

# Application of the numerical conformal bootstrap

David Poland

Yale

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Bootstat 2021



# Conformal Bootstrap

Can we use:

1. Conformal Symmetry:  $SO(D, 2)$  or  $SO(D + 1, 1)$
2. Crossing Symmetry
3. Unitarity or Reflection Positivity

to map out and solve conformal field theories?

# Crossing Symmetry

CFT 4-point functions  $\langle \sigma(x_1)\sigma(x_2)\sigma(x_3)\sigma(x_4) \rangle$  can be expanded using OPE  $\sigma \times \sigma \sim \sum_{\mathcal{O}} \lambda_{\sigma\sigma\mathcal{O}} \mathcal{O}$  in different channels, giving crossing constraint:

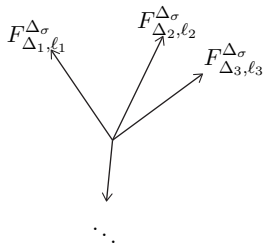
$$\Sigma \text{ (planar)} = \Sigma \text{ (non-planar)}$$

$$0 = \sum_{\Delta, \ell} \lambda_{\sigma\sigma\mathcal{O}}^2 \left[ \frac{g_{\Delta, \ell}(z, \bar{z})}{(z\bar{z})^{\Delta_\sigma}} - \frac{g_{\Delta, \ell}(1-z, 1-\bar{z})}{[(1-z)(1-\bar{z})]^{\Delta_\sigma}} \right]$$

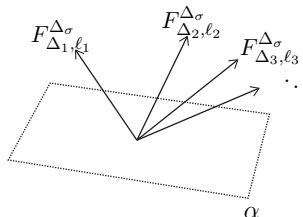
- ▶ Conformal blocks  $g_{\Delta,\ell}(z, \bar{z})$  known functions of cross ratios  $z, \bar{z}$
- ▶ Only **unknowns** are set of scaling dimensions and coefficients:  $\{\Delta, \lambda_{ijk}\}$

# Numerical Approach

Yes



No

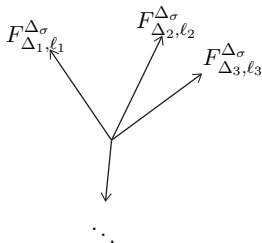


- Make some assumption on  $\{\Delta, \lambda_{ijk}\}$ , search for functional

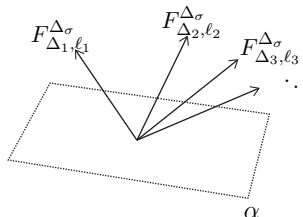
$$\alpha = \left( \sum_{m+n \leq \Lambda} \alpha_{mn} \partial_z^m \partial_{\bar{z}}^n \Big|_{1/2, 1/2} \right) \text{ implying } 0 = \sum(\text{positive})$$

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- Find them by solving semidefinite programs: SDPB

<https://github.com/davidsd/sdpb> [Simmons-Duffin '15; Landry, Simmons-Duffin '19]

# Recent Advances

A few advances in the numerical bootstrap in the past 1-2 years:

- ▶ First applications of SDPB 2.0 [Landry, Simmons-Duffin '19]
- ▶ New “Cutting Surface” algorithm for finding allowed  $\lambda$ 's in large parameter spaces [Chester, Landry, Liu, DP, Simmons-Duffin, Su, Vichi '19]
- ▶ New tiptop algorithm for maximizing some  $\Delta$  over multi-dimensional island [Chester, Landry, Liu, DP, Simmons-Duffin, Su, Vichi, '20]
- ▶ New blocks\_3d software for efficient computation of spinning 3d conformal blocks [Erramilli, Iliesiu, Kravchuk, Landry, DP, Simmons-Duffin '20]

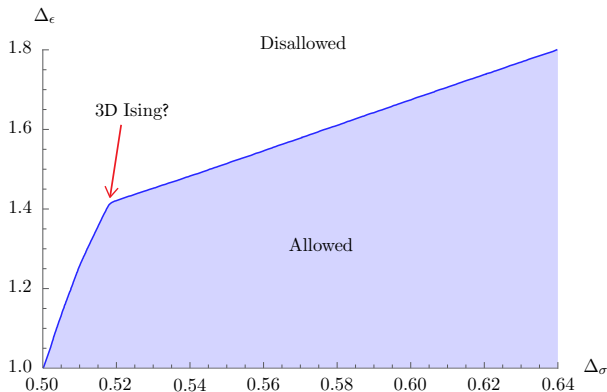
**bootstrapcollaboration** @

Group ID: 1030690

Simons Collaboration on the Nonperturbative Bootstrap (<http://bootstrapcollaboration.com>)

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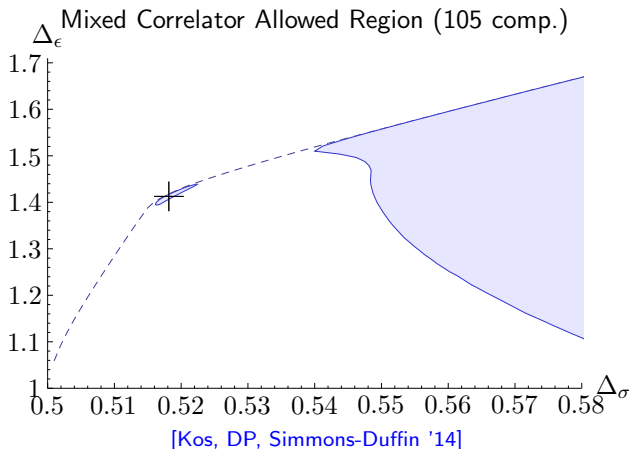
# 3D Dimension Bounds



[El-Showk, Paulos, DP, Rychkov, Simmons-Duffin, Vichi, '12; '14]

- Upper bound on first  $\mathbb{Z}_2$ -even scalar in  $\sigma \times \sigma \sim \mathbb{1} + \epsilon + \dots$  from  $\langle \sigma \sigma \sigma \sigma \rangle$

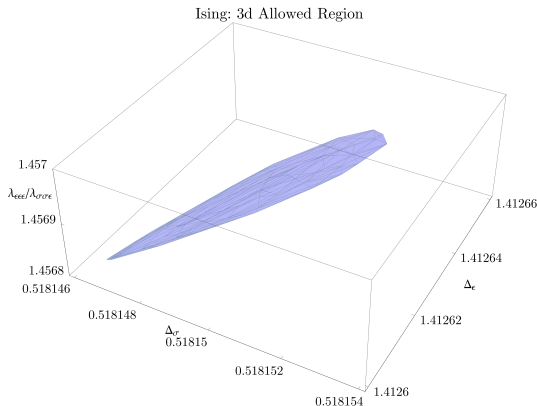
# 3D Ising Island



- Combine 5 crossing relations from  $\langle \sigma \sigma \sigma \sigma \rangle$ ,  $\langle \sigma \sigma \epsilon \epsilon \rangle$ ,  $\langle \epsilon \epsilon \epsilon \epsilon \rangle$  and impose that  $\sigma$  and  $\epsilon$  are the only **relevant** ( $\Delta < 3$ ) operators



