

# Informal Considerations on the Random Field Ising Model

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# Glasses & RFIM

## Freezing of fluctuations

Order parameter: overlap between config.

$\delta\rho(x)$  local density fluctuations

$$\langle \delta\rho(x) \delta\rho'(x) \rangle \sim q_{\text{EA}}$$

$$q_{\rho\rho'} = \int dx w(x-y) \delta\rho(x) \delta\rho'(y)$$

Order parameter: overlap between configurations

## Landau free-energy for glasses

(I) Fix a conf.  $\underline{x} = \{x_1 \dots x_n\}$

chosen with  $\mu(\underline{x}) = e^{-\beta H(\underline{x})}/Z$

Space  
↓  
 $q_x(X, Y)$   
local measure  
of similarity

(II) Consider  $\mu(Y|X) = \frac{e^{-\beta H(Y)} + \beta \int dx \epsilon_x q_x(X, Y)}{Z(X, \epsilon)}$

(III) Large deviation function

$$V(q) = -T \int dx \mu(x) \log Z(x, \epsilon)$$

$$\sum_x e^{-\beta H(x)} \log Z(x, q) \longrightarrow \sum_x e^{-\beta \sum_{a=1}^n h(x_a) + \beta \sum_{a>1} q(x_a, x_a)}$$

↓

$$Z^{n-1}(x, q)$$

Glasses:  $n \rightarrow 1$  invariance  $S_{n-1}$

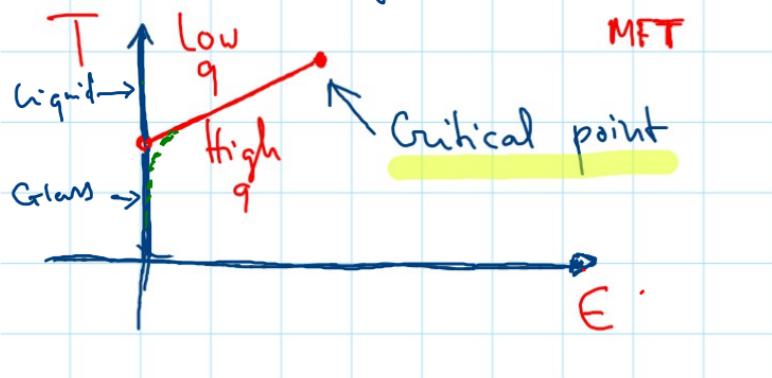
$\times$ : self-induced disorder

Fluctuations of overlap

thermal

sample-to-sample (choice of  $x$ )

# Phase diagram



Critical point :

Critical overlap  
fluctuations  
in the Univ. Class  
of RFIM  
in pert. theory

Starting point :

special Sm symmetric point  
with  $\phi_{ab}$

# RFIM (Random Field Ising Model)

## Definition of the model

$$H = -J \sum_{\langle i,j \rangle} S_i S_j - \sum_i h_i S_i$$

$h_i$ : random  $\overline{h_i h_j} = \Delta(i-j) = \Delta \cdot \delta_{ij}$

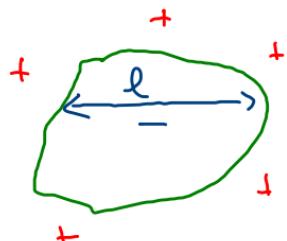
## Para-Ferro transition

$\psi$ " version  $F[\psi] = \int dx \left[ \frac{1}{2} (\nabla \psi)^2 + V(\psi) - h \psi \psi(x) \right]$

Existence of the transition

Imag-Ma argument

Stability of magnetized phase



$$\Delta E = \sigma l^{d-1} - \sqrt{\Delta} l^{d/2}$$

transition for  $d > 2$

$d = 2$  no-transition

$$G_c(r) = \overline{\langle \varphi(x)\varphi(0) \rangle} - \overline{\langle \varphi(x) \rangle} \overline{\langle \varphi(0) \rangle} \sim \frac{1}{r^{d-2+\gamma}}$$

$$G_d(r) = \overline{\langle \varphi(x)\varphi(0) \rangle} - \overline{\langle \varphi(x) \rangle} \overline{\langle \varphi(0) \rangle} \sim \frac{1}{r^{d-4+\bar{\gamma}}}$$

$$\bar{\gamma} = \gamma = 0 \quad \text{in MF}$$

Pert. theory dominated by free.

$$(1) \quad \gamma = \bar{\gamma}$$

(2) Dimensional reduction  $\rightarrow$  corr. functions (+expon)  
of the pure model in dim  $d' = d-2$

Solutions of Stochastic eq. :

$$-\Delta\varphi + V'(\varphi) = \lambda(x)$$

$$-\Delta\varphi + V'(\varphi) = h(x)$$

$$\overline{A[\{\varphi\}]} = \frac{\int d\varphi \delta(-\Delta\varphi + V' + h) |det(-\Delta + V'')| A}{\int d\varphi |det(-\Delta + V'')| A}$$

