Nabil Chouika

Irfu/DPhN, CEA Saclay - Université Paris-Saclay

GDR QCD, Orme des Merisiers, 5 Décembre, 2017









Outline

- Introduction to Generalized Parton Distributions
 - Definition and properties
 - Experimental access
- Representations of Generalized Parton Distributions
 - Overlap of Light-cone wave functions
 - Double Distributions
- Covariant extention of Generalized Parton Distributions
 - Motivation
 - Inversion of Incomplete Radon Transform
 - Results
- 4 Conclusion



Introduction to GPDs

- Introduction to Generalized Parton Distributions
 - Definition and properties
 - Experimental access
- - Overlap of Light-cone wave functions
 - Double Distributions
- - Motivation
 - Inversion of Incomplete Radon Transform
 - Results

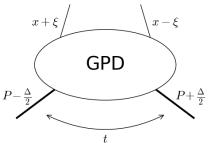


Introduction to GPDs

0000

Quark GPD (twist-2, spin-0 hadron): (Müller et al., 1994; Radyushkin, 1996; Ji, 1997)

$$H^{q}\left(x,\xi,t\right) = \frac{1}{2} \int \frac{\mathrm{d}z^{-}}{2\pi} e^{ixP^{+}z^{-}} \left\langle P + \frac{\Delta}{2} \left| \bar{q}\left(-z\right) \gamma^{+} q\left(z\right) \right| P - \frac{\Delta}{2} \right\rangle \right|_{z^{+}=0, z_{\perp}=0}.$$
with: (1)



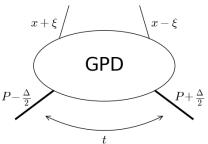
$$t=\Delta^2 \;\;,\;\; \xi=-rac{\Delta^+}{2\,P^+} \;.$$

Introduction to GPDs

0000

Quark GPD (twist-2, spin-0 hadron): (Müller et al., 1994; Radyushkin, 1996; Ji, 1997)

$$H^{q}\left(x,\xi,t\right) = \frac{1}{2} \int \frac{\mathrm{d}z^{-}}{2\pi} e^{i x P^{+}z^{-}} \left\langle P + \frac{\Delta}{2} \left| \bar{q}\left(-z\right) \gamma^{+} q\left(z\right) \right| P - \frac{\Delta}{2} \right\rangle \right|_{z^{+}=0, z_{\perp}=0}.$$
with: (1)



$$t = \Delta^2$$
, $\xi = -\frac{\Delta^+}{2P^+}$.

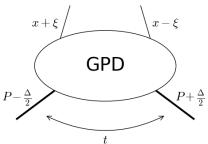
Similar matrix element for gluons.

Introduction to GPDs

0000

Quark GPD (twist-2, spin-0 hadron): (Müller et al., 1994; Radyushkin, 1996; Ji, 1997)

$$H^{q}\left(x,\xi,t\right) = \frac{1}{2} \int \frac{\mathrm{d}z^{-}}{2\pi} e^{i \times P^{+}z^{-}} \left\langle P + \frac{\Delta}{2} \left| \bar{q}\left(-z\right) \gamma^{+} q\left(z\right) \right| P - \frac{\Delta}{2} \right\rangle \right|_{z^{+}=0, z_{\perp}=0}.$$
with: (1)



$$t = \Delta^2 \;\; , \;\; \xi = -rac{\Delta^+}{2\,P^+} \; .$$

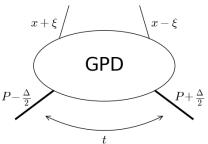
- Similar matrix element for gluons.
- More GPDs for spin-¹/₂ hadrons.

Introduction to GPDs

0000

Quark GPD (twist-2, spin-0 hadron): (Müller et al., 1994; Radyushkin, 1996; Ji, 1997)

$$H^{q}\left(x,\xi,t\right) = \frac{1}{2} \int \frac{\mathrm{d}z^{-}}{2\pi} e^{i x P^{+}z^{-}} \left\langle P + \frac{\Delta}{2} \left| \bar{q}\left(-z\right) \gamma^{+} q\left(z\right) \right| P - \frac{\Delta}{2} \right\rangle \right|_{z^{+}=0, z_{\perp}=0}.$$
with: (1)



$$t = \Delta^2$$
, $\xi = -\frac{\Delta^+}{2P^+}$.

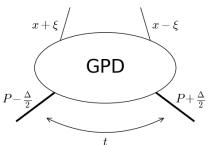
- Similar matrix element for gluons.
- More GPDs for spin-¹/₂ hadrons.
- Experimental programs at JLab, COMPASS.

Introduction to GPDs

Nabil Chouika

Quark GPD (twist-2, spin-0 hadron): (Müller et al., 1994; Radyushkin, 1996; Ji, 1997)

$$H^{q}(x,\xi,t) = \frac{1}{2} \int \frac{\mathrm{d}z^{-}}{2\pi} e^{ixP^{+}z^{-}} \left\langle P + \frac{\Delta}{2} \left| \bar{q}(-z) \gamma^{+} q(z) \right| P - \frac{\Delta}{2} \right\rangle \Big|_{z^{+}=0, z_{\perp}=0}.$$
with:
$$(1)$$



$$t=\Delta^2 \;\;,\;\; \xi=-rac{\Delta^+}{2\,P^+} \;.$$

- Similar matrix element for gluons.
- More GPDs for spin-¹/₂ hadrons.
- Experimental programs at JLab, COMPASS.

GDR QCD, 05/12/17

4 / 22

Impact parameter space GPD (at $\xi = 0$): (Burkardt, 2000)

$$q\left(x,\vec{b_{\perp}}\right) = \int \frac{\mathrm{d}^{2}\vec{\Delta_{\perp}}}{(2\pi)^{2}} e^{-i\vec{b_{\perp}}\cdot\vec{\Delta_{\perp}}} H^{q}\left(x,0,-\vec{\Delta_{\perp}}^{2}\right). \tag{2}$$

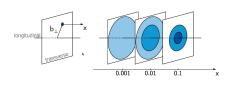
Covariant extension of GPDs

4 / 22

Introduction to GPDs

Quark GPD (twist-2, spin-0 hadron): (Müller et al., 1994; Radyushkin, 1996; Ji, 1997)

$$H^{q}(x,\xi,t) = \frac{1}{2} \int \frac{\mathrm{d}z^{-}}{2\pi} e^{i \times P^{+}z^{-}} \left\langle P + \frac{\Delta}{2} \left| \bar{q}(-z) \gamma^{+} q(z) \right| P - \frac{\Delta}{2} \right\rangle \Big|_{z^{+}=0, z_{\perp}=0}.$$
with: (1)



$$t=\Delta^2 \;\;,\;\; \xi=-rac{\Delta^+}{2\,P^+} \,.$$

- Similar matrix element for gluons.
- More GPDs for spin-¹/₂ hadrons.
- Experimental programs at JLab, COMPASS.
- Impact parameter space GPD (at $\xi=0$): (Burkardt, 2000)

$$q\left(x,\vec{b_{\perp}}\right) = \int \frac{\mathrm{d}^{2}\vec{\Delta_{\perp}}}{(2\pi)^{2}} e^{-i\vec{b_{\perp}}\cdot\vec{\Delta_{\perp}}} H^{q}\left(x,0,-\vec{\Delta_{\perp}}^{2}\right). \tag{2}$$

Nabil Chouika Covariant extension of GPDs GDR QCD, 05/12/17

Main properties:

0000

• Physical region: $(x, \xi) \in [-1, 1]^2$.

Main properties:

- Physical region: $(x,\xi) \in [-1,1]^2$.
 - ▶ DGLAP: $|x| > |\xi|$.

Main properties:

- Physical region: $(x,\xi) \in [-1,1]^2$.
 - ▶ DGLAP: $|x| > |\xi|$.
 - ▶ ERBL: $|x| < |\xi|$.