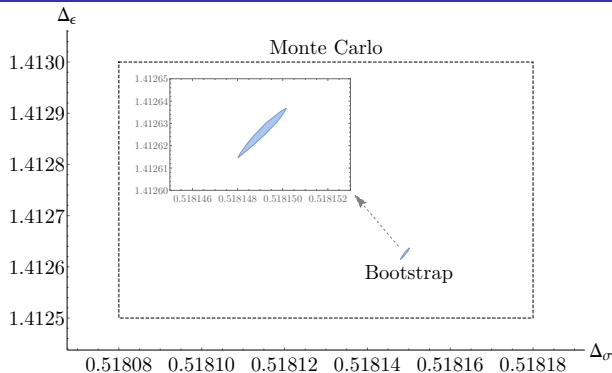
# Mixed Correlator Islands



[Kos, DP, Simmons-Duffin, Vichi '16]

- ▶ Best bounds: perform "OPE scan" over ratio  $r \equiv \lambda_{eee}/\lambda_{\sigma\sigma\epsilon} \rightarrow$  3d island
- ▶ Excludes degenerate exchanged operators at same  $\Delta_{\sigma,\epsilon}$  but different  $\lambda$ 's
- ▶ Each run imposes positivity in specific direction:  $(1\ r)\vec{\alpha} \cdot \vec{F}_{2 \times 2}(1\ r)^T > 0$

# 3D Ising Island



[Kos, DP, Simmons-Duffin, Vichi '16]

- Increase search space to  $5 \times 253 = 1265$  components ( $\Lambda = 43$ )

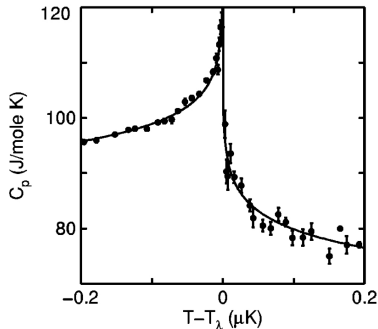
$$\begin{aligned}\{\Delta_\sigma, \Delta_\epsilon\} &= \{0.518149(1), 1.412625(10)\} \\ \{\lambda_{\sigma\sigma\epsilon}, \lambda_{\epsilon\epsilon\epsilon}\} &= \{1.0518537(41), 1.532435(19)\}\end{aligned}$$

# 3D $O(N)$ Models

Generalization to  $N$  scalars  $\phi_i$ :  $\mathcal{L}_{O(N)}^{int} \sim \lambda(\phi_i\phi^i)^2$

# 3D $O(N)$ Models

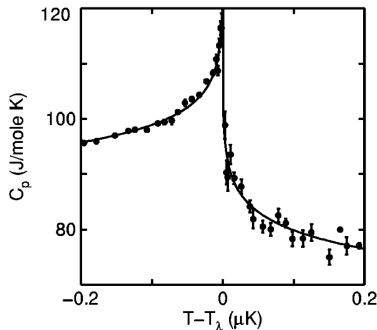
Generalization to  $N$  scalars  $\phi_i$ :  $\mathcal{L}_{O(N)}^{int} \sim \lambda(\phi_i\phi^i)^2$



►  $N = 2$ : Superfluid ( $\lambda$ ) transition in  $^4\text{He}$  [Lipa et al, '96; '03]

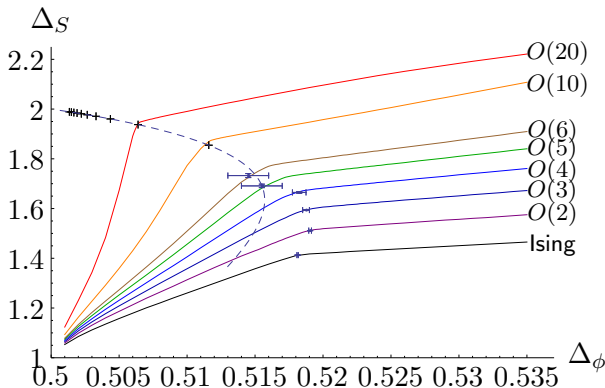
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- ▶  $N = 2$ : Superfluid ( $\lambda$ ) transition in  $^4\text{He}$  [Lipa et al, '96; '03]
- ▶  $N = 3$ : Isotropic ferromagnets (Fe, Co, Ni, ...)
- ▶ Large  $N$ : Solvable in  $1/N$  expansion

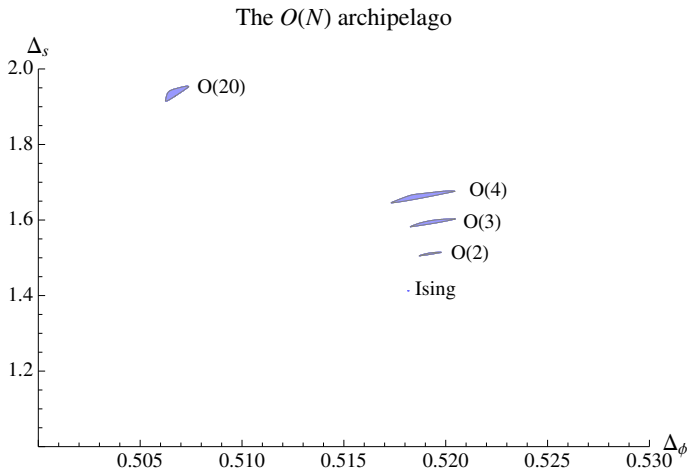
# 3D $O(N)$ Bounds



[Kos, DP, Simmons-Duffin '13]

- ▶ Extension to  $\langle \phi_i \phi_j \phi_k \phi_l \rangle$ , where  $\phi_i$  is  $O(N)$  vector
- ▶ Large  $N$ : matches  $1/N$  expansion, Small  $N$ : matches experiment!

