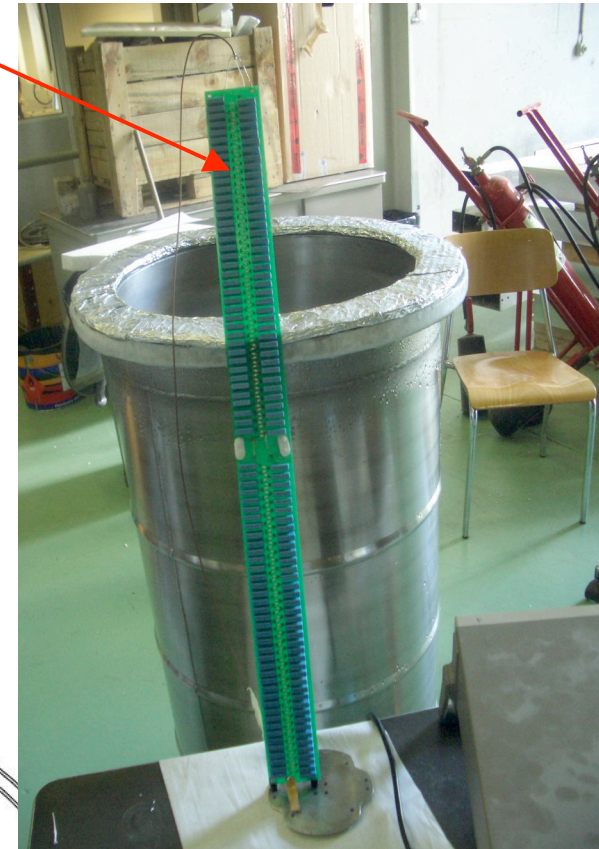
**ArDM
1-ton prototype
@ CERN**



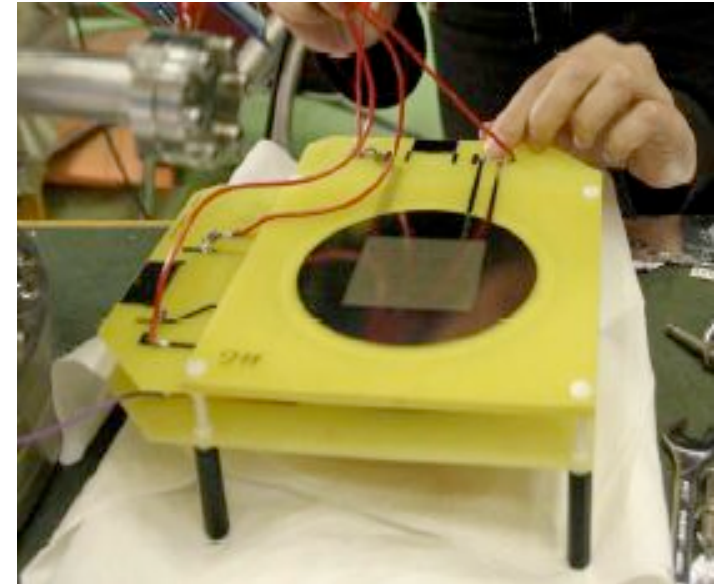
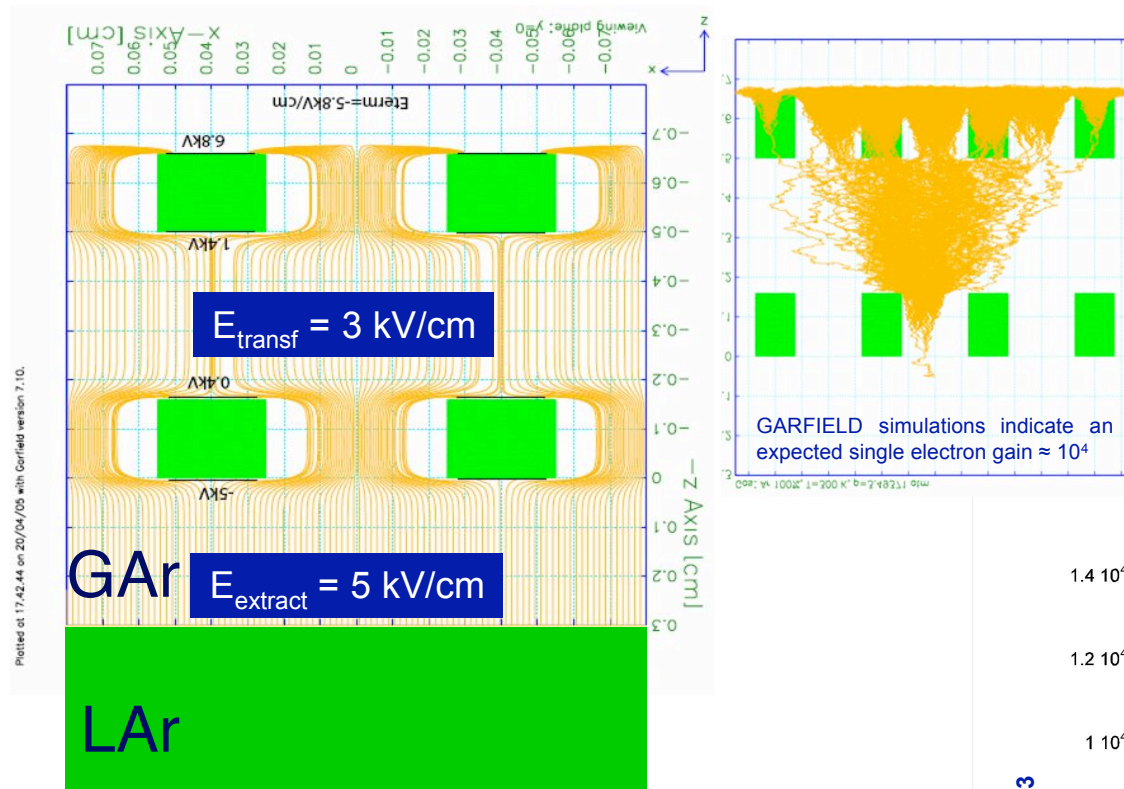
High Voltage system for drift field generation

- A cascade of rectifier cells (Greinacher/Cockroft-Walton circuit) used
- The total voltage we aim to reach is $V_{\text{tot}} = 500 \text{ kV}$, i.e. $\approx 4 \text{ kV/cm}$
- Tests in liquid nitrogen have been performed
- The largest system successfully operated consists of 80 stages and reached stable operation at up to $120 \text{ kV} \approx 2 \text{ kV/cm}$