5 / 22

Main properties:

Introduction to GPDs

0000

- Physical region: $(x, \xi) \in [-1, 1]^2$.
 - ▶ DGLAP: $|x| > |\xi|$.
 - ▶ ERBL: $|x| < |\xi|$.
- Link to PDFs and Form Factors:

$$\int \mathrm{d}x \, H^q(x,\xi,t) = F^q(t) \;, \tag{3}$$

$$H^{q}(x,0,0) = \theta(x) q(x) - \theta(-x) \bar{q}(-x). \tag{4}$$

Main properties:

Introduction to GPDs

- Physical region: $(x, \xi) \in [-1, 1]^2$.
 - ▶ DGLAP: $|x| > |\xi|$.
 - ▶ ERBL: $|x| < |\xi|$.
- Link to PDFs and Form Factors:

$$\int \mathrm{d}x \, H^q(x,\xi,t) = F^q(t) \;, \tag{3}$$

$$H^{q}(x,0,0) = \theta(x) q(x) - \theta(-x) \bar{q}(-x)$$
 (4)

Polynomiality:

$$\int_{-1}^{1} \mathrm{d}x \, x^{m} \, H(x, \xi, t) = \text{Polynomial in } \xi. \tag{5}$$

Main properties:

Introduction to GPDs

- Physical region: $(x, \xi) \in [-1, 1]^2$.
 - ▶ DGLAP: $|x| > |\xi|$.
 - ▶ ERBL: $|x| < |\xi|$.
- Link to PDFs and Form Factors:

$$\int dx H^{q}(x,\xi,t) = F^{q}(t) , \qquad (3)$$

$$H^{q}(x,0,0) = \theta(x) q(x) - \theta(-x) \bar{q}(-x).$$

$$(4)$$

Polynomiality:

$$\int_{-1}^{1} \mathrm{d}x \, x^{m} \, H(x, \xi, t) = \text{Polynomial in } \xi.$$
 (5)

From Lorentz invariance.

Representations of GPDs

Main properties:

- Physical region: $(x, \xi) \in [-1, 1]^2$.
 - ▶ DGLAP: $|x| > |\xi|$.
 - ▶ ERBL: $|x| < |\xi|$.
- Link to PDFs and Form Factors:

$$\int dx H^{q}(x,\xi,t) = F^{q}(t) , \qquad (3)$$

$$H^{q}(x,0,0) = \theta(x) q(x) - \theta(-x) \bar{q}(-x).$$

$$\tag{4}$$

Polynomiality:

$$\int_{-1}^{1} dx \, x^{m} H(x, \xi, t) = \text{Polynomial in } \xi.$$
 (5)

- From Lorentz invariance.
- Positivity (in DGLAP): (Pire et al., 1999; Radyushkin, 1999)

$$|H^{q}(x,\xi,t)| \leq \sqrt{q\left(\frac{x-\xi}{1-\xi}\right)q\left(\frac{x+\xi}{1+\xi}\right)}.$$
 (6)

Main properties:

- Physical region: $(x, \xi) \in [-1, 1]^2$.
 - ▶ DGLAP: $|x| > |\xi|$.
 - ▶ ERBL: $|x| < |\xi|$.
- Link to PDFs and Form Factors:

$$\int dx H^{q}(x,\xi,t) = F^{q}(t) , \qquad (3)$$

$$H^{q}(x,0,0) = \theta(x) q(x) - \theta(-x) \bar{q}(-x).$$
(4)

Polynomiality:

$$\int_{-1}^{1} dx \, x^{m} H(x, \xi, t) = \text{Polynomial in } \xi.$$
 (5)

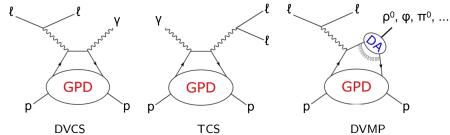
- From Lorentz invariance.
- Positivity (in DGLAP): (Pire et al., 1999; Radyushkin, 1999)

$$|H^{q}(x,\xi,t)| \leq \sqrt{q\left(\frac{x-\xi}{1-\xi}\right)q\left(\frac{x+\xi}{1+\xi}\right)}.$$
 (6)

Cauchy-Schwarz theorem in Hilbert space.

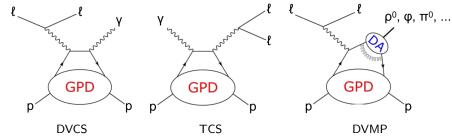
Nabil Chouika Covariant extension of GPDs GDR QCD, 05/12/17 5 / 22

Exclusive processes:



Introduction to GPDs

Exclusive processes:

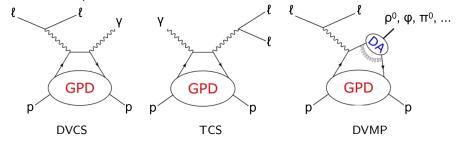


Compton Form Factors: (Belitsky et al., 2002)

$$\mathcal{F}\left(\xi, t, Q^{2}\right) = \int_{-1}^{1} dx \, C\left(x, \xi, \alpha_{S}\left(\mu_{F}\right), \frac{Q}{\mu_{F}}\right) F\left(x, \xi, t, \mu_{F}\right). \tag{7}$$

Introduction to GPDs

Exclusive processes:



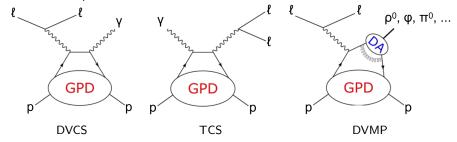
Compton Form Factors: (Belitsky et al., 2002)

$$\mathcal{F}\left(\xi, t, Q^{2}\right) = \int_{-1}^{1} dx \, C\left(x, \xi, \alpha_{S}\left(\mu_{F}\right), \frac{Q}{\mu_{F}}\right) F\left(x, \xi, t, \mu_{F}\right). \tag{7}$$

Observables are convolutions of:

Introduction to GPDs

Exclusive processes:

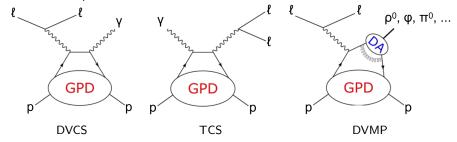


Compton Form Factors: (Belitsky et al., 2002)

$$\mathcal{F}\left(\xi, t, Q^{2}\right) = \int_{-1}^{1} dx \, C\left(x, \xi, \alpha_{S}\left(\mu_{F}\right), \frac{Q}{\mu_{F}}\right) F\left(x, \xi, t, \mu_{F}\right). \tag{9}$$

- Observables are convolutions of:
 - a soft part, i.e. the GPD, with long distance interactions encoded (non-perturbative QCD).

Exclusive processes:



Compton Form Factors: (Belitsky et al., 2002)

$$\mathcal{F}\left(\xi, t, Q^{2}\right) = \int_{-1}^{1} dx \, \mathbf{C}\left(x, \xi, \alpha_{S}\left(\mu_{F}\right), \frac{Q}{\mu_{F}}\right) F\left(x, \xi, t, \mu_{F}\right). \tag{8}$$

- Observables are convolutions of:
 - a soft part, i.e. the GPD, with long distance interactions encoded (non-perturbative QCD).
 - a hard-scattering kernel, calculated with perturbative QCD (short distance interactions).

Nabil Chouika

Outline

Representations of GPDs

- Definition and properties
- Experimental access
- Representations of Generalized Parton Distributions
 - Overlap of Light-cone wave functions
 - Double Distributions
- - Motivation
 - Inversion of Incomplete Radon Transform
 - Results



A given hadronic state is decomposed in a Fock basis: (Brodsky et al., 1981)

$$|H; P, \lambda\rangle = \sum_{N,\beta} \int [\mathrm{d}x]_N \left[\mathrm{d}^2\mathbf{k}_\perp\right]_N \Psi_{N,\beta}^{\lambda} \left(x_1, \mathbf{k}_{\perp 1}, ...\right) |N, \beta; k_1, ..., k_N\rangle , \qquad (10)$$

where the $\Psi_{N,\beta}^{\lambda}$ are the Light-front wave functions (LFWF).