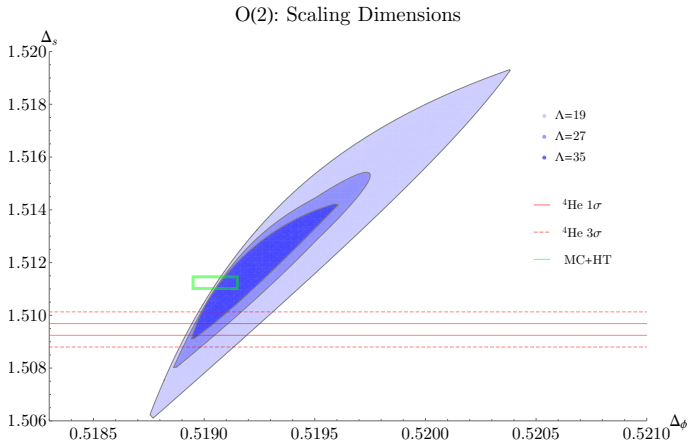
# $O(N)$ Archipelago from Mixed Correlators



[Kos, DP, Simmons-Duffin, Vichi '15; '16]

- With  $\{\phi_i, s\}$  correlators, again find islands (Input:  $\Delta_{\phi'} > 3, \Delta_s > 3$ )

# $O(2)$ from $\{\phi_i, s\}$ System

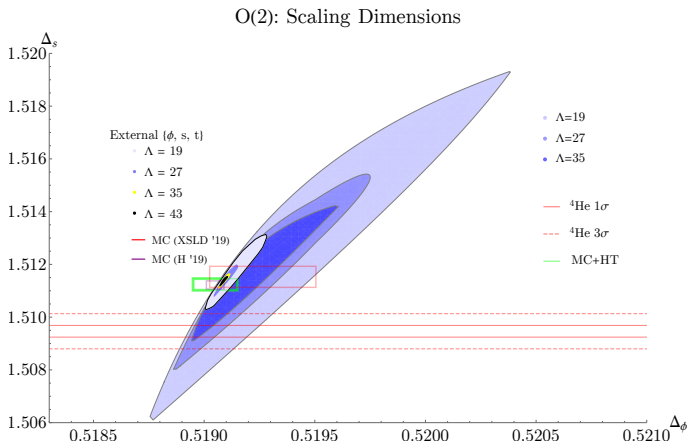


[Kos, DP, Simmons-Duffin, Vichi '15; '16]

- ▶  $\{\Delta_\phi, \Delta_s, \lambda_{\phi\phi s}, \lambda_{sss}\} = \{.51926(32), 1.5117(25), .68726(65), .8286(60)\}$
- ▶  **$8\sigma$  discrepancy** between lattice and expt ( ${}^4\text{He}$ )



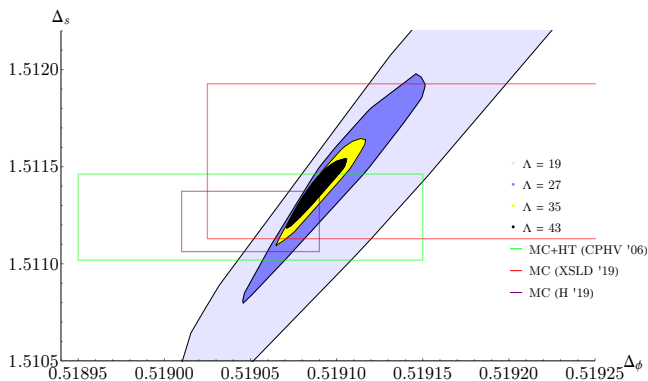
# $O(2)$ from $\{\phi_i, s, t_{ij}\}$ System



[Chester, Landry, Liu, DP, Simmons-Duffin, Su, Vichi '19]

► Result from  $\{\phi_i, s, t_{ij}\}$  system (22 crossing equations)

# $O(2)$ from $\{\phi_i, s, t_{ij}\}$ System



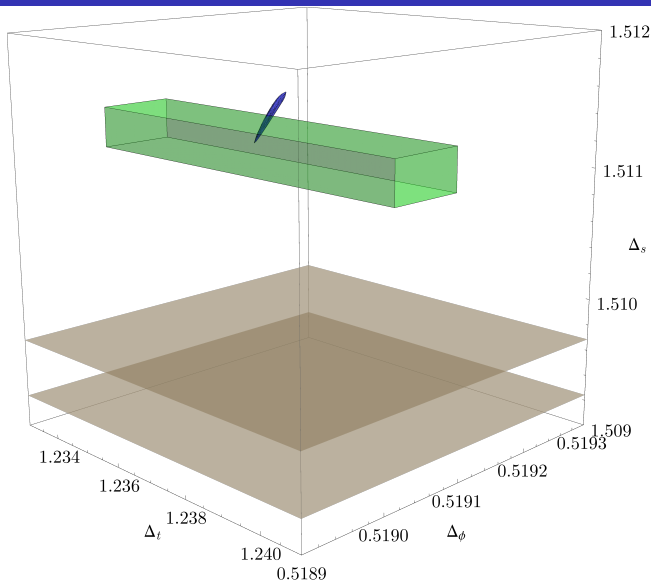
[Chester, Landry, Liu, DP, Simmons-Duffin, Su, Vichi '19]

Bootstrap:  $\alpha = -0.01527(21)$

MC:  $\alpha = -0.01507(21)$  [Hasenbusch '19]

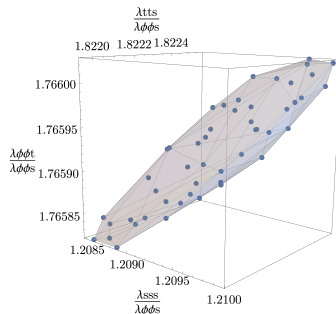
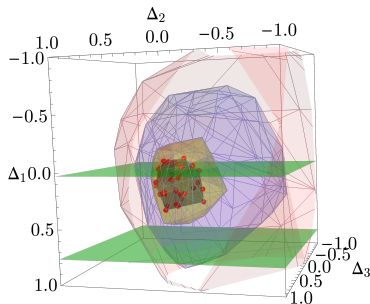
${}^4\text{He}$ :  $\alpha = -0.0127(3)$  [Lipa et al, '96; '03]

# $O(2)$ from $\{\phi_i, s, t_{ij}\}$ System



[Chester, Landry, Liu, DP, Simmons-Duffin, Su, Vichi '19]

# $O(2)$ from $\{\phi_i, s, t_{ij}\}$ System



[Chester, Landry, Liu, DP, Simmons-Duffin, Su, Vichi, '19]

- Best results require 6d search over  $\{\Delta_\phi, \Delta_s, \Delta_t, \frac{\lambda_{sss}}{\lambda_{\phi\phi s}}, \frac{\lambda_{tts}}{\lambda_{\phi\phi s}}, \frac{\lambda_{\phi\phi t}}{\lambda_{\phi\phi s}}\}$

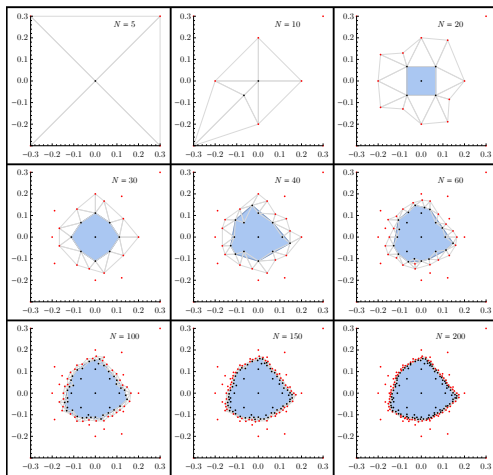
$$\begin{aligned} \{\Delta_\phi, \Delta_s, \Delta_t\} &= \{0.519088(17), 1.51136(18), 1.23629(9)\} \\ \left\{ \frac{\lambda_{sss}}{\lambda_{\phi\phi s}}, \frac{\lambda_{tts}}{\lambda_{\phi\phi s}}, \frac{\lambda_{\phi\phi t}}{\lambda_{\phi\phi s}} \right\} &= \{1.20926(46), 1.82227(19), 1.765918(64)\} \end{aligned}$$