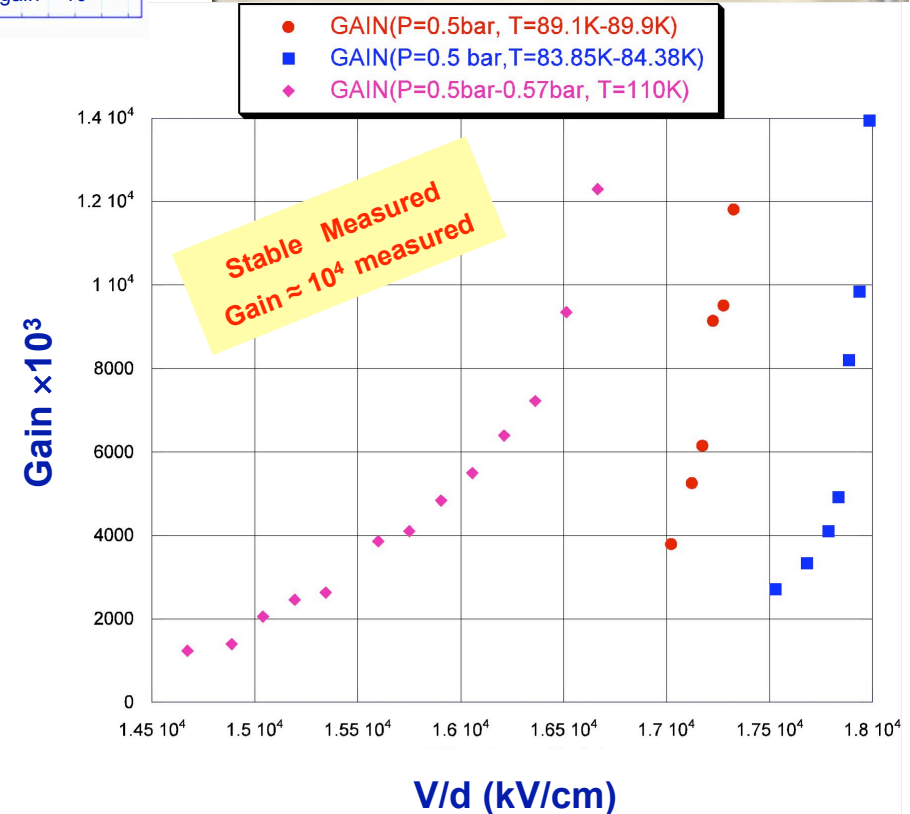
Mounted on field shaper rings



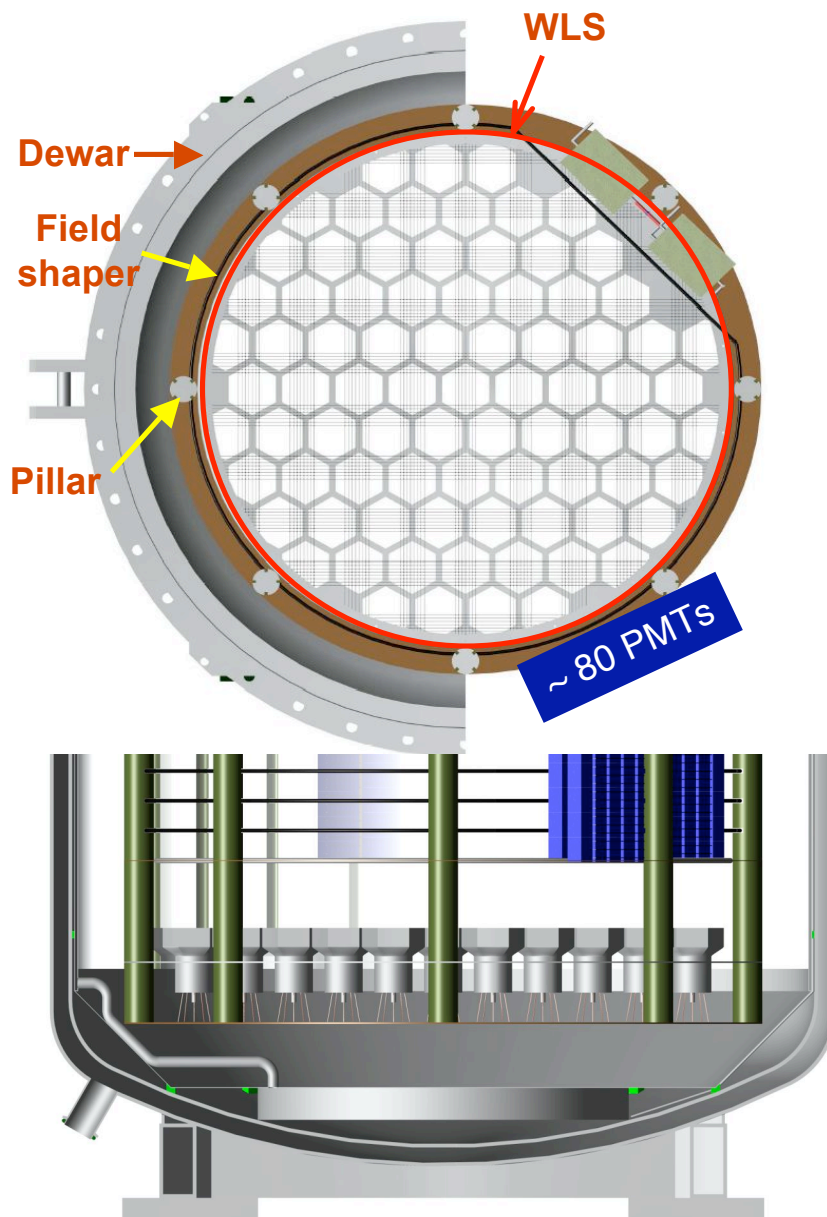
Layout of the charge readout system



- Distance between stages: 3 mm
- Avalanche spreads into several holes at second stage
- Higher Gain reached as with one stage, with good stability
- Hole dimension: 500 μm diameter, 800 μm distance.
- Thickness of PCB: 1.6 mm



Layout of Light Readout system and PMT



1. Scintillation light detection via PMTs and WLS reflector in mirror walls

- ✓ Advantage: easier to find materials with high reflectivity for visible light than for UV light
- ✓ 128 nm γ produced in fiducial volume, hits the WLS mirror \rightarrow diffuse reflection at 430 nm.
- ✓ Two layer WLS cylinder: outer to reflect visible light (99%), inner thin WLS added. Mylar foils coated with thin Al+MgF₂ layers (CERN)

2. Scintillation light detection via PMTs coated with WLS



- Polymer and Tetra-Phenyl-Butadiene (TPB) compound coated on PMT window shifts the DUV light (128 nm) to 430 nm
- Efficiency of wavelength shifting: 20% to 30%
- PMT's: array of ~80 photosensors at bottom of detector

Slow Control Devices

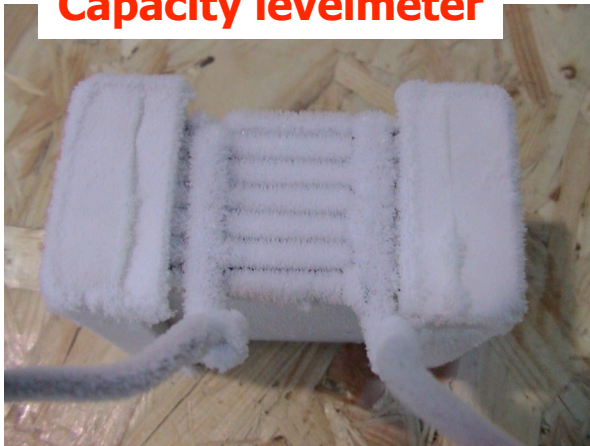
- A series of **custom designed Slow Control devices** have been built, tested and installed to monitor temp., level, purity ...