Introduction to GPDs

• A given hadronic state is decomposed in a Fock basis: (Brodsky et al., 1981)

$$|H; P, \lambda\rangle = \sum_{N,\beta} \int [\mathrm{d}x]_N \left[\mathrm{d}^2 \mathbf{k}_\perp \right]_N \Psi_{N,\beta}^{\lambda} \left(x_1, \mathbf{k}_{\perp 1}, ... \right) |N, \beta; k_1, ..., k_N \rangle , \qquad (10)$$

where the $\Psi^{\lambda}_{N,\beta}$ are the Light-front wave functions (LFWF).

• For example, for the pion:

$$\left|\pi^{+}\right\rangle = \psi_{u\bar{d}}^{\pi} \left|u\bar{d}\right\rangle + \psi_{u\bar{d}g}^{\pi} \left|u\bar{d}g\right\rangle + \dots \tag{11}$$

Conclusion

Overlap of Light-cone wave functions

A given hadronic state is decomposed in a Fock basis: (Brodsky et al., 1981)

$$|H;P,\lambda\rangle = \sum_{N,\beta} \int [\mathrm{d}x]_N \left[\mathrm{d}^2\mathbf{k}_\perp\right]_N \Psi_{N,\beta}^{\lambda} \left(x_1,\mathbf{k}_{\perp 1},...\right) |N,\beta;k_1,...,k_N\rangle , \qquad (10)$$

where the $\Psi_{N.\beta}^{\lambda}$ are the Light-front wave functions (LFWF).

• For example, for the pion:

$$\left|\pi^{+}\right\rangle = \psi_{u\bar{d}}^{\pi} \left|u\bar{d}\right\rangle + \psi_{u\bar{d}g}^{\pi} \left|u\bar{d}g\right\rangle + \dots \tag{11}$$

GPD as an overlap of LFWFs: (Diehl et al., 2001; Diehl, 2003)

$$H^{q}(x,\xi,t) = \sum_{N,\beta} \sqrt{1-\xi^{2-N}} \sqrt{1+\xi^{2-N}} \sum_{a} \delta_{a,q}$$
 (12)

$$\times \quad \int \left[\mathrm{d}\bar{x}\right]_{\textit{N}} \left[\mathrm{d}^{2}\bar{\boldsymbol{k}}_{\perp}\right]_{\textit{N}} \delta\left(x-\bar{x}_{\textit{a}}\right) \Psi_{\textit{N},\beta}^{*} \left(\hat{x}_{1}^{'},\hat{\boldsymbol{k}}_{\perp1}^{'},...\right) \Psi_{\textit{N},\beta} \left(\tilde{x}_{1},\tilde{\boldsymbol{k}}_{\perp1},...\right) \,,$$

in the DGLAP region $\xi < x < 1$ (pion case).

Introduction to GPDs

Conclusion

Overlap of Light-cone wave functions

A given hadronic state is decomposed in a Fock basis: (Brodsky et al., 1981)

$$|H;P,\lambda\rangle = \sum_{N,\beta} \int [\mathrm{d}x]_N \left[\mathrm{d}^2\mathbf{k}_\perp\right]_N \Psi_{N,\beta}^{\lambda} \left(x_1,\mathbf{k}_{\perp 1},...\right) |N,\beta;k_1,...,k_N\rangle , \qquad (10)$$

where the $\Psi_{N,\beta}^{\lambda}$ are the Light-front wave functions (**LFWF**).

• For example, for the pion:

$$\left|\pi^{+}\right\rangle = \psi_{u\bar{d}}^{\pi} \left|u\bar{d}\right\rangle + \psi_{u\bar{d}g}^{\pi} \left|u\bar{d}g\right\rangle + \dots \tag{11}$$

GPD as an overlap of LFWFs: (Diehl et al., 2001; Diehl, 2003)

$$H^{q}(x,\xi,t) = \sum_{N,\beta} \sqrt{1-\xi^{2-N}} \sqrt{1+\xi^{2-N}} \sum_{a} \delta_{a,q}$$
 (12)

$$\times \int [\mathrm{d}\bar{\mathbf{x}}]_{N} \left[\mathrm{d}^{2}\bar{\mathbf{k}}_{\perp}\right]_{N} \delta\left(\mathbf{x} - \bar{\mathbf{x}}_{\mathsf{a}}\right) \Psi_{N,\beta}^{*} \left(\hat{\mathbf{x}}_{1}^{'}, \hat{\mathbf{k}}_{\perp 1}^{'}, ...\right) \Psi_{N,\beta} \left(\tilde{\mathbf{x}}_{1}, \tilde{\mathbf{k}}_{\perp 1}, ...\right),$$

in the DGLAP region $\xi < x < 1$ (pion case).

Similar result in ERBL $(-\xi < x < \xi)$, but with N and N + 2...

Introduction to GPDs

• A given hadronic state is decomposed in a Fock basis: (Brodsky et al., 1981)

$$|H;P,\lambda\rangle = \sum_{N,\beta} \int [\mathrm{d}x]_N \left[\mathrm{d}^2\mathbf{k}_\perp\right]_N \Psi_{N,\beta}^{\lambda} \left(x_1,\mathbf{k}_{\perp 1},...\right) |N,\beta;k_1,...,k_N\rangle , \qquad (10)$$

where the $\Psi_{N,\beta}^{\lambda}$ are the *Light-front wave functions* (**LFWF**).

• For example, for the pion:

$$\left|\pi^{+}\right\rangle = \psi_{u\bar{d}}^{\pi} \left|u\bar{d}\right\rangle + \psi_{u\bar{d}g}^{\pi} \left|u\bar{d}g\right\rangle + \dots \tag{11}$$

GPD as an overlap of LFWFs: (Diehl et al., 2001; Diehl, 2003)

$$H^{q}(x,\xi,t) = \sum_{N,\beta} \sqrt{1-\xi^{2-N}} \sqrt{1+\xi^{2-N}} \sum_{a} \delta_{a,q}$$
 (12)

$$\times \int [\mathrm{d}\bar{\mathbf{x}}]_{N} \left[\mathrm{d}^{2}\bar{\mathbf{k}}_{\perp}\right]_{N} \delta\left(\mathbf{x} - \bar{\mathbf{x}}_{a}\right) \Psi_{N,\beta}^{*} \left(\hat{\mathbf{x}}_{1}^{'}, \hat{\mathbf{k}}_{\perp 1}^{'}, ...\right) \Psi_{N,\beta} \left(\tilde{\mathbf{x}}_{1}, \tilde{\mathbf{k}}_{\perp 1}, ...\right),$$

in the DGLAP region $\xi < x < 1$ (pion case).

- Similar result in ERBL $(-\xi < x < \xi)$, but with N and N + 2...
- GPD is a scalar product of LFWFs:

4 D > 4 A > 4 B > 4 B > B | E | 9 Q ()

• A given hadronic state is decomposed in a Fock basis: (Brodsky et al., 1981)

$$|H;P,\lambda\rangle = \sum_{N,\beta} \int [\mathrm{d}x]_N \left[\mathrm{d}^2\mathbf{k}_\perp\right]_N \Psi_{N,\beta}^{\lambda} \left(x_1,\mathbf{k}_{\perp 1},...\right) |N,\beta;k_1,...,k_N\rangle , \qquad (10)$$

where the $\Psi_{N,\beta}^{\lambda}$ are the *Light-front wave functions* (**LFWF**).