## $O(2)$ from $\{\phi_i, s, t_{ij}\}$ System

To carry out this 6d search, we employed the following strategy:

1. Use SDPB 2.0 [Landry, Simmons-Duffin '19], take advantage of parallelization
2. Use hotstarting [Go, Tachikawa '19] to run SDPB for fewer iterations
3. Search over  $\Delta$ 's carried out using “Delaunay triangulation” search
4. Search over  $\lambda$ 's carried out using “Cutting Surface” algorithm

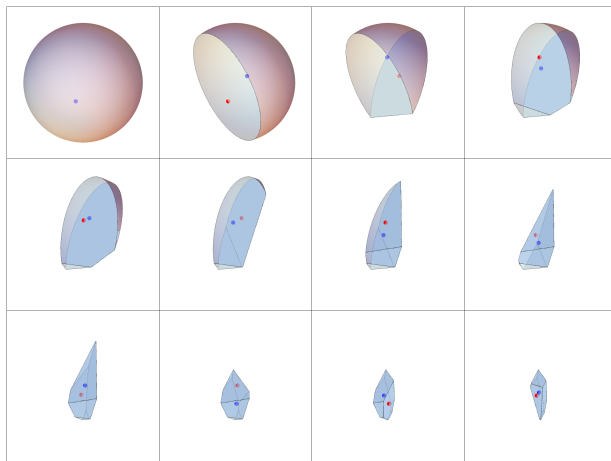
# Delaunay Triangulation Search



[Chester, Landry, Liu, DP, Simmons-Duffin, Su, Vichi '19]

- Compute Delaunay triangulation of all tested points, pick midpoint of “biggest” triangle connecting disallowed to allowed

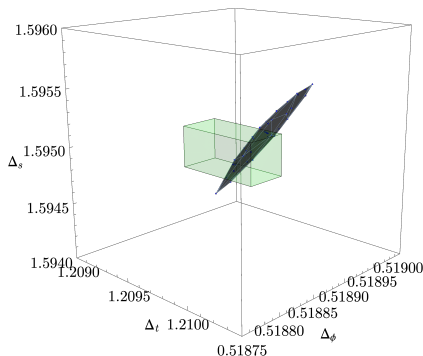
# Cutting Surface Algorithm



[Chester, Landry, Liu, DP, Simmons-Duffin, Su, Vichi '19]

- ▶ Each computation excludes a *region* of  $\vec{\lambda}$ -space:  $\vec{\lambda} \cdot \alpha[F_{\vec{\Delta}}] \cdot \vec{\lambda} > 0$
- ▶ After  $\sim 10 - 30$  tests either find allowed point or rule out entire region

# $O(3)$ from $\{\phi_i, s, t_{ij}\}$ System



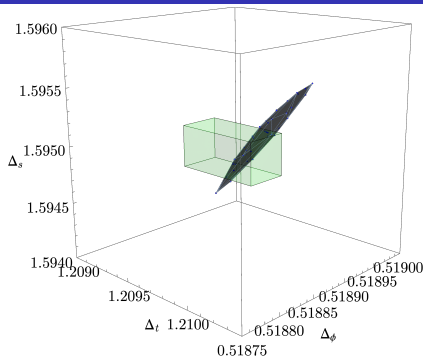
Black (Bootstrap): [Chester, Landry, Liu, DP, Simmons-Duffin, Su, Vichi '20]

Green (Monte Carlo): [Hasenbusch, Vicari '11; Hasenbusch '20]

- ▶ New  $O(3)$  island:  $\{\Delta_\phi, \Delta_s, \Delta_t\} = \{0.51894(5), 1.5949(6), 1.2095(2)\}$
- ▶ Open question: is  $\phi^i \phi^j \phi^k \phi^l$  relevant or irrelevant in  $O(3)$  model?



# $O(3)$ from $\{\phi_i, s, t_{ij}\}$ System

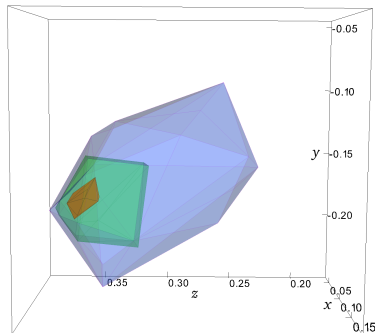
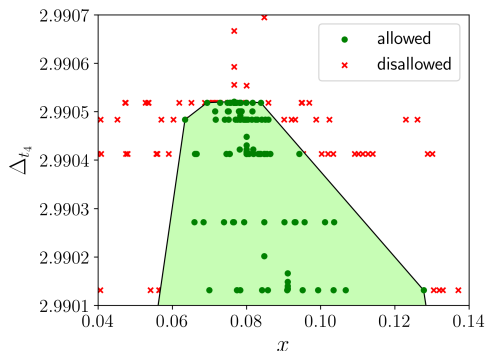


Black (Bootstrap): [Chester, Landry, Liu, DP, Simmons-Duffin, Su, Vichi, '20]

Green (Monte Carlo): [Hasenbusch, Vicari '11; Hasenbusch '20]

- ▶ Using tiptop search, find it is relevant:  $\Delta_{\phi^{\{i\phi^j\phi^k\phi^l\}}} < 2.99056!$
- ▶ Proof that critical Heisenberg magnets are unstable to cubic anisotropy, should flow to fixed point with cubic symmetry  $C_3$  rather than  $O(3)$

# $O(3)$ from $\{\phi_i, s, t_{ij}\}$ System

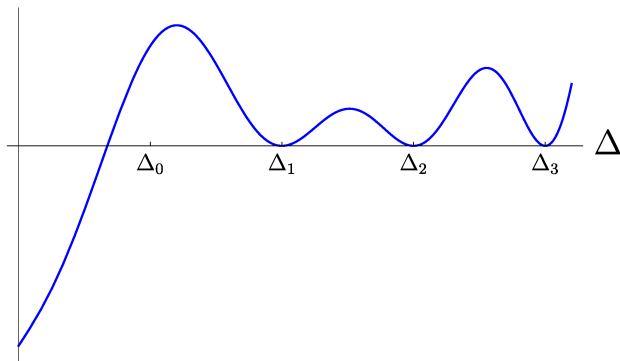


[Chester, Landry, Liu, DP, Simmons-Duffin, Su, Vichi, '20]

- ▶ tiptop search: map out approximate island at a given gap  $\Delta t_4 - 3$ , shrink bounding box of island by factor of 2, increase gap (via binary search) until it can no longer accommodate smaller box, then iterate