If the StochEq has only one solution

$$|det| \rightarrow det \rightarrow$$

$$\delta(-\Delta\varphi + V' + h) |det(-\Delta + V'')| = \int d\hat{\varphi} d\bar{x} dx e^{i\hat{\varphi}(-\Delta\varphi + V' + h)} \times e^{-\frac{1}{2}\bar{x}(-\Delta + V'')x}$$

$$S[\varphi, \hat{\varphi}, \bar{x}, x] \rightarrow \text{BRST Symmetric}$$

$$i\hat{\varphi} = \frac{\partial}{\partial h}$$

$$\overline{e^{\int i\hat{\varphi}(x)h(x)}} = e^{\frac{\Delta}{2} \int (i\hat{\varphi})^2} \Rightarrow \text{Susy}$$

$$\underline{\Phi} = \varphi + \theta \bar{x} + \bar{\theta} x + \theta \bar{\theta} i \hat{\psi} \quad \theta, \bar{\theta}$$

Grassmann vars.

$$i \hat{\psi}^2 = \int d\theta d\bar{\theta} \quad \underline{\Phi} \frac{\partial}{\partial \theta \partial \bar{\theta}} \underline{\Phi}$$

$$\mathcal{S} = \int dx d\theta d\bar{\theta} \left[ \underline{\Phi} \left( \partial_x^2 + \Delta \partial_\theta \partial_{\bar{\theta}} \right) \underline{\Phi} + V(\underline{\Phi}) \right]$$

$$\langle \underline{\Phi}(x, \theta \bar{\theta}) \underline{\Phi}(0) \rangle = G(x^2 + \theta \bar{\theta} / \Delta) = G(x^2) + \frac{\partial \bar{\theta}}{\Delta} G'(x^2)$$

||

$$\langle \varphi \varphi \rangle + \theta \langle \bar{x} \varphi \rangle + \bar{\theta} \langle x \varphi \rangle + \theta \bar{\theta} \langle \hat{\psi} \varphi \rangle$$

$\downarrow$   
Gd

$\downarrow$   
Gc

$$\begin{aligned} \langle \hat{\Phi}(x, \theta, \bar{\theta}) \hat{\Phi}(0) \rangle &= G(x^2 + \frac{\theta\bar{\theta}}{d}) = G(x^2) + \frac{\theta\bar{\theta}}{d} G'(x^2) \\ &= \langle \varphi(x) \varphi(0) \rangle + \theta \cancel{\langle \varphi(x) \tilde{x}(x) \rangle} + \bar{\theta} \cancel{\langle \varphi(0) \tilde{x}(x) \rangle} + \theta \bar{\theta} \langle i \hat{\varphi}(x) \varphi(0) \rangle \end{aligned}$$

$\downarrow$      $\downarrow$

$$\frac{1}{|x|^{d-4+\eta}} \longrightarrow \boxed{\eta = \bar{\eta}} \longleftarrow \frac{1}{|x|^{d-2+\eta}}$$

Dimensional Reduction:  $d' = d-2$  : integrals.

$$\int d^d x f'(x^2) \propto \int d^{d-2} x f(x^2)$$

RFIM in dim  $d$   $\sim$  IM in dim  $d-2$

D<sub>uc</sub> = 6

Evidently wrong in low  $d$ .

Failure of Sisy  $-\Delta\varphi + V(\varphi) = h(x)$

- (I) Multiple solutions  $| \text{det} | \neq \text{det}$
- (II) Discontinuous jumps :  $h \rightarrow h + \delta h$  change of dominant solution

(III) Marginality : Solutions cannot be followed  
when  $h(x) \rightarrow h(x) + \delta h(x)$

$$\frac{\delta\varphi}{\delta h} \neq (-\Delta + V'')^{-1} \rightarrow \begin{array}{l} \text{Breaking of FDT} \\ (\text{Word Id of BRST}) \end{array}$$

Avalanches Large rearrangements

$$\langle i\hat{\varphi}\psi \rangle = \langle \bar{X}X \rangle \Leftrightarrow \frac{\delta\varphi}{\delta h} = (-\Delta + V'')^{-1} \quad \begin{array}{l} \text{min. can} \\ \text{be followed} \end{array}$$

## Replicars

$$\overline{e^{-\beta \sum_{a=1}^n H(\varphi_a)}} = e^{S(\varphi_a)}$$

$$S[\varphi_a] = \int dx \sum_{a=1}^n -\varphi_a \Delta \varphi_a + V(\varphi_a) + \frac{\Delta}{2} \left( \sum_a \varphi_a \right)^2$$