PT10K resistors

10 K Ω at 0°C
Range: -200 to 400°C

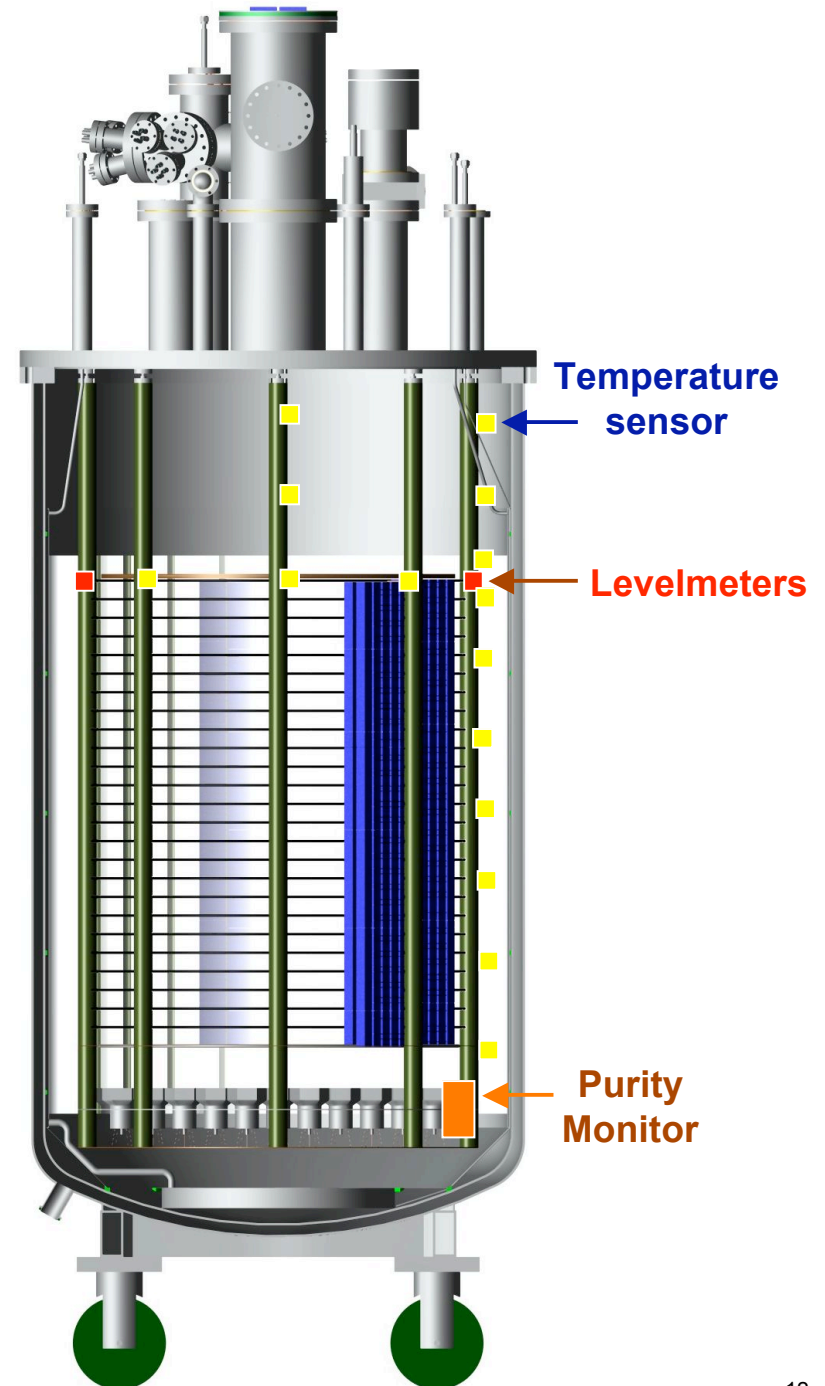


Capacity levelmeter



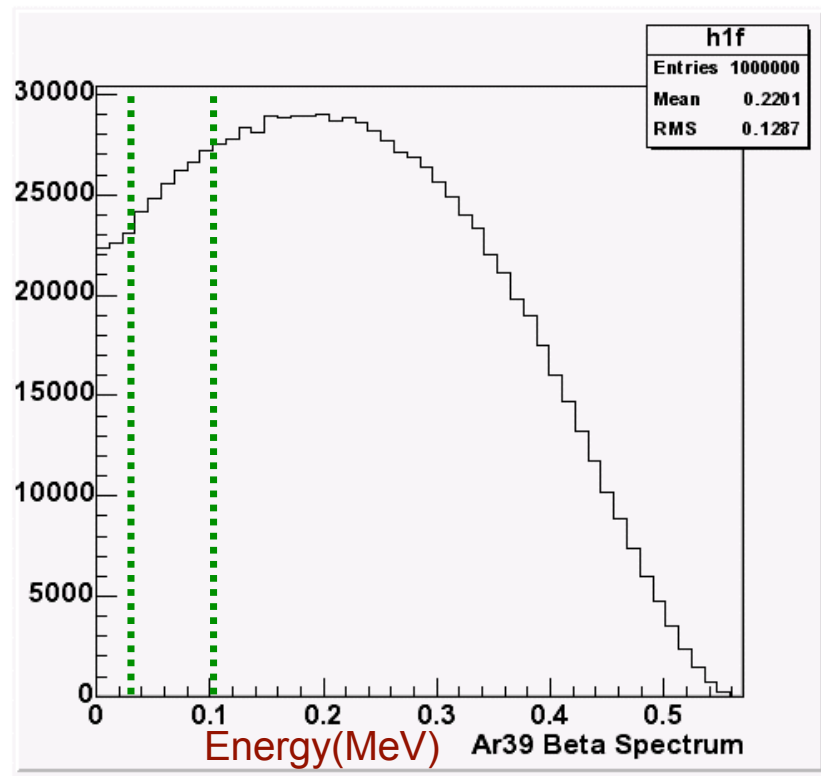
~ 0.7 pF/mm
precision of ~ 0.03 mm

Argon Purity monitor



Intrinsic background from Argon 39 isotope

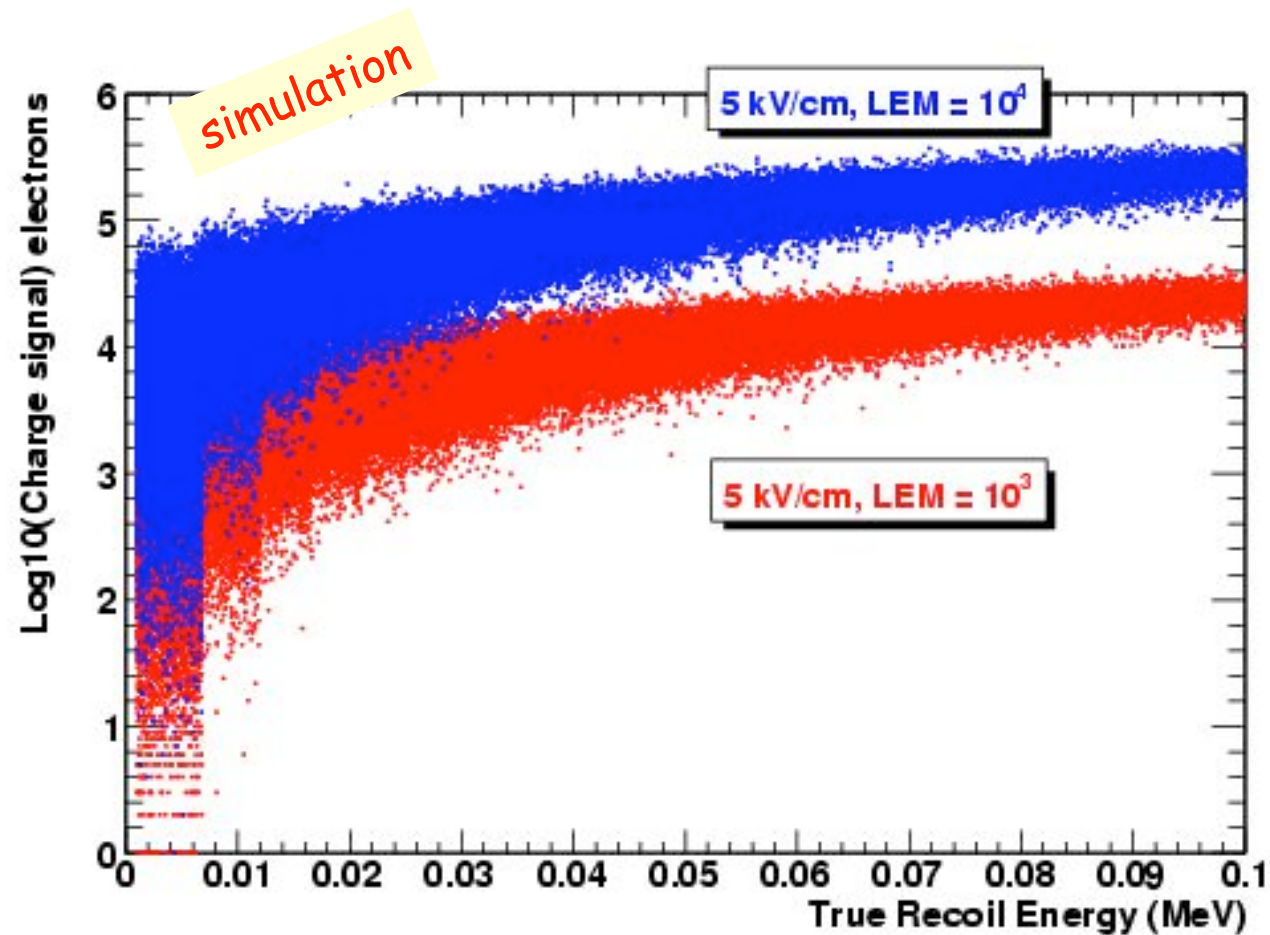
Natural argon from liquefaction of air contains small fractions of ^{39}Ar radioactive isotope (well known to geophysicists)



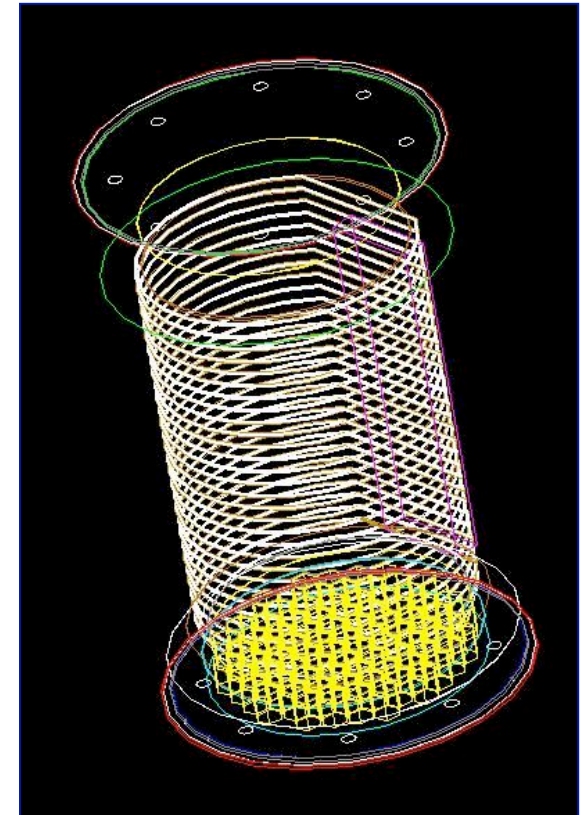
- Induced in atmospheric argon by cosmic rays
- Concentration in natural Ar: 8.1×10^{-16} $^{39}\text{Ar}/\text{Ar}$ [H.H. Loosli, Earth and Planetary Science Letters, 63 (1983) 51 and "Nachweis von ^{39}Ar in atmosphärischem Argon" PhD thesis University Bern 1968]
- Half life: 269 years, $Q=565$ KeV
- Mean Energy: 218 keV
- Integrated rate in 1 ton LAr $\sim 1\text{kHz}$ [WARP Coll.] astro-ph/0603131

To suppress ^{39}Ar fraction we consider using Ar extracted from well gases (extracted from underground natural gas). On the other hand, this source, evenly distributed in the target, provides precise calibration and monitoring of the detector response.