# Extremal Functional

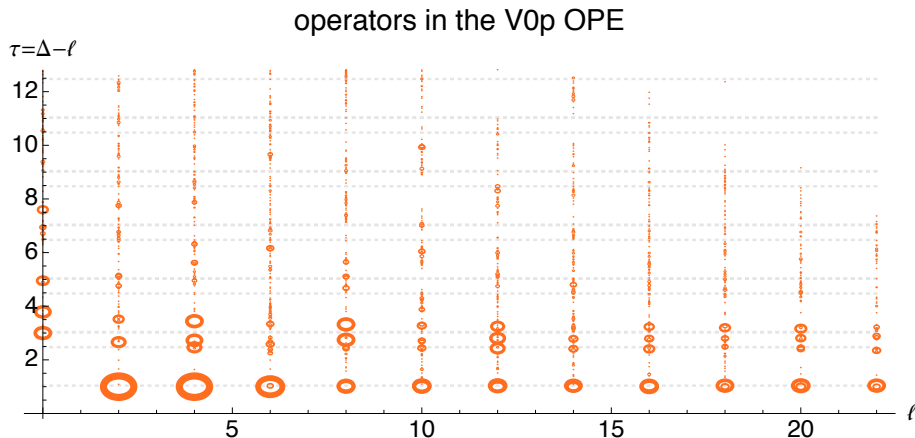
$$\alpha \cdot F_{\Delta,\ell}$$



[DP, Simmons-Duffin '10; Paulos, El-Showk '12; plot from Paulos, Zan '20]

- By going to a boundary of the allowed region (e.g., extremizing  $\lambda_{\phi\phi_s}$ ), we can extract extremal spectra corresponding to the zeros of  $\alpha \cdot F_{\Delta,\ell}$

# $O(2)$ from $\{\phi_i, s, t_{ij}\}$ System



[Chester, Landry, Liu, DP, Simmons-Duffin, Su, Vichi, '19]

- Can be done in practice using spectrum-extraction Python script (`spectrum.py`), which uses `sdpb` output

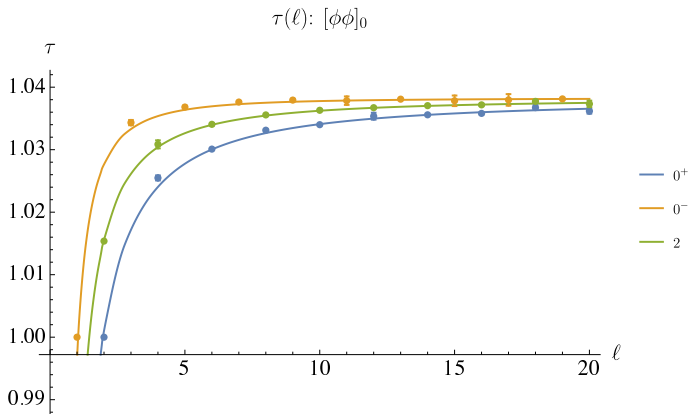
# Analytical Bootstrap

Extremal spectra can be compared with analytical bootstrap predictions:

- ▶ Lightcone Bootstrap ( $z \rightarrow 0$ ):  $\exists$  trajectories of “double-twist” operators  $\sim \sigma \partial^\ell \sigma$  with twist asymptoting to  $\tau(\ell \rightarrow \infty) = 2\Delta_\sigma - \frac{\#}{\ell} - \frac{\#}{\ell \Delta_\epsilon} + \dots$   
[Fitzpatrick, Kaplan, DP, Simmons-Duffin '12; Komargodski, Zhiboedov '12]
- ▶ Lorentzian Inversion  $\rightarrow$  All-orders analytic function [Caron-Huot '17]

$$\tau(\ell) \sim \int d\text{Disc}[g] \sim \sum_{\mathcal{O}} 4F_3(\dots) \text{ [see Henriksson talk...]}$$

# $O(2)$ : Comparison with Analytics



[Albayrak, Meltzer, DP '19; J. Liu, Meltzer, DP, Simmons-Duffin '20]

- Excellent agreement between leading-twist extremal spectra and analytics after including exchange of  $\{s, t, J^\mu, T^{\mu\nu}\}$

# The Spinning Frontier

- General 3d conformal blocks can be expanded recursively in poles:

$$g_{\Delta,j,I}^{ab} \sim \frac{1}{\Delta - \Delta_{j,i}} (\mathcal{L}_{j,i})_{a'}^a (\mathcal{R}_{j,i})_{b'}^b g_{\Delta'_{j,i},j'_{j,i},I}^{a'b'}(z, \bar{z})$$

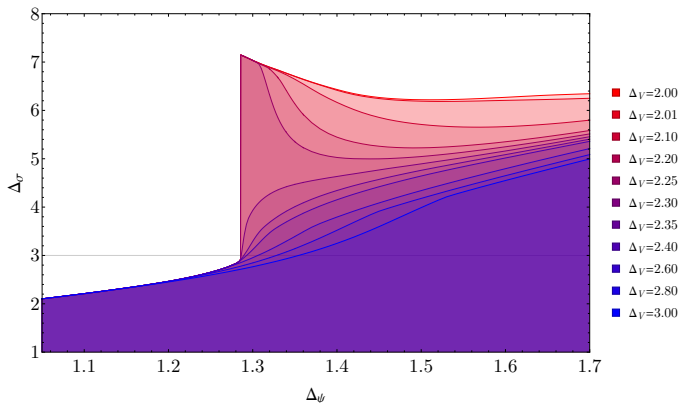
- blocks\_3d is efficient, multithreaded, C++ implementation  
[Erramilli, Iliesiu, Kravchuk '19; Erramilli, Iliesiu, Kravchuk, Landry, DP, Simmons-Duffin '20]
- Practical for external fermions  $\psi$ , currents  $J$ , stress-tensors  $T$ , ...

block ( $\Lambda = 25$ )	$j_{12}$	$j_{43}$	Memory (GB)	Time (hr)
$\langle \phi \phi \phi \phi \rangle$	0	0	4	0.014
$\langle \phi \psi \phi \psi \rangle$	$\frac{1}{2}$	$\frac{1}{2}$	7	0.025
$\langle T \phi \phi \phi \rangle$	2	0	11	0.045
$\langle \psi \psi \psi \psi \rangle$	1	1	15	0.068
$\langle T \phi T \phi \rangle$	2	2	36	0.20
$\langle T \psi T \psi \rangle$	$\frac{5}{2}$	$\frac{5}{2}$	48	0.62
$\langle T T T \phi \rangle$	4	2	62	0.94
$\langle T T T T \rangle$	4	4	106	6.9

(See CFTs4D package for spinning 4d blocks [Cuomo, Karateev, Kravchuk '17])



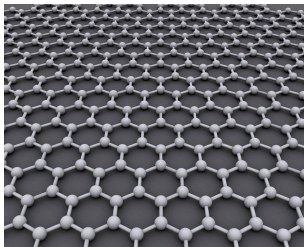
# Mysterious jump?