$n \rightarrow \infty$

## Quadratic fluctuations

$$\langle \varphi_a^2 \rangle_c = \frac{1}{k^2 + m^2} \quad ; \quad \langle \varphi_a \varphi_b \rangle = \frac{1}{k^4 + m^4}$$

$\varphi_a$ : not well defined scaling dimension

## Cardy transformation

$$\varphi_a = \varphi + \delta_{a,1} i\hat{\psi} + \chi_a \quad ; \quad \chi_i = 0, \quad , \quad \sum_a \chi_a = 0$$

$\varphi, i\hat{\psi}, \chi_a$  have well defined Scaling dimension

$$\Delta_\varphi = D_L/2 - 2; \quad \Delta_{i\hat{\psi}} = D_L/2; \quad \Delta_\chi = D_L/2 - 1$$

keeping only the relevant terms leads .

$$S = \int dx i\hat{\psi} \left[ -\Delta \varphi + V'(\varphi) \right] + \frac{1}{2} \sum_a \chi_a (-\Delta + V'') \chi_a$$

$$\chi_a \rightarrow n-2 \text{ indep. components} \rightarrow \left( \det(-\Delta + V'') \right)^{(n-2)/2}$$

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## The cusp

Avalanches lead to singular contributions to correlation functions

Copies with different external fields  $h$ ,  $h + \delta h$        $\delta h \ll 1$

- Probability of meeting an avalanche of size  $S$

$$\rho(S) dh$$

$$\frac{(m(h) - m(h + dh))^c}{(m(h) - m(h + dh))^{c'}} = 18h \int dS \rho(S) \frac{S^c}{L^{2d}} + \text{reg.}$$

  
cusp

- Fractal dimension of  
avalanches

Tarjus-Tressler  
2012

$$S \sim \ell^{d_f}$$

If  $d_f < \frac{1}{2}(d+4 + \bar{\eta})$  dim. reduction

$\rightarrow d_f = \frac{1}{2}(d+4 + \bar{\eta})$  Susy Breaking

In  $d > 6$   $d_f = 4$   $4 < \frac{d+4}{2} \Rightarrow$  Room for  
 $\bar{\eta} = 0$  Susy below  $d = 6$

## The Cusp

Failure of FDT and avalanches

$G_d(x) = \overline{\langle \varphi(x) \varphi(0) \rangle} \rightarrow$  Generating function of disconnected correlations

$$e^{W[h]} = \int d\varphi e^{-S(\varphi) + \int \varphi h d\chi}$$
$$\log e^{\frac{W[h+\delta h_1] + W[h+\delta h_2]}{W[h]}} = W_2[\delta h_1, \delta h_2]$$

$$\frac{\partial^2}{\partial h_1 \partial h_2} \Big|_{h_1=h_2=0} \frac{W_2}{\delta h_1(x) \delta h_2(x)} \Big|_{\delta h_1^2 \delta h_2 = 0} = \underbrace{\langle \varphi(x) \rangle \langle \varphi(0) \rangle}_{\text{singular if avalanches}}$$

## Non perturbative RG

Uncontrolled closure of exact RG eqns

$$R_k(x) \quad \text{IR regulator} \quad \Delta S_k[\phi] = \int dx \frac{1}{2} \phi(x) R_k(x-y) \phi(y)$$

$$R_0(x) = 0 \quad R_\Lambda(x) = \delta(x-y) \Lambda^2$$

$$e^{W[h]} = \int d\phi e^{-S - \Delta S + (h, \phi)}$$

$$\Gamma(\phi) = -W[h] + h\phi - \Delta S(\phi)$$

$$\frac{\partial \Gamma}{\partial \log k} = \frac{1}{2} \text{Tr} \left\{ \frac{\partial \log k}{\partial \log k} R \left[ \Gamma_k^{(0)}(\phi) + R \right]^{-1} \right\}$$

$$\frac{\partial \Gamma}{\partial \log k} = \frac{1}{2} \text{Tr} \left\{ \frac{\partial \log k}{R} \left[ \Gamma_k^{(r)}(\phi) + R \right]^{-1} \right\} \quad \star \star$$

$$\Gamma^{(r)} = \frac{\sum \Gamma}{\delta \phi \delta \phi}$$

- Assume a given form for  $\Gamma$
- Use the eq. under the restrict form

e.g.

$$\Gamma(\phi) = \int d\mathbf{x} \frac{1}{2} (\nabla \phi)^2 + U(\phi)$$

Impose the validity of  $\star \star$  on  
spatially homogeneous config. of  $\phi$

## Combining

Non Perturbative + Functional RG  
replicas

Tissier-Tarjus + could find that

The RG flow gives rise to a cusp  
at a finite time below  $d_c \simeq 5.1$  ← dim.  
reduction