• For example, for the pion:

$$\left|\pi^{+}\right\rangle = \psi_{u\bar{d}}^{\pi} \left|u\bar{d}\right\rangle + \psi_{u\bar{d}g}^{\pi} \left|u\bar{d}g\right\rangle + \dots \tag{11}$$

GPD as an overlap of LFWFs: (Diehl et al., 2001; Diehl, 2003)

$$H^{q}(x,\xi,t) = \sum_{N,\beta} \sqrt{1-\xi^{2-N}} \sqrt{1+\xi^{2-N}} \sum_{a} \delta_{a,q}$$
 (12)

$$\times \int [\mathrm{d}\bar{x}]_{N} \left[\mathrm{d}^{2}\bar{\mathbf{k}}_{\perp}\right]_{N} \delta\left(x-\bar{x}_{a}\right) \Psi_{N,\beta}^{*}\left(\hat{x}_{1}^{'},\hat{\mathbf{k}}_{\perp 1}^{'},...\right) \Psi_{N,\beta}\left(\tilde{x}_{1},\tilde{\mathbf{k}}_{\perp 1},...\right),$$

in the DGLAP region $\xi < x < 1$ (pion case).

- Similar result in ERBL $(-\xi < x < \xi)$, but with N and N + 2...
- GPD is a scalar product of LFWFs:
 - ► Cauchy-Schwarz theorem ⇒ **Positivity** fulfilled!

- ◀ □ ▶ ◀♬ ▶ ◀ 툴 ▶ 세 툴 ▶ 포 [표 150 Q C

8 / 22

A given hadronic state is decomposed in a Fock basis: (Brodsky et al., 1981)

$$|H; P, \lambda\rangle = \sum_{N,\beta} \int [\mathrm{d}x]_N \left[\mathrm{d}^2\mathbf{k}_\perp\right]_N \Psi_{N,\beta}^{\lambda} \left(x_1, \mathbf{k}_{\perp 1}, ...\right) |N, \beta; k_1, ..., k_N\rangle , \qquad (10)$$

where the $\Psi_{N.\beta}^{\lambda}$ are the *Light-front wave functions* (**LFWF**).

• For example, for the pion:

$$\left|\pi^{+}\right\rangle = \psi_{u\bar{d}}^{\pi} \left|u\bar{d}\right\rangle + \psi_{u\bar{d}g}^{\pi} \left|u\bar{d}g\right\rangle + \dots \tag{11}$$

GPD as an overlap of LFWFs: (Diehl et al., 2001; Diehl, 2003)

$$H^{q}(x,\xi,t) = \sum_{N,\beta} \sqrt{1-\xi^{2-N}} \sqrt{1+\xi^{2-N}} \sum_{a} \delta_{a,q}$$
 (12)

$$\times \int [\mathrm{d}\bar{x}]_{N} \left[\mathrm{d}^{2}\bar{\mathbf{k}}_{\perp}\right]_{N} \delta\left(x-\bar{x}_{a}\right) \Psi_{N,\beta}^{*}\left(\hat{x}_{1}^{'},\hat{\mathbf{k}}_{\perp 1}^{'},...\right) \Psi_{N,\beta}\left(\tilde{x}_{1},\tilde{\mathbf{k}}_{\perp 1},...\right),$$

in the DGLAP region $\xi < x < 1$ (pion case).

- Similar result in ERBL $(-\xi < x < \xi)$, but with N and N + 2...
- GPD is a scalar product of LFWFs:
 - ► Cauchy-Schwarz theorem ⇒ Positivity fulfilled!
 - Polynomiality not manifest...

DD representation of GPDs:

$$H(x,\xi,t) \propto \int_{\Omega} \mathrm{d}\beta \,\mathrm{d}\alpha \,h(\beta,\alpha,t) \,\delta(x-\beta-\alpha\xi)$$
 (13)

DD representation of GPDs:

$$H(x,\xi,t) \propto \int_{\Omega} \mathrm{d}\beta \,\mathrm{d}\alpha \,h(\beta,\alpha,t) \,\delta(x-\beta-\alpha\xi) \;.$$
 (13)

• DD *h* is defined on the support $\Omega = \{|\beta| + |\alpha| \le 1\}$.

DD representation of GPDs:

$$H(x,\xi,t) \propto \int_{\Omega} \mathrm{d}\beta \,\mathrm{d}\alpha \,h(\beta,\alpha,t) \,\delta(x-\beta-\alpha\xi) \ .$$
 (13)

- DD h is defined on the support $\Omega = \{|\beta| + |\alpha| \le 1\}$.
- **Polynomial** in ξ :

$$\int_{-1}^{1} dx \, x^{m} \, H(x, \xi, t) \quad \propto \quad \int dx \, x^{m} \int_{\Omega} d\beta \, d\alpha \, h(\beta, \alpha, t) \, \delta(x - \beta - \alpha \xi)$$

$$\propto \quad \int_{\Omega} d\beta \, d\alpha \, (\beta + \xi \alpha)^{m} \, h(\beta, \alpha, t) . \tag{14}$$

• DD representation of GPDs:

$$H(x,\xi,t) \propto \int_{\Omega} \mathrm{d}\beta \,\mathrm{d}\alpha \,h(\beta,\alpha,t) \,\delta(x-\beta-\alpha\xi) \ .$$
 (13)

- DD h is defined on the support $\Omega = \{|\beta| + |\alpha| \le 1\}$.
- **Polynomial** in *ξ*:

$$\int_{-1}^{1} dx \, x^{m} \, H(x, \xi, t) \quad \propto \quad \int dx \, x^{m} \int_{\Omega} d\beta \, d\alpha \, h(\beta, \alpha, t) \, \delta(x - \beta - \alpha \xi)$$

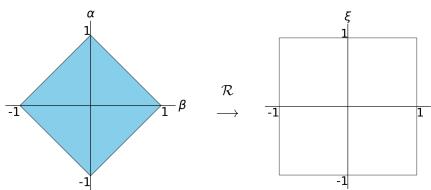
$$\propto \quad \int_{\Omega} d\beta \, d\alpha \, (\beta + \xi \alpha)^{m} \, h(\beta, \alpha, t) \, . \tag{14}$$

Positivity not manifest...

• DD representation of GPDs:

$$H(x,\xi,t) \propto \int_{\Omega} \mathrm{d}\beta \,\mathrm{d}\alpha \,h(\beta,\alpha,t) \,\delta(x-\beta-\alpha\xi)$$
 (13)

• Radon Transform: (Deans, 1983; Teryaev, 2001)

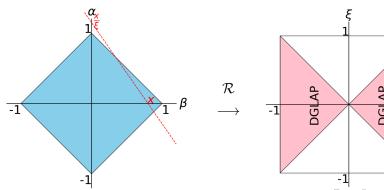


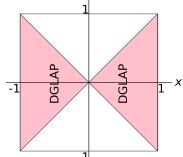
Nabil Chouika

DD representation of GPDs:

$$H(x,\xi,t) \propto \int_{\Omega} \mathrm{d}\beta \,\mathrm{d}\alpha \,h(\beta,\alpha,t) \,\delta(x-\beta-\alpha\xi) \ .$$
 (13)

Radon Transform: (Deans, 1983; Teryaev, 2001)



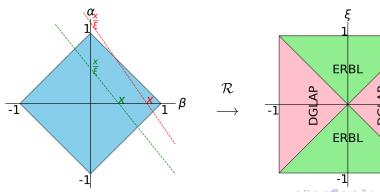


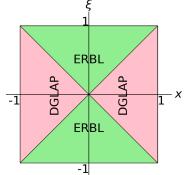
Double Distributions (DDs)

DD representation of GPDs:

$$H(x,\xi,t) \propto \int_{\Omega} \mathrm{d}\beta \,\mathrm{d}\alpha \,h(\beta,\alpha,t) \,\delta(x-\beta-\alpha\xi) \ .$$
 (13)

Radon Transform: (Deans, 1983; Teryaev, 2001)



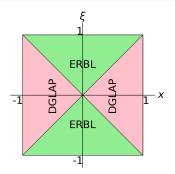


Outline

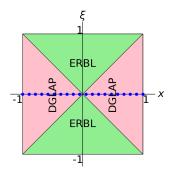
- Introduction to Generalized Parton Distributions
 - Definition and properties
 - Experimental access
- 2 Representations of Generalized Parton Distributions
 - Overlap of Light-cone wave functions
 - Double Distributions
- Covariant extention of Generalized Parton Distributions
 - Motivation
 - Inversion of Incomplete Radon Transform
 - Results
- 4 Conclusion



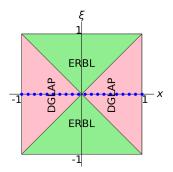
For the spatial tomography of hadrons, we need:



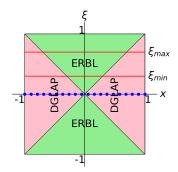
- For the spatial tomography of hadrons, we need:
 - ► GPD at ξ = 0.