[Erramilli, Iliesiu, Kravchuk, Landry, DP, Simmons-Duffin '20]

- ▶ Sharp jump in parity-odd scalar bound from  $\langle \psi\psi\psi\psi \rangle$   
[Iliesiu, Kos, DP, Pufu, Simmons-Duffin '17]
- ▶ Seems to persist after removing “fake primary effect”  
([Karateev, Kravchuk, Serone, Vichi '19]: spin-1  $V^\mu$  mimics  $\Delta = 3$  scalar)
- ▶ Could be evidence for new fermionic CFT w/ no relevant scalars?

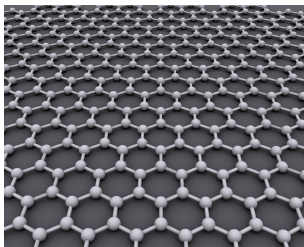
# 3D Fermion Models (Gross-Neveu-Yukawa)



Interesting CFTs obtained from fixed points involving  $N$  fermions:

$$\mathcal{L}_{GNY} \sim \frac{g}{2} \sigma \bar{\psi}^i \psi_i + \lambda \sigma^4 \text{ (and variations with multiple scalars)}$$

# 3D Fermion Models (Gross-Neveu-Yukawa)

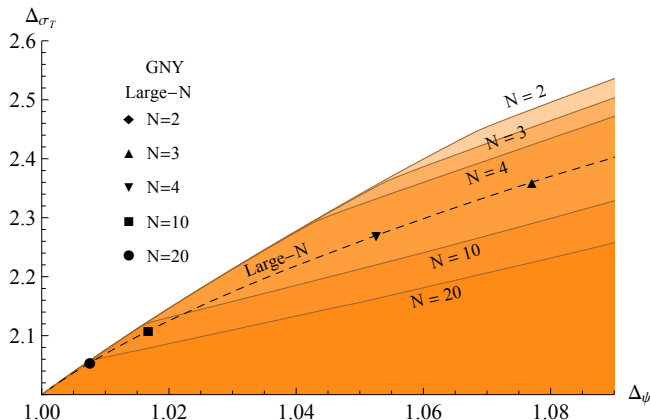


Interesting CFTs obtained from fixed points involving  $N$  fermions:

$$\mathcal{L}_{GNY} \sim \frac{g}{2} \sigma \bar{\psi}^i \psi_i + \lambda \sigma^4 \text{ (and variations with multiple scalars)}$$

- ▶ Large  $N$ : Solvable in  $1/N$  expansion [Gracey '92, '93; ...]
- ▶  $N = 8$ : Possible QCPs in D-wave superconductors or graphene [Vojta, Zhang, Sachdev '00; Herbut '06; Classen, Herbut, Scherer '17]
- ▶  $N = 4$ : Spinless fermions on honeycomb lattice, gapless semiconductors [Raghu, Qi, Honerkamp, Zhang '07; Moon, Xu, Kim, Balents '12; Herbut, Janssen '14]

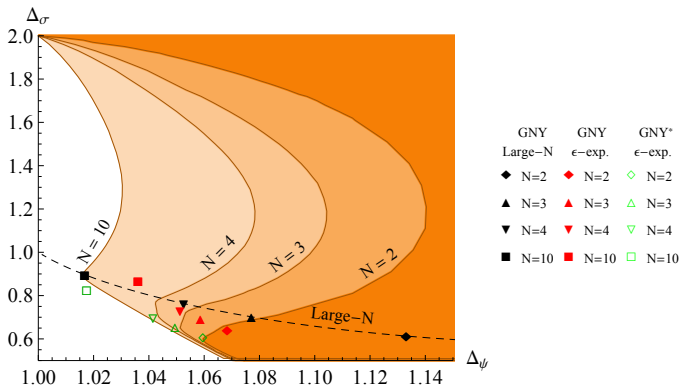
# 3D $O(N)$ Fermion Bootstrap



[Iliesiu, Kos, DP, Pufu, Simmons-Duffin '17]

- ▶ Bootstrap for fermion 4-point functions  $\langle \psi_i \psi_j \psi_k \psi_l \rangle$
- ▶ Kinks in symmetric tensor bounds match GNY models at large  $N$

# 3D $O(N)$ Fermion Bootstrap

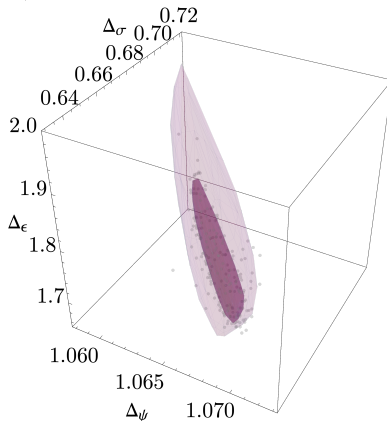


[Iliesiu, Kos, DP, Pufu, Simmons-Duffin '17]

- ▶ Intricate structure in  $\{\Delta_\psi, \Delta_\sigma\}$  plane assuming  $\sigma'$  irrelevant
- ▶ Upper kinks plausibly related to GNV models

# Preliminary Island for $N = 2$ Gross-Neveu-Yukawa Model

$$\Delta_{\psi'} > 2, \Delta_{\sigma'} > 2.5, \Delta_{\epsilon'} > 3, \Delta_{\sigma_T} > 2, \Delta_{\chi} > 3.5$$

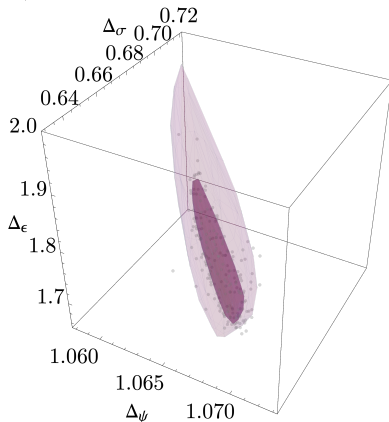


[Erramilli, Iliesiu, Kravchuk, A. Liu, DP, Simmons-Duffin, in progress]

- ▶ Preliminary island from system containing  $\{\sigma, \psi_i, \epsilon\}$  at  $\Lambda = 11, 15$
- ▶ Gaps motivated by large  $N$  estimates and E.O.M.  $\partial\psi \sim \sigma\psi$

# Preliminary Island for $N = 2$ Gross-Neveu-Yukawa Model

$$\Delta_{\psi'} > 2, \Delta_{\sigma'} > 2.5, \Delta_{\epsilon'} > 3, \Delta_{\sigma_T} > 2, \Delta_{\chi} > 3.5$$

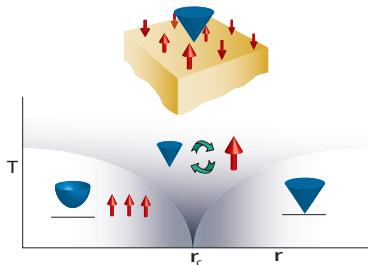


[Erramilli, Iliesiu, Kravchuk, A. Liu, DP, Simmons-Duffin, in progress]

$$\{\Delta_\psi, \Delta_\sigma, \Delta_\epsilon\} = \{1.066(4), 0.666(18), 1.77(7)\}$$

► Now preparing to run at high derivative order and at other  $N$ ...