Charged (WIMP) signal readout



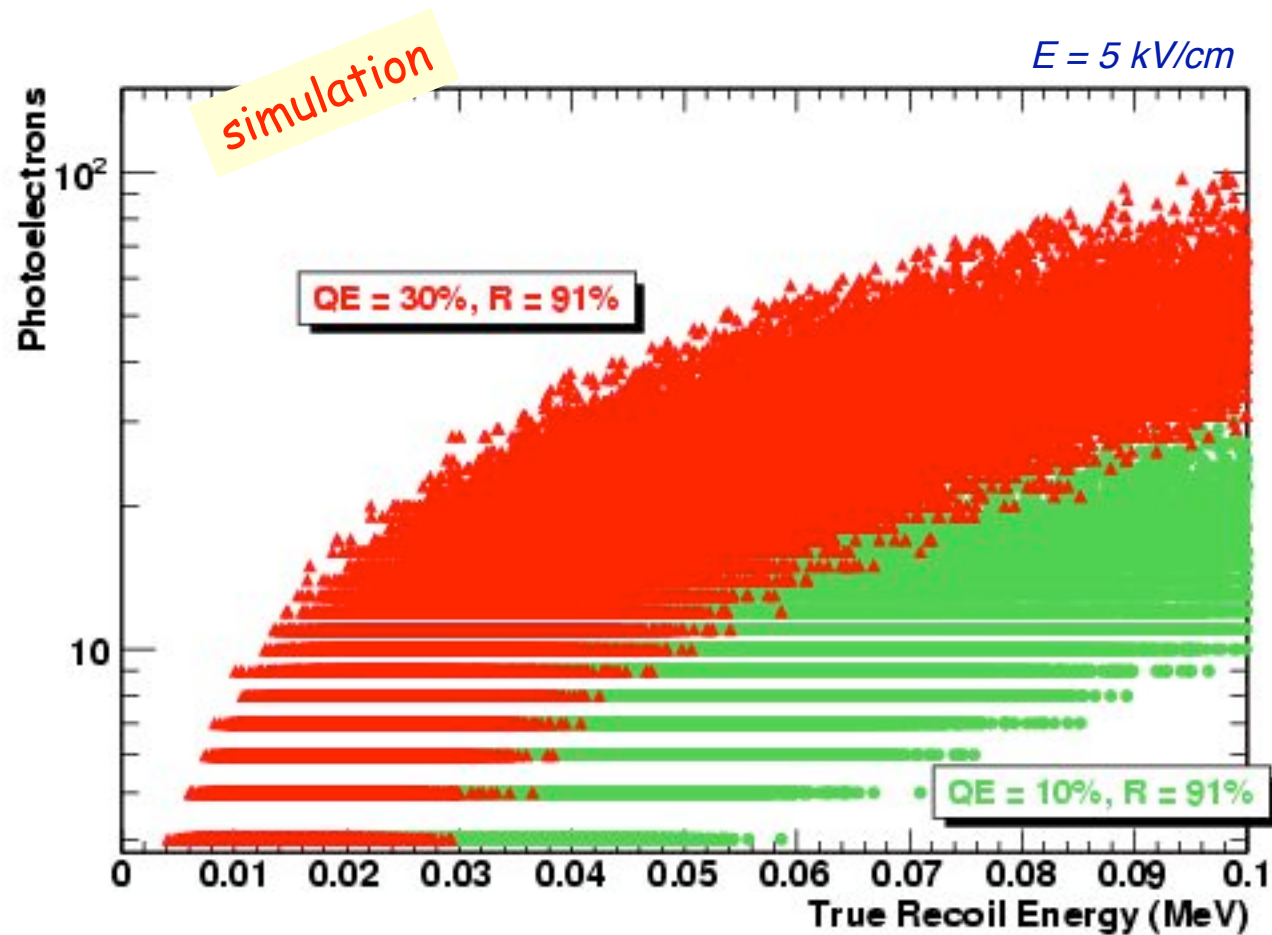
Detailed GEANT4 simulation of charge readout for various configurations



- Drift field ≈ 5 kV/cm
- Charge LEM readout: Single electron gain $\approx 10^3 \div 10^4$
- Electron lifetime included (1.4 ms \sim 0.2 ppb O₂ eq.)

Light (WIMP) signal readout

Detailed simulation of charge readout for various configurations



- Surface wall reflectivity $R \approx 90\%$
- Cathode transparency = 80%
- PMT coverage = 70%
- PMT quantum efficiency $QE \approx 10\% \div 30\%$



Average light yield

$0.15 \div 0.43$ phe/keV

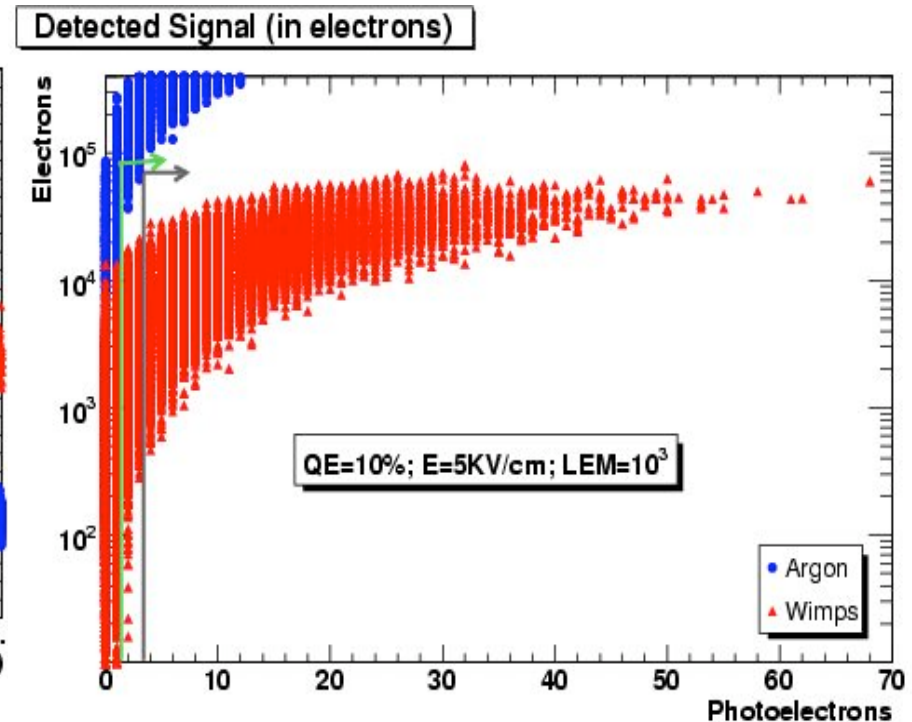
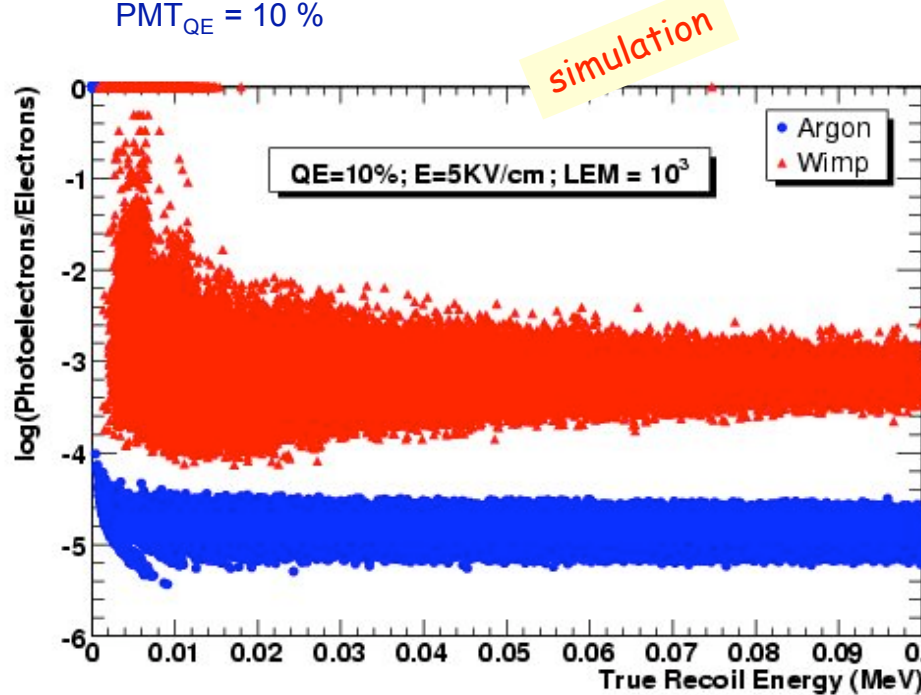
WIMPs vs. ^{39}Ar background discrimination

This is MONTE CARLO, this relies heavily on MC, there is no reason to believe this is OK, this is exactly what the 1 ton test at CERN should prove.