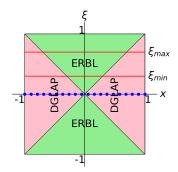
- For the spatial tomography of hadrons, we need:
 - ► GPD at ξ = 0.
- Experimental access through exclusive processes:



- For the spatial tomography of hadrons, we need:
 - ► GPD at ξ = 0.
 - Experimental access through exclusive processes:
 - ▶ Integrals over x of GPD at $\xi > 0$.

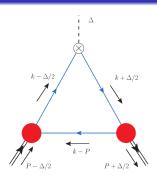


- For the spatial tomography of hadrons, we need:
 - ► GPD at ξ = 0.
 - Experimental access through exclusive processes:
 - Integrals over x of GPD at $\xi > 0$.
 - Extrapolation?



0000000

- For the spatial tomography of hadrons, we need:
 - ► GPD at ξ = 0.
 - Experimental access through exclusive processes:
 - Integrals over x of GPD at $\xi > 0$.
 - Extrapolation?
 - Covariant approach?



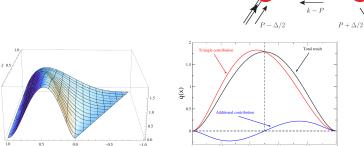
 $k - \Delta/2$

 $k + \Delta/2$

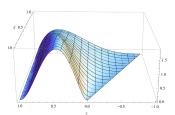
Covariant extension

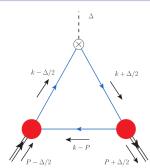
Introduction to GPDs

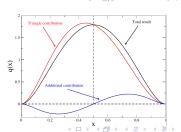
- For the spatial tomography of hadrons, we need:
 - ► GPD at ξ = 0.
 - Experimental access through exclusive processes:
 - Integrals over x of GPD at $\xi > 0$.
 - Extrapolation?
 - Covariant approach?
 - Positivity?



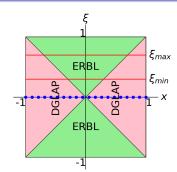
- For the spatial tomography of hadrons, we need:
 - ► GPD at ξ = 0.
 - Experimental access through exclusive processes:
 - Integrals over x of GPD at $\xi > 0$.
 - Extrapolation?
 - Covariant approach?
 - Positivity?
 - Loss of symmetries...



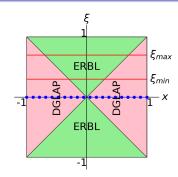


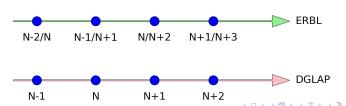


- For the spatial tomography of hadrons, we need:
 - ► GPD at ξ = 0.
 - Experimental access through exclusive processes:
 - Integrals over x of GPD at $\xi > 0$.
 - Extrapolation?
 - Light-front wave functions approach:

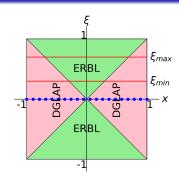


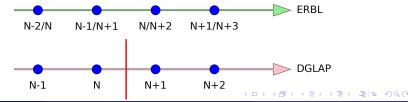
- For the spatial tomography of hadrons, we need:
 - ▶ GPD at $\xi = 0$.
 - Experimental access through exclusive processes:
 - Integrals over x of GPD at $\xi > 0$.
 - Extrapolation?
 - Light-front wave functions approach:
 - Derive all parton distributions (not just GPDs)!



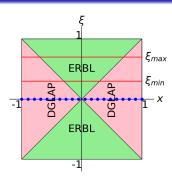


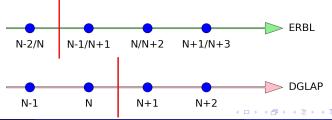
- For the spatial tomography of hadrons, we need:
 - ► GPD at *ξ* = 0.
 - Experimental access through exclusive processes:
 - Integrals over x of GPD at $\xi > 0$.
 - Extrapolation?
 - **Light-front wave functions** approach:
 - Derive all parton distributions (not just GPDs)!
 - But how to truncate?



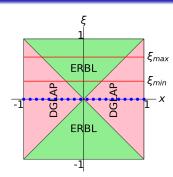


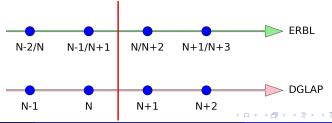
- For the spatial tomography of hadrons, we need:
 - ► GPD at *ξ* = 0.
 - Experimental access through exclusive processes:
 - Integrals over x of GPD at $\xi > 0$.
 - Extrapolation?
 - **Light-front wave functions** approach:
 - Derive all parton distributions (not just GPDs)!
 - But how to truncate?



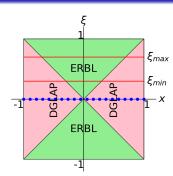


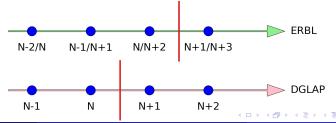
- For the spatial tomography of hadrons, we need:
 - ► GPD at *ξ* = 0.
 - Experimental access through exclusive processes:
 - Integrals over x of GPD at $\xi > 0$.
 - Extrapolation?
 - **Light-front wave functions** approach:
 - Derive all parton distributions (not just GPDs)!
 - But how to truncate?





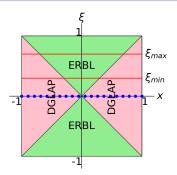
- For the spatial tomography of hadrons, we need:
 - ► GPD at *ξ* = 0.
 - Experimental access through exclusive processes:
 - Integrals over x of GPD at $\xi > 0$.
 - Extrapolation?
 - **Light-front wave functions** approach:
 - Derive all parton distributions (not just GPDs)!
 - But how to truncate?

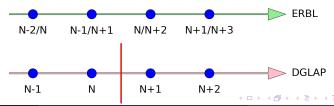




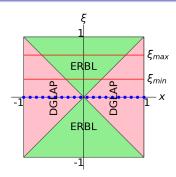
0000000

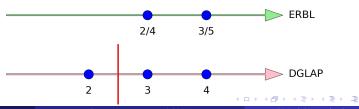
- For the spatial tomography of hadrons, we need:
 - ► GPD at *ξ* = 0.
 - Experimental access through exclusive processes:
 - ▶ Integrals over x of GPD at $\xi > 0$.
 - Extrapolation?
 - **Light-front wave functions** approach:
 - Derive all parton distributions (not just GPDs)!
 - But how to truncate?
 - Use Lorentz invariance to extend from DGLAP! (Hwang and Mueller, 2008; Chouika et al., 2017)





- For the spatial tomography of hadrons, we need:
 - ► GPD at *ξ* = 0.
 - Experimental access through exclusive processes:
 - ▶ Integrals over x of GPD at $\xi > 0$.
 - Extrapolation?
 - **Light-front wave functions** approach:
 - Derive all parton distributions (not just GPDs)!
 - But how to truncate?
- Use Lorentz invariance to extend from DGLAP! (Hwang and Mueller, 2008; Chouika et al., 2017)





Problem

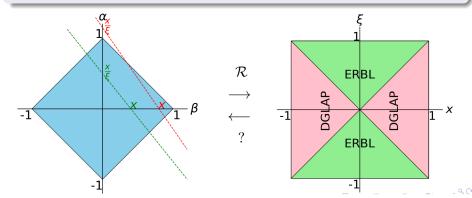
Find $h(\beta, \alpha)$ on square $\{|\alpha| + |\beta| \le 1\}$ such that

$$H(x,\xi)|_{\mathrm{DGLAP}} \propto \int \mathrm{d}\beta\,\mathrm{d}\alpha\,h(\beta,\alpha)\,\delta(x-\beta-\alpha\xi)$$
.