# Minimal 3D SCFT ( $N = 1$ Gross-Neveu-Yukawa)

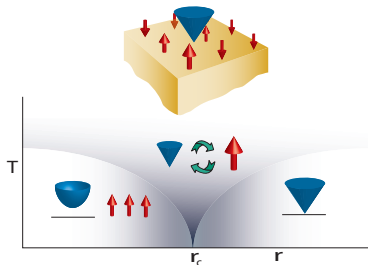


- The minimal SUSY extension of 3D Ising has  $\mathcal{N} = 1$  supersymmetry:

$$V = \frac{g}{2}\sigma\bar{\psi}\psi + \frac{g}{8}\sigma^4 \quad \leftrightarrow \quad W = \frac{g}{3}\Sigma^3, \quad \Sigma = \sigma + \theta\psi + \theta^2\epsilon$$



# Minimal 3D SCFT ( $N = 1$ Gross-Neveu-Yukawa)

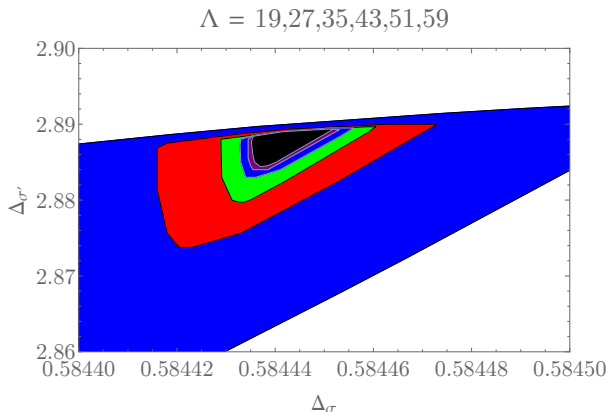


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- May be realizable in (3+1)D topological superconductors, with (2+1)D boundary supporting Majorana fermions [Grover, Sheng, Vishwanath '13]

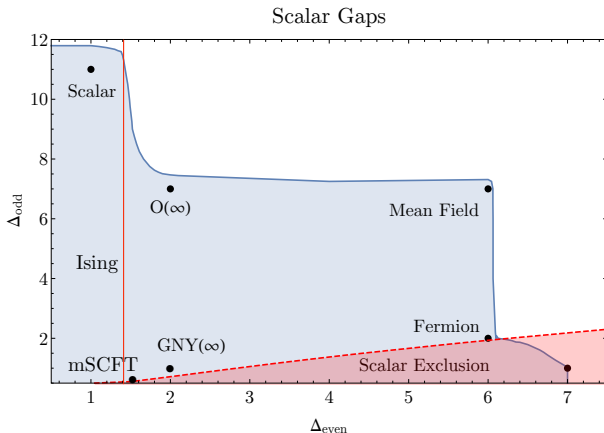
# Supersymmetric Island



[Rong, Su '18; Atanasov, Hillman, DP, Rong, Su, in progress]

- $\{\sigma, \epsilon\}$  SUSY system  $\rightarrow \Delta_{\sigma} = .584443(8), \Delta_{\sigma'} = 2.8869(25)$   
Compare to  $\epsilon$ -expansion:  $\Delta_{\sigma} = .5837(14)$  [Ihrig, Mihaila, Scherer '18]  
(Assumption:  $\mathcal{N} = 1$  SUSY,  $\Delta_{\epsilon'} \geq 3$ )

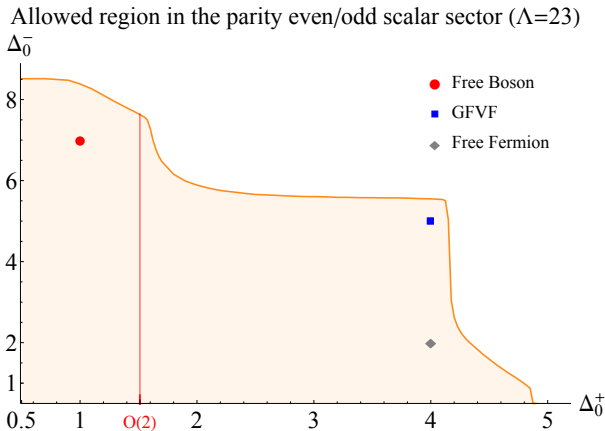
# Map of Allowed Scalar Gaps from $\langle T^{\mu\nu} T^{\rho\sigma} T^{\alpha\beta} T^{\gamma\delta} \rangle$



[Dymarsky, Kos, Kravchuk, DP, Simmons-Duffin, '17]

- Allowed {scalar, pseudoscalar} gaps from stress tensor 4-point functions

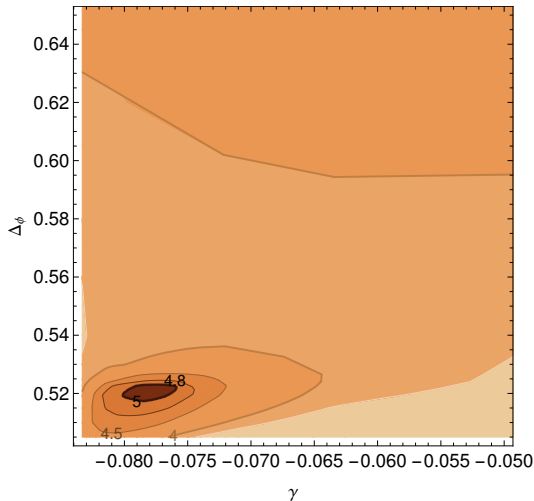
# Map of Allowed Scalar Gaps from $\langle J^\mu J^\nu J^\rho J^\sigma \rangle$



[Dymarsky, Penedones, Trevisani, Vichi '17]

- Allowed {scalar, pseudoscalar} gaps from current 4-point functions

# Map of Allowed Couplings from $\{J^\mu, \phi\}$ System



[Reehorst, Trevisani, Vichi '19]

- Allowed couplings  $\langle JJT \rangle \propto \gamma$  in  $O(2)$  model after imposing  $T'_{\mu\nu}$  gap

Where do we go from here?