Drift field $E = 5 \text{ kv/cm}$

$\text{LEM} = 10^3$

$\text{PMT}_{\text{QE}} = 10 \%$



CUTS:

True recoil energy $> 30 \text{ keV}$

$Q > 2000 \text{ electrons}$

light $\geq \text{phe}$

phe = 2 : $\sim 91 \text{ WIMP evts/day}$

phe = 4 : $\sim 85 \text{ WIMP evts/day}$

(If Quenching = 0.28)

phe = 2 : $\sim 39 \text{ WIMP evts/day}$

phe = 4 : $\sim 9 \text{ WIMP evts/day}$

Neutron Background from detector components

Neutron sources:

- **Uranium and Thorium** contamination (spontaneous fission) of the detector components and the surrounding rock:
 - flux about $3.8 \times 10^{-6} \text{ cm}^{-2} \text{ s}^{-1}$ (at 2450 m.w.e.)
 - can be shielded, e.g. by a hydrocarbon shield
- **Muon-induced neutrons** from surrounding rock, shielding and detector components

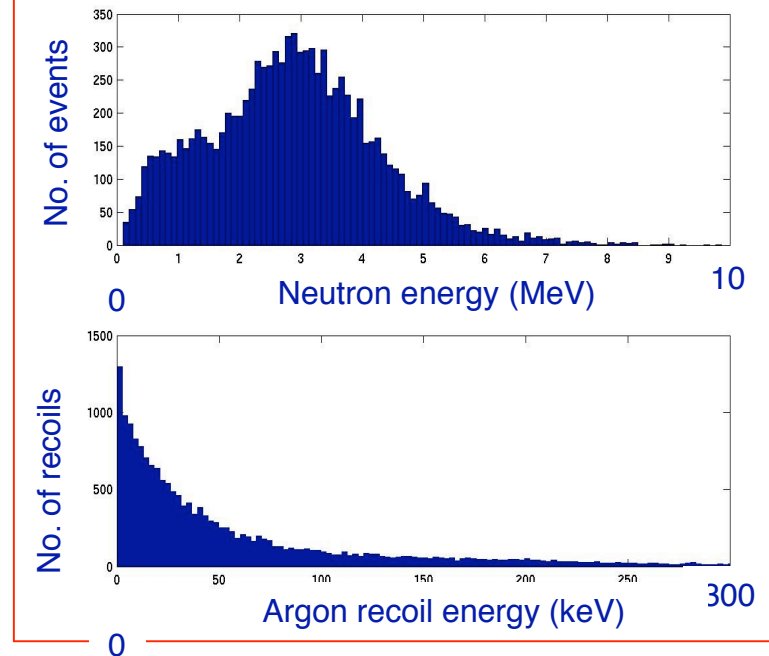
High energy neutrons penetrate shielding, are thereby moderated and can cause WIMP-like events.

Event numbers per year

Component	n per year	WIMP-like recoils
Container	~ 400	~ 50
LEM (std. materials)	~ 10000	~ 1300
LEM (PEEK)	< 18	< 1
81 PMTs (std. materials)	~ 12000	~ 1350
81 PMTs (screened)	~ 600	~ 70

Compared with ~ 3500 WIMP events at $\sigma = 10^{-43} \text{ cm}^{-2}$

Geant4 simulation



Nuclear recoils:

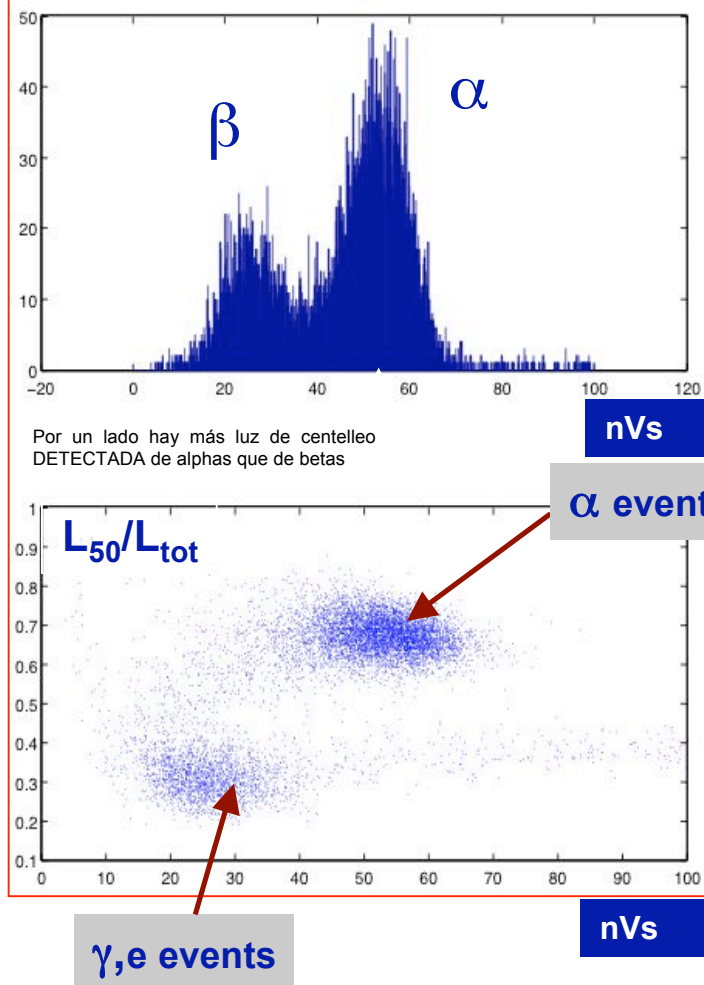
- 70% scatter more than once within the fiducial volume → advantage of large detectors
- 10% produce a WIMP-like event (single scattering, recoil energy $\in [30, 100] \text{ keV}$)

Light measurements in Liquid Argon (preliminary)

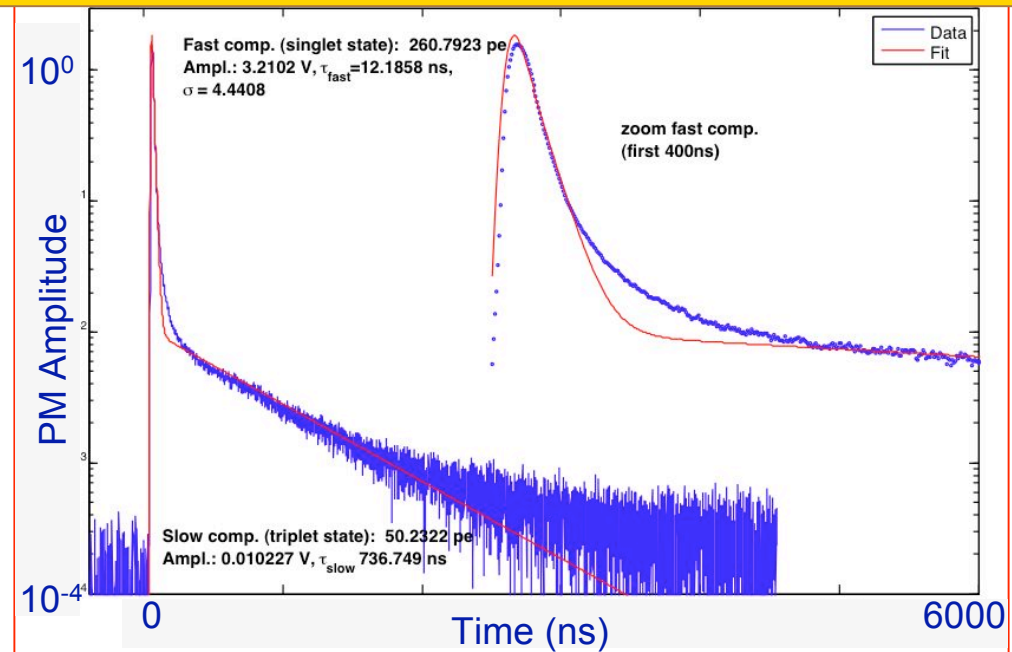
real data

Radioactive source: α (5.4 MeV) + β (Q = 1.163 MeV)

Event separation in liquid argon



Scintillation light from α in 1200 mbar liquid argon



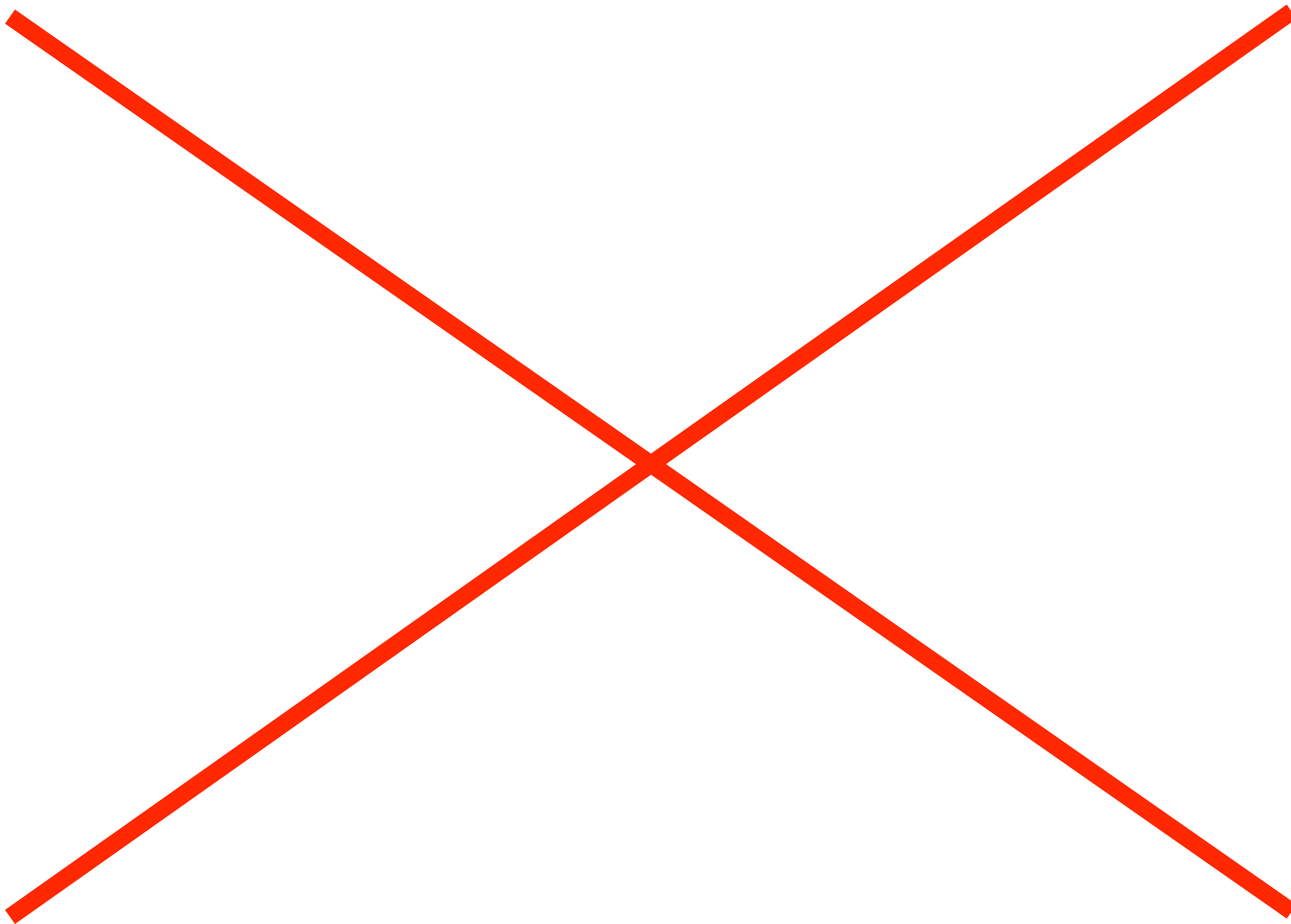
- α events separate well from γ, e events
- Fast and slow light components distinguishable