Problem

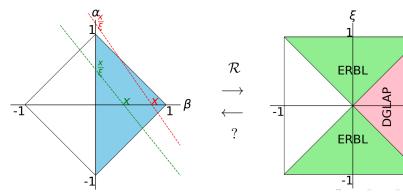
Find $h(\beta, \alpha)$ on square $\{|\alpha| + |\beta| \le 1\}$ such that

$$H(x,\xi)|_{\mathrm{DGLAP}} \propto \int \mathrm{d}\beta\,\mathrm{d}\alpha\,h(\beta,\alpha)\,\delta(x-\beta-\alpha\xi)$$
.



Introduction to GPDs

• Quark GPD: $H(x,\xi) = 0$ for $-1 < x < -|\xi| \Longrightarrow h(\beta,\alpha) = 0$ for $\beta < 0$.

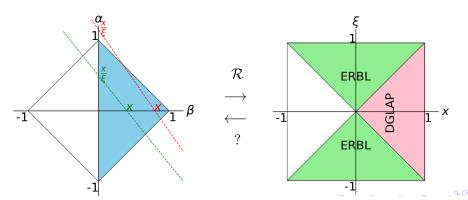


Nabil Chouika

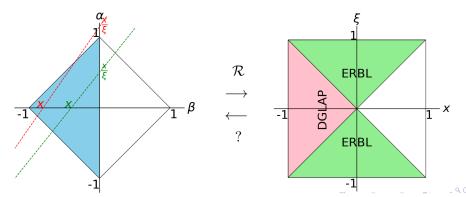
Covariant extension of GPDs

- X

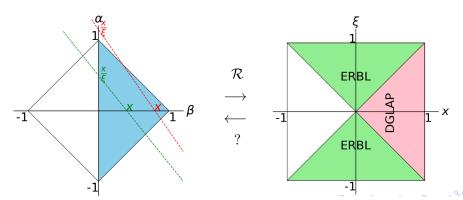
- Quark GPD: $H(x,\xi) = 0$ for $-1 < x < -|\xi| \Longrightarrow h(\beta,\alpha) = 0$ for $\beta < 0$.
- Domains $\beta < 0$ and $\beta > 0$ are uncorrelated in the DGLAP region.



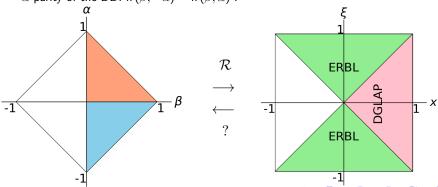
- Quark GPD: $H(x,\xi) = 0$ for $-1 < x < -|\xi| \Longrightarrow h(\beta,\alpha) = 0$ for $\beta < 0$.
- Domains $\beta < 0$ and $\beta > 0$ are uncorrelated in the DGLAP region.



- Quark GPD: $H(x,\xi) = 0$ for $-1 < x < -|\xi| \Longrightarrow h(\beta,\alpha) = 0$ for $\beta < 0$.
- Domains β < 0 and β > 0 are uncorrelated in the DGLAP region.
- Divide and conquer:
 - Better numerical stability.
 - Lesser complexity: $O(N^p + N^p) \ll O((N + N)^p)$.



- Quark GPD: $H(x,\xi) = 0$ for $-1 < x < -|\xi| \Longrightarrow h(\beta,\alpha) = 0$ for $\beta < 0$.
- Domains $\beta < 0$ and $\beta > 0$ are uncorrelated in the DGLAP region.
- Divide and conquer:
 - Better numerical stability.
 - ▶ Lesser complexity: $O(N^p + N^p) \ll O((N + N)^p)$.
- α -parity of the DD: $h(\beta, -\alpha) = h(\beta, \alpha)$.



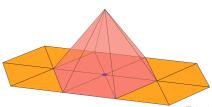
• Expansion of the DD into basis functions $\{v_i\}$:

$$h(\beta,\alpha) = \sum_{j} h_{j} v_{j}(\beta,\alpha) , \qquad (15)$$

• Expansion of the DD into basis functions $\{v_i\}$:

$$h(\beta,\alpha) = \sum_{i} h_{i} v_{i}(\beta,\alpha) , \qquad (15)$$

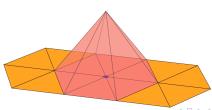
▶ Piece-wise constant, piece-wise linear, etc.



• Expansion of the DD into basis functions $\{v_j\}$:

$$h(\beta,\alpha) = \sum_{i} h_{i} v_{i}(\beta,\alpha) , \qquad (15)$$

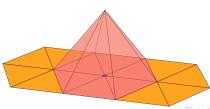
- ▶ Piece-wise constant, piece-wise linear, etc.
- n columns of the matrix.



• Expansion of the DD into basis functions $\{v_j\}$:

$$h(\beta,\alpha) = \sum_{i} h_{i} v_{j}(\beta,\alpha) , \qquad (15)$$

- ▶ Piece-wise constant, piece-wise linear, etc.
- n columns of the matrix.
- Sampling:

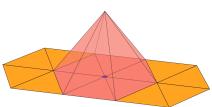


713CTCTIZACIOTI

• Expansion of the DD into basis functions $\{v_j\}$:

$$h(\beta,\alpha) = \sum_{j} h_{j} v_{j}(\beta,\alpha) , \qquad (15)$$

- ▶ Piece-wise constant, piece-wise linear, etc.
- n columns of the matrix.
- Sampling:
 - ▶ Random couples $(x, \xi) \longrightarrow m \ge n$ lines of the matrix.

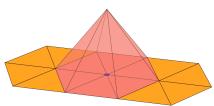


Expansion of the DD into basis functions $\{v_i\}$:

$$h(\beta,\alpha) = \sum_{j} h_{j} v_{j}(\beta,\alpha) , \qquad (15)$$

Covariant extention of GPDs

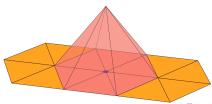
- Piece-wise constant, piece-wise linear, etc.
- n columns of the matrix.
- Sampling:
 - ▶ Random couples $(x,\xi) \longrightarrow m \ge n$ lines of the matrix.
- Linear problem: AX = B where $B_i = H(x_i, \xi_i)$ and $A_{ii} = \mathcal{R}v_i(x_i, \xi_i)$.



Expansion of the DD into basis functions $\{v_i\}$:

$$h(\beta,\alpha) = \sum_{j} h_{j} v_{j}(\beta,\alpha) , \qquad (15)$$

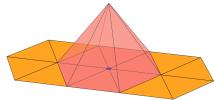
- Piece-wise constant, piece-wise linear, etc.
- n columns of the matrix.
- Sampling:
 - ▶ Random couples $(x, \xi) \longrightarrow m \ge n$ lines of the matrix.
- Linear problem: AX = B where $B_i = H(x_i, \xi_i)$ and $A_{ij} = \mathcal{R}v_i(x_i, \xi_i)$.
- Regularization necessary: discrete ill-posed problem.



Expansion of the DD into basis functions $\{v_i\}$:

$$h(\beta,\alpha) = \sum_{i} h_{i} v_{i}(\beta,\alpha) , \qquad (15)$$

- Piece-wise constant, piece-wise linear, etc.
- n columns of the matrix.
- Sampling:
 - ▶ Random couples $(x, \xi) \longrightarrow m \ge n$ lines of the matrix.
- Linear problem: AX = B where $B_i = H(x_i, \xi_i)$ and $A_{ii} = \mathcal{R}v_i(x_i, \xi_i)$.
- Regularization necessary: discrete ill-posed problem.
 - Trade-off between noise and convergence.