- ▶ Find **islands** for other interesting CFTs
  - ▶ 3d Gross-Neveu-Yukawa Models:  $\mathcal{L} \sim \lambda \sigma \psi_i \psi^i + \text{variations}$
  - ▶ Conformal windows of gauge theories (3d QED, 4d QCD, ...)
  - ▶ Superconformal zoo
  - ▶ Other interesting theories described in this workshop!
- ▶ Study **larger systems** of bootstrap equations
  - ▶ Mixed correlators with spinning operators ( $\psi$ ,  $J^\mu$ ,  $T^{\mu\nu}$ )
  - ▶ Improve algorithms and software tools
- ▶ Improve **analytical** understanding of bootstrap equations
  - ▶ Match to Lorentzian Inversion formula, conformal dispersion relations
  - ▶ Incorporate analytical insights into numerical algorithms

# Backup Slides

# Lorentzian Inversion Review

- ▶ The basic idea is to decompose CFT 4-pt functions  $\langle \sigma \sigma \sigma \sigma \rangle \propto g(z, \bar{z})$  in a basis of “principal series” ( $\Delta = d/2 + i\alpha$ ) partial waves

$$g(z, \bar{z}) = \sum_{\ell=0}^{\infty} \int_{d/2-i\infty}^{d/2+i\infty} \frac{d\Delta}{2\pi i} c(\Delta, \ell) g_{\Delta, \ell}(z, \bar{z}) + (\text{non-norm.})$$

where the physical spectrum is encoded in the poles of  $c(\Delta, \ell)$ .



# Lorentzian Inversion Review

- Using orthogonality and Lorentzian continuation one “inverts” the formula to obtain  $c(\Delta, \ell) = c^t(\Delta, \ell) + (-1)^\ell c^u(\Delta, \ell)$ :

$$c^t(\Delta, \ell) = \frac{\kappa_{\Delta+\ell}}{4} \int_0^1 dz d\bar{z} \mu(z, \bar{z}) g_{\ell+d-1, \Delta+1-d}(z, \bar{z}) d\text{Disc} [g(z, \bar{z})]$$

with

$$\begin{aligned} \mu(z, \bar{z}) &= \left| \frac{z - \bar{z}}{z\bar{z}} \right|^{d-2} \frac{1}{(z\bar{z})^2} \\ d\text{Disc} [g(z, \bar{z})] &= g(z, \bar{z}) - \frac{1}{2} g(z, \bar{z} e^{2\pi i}) - \frac{1}{2} g(z, \bar{z} e^{-2\pi i}) \end{aligned}$$

See [\[Caron-Huot '17; Simmons-Duffin, Stanford, Witten '17; Kravchuk, Simmons-Duffin '18\]](#)

# Lorentzian Inversion Review

- ▶ Expanding  $g(z, \bar{z}) = \left( \frac{z\bar{z}}{(1-z)(1-\bar{z})} \right)^{\Delta_\sigma} \sum \lambda_{\sigma\sigma\mathcal{O}}^2 g_{\Delta,\ell}(1-z, 1-\bar{z})$  in a finite number of known contributions, we can compute the integrals
- ▶ Matching identity operator reveals poles  $\frac{1}{\Delta - (2\Delta_\sigma + \ell)}$  corresponding to “double-twist” operators:  $\sigma \partial_{\mu_1} \dots \partial_{\mu_\ell} \sigma$
- ▶ Other exchanged operators give anomalous dimensions ( $\log(z)$  terms) and correct their OPE coefficients (regular terms)

Approach developed in various works: [Sleight, Taronna '18; Kravchuk, Simmons-Duffin '18; Cardona, Sen '18; Karateev, Kravchuk, Simmons-Duffin '18; Cardona, Guha, Kanumilli, Sen '18; Albayrak, Meltzer, DP '19, '20; Caron-Huot, Gobeil, Zahraee '20]

# Anomalous Dimensions from Scalar Exchange

$$\delta\tau_{[\sigma\sigma]_0}(\bar{h})_{\text{pert}} = -\frac{\lambda_{\sigma\sigma\epsilon}^2}{1 + \delta P_{[\sigma\sigma]_0}(\bar{h})} \frac{2\Gamma(\Delta_\sigma)^2 \Gamma\left(1 + \frac{\Delta_\epsilon - 2\Delta_\sigma}{2}\right)^2 \Gamma(\Delta_\epsilon)}{\Gamma\left(1 - \frac{\Delta_\epsilon - 2\Delta_\sigma}{2}\right)^2 \Gamma\left(\frac{\Delta_\epsilon - 2\Delta_\sigma}{2}\right)^2 \Gamma\left(\frac{\Delta_\epsilon}{2}\right)^2} \\ \times \frac{\Gamma(\bar{h} - \Delta_\sigma + 1) \Gamma\left(\bar{h} - \frac{\Delta_\epsilon - 2\Delta_\sigma}{2} - 1\right)}{\Gamma(\bar{h} + \Delta_\sigma - 1) \Gamma\left(\bar{h} + \frac{\Delta_\epsilon - 2\Delta_\sigma}{2} + 1\right)}$$

$$\times {}_4F_3\left(\frac{\Delta_\epsilon - d + 2}{2}, \frac{\Delta_\epsilon - 2\Delta_\sigma}{2} + 1, \frac{\Delta_\epsilon - 2\Delta_\sigma}{2} + 1, \frac{\Delta_\epsilon}{2}; 1\right)$$

$$\delta\tau_{[\sigma\sigma]_0}(\bar{h})_{\text{np}} = -\frac{\lambda_{\sigma\sigma\epsilon}^2}{1 + \delta P_{[\phi\phi]_0}(\bar{h})} \frac{2\Gamma(\Delta_\sigma)^2 \Gamma\left(\Delta_\epsilon - \frac{d-2}{2}\right) \Gamma(\Delta_\epsilon)}{\Gamma\left(1 - \frac{\Delta_\epsilon - 2\Delta_\sigma}{2}\right)^2 \Gamma\left(\frac{\Delta_\epsilon - 2\Delta_\sigma}{2}\right)^2 \Gamma\left(\frac{\Delta_\epsilon}{2}\right)^3 \Gamma\left(\frac{\Delta_\epsilon - d - 2}{2}\right)} \\ \times \frac{\Gamma(\bar{h})^2 \Gamma(\bar{h} - \Delta_\sigma + 1) \Gamma\left(\bar{h} + \Delta_\sigma - \frac{d}{2}\right) \Gamma\left(\frac{\Delta_\epsilon - 2\Delta_\sigma}{2} - \bar{h} + 1\right)}{\Gamma(2\bar{h}) \Gamma\left(\bar{h} + \Delta_\sigma + \frac{\Delta_\epsilon}{2} - \frac{d}{2}\right)} \\ \times {}_4F_3\left(\bar{h}, \bar{h}, \bar{h} + \Delta_\sigma - 1, \bar{h} + \Delta_\sigma - \frac{d}{2}; 2\bar{h}, \bar{h} - \frac{\Delta_\epsilon - 2\Delta_\sigma}{2}, \bar{h} + \Delta_\sigma + \frac{\Delta_\epsilon}{2} - \frac{d}{2}; 1\right)$$

► E.g., for exchange of a scalar  $\epsilon$ , integral yields two pieces:

$$\delta\tau_{[\sigma\sigma]_0}(\bar{h}) = \delta\tau_{[\sigma\sigma]_0}(\bar{h})_{\text{pert}} + \delta\tau_{[\sigma\sigma]_0}(\bar{h})_{\text{np}} \quad (\text{here } \bar{h} \equiv \frac{\Delta + \ell}{2})$$