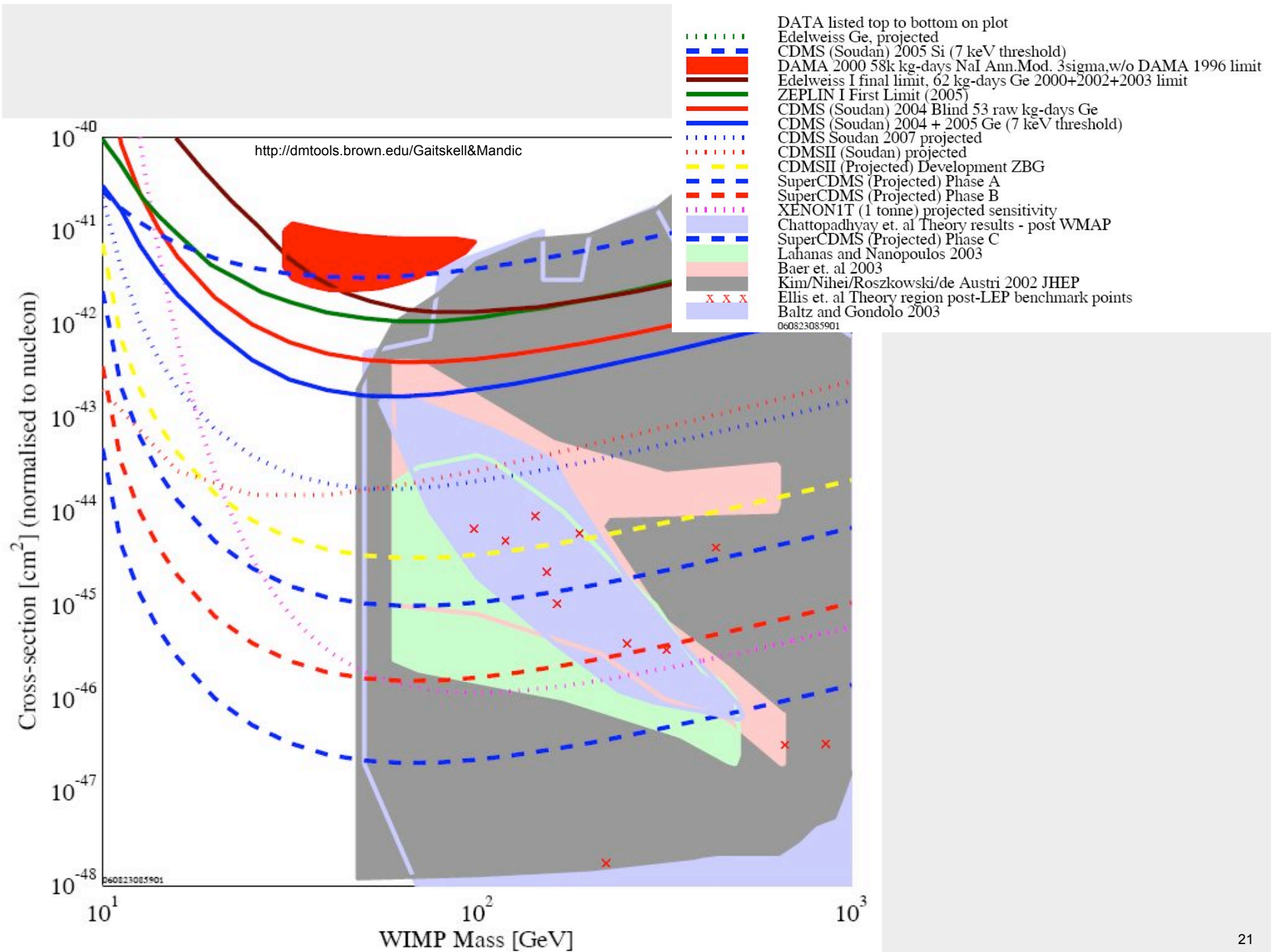
$$\frac{I_S}{I_T} = \begin{matrix} \sim 0.3 & \text{for } \beta \text{ particles} \\ \sim 1.3 & \text{for } \alpha \text{ particles} \\ \sim 3 & \text{for fission fragments} \end{matrix}$$