Some examples (Dyson-Schwinger model)

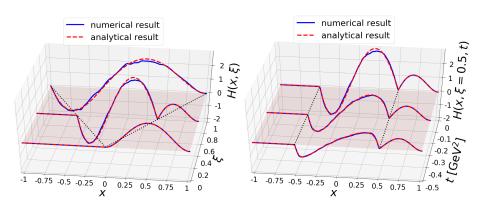


Figure: Extension of GPDs for the pion DSE model of Refs. (Mezrag, 2015; Mezrag et al., 2016). Comparison to the analytical result. Left: Plot for fixed ξ values 0, 0.5 and 1, at t = 0 GeV². Right: Plot for fixed t values 0, -0.25 and -0.5 GeV², at $\xi = 0.5$.

Some examples (Spectator model)

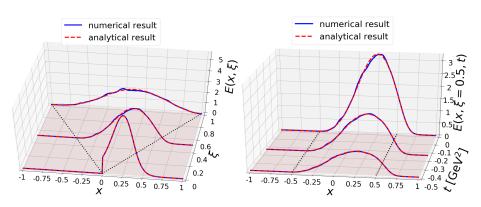


Figure: Extension of GPD E for the nucleon model of Ref. (Hwang and Mueller, 2008). Comparison to the analytical result of the authors. Left: Plot for fixed ξ values 0, 0.5 and 1, at t = 0 GeV². Right: Plot for fixed t values 0, -0.25 and -0.5 GeV², at $\xi = 0.5$.

Some examples (gaussian model)

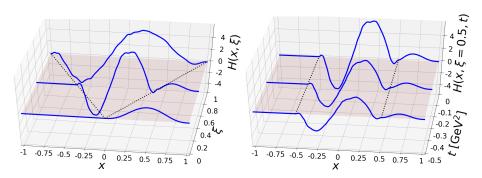


Figure: Extension of GPD for a gaussian pion model (in the vein of AdS/QCD). Left: Plot for fixed ξ values 0, 0.5 and 1, at t=0 GeV². Right: Plot for fixed t values 0, -0.25 and -0.5 GeV², at $\xi=0.5$.

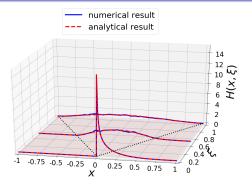


Figure: Extension of GPD for a nucleon toy model with Regge behavior. Plot for fixed ξ values 0, 0.5 and 1.

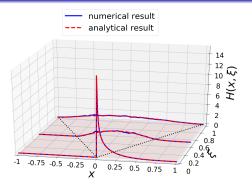


Figure: Extension of GPD for a nucleon toy model with Regge behavior. Plot for fixed ξ values 0, 0.5 and 1.

Integrable singularity for the GPD at $x \sim 0$: $H(x,\xi) \propto \frac{1}{\sqrt{x}}$.

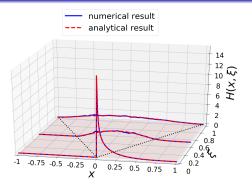


Figure: Extension of GPD for a nucleon toy model with Regge behavior. Plot for fixed ξ values 0, 0.5 and 1.

- Integrable singularity for the GPD at $x \sim 0$: $H(x,\xi) \propto \frac{1}{\sqrt{x}}$.
- Equivalent to an integrable singularity for the DD at $\beta \sim$ 0: $h(\beta, \alpha) \propto \frac{1}{\sqrt{\beta}}$.

4 D > 4 P > 4 E > 4 E > E E 9 9 0

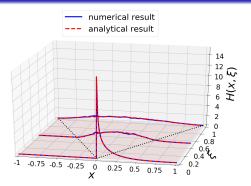


Figure: Extension of GPD for a nucleon toy model with Regge behavior. Plot for fixed ξ values 0, 0.5 and 1.

- Integrable singularity for the GPD at $x \sim 0$: $H(x,\xi) \propto \frac{1}{\sqrt{x}}$.
- Equivalent to an integrable singularity for the DD at $\beta \sim$ 0: $h(\beta, \alpha) \propto \frac{1}{\sqrt{\beta}}$.
 - We solve for $\sqrt{\beta} h(\beta, \alpha)$ instead of $h(\beta, \alpha)$!

1 > ◆@ > ◆분 > ◆분 > 분 | 트 / 의 < ©

Nabil Chouika

Generalized Parton Distributions

Representations of GPDs

- Generalized Parton Distributions
 - encode information about the 3D structure of a hadron.

- Generalized Parton Distributions
 - encode information about the 3D structure of a hadron.
 - are accessible with exclusive processes in experiments: JLab, COMPASS, etc.

18 / 22

- Generalized Parton Distributions
 - encode information about the 3D structure of a hadron.
 - are accessible with exclusive processes in experiments: JLab, COMPASS, etc.
- Systematic procedure for GPD modeling from first principles:

- Generalized Parton Distributions
 - encode information about the 3D structure of a hadron.
 - are accessible with exclusive processes in experiments: JLab, COMPASS, etc.
- Systematic procedure for GPD modeling from first principles:
 - $\qquad \qquad \mathsf{LFWFs} \underset{\mathrm{Overlap}}{\longrightarrow} \mathsf{GPD} \ \mathsf{in} \ \mathsf{DGLAP} \underset{\mathrm{Inverse \ Radon \ Transform}}{\longrightarrow} \mathsf{DD} \underset{\mathrm{RT}}{\longrightarrow} \mathsf{GPD}.$

- Generalized Parton Distributions
 - encode information about the 3D structure of a hadron.
 - are accessible with exclusive processes in experiments: JLab, COMPASS, etc.
- Systematic procedure for GPD modeling from first principles:
 - $\qquad \qquad \mathsf{LFWFs} \overset{\longrightarrow}{\longrightarrow} \mathsf{GPD} \; \mathsf{in} \; \mathsf{DGLAP} \overset{\longrightarrow}{\underset{\mathrm{Inverse \; Radon \; Transform}}{\longrightarrow}} \mathsf{DD} \overset{\longrightarrow}{\underset{\mathrm{RT}}{\longrightarrow}} \mathsf{GPD}.$
 - Both polynomiality and positivity!

- Generalized Parton Distributions
 - encode information about the 3D structure of a hadron.
 - are accessible with exclusive processes in experiments: JLab, COMPASS, etc.
- **Systematic procedure** for GPD modeling from **first principles**:
 - ightharpoonup LFWFs $\underset{\mathrm{Overlap}}{\longrightarrow}$ GPD in DGLAP $DD \xrightarrow{RT} GPD.$ —→
 Inverse Radon Transform
 - Both polynomiality and positivity!
 - Compromise with respect to noise and convergence.

- Generalized Parton Distributions
 - encode information about the 3D structure of a hadron.
 - are accessible with exclusive processes in experiments: JLab, COMPASS, etc.
- Systematic procedure for GPD modeling from first principles:
 - $\qquad \qquad \mathsf{LFWFs} \underset{\mathrm{Overlap}}{\longrightarrow} \mathsf{GPD} \ \mathsf{in} \ \mathsf{DGLAP} \underset{\mathrm{Inverse \ Radon \ Transform}}{\longrightarrow} \mathsf{DD} \underset{\mathrm{RT}}{\longrightarrow} \mathsf{GPD}.$
 - Both polynomiality and positivity!
 - Compromise with respect to noise and convergence.
 - arxiv:1711.05108.

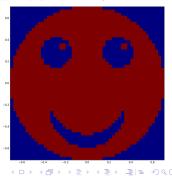
- Generalized Parton Distributions
 - encode information about the 3D structure of a hadron.
 - are accessible with exclusive processes in experiments: JLab, COMPASS, etc.
- Systematic procedure for GPD modeling from first principles:
 - $\qquad \qquad \mathsf{LFWFs} \overset{\longrightarrow}{\longrightarrow} \mathsf{GPD} \; \mathsf{in} \; \mathsf{DGLAP} \; \underset{\mathrm{Inverse \; Radon \; Transform}}{\longrightarrow} \; \mathsf{DD} \overset{\longrightarrow}{\longrightarrow} \; \mathsf{GPD}.$
 - Both polynomiality and positivity!
 - Compromise with respect to noise and convergence.
 - arxiv:1711.05108.
- Unified phenomonelogy of GPDs and TMDs at the level of LFWFs?

Generalized Parton Distributions

- encode information about the 3D structure of a hadron.
- are accessible with exclusive processes in experiments: JLab, COMPASS, etc.
- Systematic procedure for GPD modeling from first principles:

 $\qquad \qquad \mathsf{LFWFs} \overset{\longrightarrow}{\longrightarrow} \mathsf{GPD} \ \mathsf{in} \ \mathsf{DGLAP} \overset{\longrightarrow}{\underset{\mathrm{Inverse \ Radon \ Transform}}{\longrightarrow}} \mathsf{DD} \overset{\longrightarrow}{\underset{\mathrm{RT}}{\longrightarrow}} \mathsf{GPD}.$

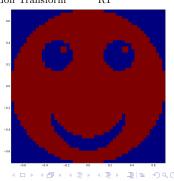
- Both polynomiality and positivity!
- Compromise with respect to noise and convergence.
- ► arxiv:1711.05108.
- Unified phenomonelogy of GPDs and TMDs at the level of LFWFs?
- Thank you!



- Generalized Parton Distributions
 - encode information about the 3D structure of a hadron.
 - are accessible with exclusive processes in experiments: JLab, COMPASS, etc.
- **Systematic procedure** for GPD modeling from **first principles**:

ightharpoonup LFWFs $\underset{\mathrm{Overlap}}{\longrightarrow}$ GPD in DGLAP $\mathsf{DD} \xrightarrow[\mathrm{RT}]{} \mathsf{GPD}.$ → Inverse Radon Transform

- Both polynomiality and positivity!
- Compromise with respect to noise and convergence.
- arxiv:1711.05108.
- Unified phenomonelogy of GPDs and TMDs at the level of LFWFs?
- Thank you!
 - Any questions?

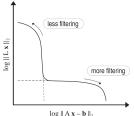


III-posed problems and Regularization

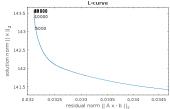
- Ill-posed problems?
 - For example the inversion of a Fredholm equation of the first kind:

$$\int K(x,y) f(y) dy = g(x).$$
 (16)

- ▶ The inverse is not continuous: an arbitrarily small variation Δg of the rhs can lead to an arbitrarily large variation Δf of the solution.
- The corresponding discrete problem needs to be regularized.
 - ▶ E.g Tikhonov regularization: $\min \{ \|AX B\|^2 + \epsilon \|X\|^2 \}$.



Theoretical "L-curve": curve parameterized by the regularization factor.



L-curve with the iteration number as regularization factor.

(fig. taken from Ref. (Hansen, 2007))

D-term considerations

Polynomiality property:

$$\int_{-1}^{1} \mathrm{d}x \, x^{m} H(x, \xi, t) = \sum_{\substack{k=0\\k \text{ even}}}^{m+1} c_{k}^{(m)}(t) \xi^{k} \,. \tag{17}$$

■ Recast polynomiality property for H − D:

$$\int_{-1}^{1} \mathrm{d}x \, x^{m} \left(H(x, \xi, t) - D\left(\frac{x}{\xi}, t\right) \right) = \sum_{\substack{k=0\\k \text{ even}}}^{m} c_{k}^{(m)}(t) \xi^{k} \,, \tag{18}$$

where $D\left(\frac{x}{\xi},t\right)$ is the so-called D-term with support on $-\xi < x < \xi$.

● H — D is a Radon Transform:

$$H(x,\xi,t) - D\left(\frac{x}{\xi},t\right) = \int_{\Omega} d\beta \,d\alpha \,h_{PW}(\beta,\alpha) \,\delta(x-\beta-\alpha\xi) \ . \tag{19}$$

- ► The DGLAP region gives no information on the D-term.
- With other DD representations, we can generate intrinsic D-terms, e.g. Pobylitsa representation:

$$H(x,\xi,t) = (1-x) \int_{0}^{\infty} d\beta \, d\alpha \, h_{P}(\beta,\alpha) \, \delta(x-\beta-\alpha\xi) . \tag{20}$$

Still freedom of extra D-term.

Bibliography I

- D. Müller. D. Robaschik, B. Geyer, F. M. Dittes, and J. Hořejši, "Wave functions, evolution equations and evolution kernels from light ray operators of QCD", Fortsch. Phys. 42 (1994) 101-141, arXiv:hep-ph/9812448 [hep-ph].
- A. V. Radyushkin, "Scaling limit of deeply virtual Compton scattering", Phys. Lett. B380 (1996) 417-425, arXiv:hep-ph/9604317 [hep-ph].
- X.-D. Ji, "Deeply virtual Compton scattering", Phys. Rev. D55 (1997) 7114-7125, arXiv:hep-ph/9609381 [hep-ph].
- M. Burkardt, "Impact parameter dependent parton distributions and off forward parton distributions for zeta —> 0", Phys. Rev. **D62** (2000) 071503, arXiv:hep-ph/0005108 [hep-ph], [Erratum: Phys. Rev.D66,119903(2002)].
- B. Pire, J. Soffer, and O. Teryaev, "Positivity constraints for off forward parton distributions", Eur. Phys. J. C8 (1999) 103-106, arXiv:hep-ph/9804284 [hep-ph].
- A. V. Radyushkin, "Double distributions and evolution equations", Phys. Rev. D59 (1999) 014030, arXiv:hep-ph/9805342 [hep-ph].
- A. V. Belitsky, D. Mueller, and A. Kirchner, "Theory of deeply virtual Compton scattering on the nucleon", Nucl. Phys. **B629** (2002) 323-392, arXiv:hep-ph/0112108 [hep-ph].
- S. J. Brodsky, T. Huang, and G. P. Lepage, "Hadronic wave functions and high momentum transfer interactions in quantum chromodynamics", Conf. Proc. C810816 (1981) 143-199.

Bibliography II

- M. Diehl, T. Feldmann, R. Jakob, and P. Kroll, "The overlap representation of skewed quark and gluon distributions", Nucl. Phys. B596 (2001) 33-65, arXiv:hep-ph/0009255 [hep-ph]. [Erratum: Nucl. Phys.B605,647(2001)].
- M. Diehl. "Generalized parton distributions". Phys. Rept. 388 (2003) 41-277. arXiv:hep-ph/0307382 [hep-ph].
- S. R. Deans, "The Radon Transform and Some of Its Applications", Wiley-Interscience, 1983.
- O. V. Teryaev, "Crossing and radon tomography for generalized parton distributions", Phys. Lett. **B510** (2001) 125-132, arXiv:hep-ph/0102303 [hep-ph].
- D. S. Hwang and D. Mueller, "Implication of the overlap representation for modelling generalized parton distributions", Phys. Lett. **B660** (2008) 350-359, arXiv:0710.1567 [hep-ph].
- N. Chouika, C. Mezrag, H. Moutarde, and J. Rodríguez-Quintero, "Covariant Extension of the GPD overlap representation at low Fock states", arXiv:1711.05108 [hep-ph].
- C. Mezrag, "Generalised Parton Distributions: from phenomenological approaches to Dyson-Schwinger equations", PhD thesis, IRFU, SPhN, Saclay, 2015.
- C. Mezrag, H. Moutarde, and J. Rodriguez-Quintero, "From Bethe-Salpeter Wave functions to Generalised Parton Distributions", Few Body Syst. 57 (2016), no. 9, 729-772, arXiv:1602.07722 [nucl-th].
- P. C. Hansen, "Regularization Tools version 4.0 for Matlab 7.3", Numerical Algorithms 46 (2007), no. 2, 189-194.