#photo-electrons	Rejection (@E=0V/cm)
>10	> 10^2
>20	> 5×10^3
>30	> 10^5

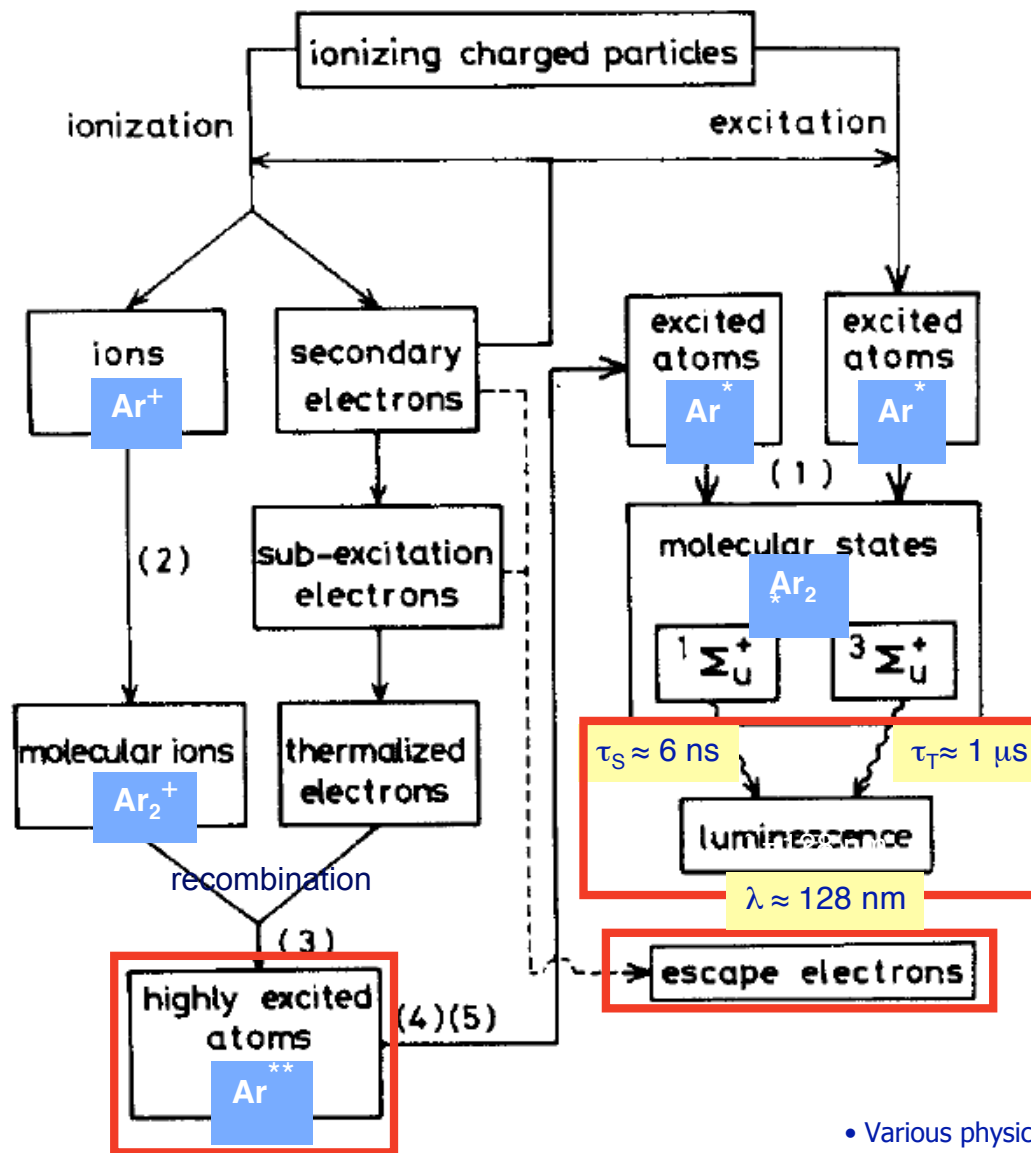
Outlook

- ✓ We are constructing a 1 ton prototype at CERN whose goal is to demonstrate the validity of the design.
- ✓ Our first milestone is a **proof of principle** on γ -rays and beta electron (^{39}Ar) rejection vs nuclear recoils above ground (1st phase, 2006-2007) .
- ✓ The goal requires a successful implementation and operation of:
 - ✓ Purified Liquid Argon target
 - ✓ High drift field
 - ✓ Independent charge (LEM-based charge readout) and
 - ✓ Scintillation light (PMTs) detection systems
- ✓ Following its successful operation, we will consider a **deep underground operation** (2st phase, 2007?). An EOI to the Canfranc Laboratory has been submitted.
- ✓ The expected exclusion sensitivity of the prototype is **$10^{-42} - 10^{-44} \text{ cm}^2 \text{ SI}$** ($10^{-6} - 10^{-8} \text{ pb}$) depending on the background rejection power.
- ✓ This technology could provide the means to develop larger detectors to reach sensitivities below $10^{-44} \text{ cm}^2 \text{ SI}$ cross section.





Processes induced by charged particles in Argon



The electron-ion density is so high in the ionization track produced by alpha particles that the recombination process proceeds much faster than the electron drift, and produces "singlet" excitons abundantly, resulting in the fast recombination luminescence.

Since the specific ionization density significantly affects the time dependence of the scintillation light of liquid Argon, it is possible to distinguish different ionizing particles by using a pulse shape discriminator. This will provide an effective background rejection mechanism.

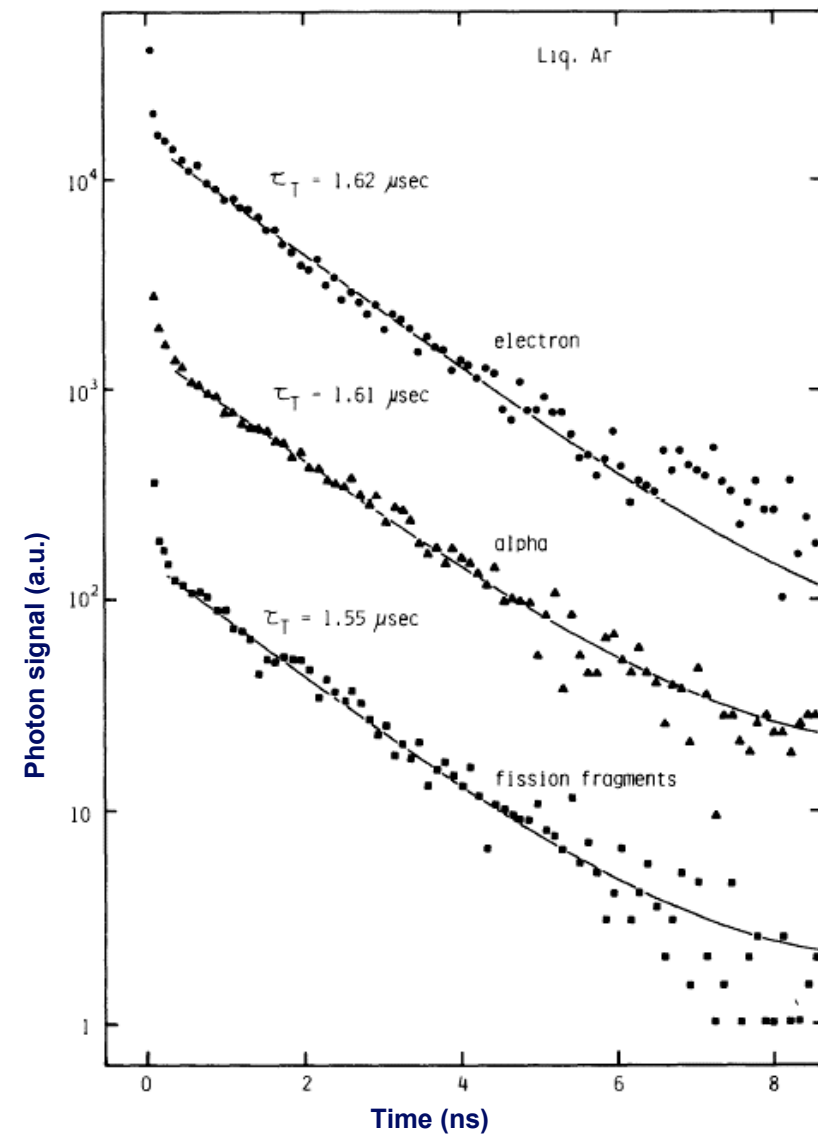
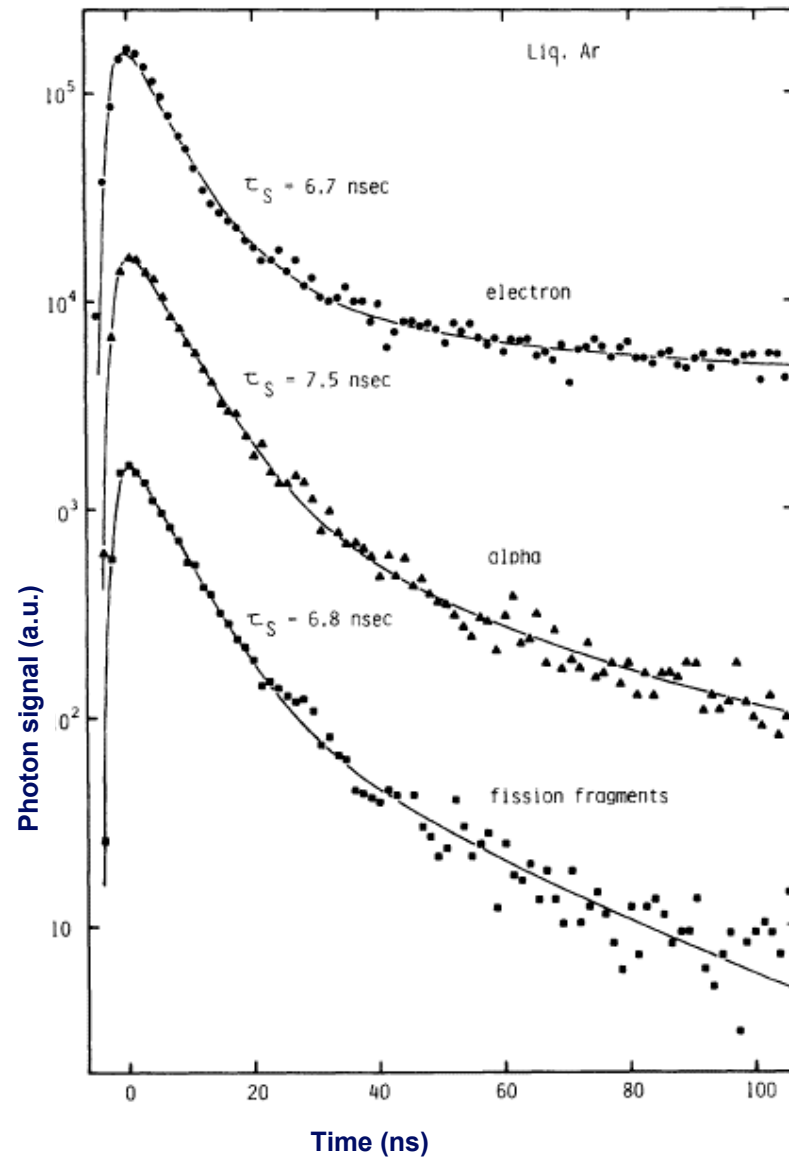
En Xenon líquido la emisión de centelleo tiene:
 $\lambda \approx 170 \text{ nm}$ $\tau_S \approx 5 \text{ ns}$ $\tau_T \approx 30 \text{ ns}$

UV Light (two components)
 +
 Charge

- Various physical processes leading to scintillation & ionization
- Yields are particle, energy and drift field dependent
- Simulation describes different response to WIMP and MIPs

Time dependence of scintillation light

A. Hitachi et al., Phys. Rev. B 27 (1983) 5279

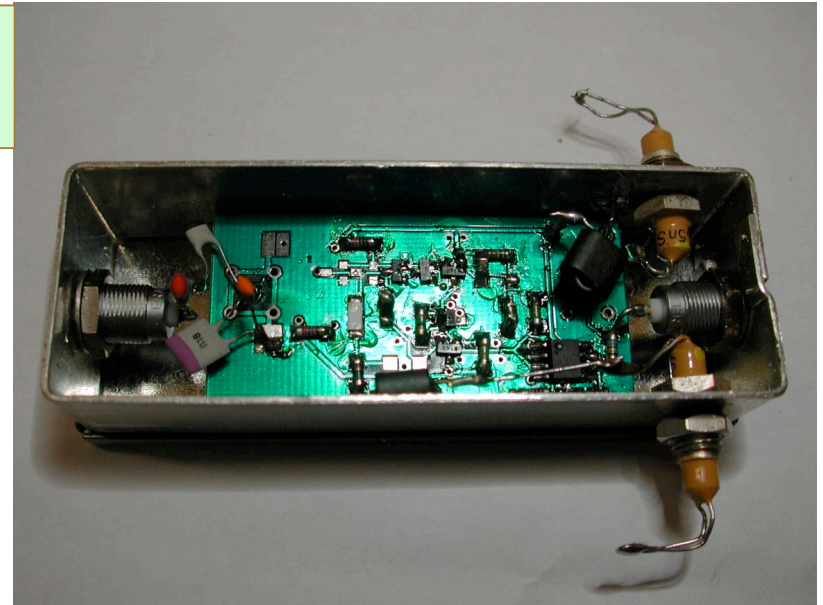


Effect of ionization density on the time dependence of luminescence !

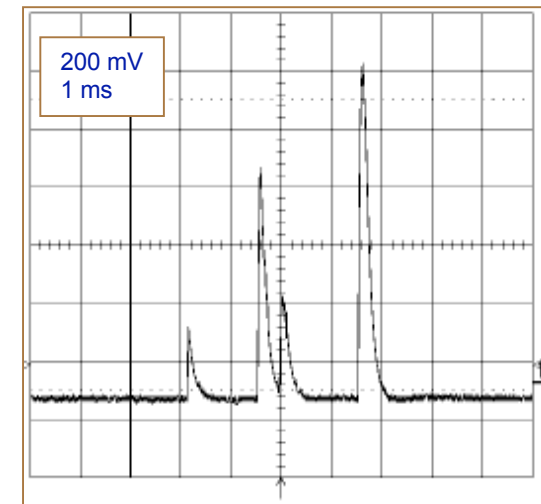
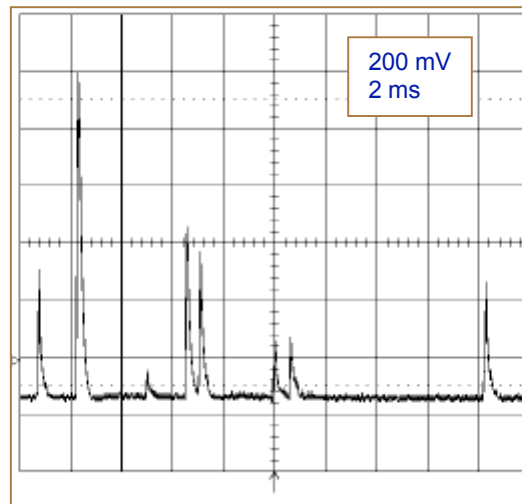
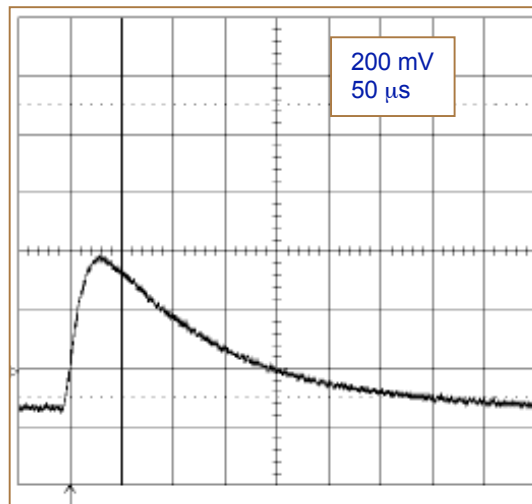
Two-stage LEM: measurements (preliminary)

Custom-made front-end charge preamplifier (3.3 mV/fC)

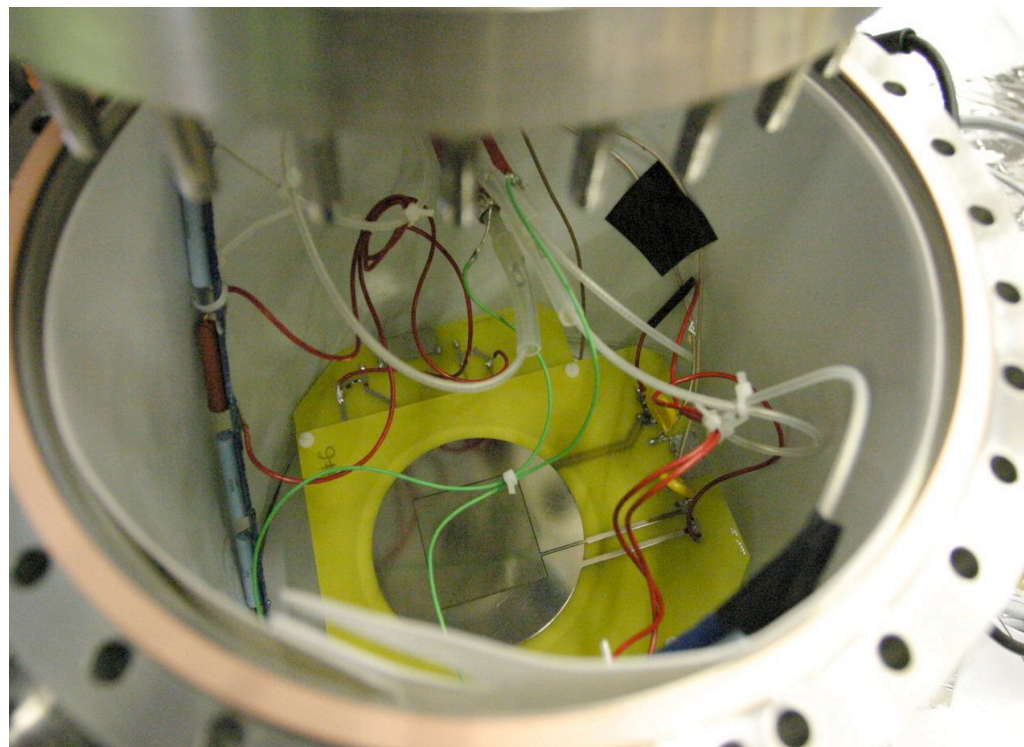
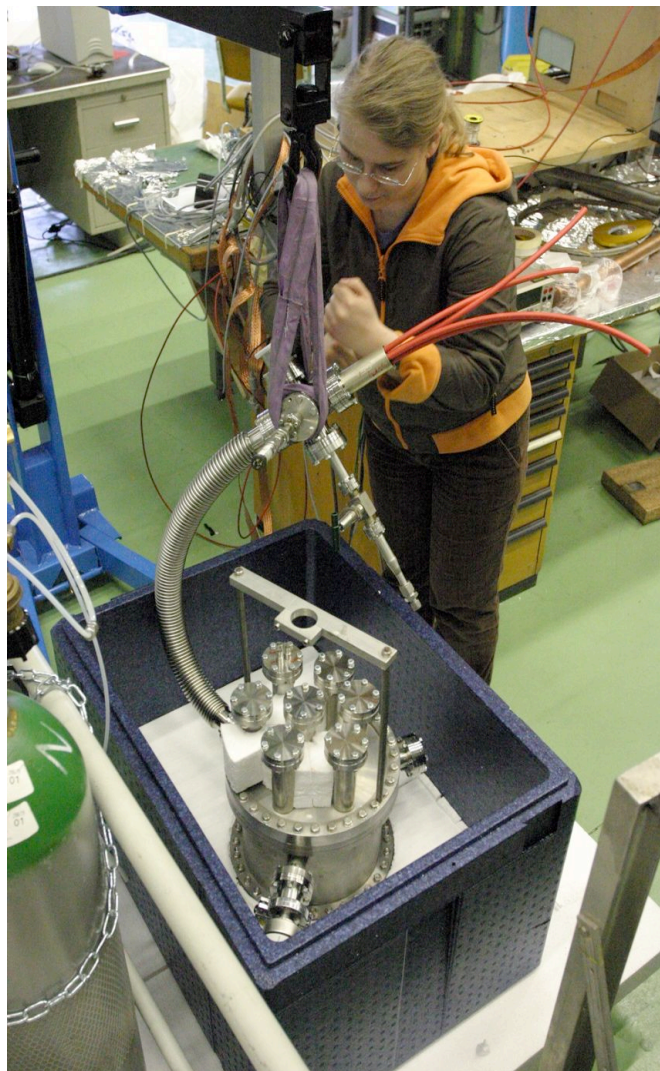
inspired from C.Boiano et al. IEEE Transact. on Nucl. Science, Vol. 51 N°5 2004



- Shapes of the signals from double-stage LEM system
- Average signal rise time: 12 μ s.
- Signals from Fe^{55} radioactive source (5.8 keV), event rate about 1 kHz.



Two-stage LEM test setup



Considered PMT	Hamamatsu R6237-01MOD*
Shape	Square
Size	$7.6 \times 7.6 \text{ cm}^2$
Peak sensitivity	420 nm
Window	Borosilicate
Cathode type	Bialkali
Typ. Gain	3×10^5 (8 dynodes)
Quantum Eff.	$\sim 20\%$ (430 nm)

